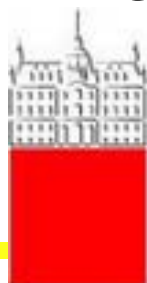


# Experiments at $e^+e^-$ flavour factories and LHCb

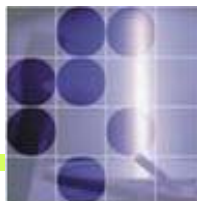
## Part 2: Belle and BaBar II

**Peter Križan**

*University of Ljubljana and J. Stefan Institute*



University  
of Ljubljana



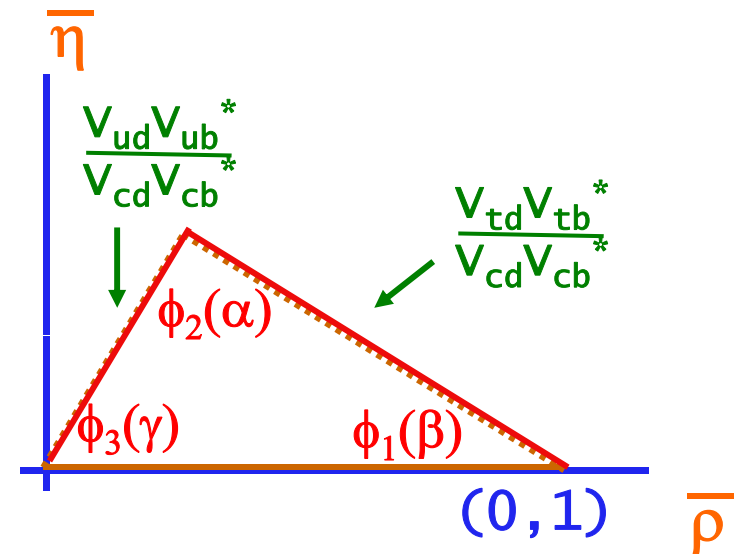
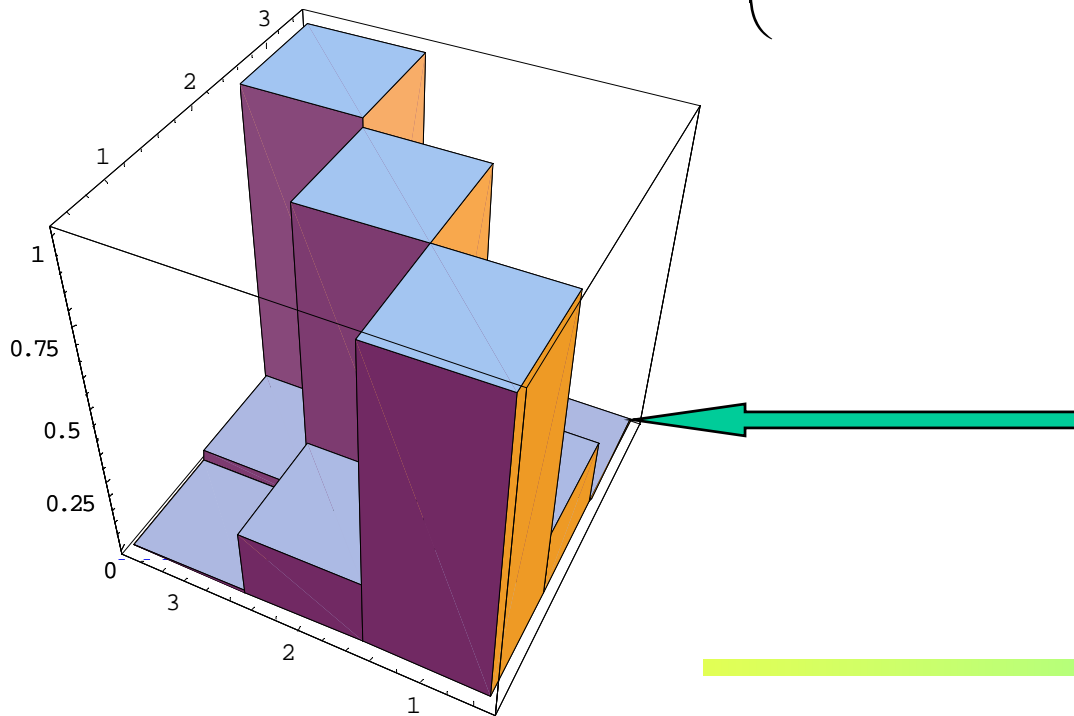
"Jožef Stefan"  
Institute

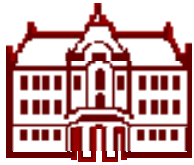




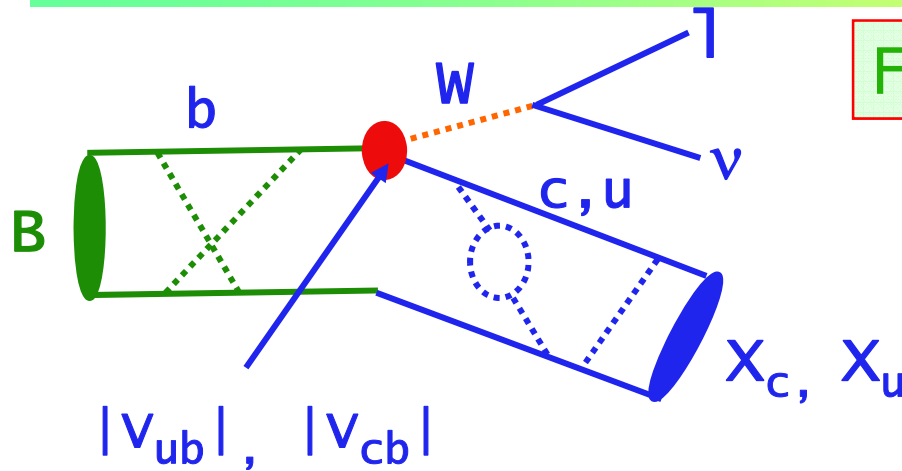
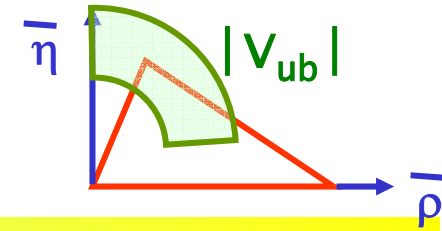
# Unitary triangle: one of the sides is determined by $V_{ub}$

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} V = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4)$$





## $|V_{ub}|$ measurements



From semileptonic B decays

$b \rightarrow c \nu$  background typically an order of magnitude larger.

Traditional inclusive method: fight the background from  $b \rightarrow c \nu$  decays by using only events with electron momentum above the  $b \rightarrow c \nu$  kinematic limit. Problem: extrapolation to the full phase space  $\rightarrow$  large theoretical uncertainty.

New method: fully reconstruct one of the B mesons, check the properties of the other (semileptonic decay, low mass of the hadronic system)

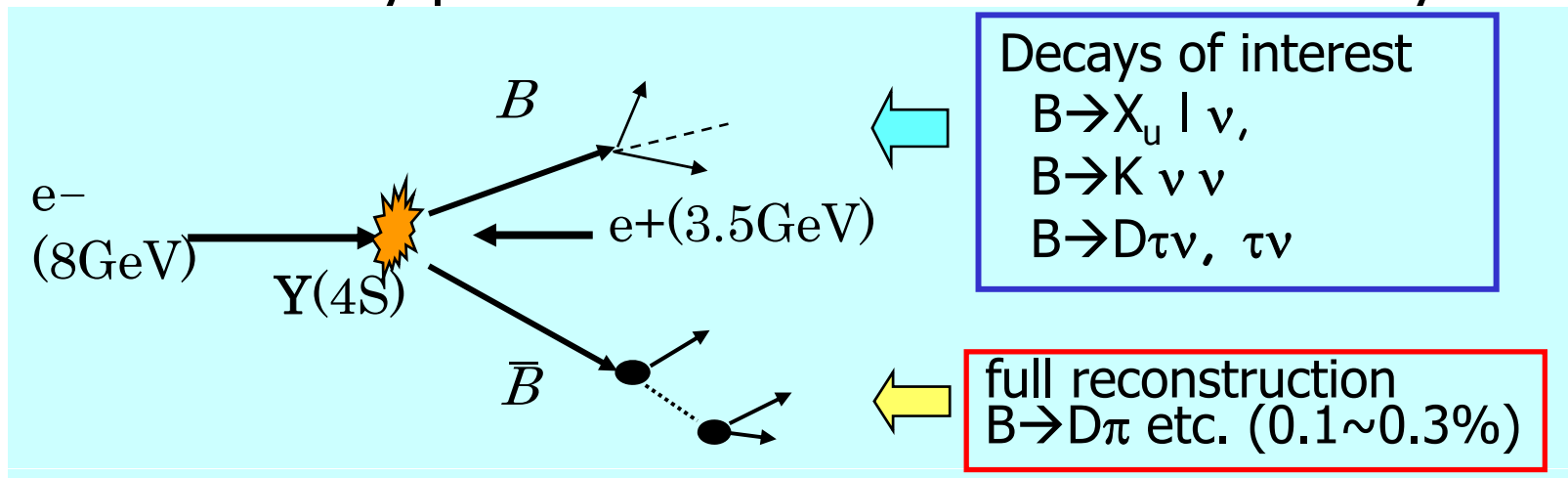
- Very good signal to noise
- Low yield (full reconstruction efficiency is 0.3-0.4%)



# Full Reconstruction Method

Fully reconstruct one of the B's to

- Tag B flavor/charge
- Determine B momentum
- Exclude decay products of one B from further analysis



→ Offline B meson beam!

Powerful tool for B decays with neutrinos



## Fully reconstructed sample

### Fully reconstructed sample

Clean environment but small sample:  $\epsilon_{\text{reco}} \approx 3 \cdot 10^{-3}$

Exclusive method: 180 decay channels

### Reconstructed channels:

$$B^0 \rightarrow D^{(*)-} \pi^+ / D^{(*)-} \rho^+ / D^{(*)-} a_1^+ / D^{(*)-} D_s^{(*)+}$$

$$B^+ \rightarrow D^{(*)0} \pi^+ / D^{(*)0} \rho^+ / D^{(*)0} a_1^+ / D^{(*)0} D_s^{(*)+}$$

$$D^{*0} \rightarrow D^0 \pi^0$$

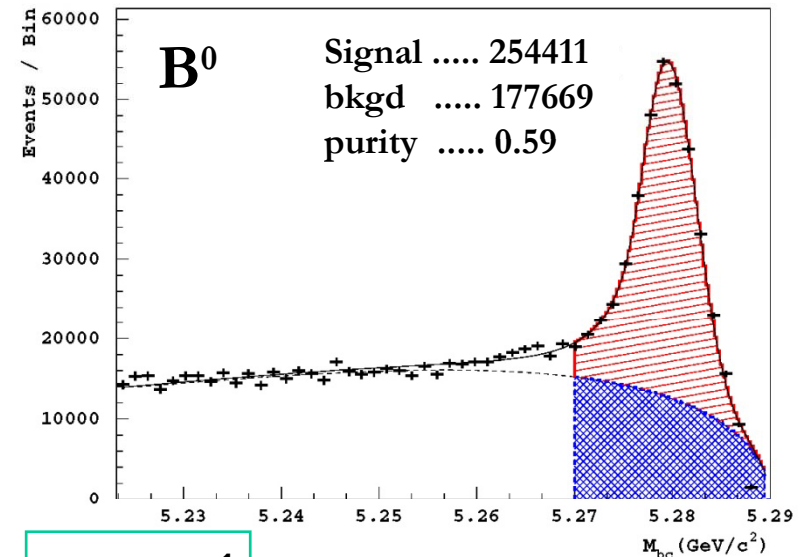
$$D^* \rightarrow D^0 \pi / D \pi^0$$

$$D_s^* \rightarrow D_s \gamma$$

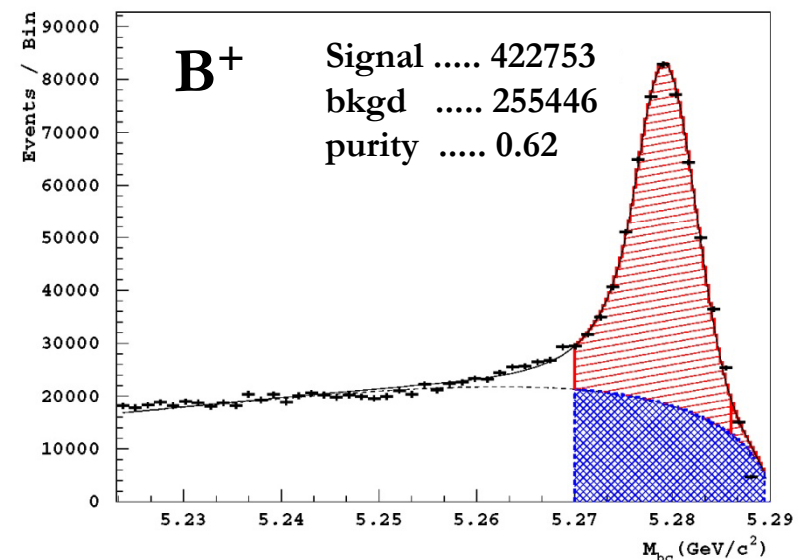
$$D^0 \rightarrow K\pi / K\pi\pi^0 / K\pi\pi\pi / K_s \pi^0 / K_s \pi\pi / K_s \pi\pi\pi^0 / KK$$

$$D \rightarrow K\pi\pi / K\pi\pi\pi^0 / K_s \pi / K_s \pi\pi^0 / K_s \pi\pi\pi / KK\pi$$

$$D_s \rightarrow K_s K\pi / KK\pi$$



253 fb<sup>-1</sup>





## $M_x$ analysis

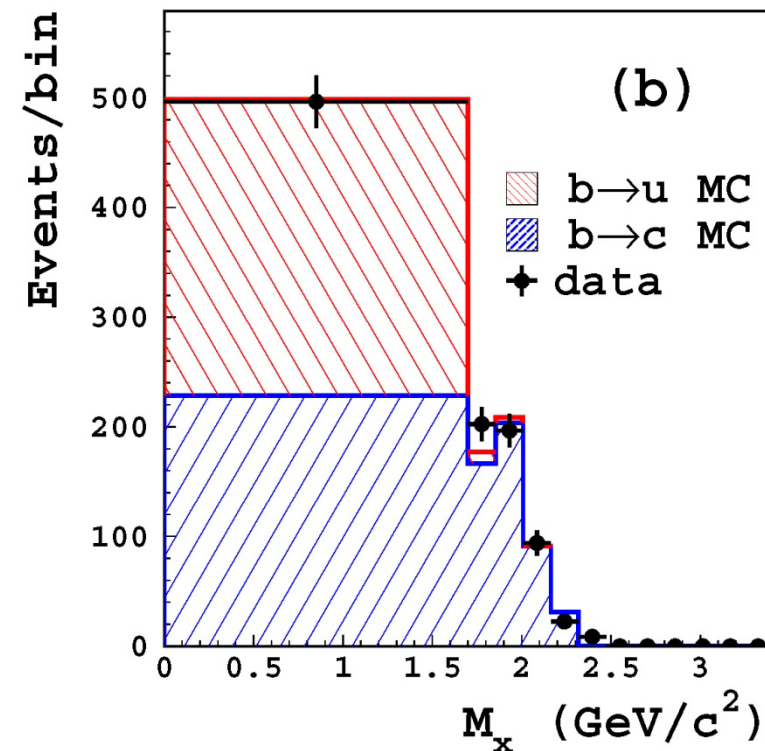
Use the mass of the hadronic system  $M_x$  as the discriminating variable against  $b \rightarrow cl\nu$

$M_x =$  mass of all hadrons from the B decay.

Expect:

- $M_x$  for  $b \rightarrow cl\nu$  to be above 1.8 GeV ( $b \rightarrow cl\nu$  results in a D meson with  $>1.8$  GeV)

- $M_x$  for  $b \rightarrow ul\nu$  to be mainly below 1.8 GeV ( $B \rightarrow \pi l\nu, \rho l\nu, \omega l\nu \dots$ )

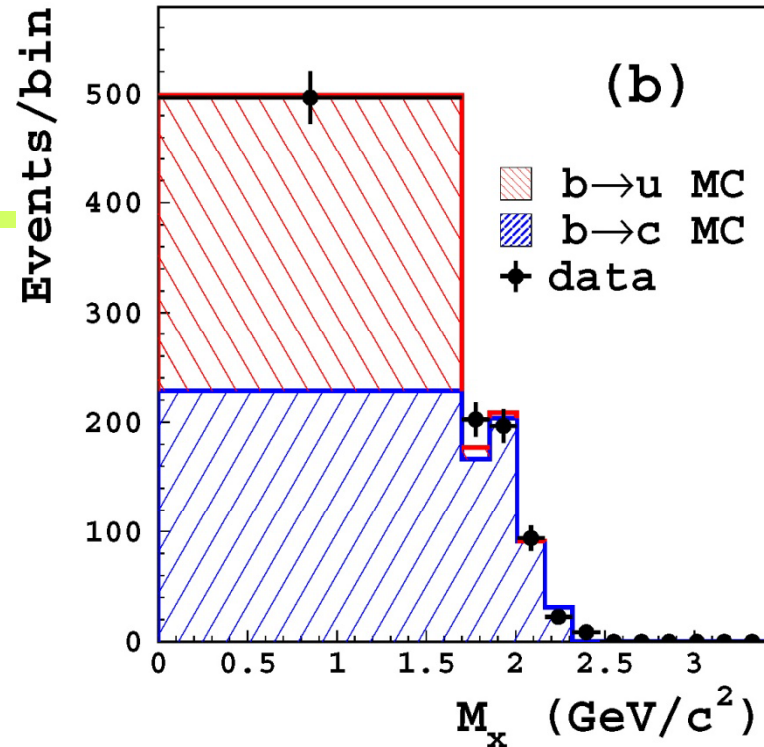




# $M_x$ analysis

$M_x < 1.7 \text{ GeV}/c^2 / q^2 > 8 \text{ GeV}^2/c^2$

Total error on  $|V_{ub}|$  ..... 12%



253 fb<sup>-1</sup>

$$|V_{ub}| = (4.93 \pm 0.25 \pm 0.22 \pm 0.15 \pm 0.13 \pm 0.46^{+0.20}_{-0.22}) \times 10^{-3}$$

stat    syst    b→u    b→c    SF    theo  
model dep.

$M_x < 1.7 \text{ GeV}/c^2 / \text{no } q^2 \text{ cut} : \text{ total error on } |V_{ub}| \text{ ..... 11\%}$

253 fb<sup>-1</sup>

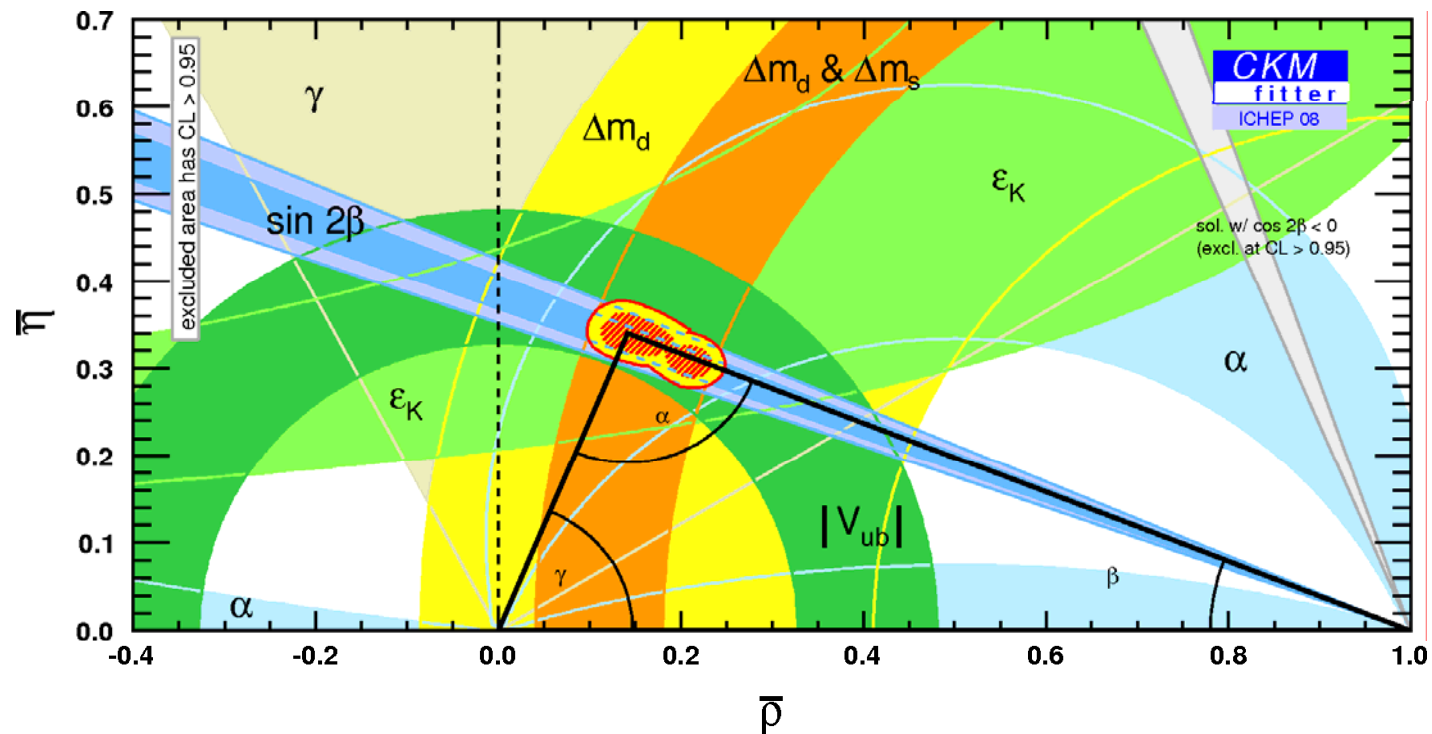
$$|V_{ub}| = (4.35 \pm 0.20 \pm 0.15 \pm 0.13 \pm 0.05 \pm 0.40^{+0.13}_{-0.14}) \times 10^{-3}$$

stat    syst    b→u    b→c    SF    theo  
model dep.



# All measurements combined...

Constraints from measurements of angles and sides of the unitarity triangle →



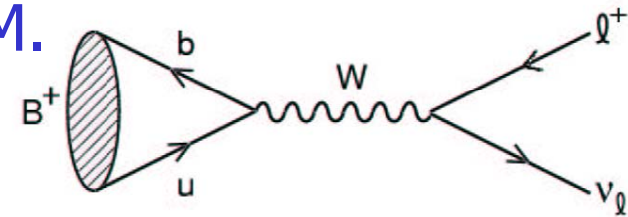
→ Remarkable agreement





## Purely leptonic decay $B \rightarrow \tau \nu$

- Challenge: B decay with at least two neutrinos
- Proceeds via W annihilation in the SM.



- Branching fraction

$$\mathcal{B}(B^- \rightarrow \ell^- \bar{\nu}) = \frac{G_F^2 m_B m_\ell^2}{8\pi} \left(1 - \frac{m_\ell^2}{m_B^2}\right)^2 f_B^2 |V_{ub}|^2 \tau_B$$

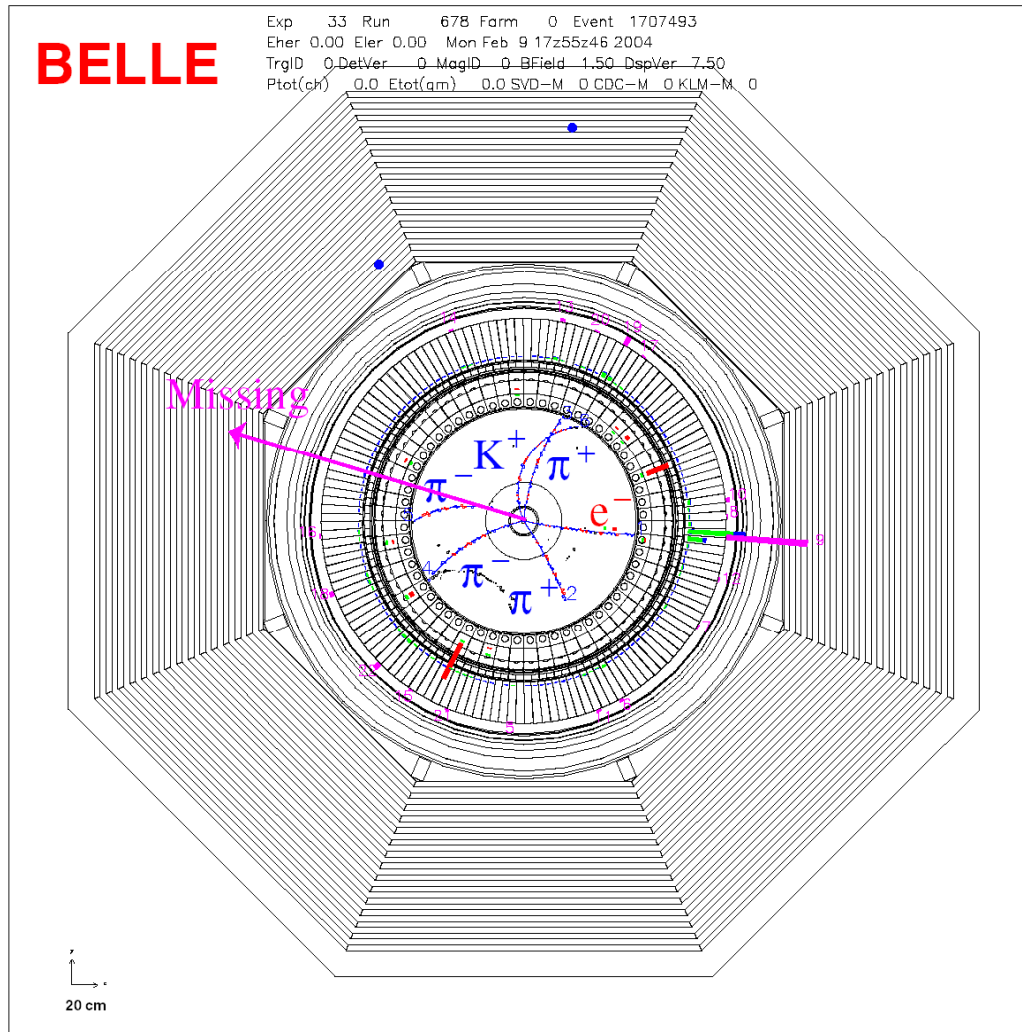
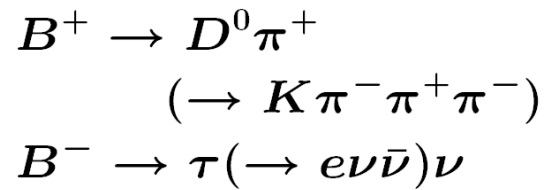
- Provide information of  $f_B |V_{ub}|$ 
  - $|V_{ub}|$  from  $B \rightarrow X_u \ell \nu \rightarrow f_B \iff$  cf) Lattice
  - $\text{Br}(B \rightarrow \tau \nu) / \Delta m_d \rightarrow |V_{ub}| / |V_{td}|$

- Limits on charged Higgs



# Event candidate $B^- \rightarrow \tau^- \nu_\tau$

Again: fully reconstruct  
one of the B's





# $B \rightarrow \tau \nu$

## $\tau$ decay modes

$$\tau^- \rightarrow \mu^- \nu \bar{\nu}, e^- \nu \bar{\nu}$$

$$\tau^- \rightarrow \pi^- \nu, \pi^- \pi^0 \nu, \pi^- \pi^+ \pi^- \nu$$

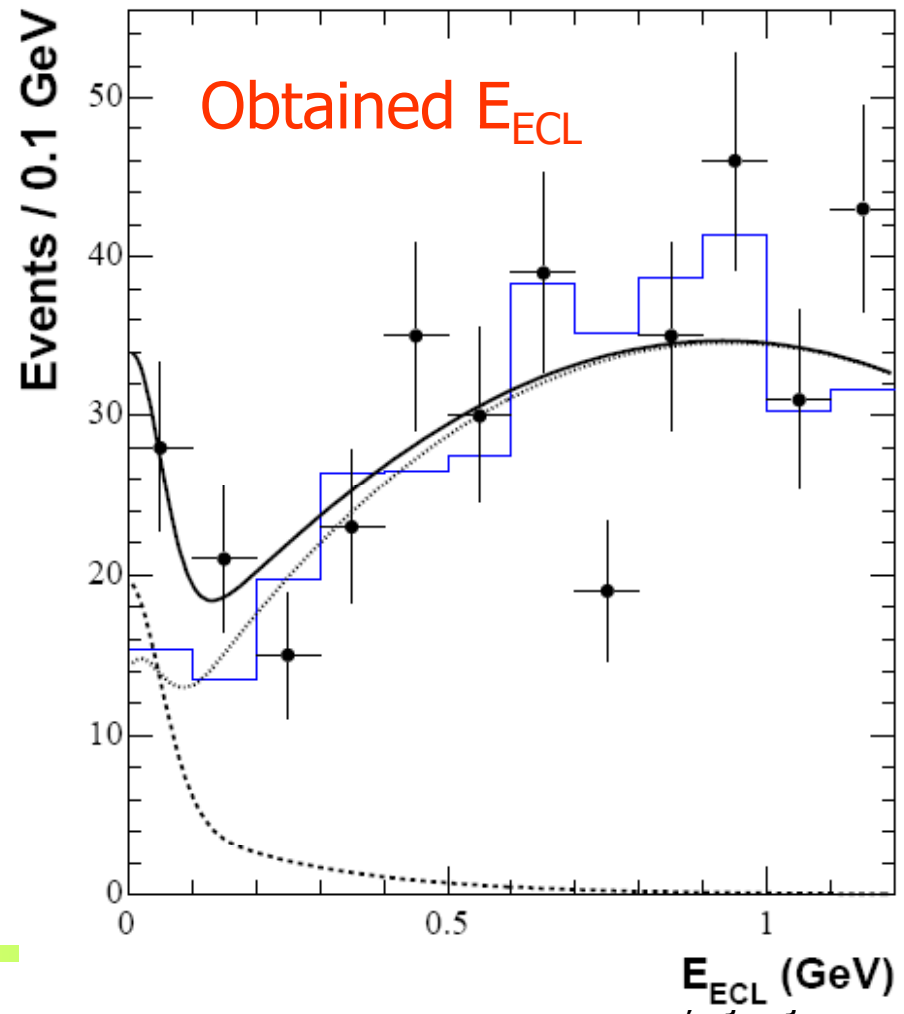
- Cover 81% of  $\tau$  decays
- Efficiency 15.8%

## Event selection

- Main discriminant: extra neutral ECL energy

Fit to  $E_{\text{residual}} \rightarrow 17.2^{+5.3}_{-4.7}$   
signal events.

$\rightarrow 3.5\sigma$  significance  
including systematics





$$B \rightarrow \tau \nu_\tau$$



$$\text{BF}(B^+ \rightarrow \tau^+ \nu_\tau) = (1.79_{-0.49-0.51}^{+0.56+0.46}) \times 10^{-4}$$

$$\Gamma^{SM}(B^+ \rightarrow \ell^+ \nu) = \frac{G_F^2}{8\pi} |V_{ub}|^2 f_B^2 m_B m_\ell^2 \left(1 - \frac{m_\ell^2}{m_B^2}\right)$$

→ Product of B meson decay constant  $f_B$  and CKM matrix element  $|V_{ub}|$

$$f_B \times V_{ub} = (10.1_{-1.4-1.4}^{+1.6+1.3}) \times 10^{-4} \text{ GeV}$$

Using  $|V_{ub}| = (4.39 \pm 0.33) \times 10^{-3}$  from HFAG

$$f_B = 229_{-31-37}^{+36+34} \text{ MeV}$$

$$\begin{array}{cc} \uparrow & \uparrow \\ 15\% & 15\% = 13\%(\text{exp.}) + 8\%(V_{ub}) \end{array}$$

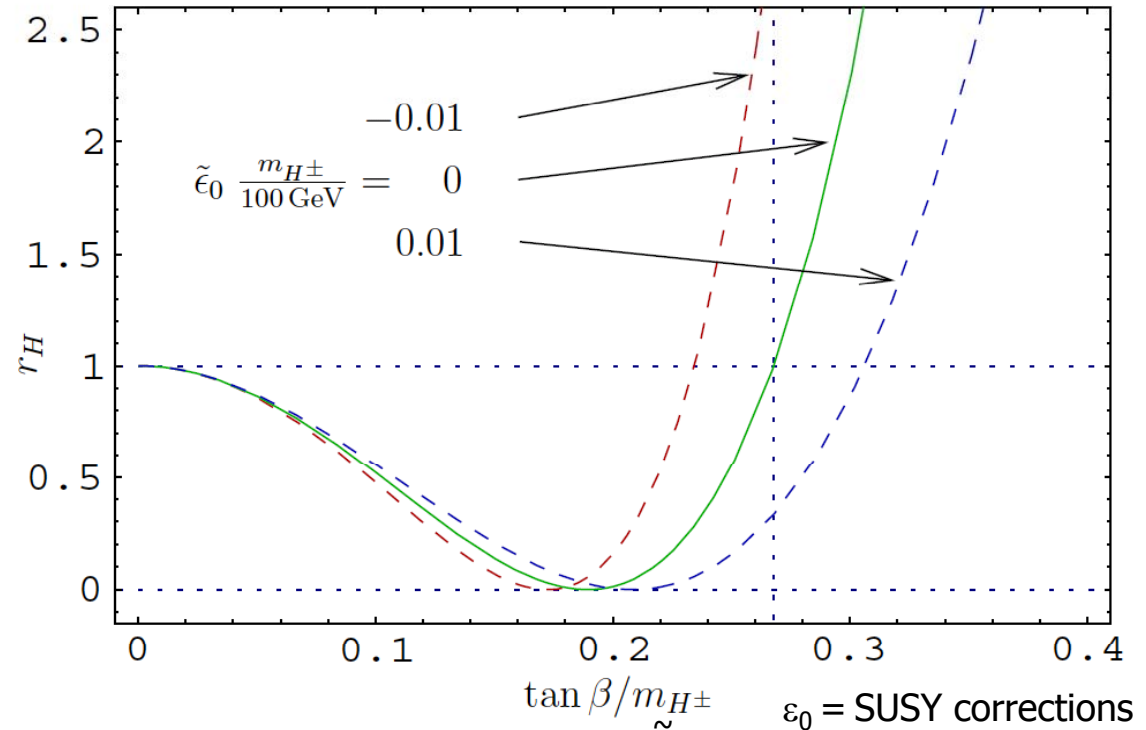
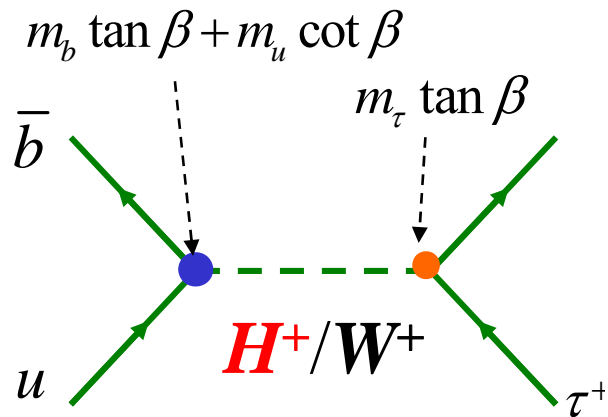
First measurement of  $f_B$ !

$f_B = (216 \pm 22) \text{ MeV}$  from unquenched lattice calculation

[HPQCD, Phys. Rev. Lett. 95, 212001 (2005) ]



# Charged Higgs contribution to $B \rightarrow \tau \nu$



$$\mathcal{B}(B \rightarrow \tau \nu) = \mathcal{B}(B \rightarrow \tau \nu)_{\text{SM}} \times r_H,$$

$$r_H = \left( 1 - \frac{m_B^2}{m_H^2} \tan^2 \beta \right)^2$$

The interference is destructive in 2HDM (type II).  $\mathcal{B} > \mathcal{B}_{\text{SM}}$  implies that  $H^+$  contribution dominates

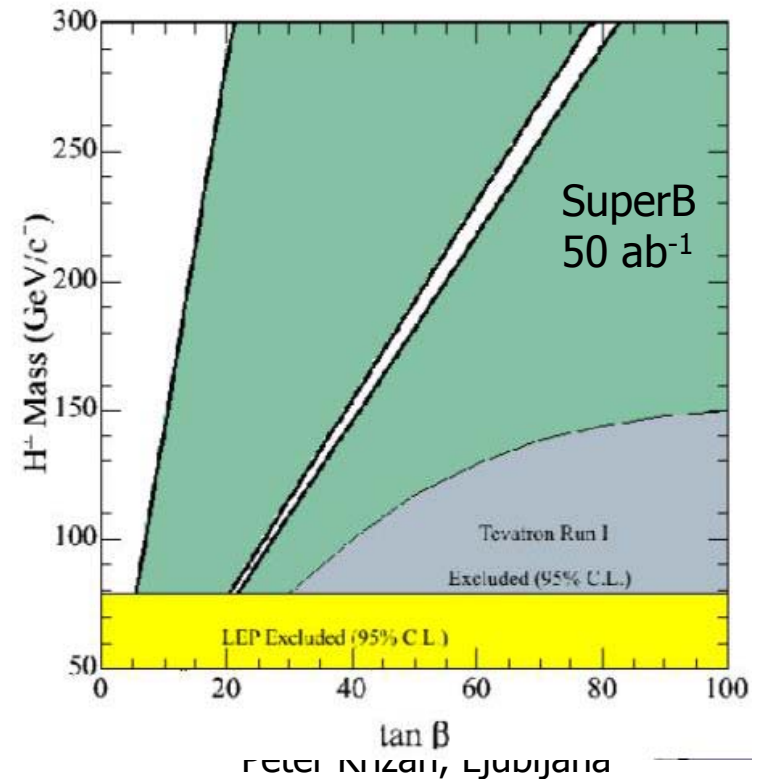
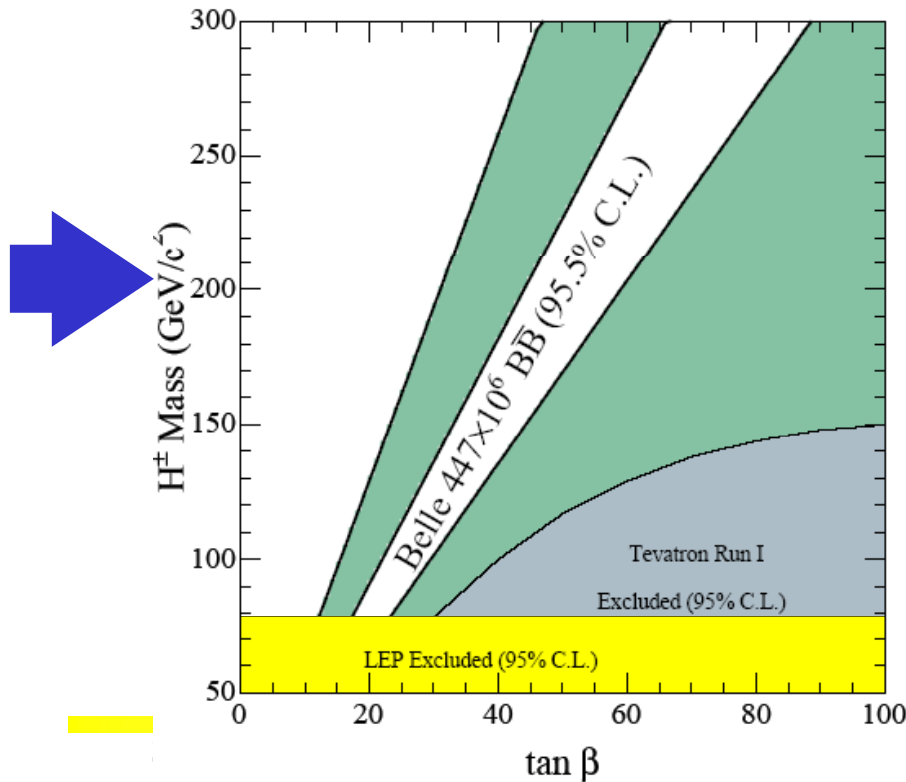
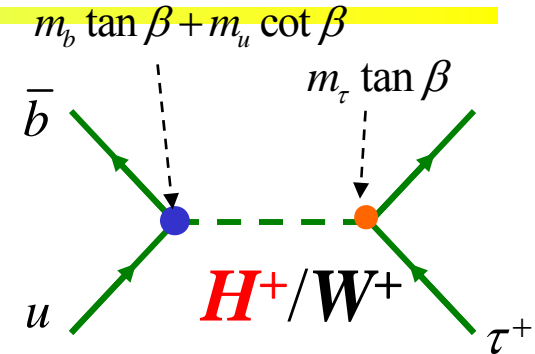
Phys. Rev. D **48**, 2342 (1993)



# Charged Higgs limits from $B^- \rightarrow \tau^- \nu_\tau$

If the theoretical prediction is taken for  $f_B$   
 $\rightarrow$  limit on charged Higgs mass vs.  $\tan\beta$

$$r_H = \frac{BF(B \rightarrow \tau\nu)}{BF(B \rightarrow \tau\nu)_{SM}} = \left(1 - \frac{m_B^2}{m_H^2} \tan^2 \beta\right)^2$$





# New Belle result on $B^+ \rightarrow \tau^+ \nu$

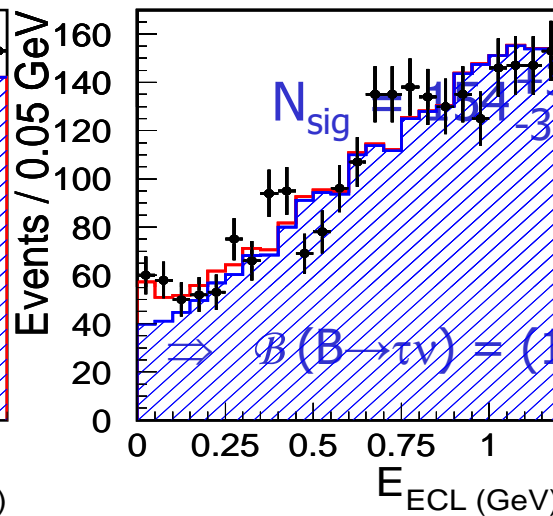
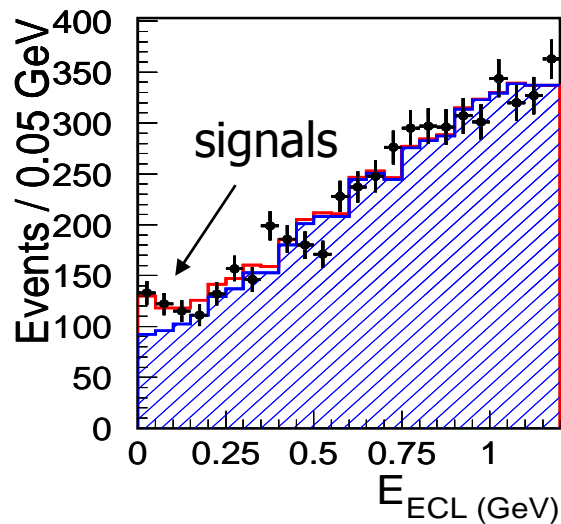
Method: Tag B on one side with the semileptonic decay  $B \rightarrow D^{(*)} l \nu$

→ Neutrino not reconstructed in the tagging B decay sequence → more background than in fully reconstructed hadronic decays

Again look for  $\tau$  signature with “extra” energy in the ECAL

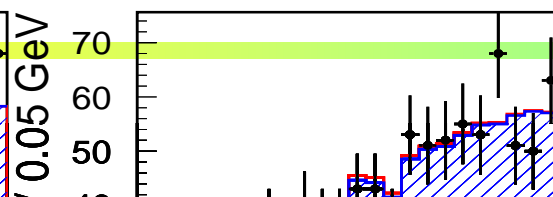
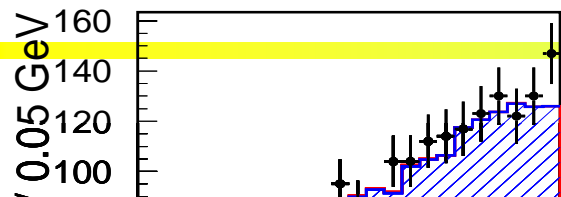


657 M  $B\bar{B}$  with  $D^{(*)} l \nu$  tag



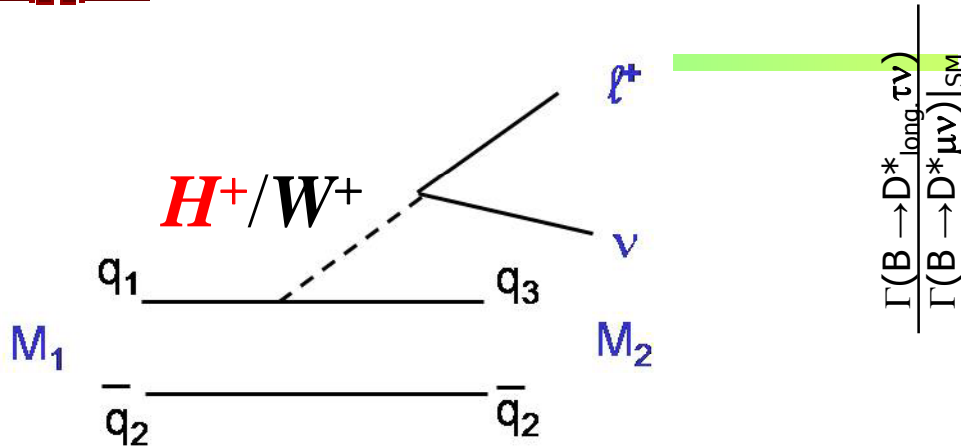
$N_{sig} = 136 \text{ (stat)} \pm 20 \text{ (syst)}$

$$\Rightarrow \mathcal{B}(B \rightarrow \tau \nu) = (1.65^{+0.38+0.35}_{-0.37-0.37}) \cdot 10^{-4}$$



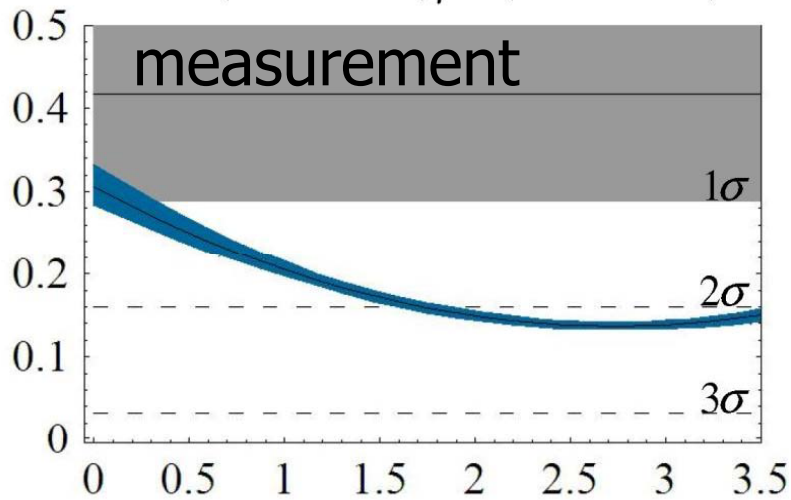


# Charged Higgs limits from $B^- \rightarrow D^{(*)} \tau^- \nu_\tau$

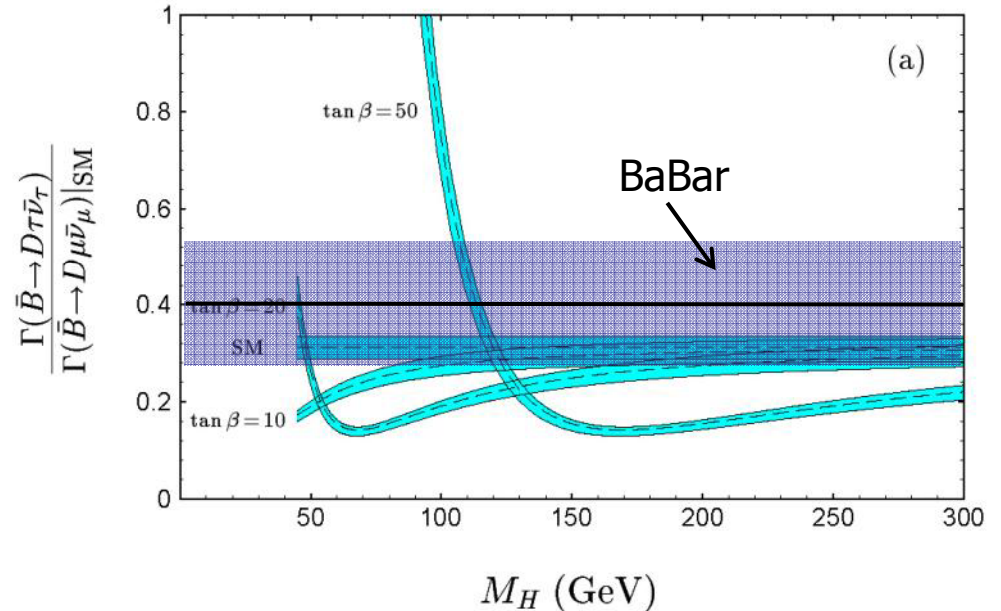
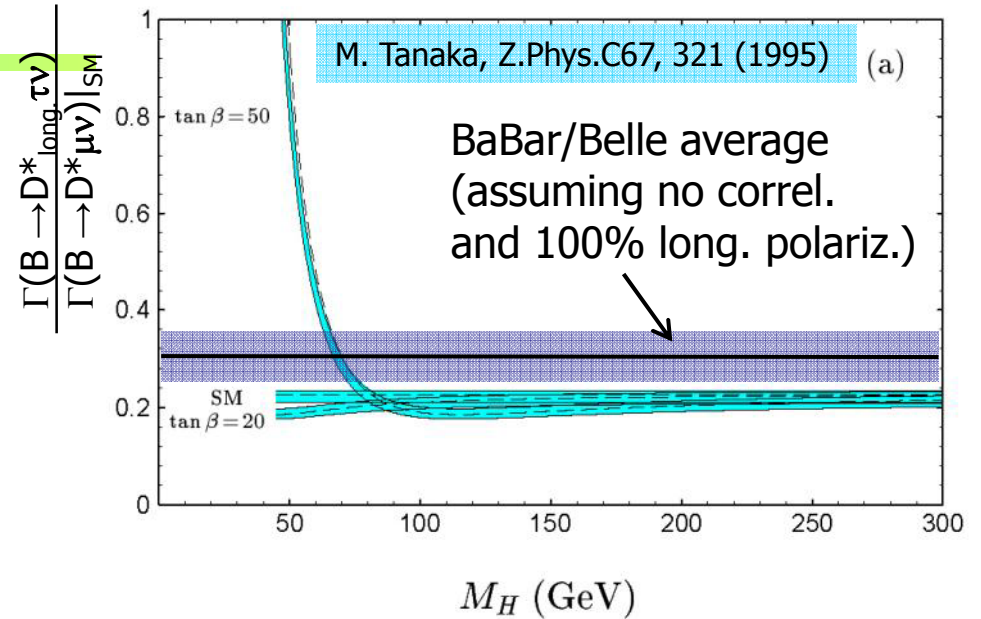


U. Nierste et al., PRD78, 015006 (2008)

$$R \equiv \mathcal{B}(B \rightarrow D\tau\nu) / \mathcal{B}(B \rightarrow D\ell\nu)$$



$m_B^2/m_H^2 \tan^2\beta$  (in 2HDM-II)

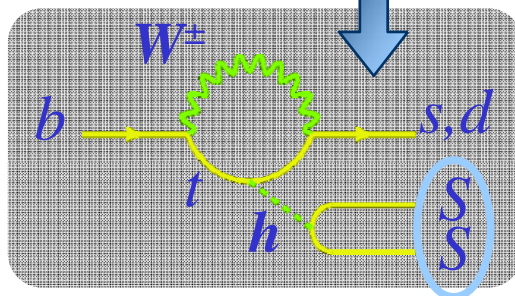
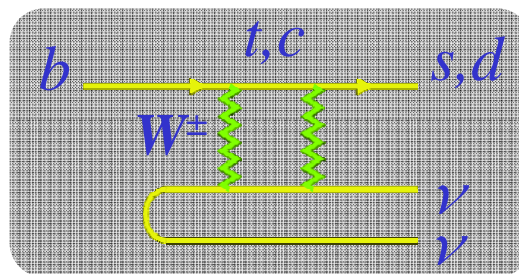
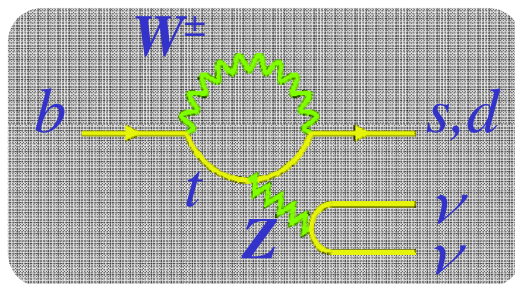






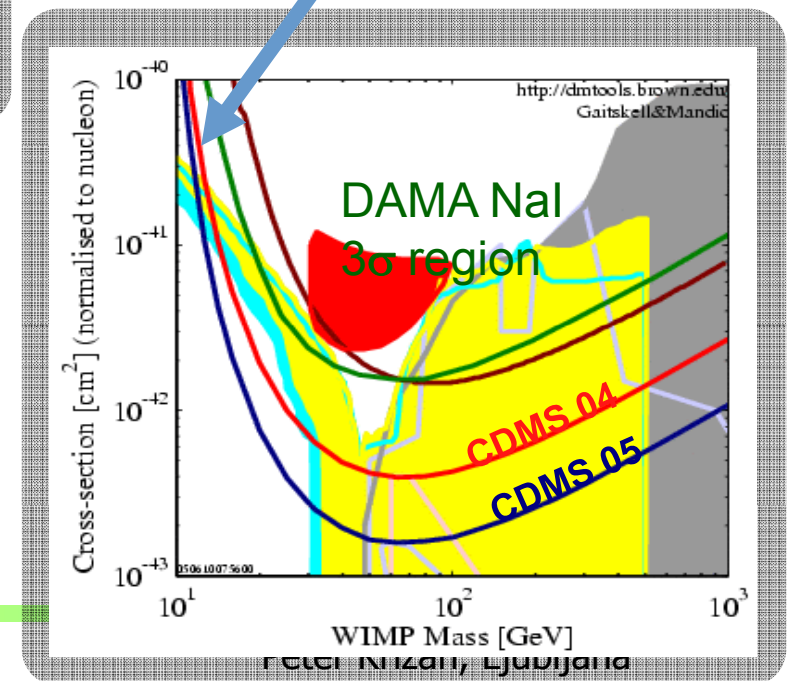
$$B \rightarrow K^{(*)} \nu \nu$$

- Proceed through electroweak penguin + box diagram.
- Sensitive to **New Physics in the loop diagram.**
- Theoretically clean: no long distance contributions.
- May be sensitive to **light dark matter** (C. Bird, PRL 93, 201803 (2004))



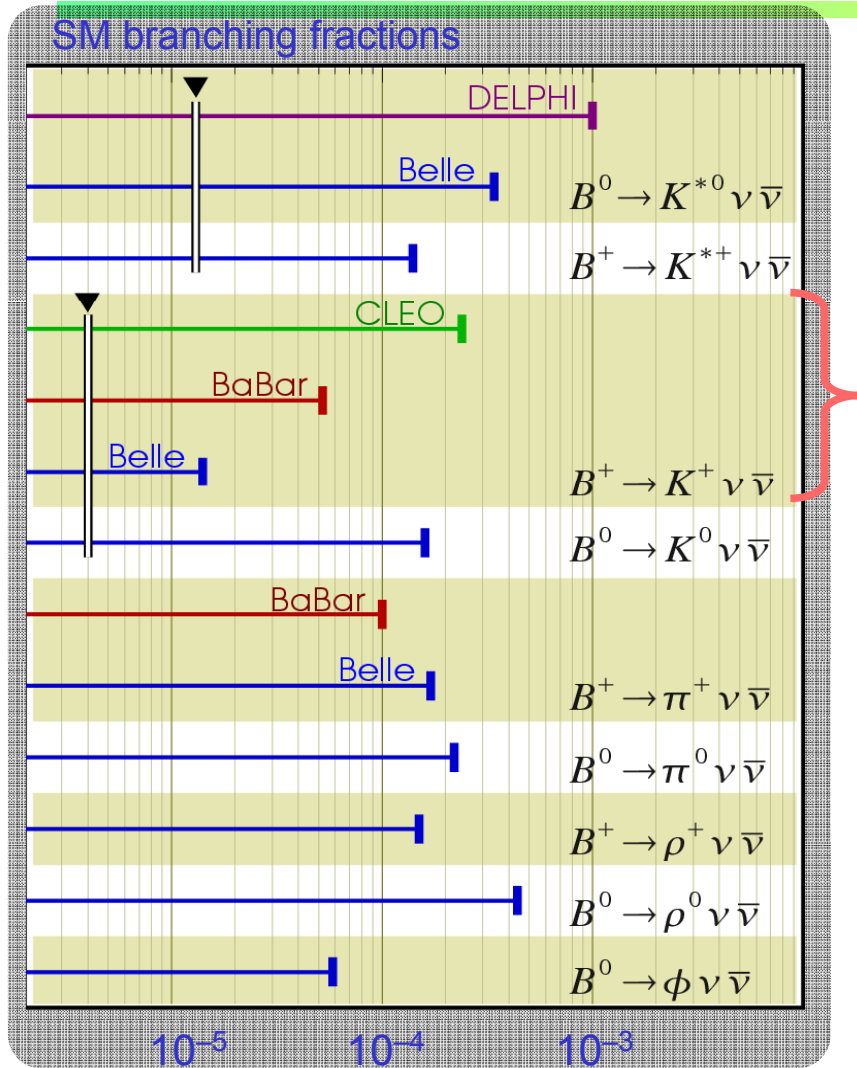
$b \rightarrow s + \text{Missing } E$   
may be enhanced by  
this extra diagram.

No sensitivity to light  
dark matter ( $M < 10$  GeV)  
in direct searches

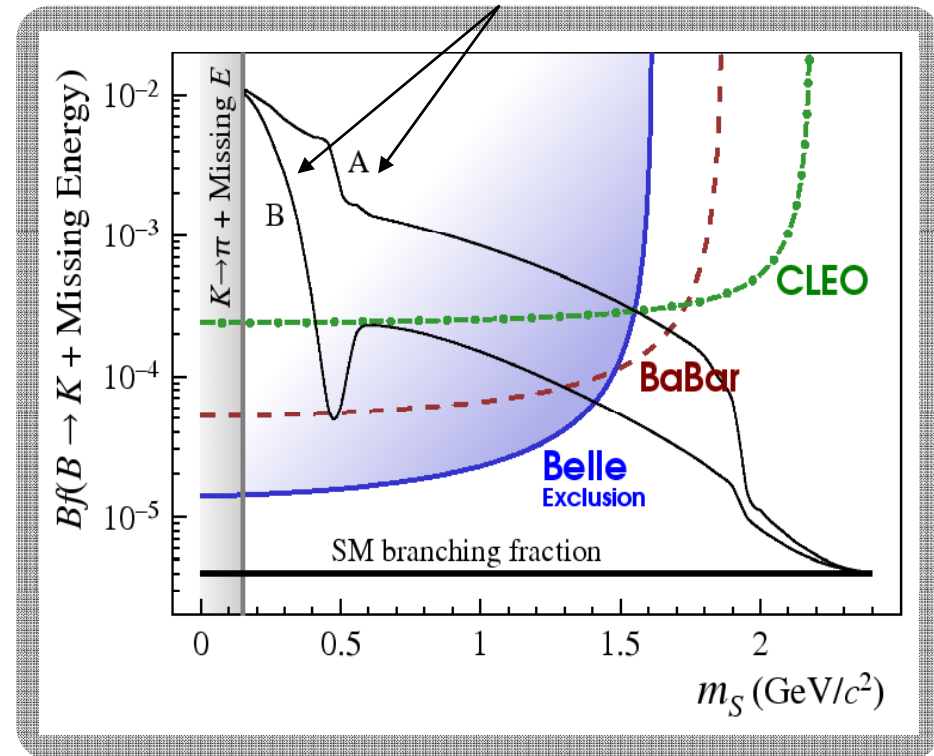




# $B \rightarrow K^{(*)} \nu \bar{\nu}$ : present limits



- Limit on light dark matter based on the  $K^+ \nu \bar{\nu}$  limits (using theory predictions, C. Bird, PRL 93, 201803 (2004))

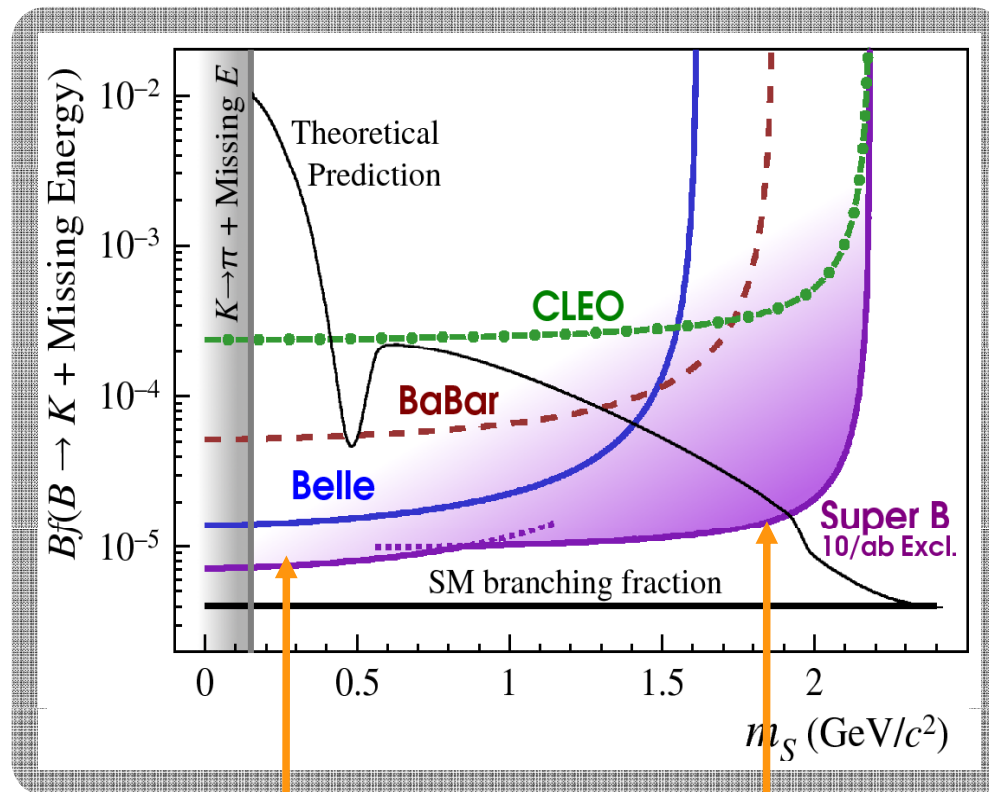


- Limit depends on  $P^*(K)$  momentum cut



# $B \rightarrow K^{(*)} \nu \nu$ : prospects for 10/ab

Assuming no changes in the analysis & detector:



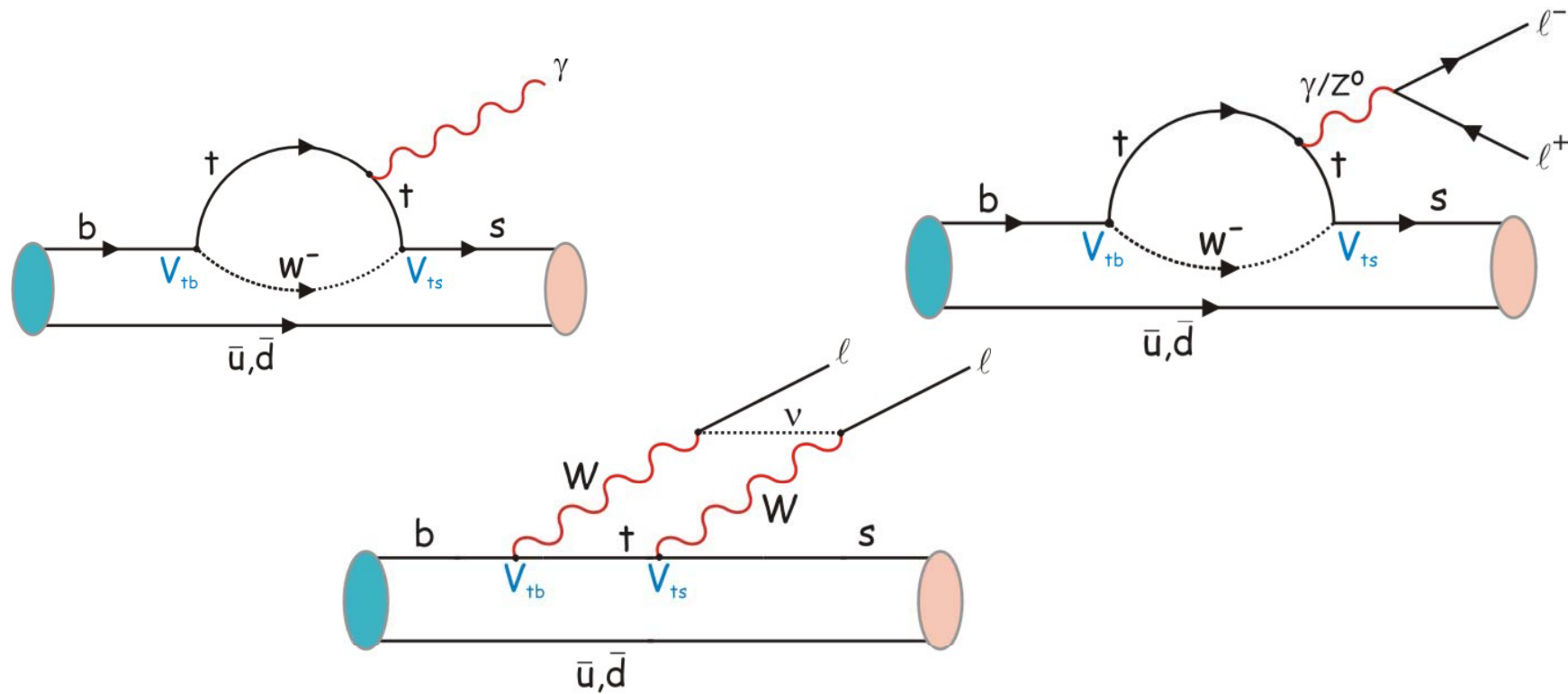
with the same  $P^*(K)$   
threshold (1.6 GeV)

with a lower  $P^*(K)$   
threshold (0.7 GeV)



# Why FCNC decays?

Flavour changing neutral current (FCNC) processes (like  $b \rightarrow s$ ,  $b \rightarrow d$ ) are forbidden at the tree level in the Standard Model. Proceed only at low rate via higher-order loop diagrams. Ideal place to search for new physics.





# How can New Physics contribute to $b \rightarrow s$ ?

For example in the process:

$$B^0 \rightarrow \eta' K^0$$

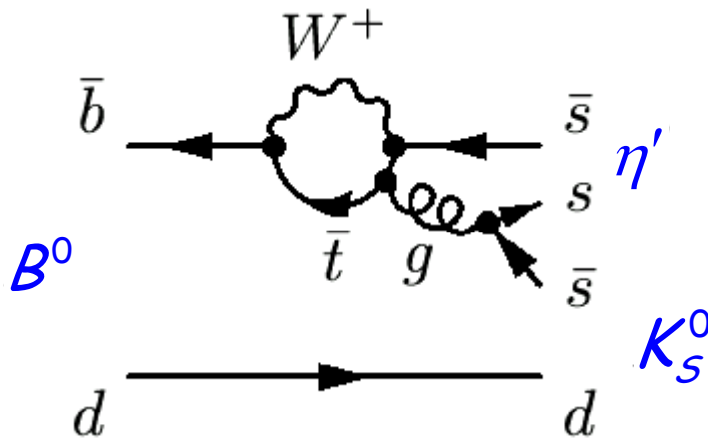
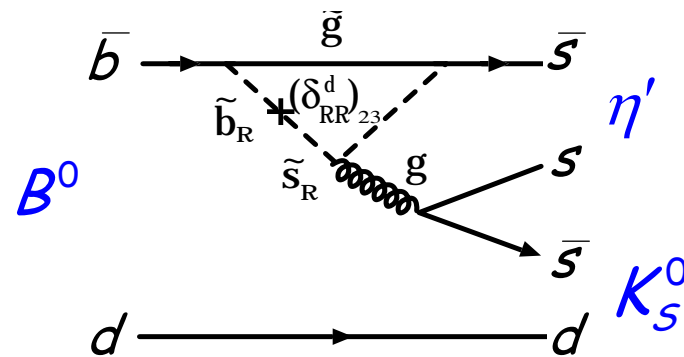


Diagram with  
supersymmetric particles

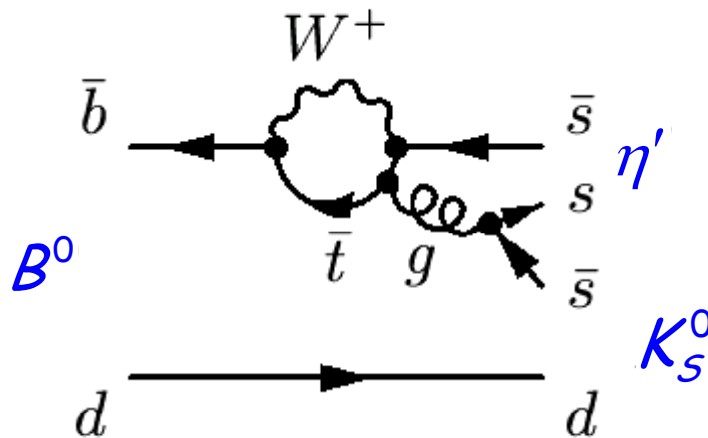
Ordinary penguin diagram with  
a t quark in the loop





# Searching for new physics phases in CP violation measurements in $b \rightarrow s$ decays

Prediction in SM:



$$a_f = -\text{Im}(\lambda_f) \sin(\Delta m t)$$

$$\text{Im}(\lambda_f) = \xi_f \sin 2\phi_1$$

The same value as in the decay  $B^0 \rightarrow J/\psi K_S^0$ !

This is only true if there are no other particles in the loop! In general the parameter can assume a different value  $\sin 2\phi_1^{\text{eff}}$

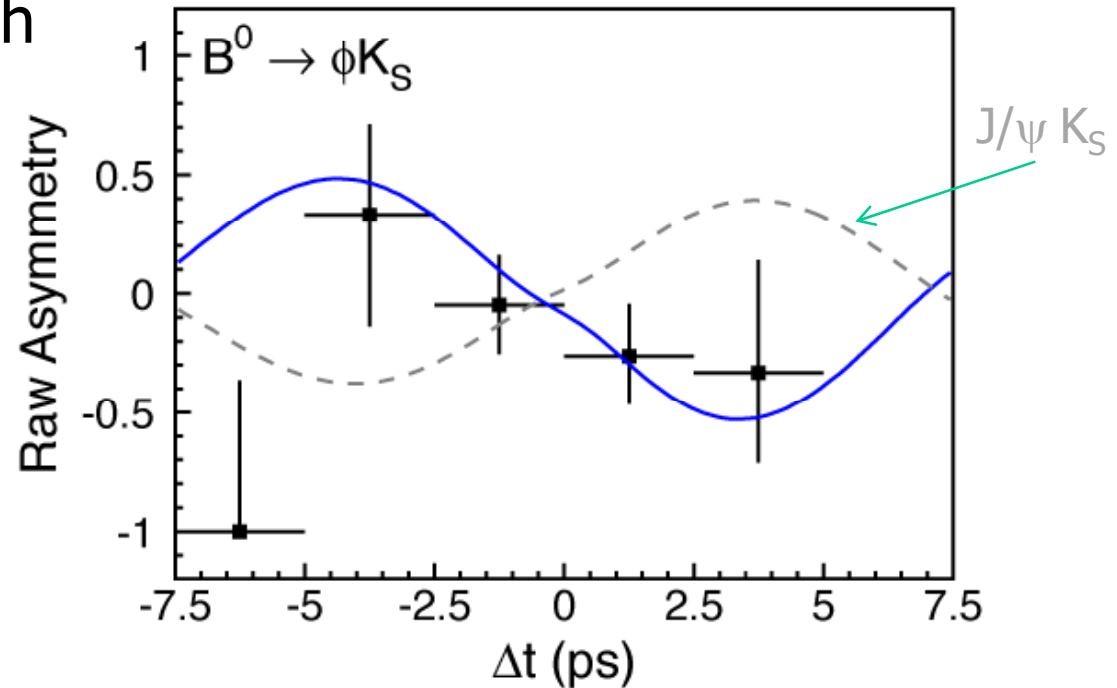


## Result of 2003 (140/fb): surprise!

Measurement: points with error bars.

Standard Model predictions: dashed

Result of the unbinned likelihood fit: blue curve



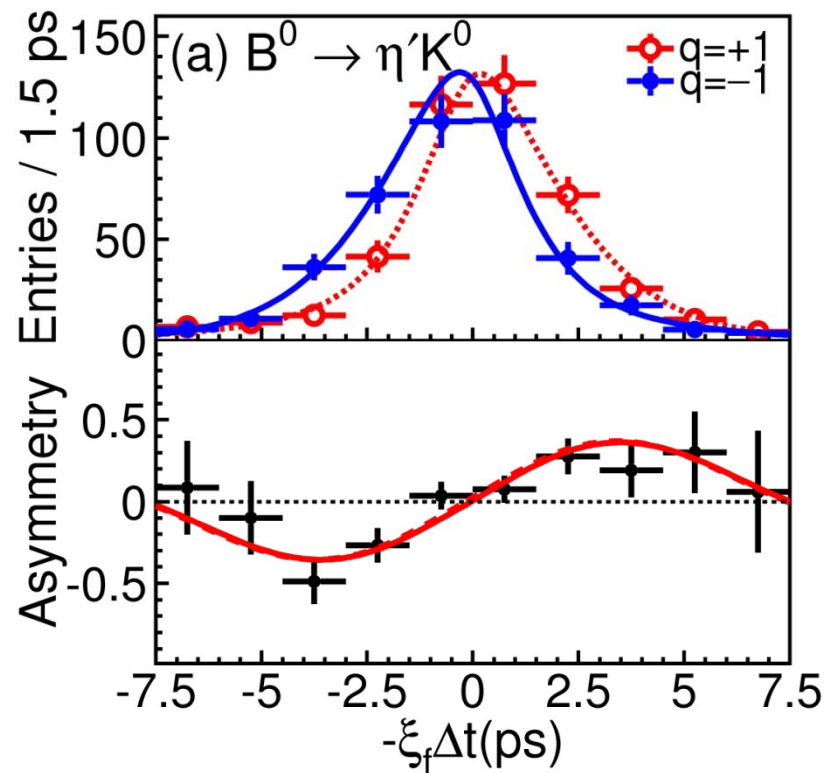
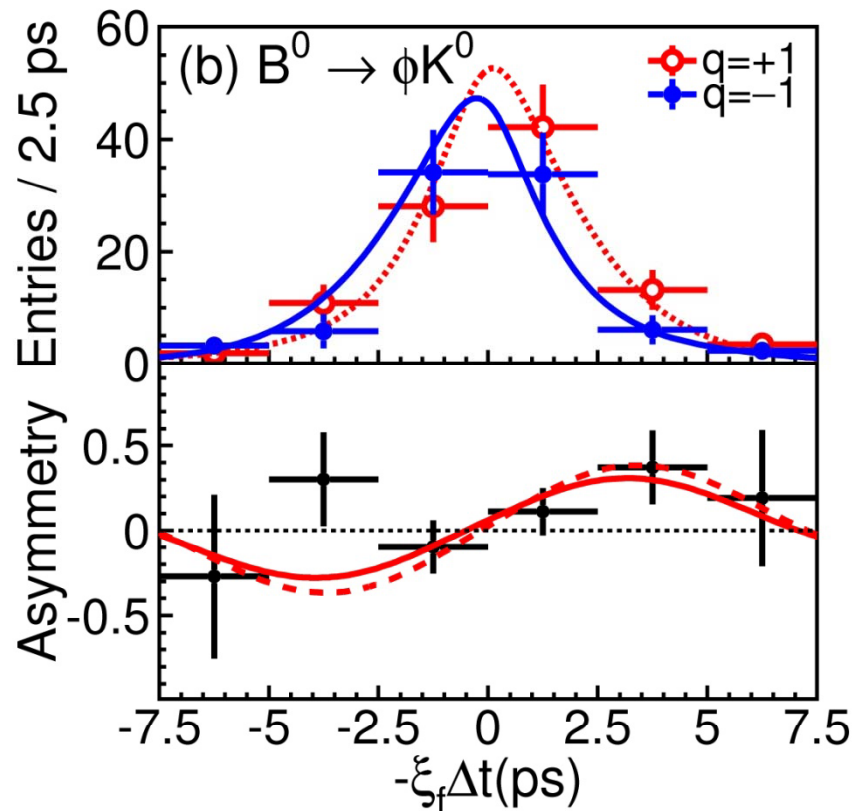
Measure:  $S = -0.96 \pm 0.50$ , expect  $S = \sin 2\phi_1 = +0.731 \pm 0.056$

not conclusive  $\rightarrow$  needed more data

... with more data ...

$B \rightarrow \phi K_S$

$B \rightarrow \eta' K_S$



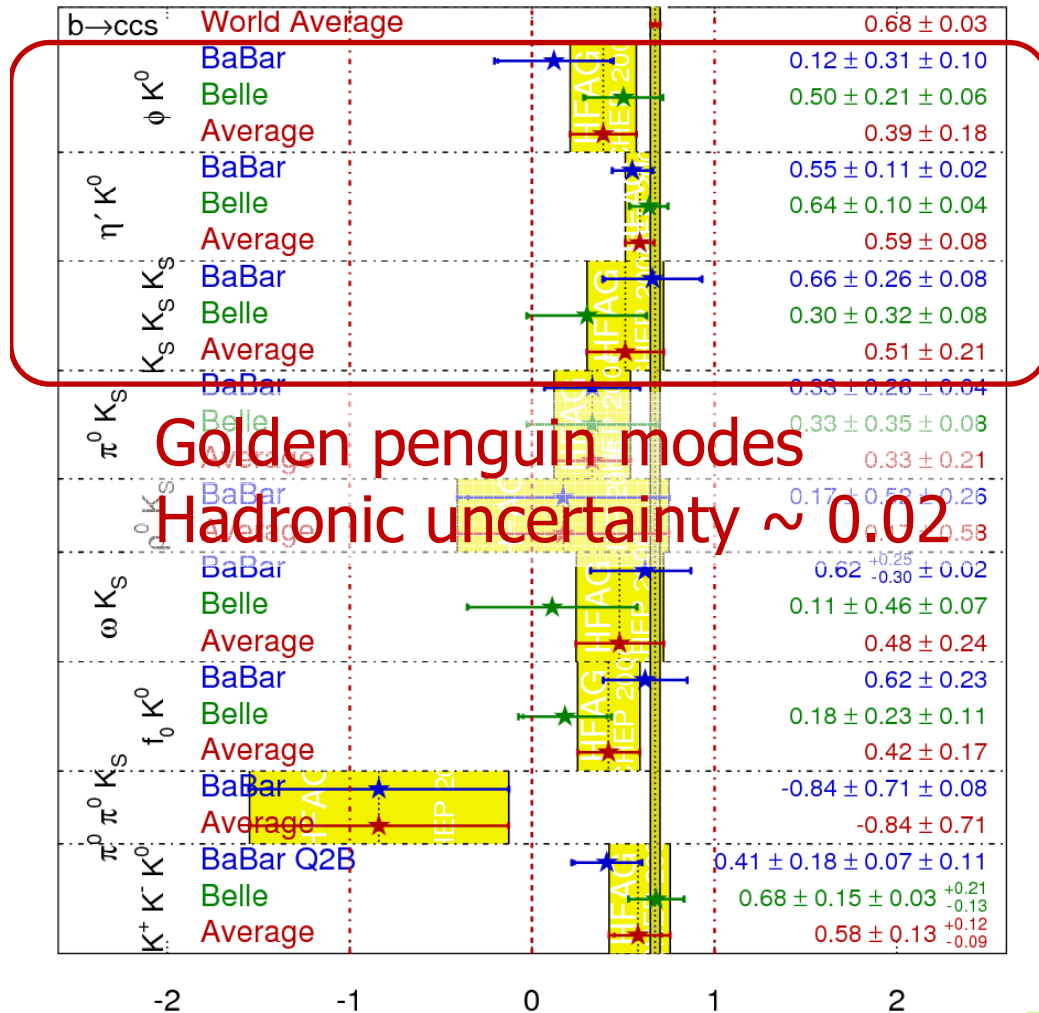




# Search for NP: $b \rightarrow sq\bar{q}$

$$\sin(2\beta^{\text{eff}}) \equiv \sin(2\phi_1^{\text{eff}})$$

**HFAG**  
ICHEP 2006  
PRELIMINARY



Golden penguin modes  
Hadronic uncertainty  $\sim 0.02$

ICHEP08

BaBar

Belle

Naïve average

$0.26 \pm 0.25 \pm 0.04$

$0.67 \pm 0.25 \pm 0.07$   
 $0.27 \pm 0.07$

$0.45 \pm 0.18$

$0.57 \pm 0.08 \pm 0.02$

$0.64 \pm 0.10 \pm 0.04$

$0.60 \pm 0.07$

$0.71 \pm 0.24 \pm 0.04$

$0.30 \pm 0.32 \pm 0.08$

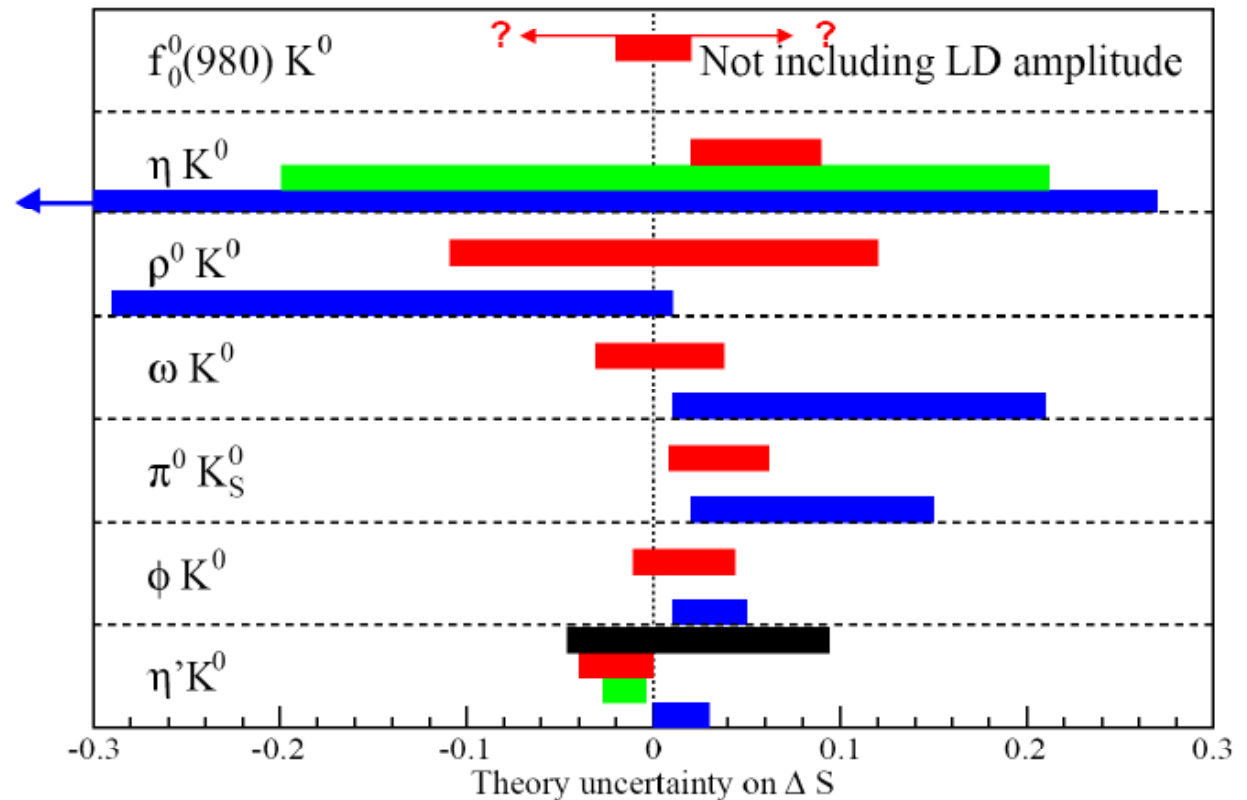
$0.57 \pm 0.20$

Need much more data  
to clarify the issue



- To find NP we need to understand the SM contributions to a process.
  - Leading order term is expected to be the same as a SM weak phase.
  - Higher order terms including re-scattering, suppressed amplitudes, final state radiation and so on can modify our expectations.

- Some channels are better understood than others.
- Sign of  $\Delta S$  correction is mode dependent.
- Most precise  $\Delta S$  correction is for  $B^0 \rightarrow \eta' K^0$ , where  $\Delta S_{\text{theory}} \sim \pm 0.01$ .
- Concentrate efforts on well understood channels.



- QCDf Beneke, PLB**620**, 143 (2005)
- SCET/QCDF, Williamson and Zupan, PRD**74**, 014003 (2006)
- QCDf Cheng, Chua and Soni, PRD**72**, 014006 (2005)
- SU(3) Gronau, Rosner and Zupan, PRD**74**, 093003 (2006)



## CP asymmetry in time integrated rates (‘direct CP’, also for charged B)

$$a_f = \frac{\Gamma(B \rightarrow f) - \Gamma(\bar{B} \rightarrow \bar{f})}{\Gamma(B \rightarrow f) + \Gamma(\bar{B} \rightarrow \bar{f})} = \frac{1 - |\bar{A}/A|^2}{1 + |\bar{A}/A|^2}$$

Need  $|\bar{A}/A| \neq 1$ : how do we get there?

In general, A is a sum of amplitudes with strong phases  $\delta_i$  and weak phases  $\phi_i$ . The amplitudes for anti-particles have the same strong phases and opposite weak phases  $\rightarrow$

$$A_f = \sum_i A_i e^{i(\delta_i + \phi_i)}$$

$$\bar{A}_{\bar{f}} = \sum_i A_i e^{i(\delta_i - \phi_i)}$$

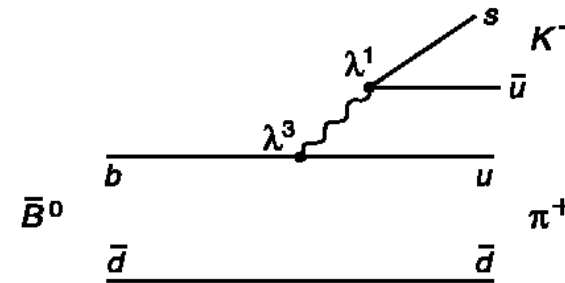
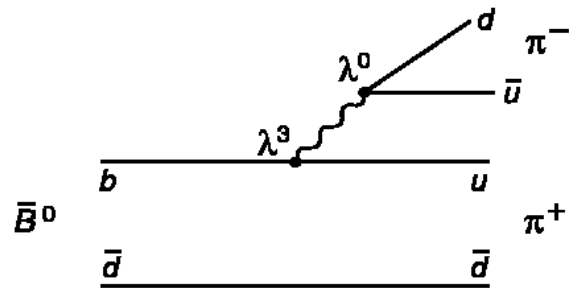
$$|A_f|^2 - |\bar{A}_{\bar{f}}|^2 = \sum_{i,j} A_i A_j \sin(\phi_i - \phi_j) \sin(\delta_i - \delta_j)$$

$\rightarrow$  Need at least two interfering amplitudes with different weak and strong phases.



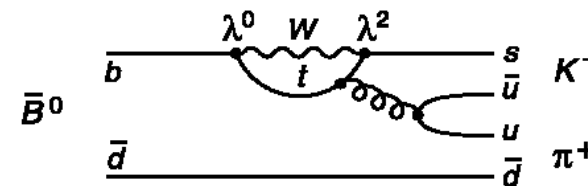
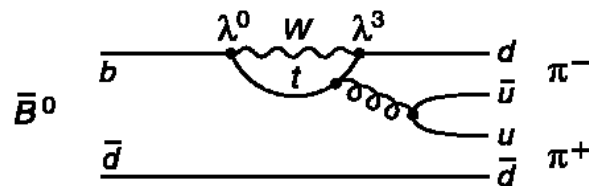
# Diagrams for $B \rightarrow \pi\pi, K\pi$ decays

$\pi\pi$



$\lambda^3 \leftarrow \lambda^4$

$K\pi$



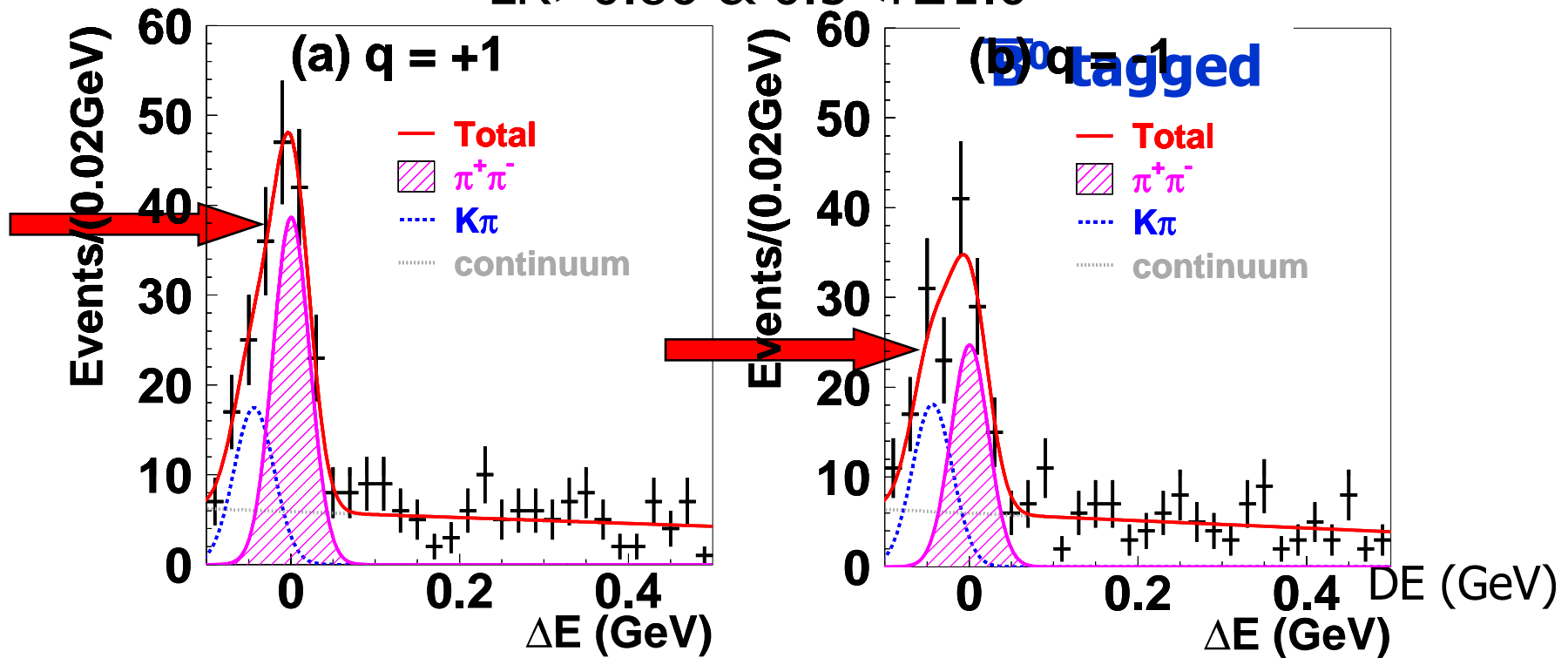
$\lambda^3 \longrightarrow \lambda^2$

- Penguin amplitudes are sizeable in both decays



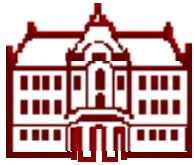
# Direct CP violation in $\pi^+\pi^-$

LR > 0.86 & 0.5 < r ≤ 1.0



Visible indication of direct CP violation.

Counting experiment consistent with the time-dependent fit (see lecture 1).



# A difference in the direct violation of CP symmetry in $B^+$ and $B^0$ decays

## CP asymmetry

$$\mathcal{A}_f = \frac{N(\bar{B} \rightarrow \bar{f}) - N(B \rightarrow f)}{N(\bar{B} \rightarrow \bar{f}) + N(B \rightarrow f)}$$

## Difference between $B^+$ and $B^0$ decays

In SM expect  $\mathcal{A}_{K^\pm \pi^\mp} \approx \mathcal{A}_{K^\pm \pi^0}$

### Measure:

$$\mathcal{A}_{K^\pm \pi^\mp} = -0.094 \pm 0.018 \pm 0.008$$

$$\mathcal{A}_{K^\pm \pi^0} = +0.07 \pm 0.03 \pm 0.01$$

$$\Delta\mathcal{A} = +0.164 \pm 0.037$$

A problem for a SM explanation  
(in particular when combined with other measurements)

A hint for new sources of CP violation?

nature

International weekly journal of science

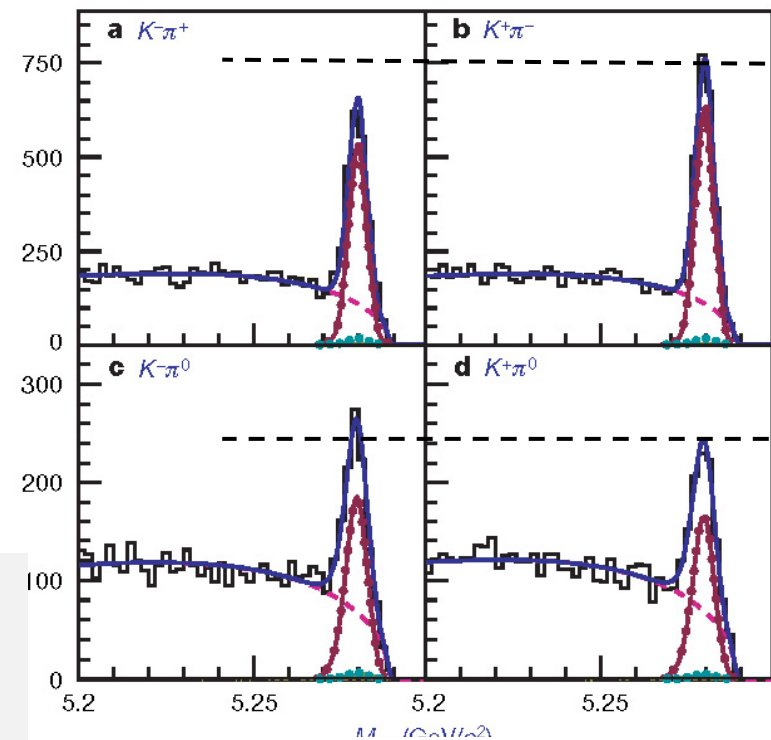
nature

Vol 452/20 March 2008 doi:10.1038/nature06827

LETTERS

**Difference in direct charge-parity violation between charged and neutral  $B$  meson decays**

The Belle Collaboration\*

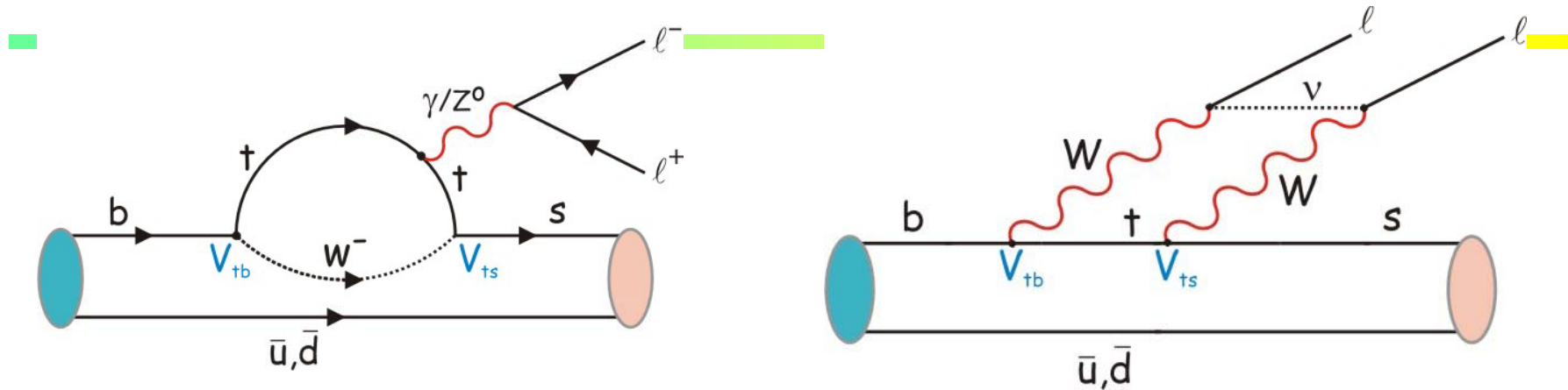


$\sim 1$  in  $10^5$   $B$  mesons decays in this decay mode

Belle, Nature 452, 332 (2008)



# Another FCNC decay: $B \rightarrow K^* l^+ l^-$



$b \rightarrow s l^+ l^-$  was first measured in  $B \rightarrow K l^+ l^-$  by Belle (2001).

Important for further searches for the physics beyond SM

Particularly sensitive: **backward-forward asymmetry in  $K^* l^+ l^-$**

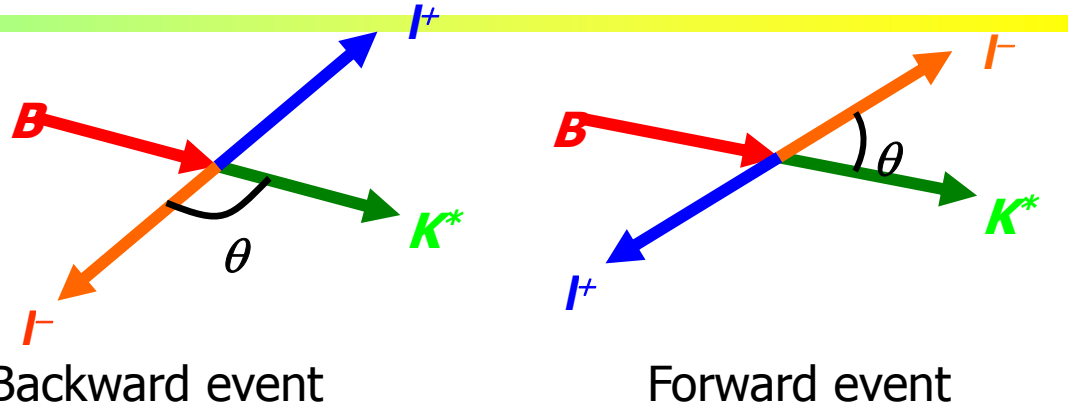
$$A_{FB} \propto \Re \left[ C_{10}^* (s C_9^{eff}(s) + r(s) C_7) \right]$$

$C_i$  : Wilson coefficients, abs. value of  $C_7$  from  $b \rightarrow s \gamma$   
 $s = \text{lepton pair mass squared}$

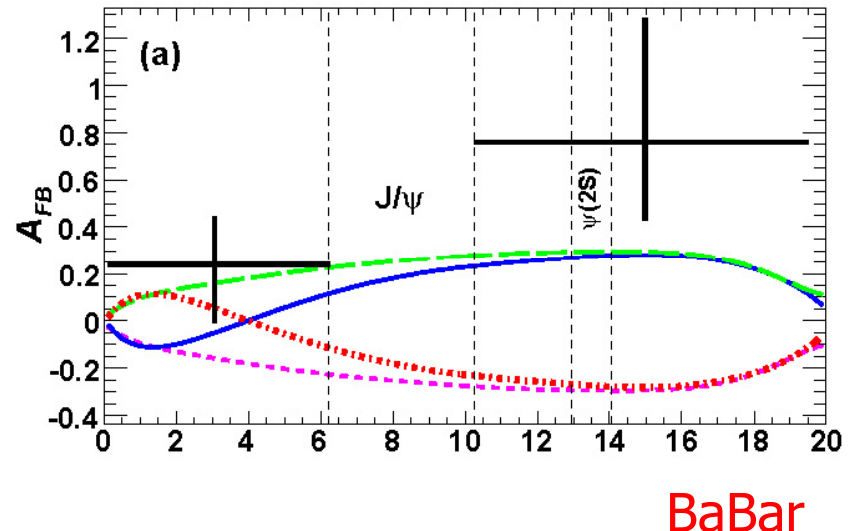
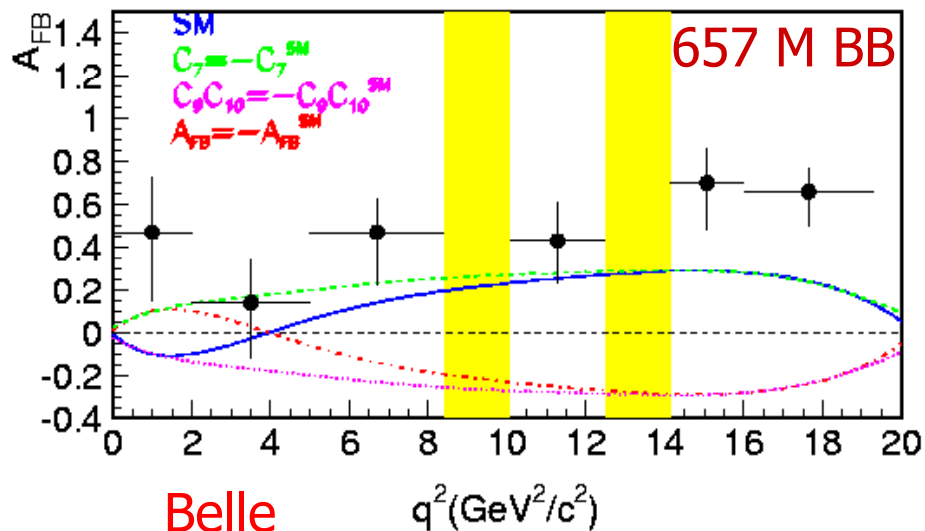


# Backward-forward asymmetry in $K^* I^+ I^-$

$$A_{FB} \propto \Re \left[ C_{10}^* (sC_9^{eff}(s) + r(s)C_7) \right]$$



$[\gamma^*$  and  $Z^*$  contributions in  $B \rightarrow K^* I^+ I^-$  interfere and give rise to forward-backward asymmetries c.f.  $e^+e^- \rightarrow \mu^+ \mu^-$ ]

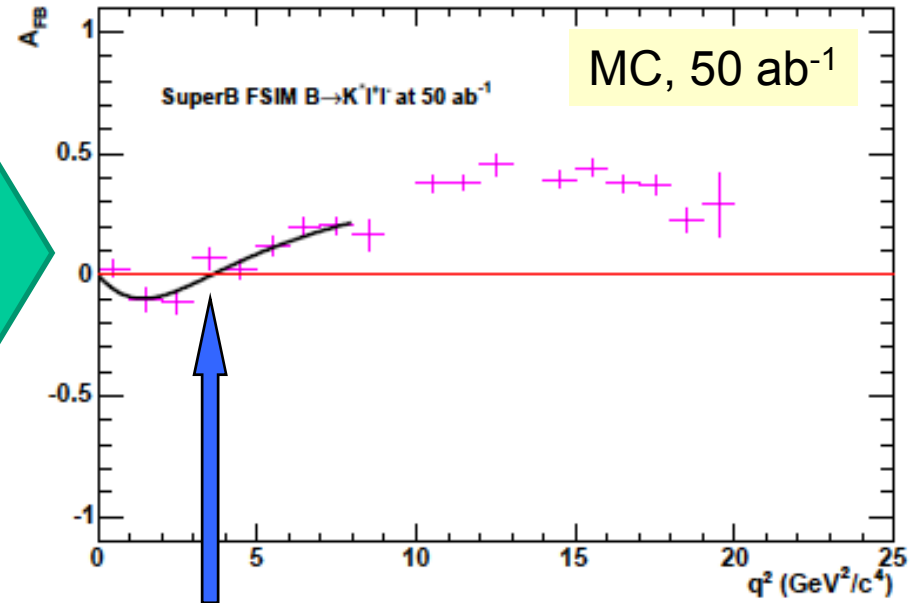
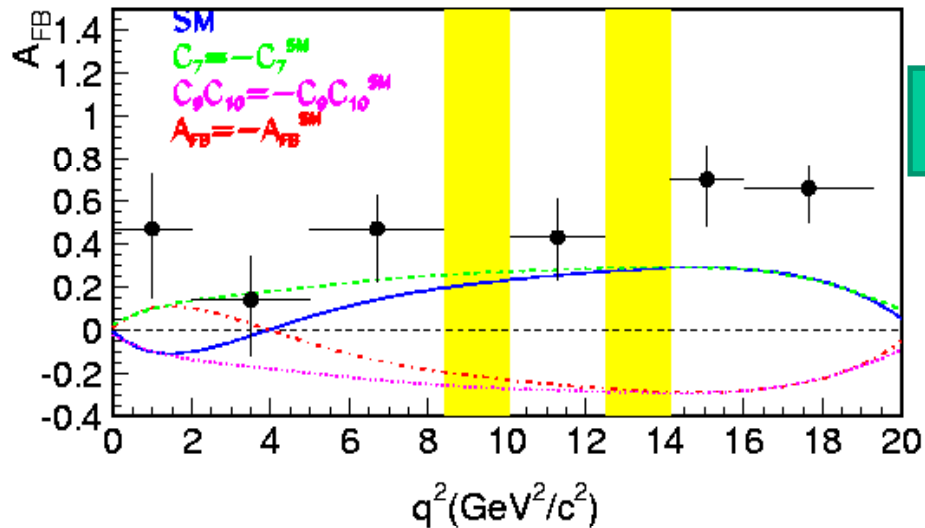






# $A_{FB}(B \rightarrow K^* l^+ l^-)[q^2]$ at a Super B Factory

657 M BB



- ▶ Zero-crossing  $q^2$  for  $A_{FB}$  will be determined with a 5% error with  $50 \text{ ab}^{-1}$ .

Strong competition from LHCb and ATLAS/CMS

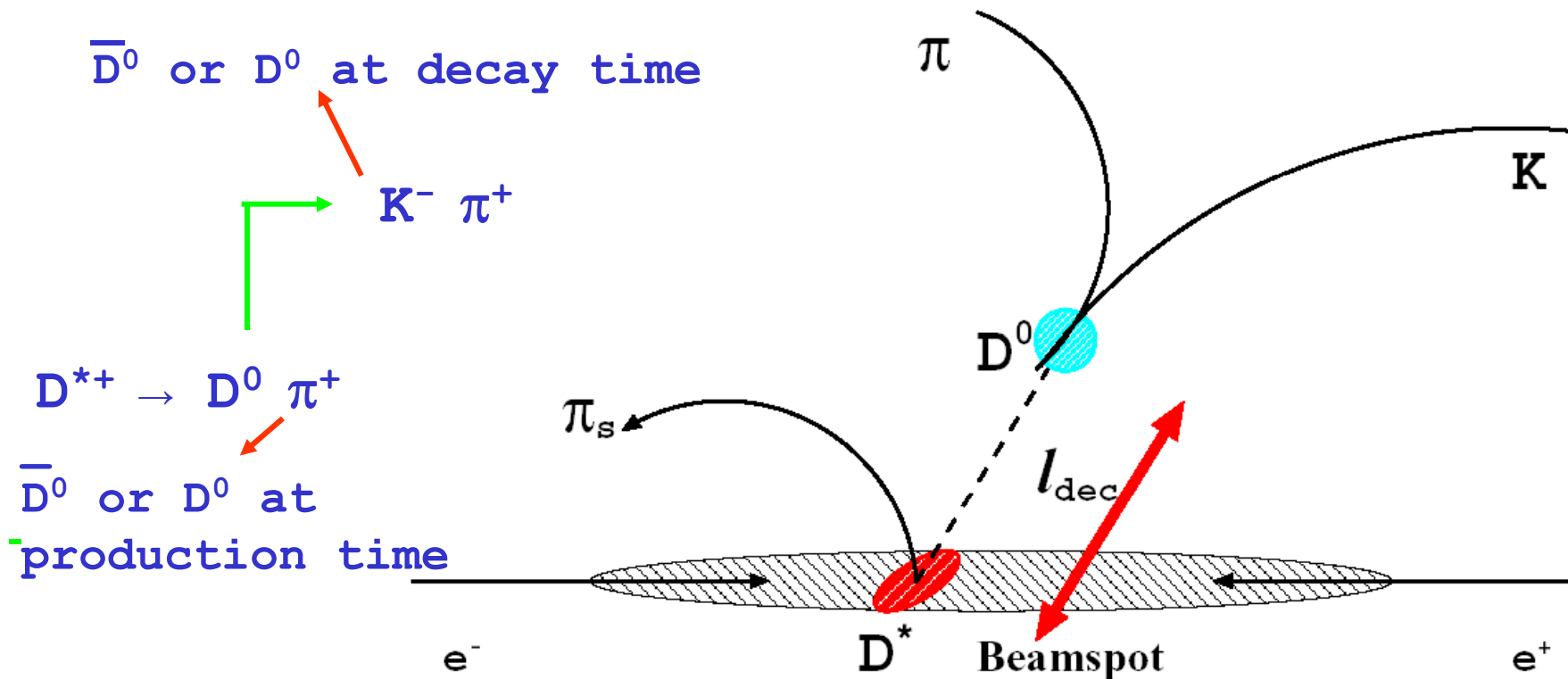


## Experimental methods in $D^0$ mixing searches

The method: investigate D decays in the decay sequence:



Used for tagging the **initial flavour** and for **background reduction**



$p_{\text{cms}}(D^*) > 2.5 \text{ GeV}/c$  eliminates D meson production from  $b \rightarrow c$



# D<sup>0</sup> mixing in K<sup>+</sup>K<sup>-</sup>, π<sup>+</sup>π<sup>-</sup>

D<sup>0</sup> → K<sup>+</sup>K<sup>-</sup> / π<sup>+</sup>π<sup>-</sup>

CP even final state;  
 in the limit of no CPV: CP|D<sub>1</sub>> = |D<sub>1</sub>>  
 ⇒ measure 1/Γ<sub>1</sub>

$$y_{CP} \equiv \frac{\tau(K^- \pi^+)}{\tau(K^- K^+)} - 1 = y \cos \varphi - \frac{1}{2} A_M x \sin \varphi =$$

$$\stackrel{\text{no CPV}}{=} y$$

S. Bergman et al., PLB486, 418 (2000)

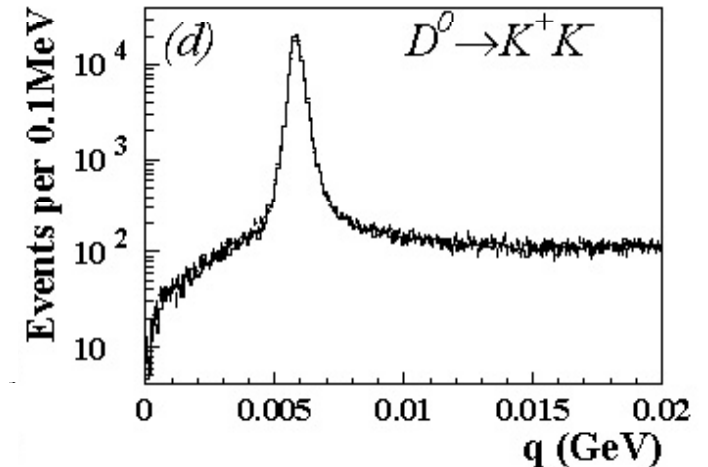
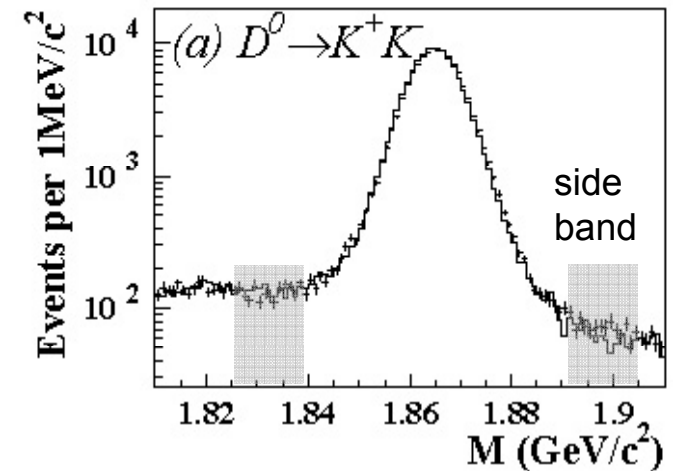
A<sub>M</sub>, φ: CPV in mixing and interference

Signal: D<sup>0</sup> → K<sup>+</sup>K<sup>-</sup> / π<sup>+</sup>π<sup>-</sup> from D<sup>\*</sup>

M, Q, σ<sub>t</sub> selection optimized in MC

	K <sup>+</sup> K <sup>-</sup>	K <sup>-</sup> π <sup>+</sup>	π <sup>+</sup> π <sup>-</sup>
N <sub>sig</sub>	111×10 <sup>3</sup>	1.22×10 <sup>6</sup>	49×10 <sup>3</sup>
purity	98%	99%	92%

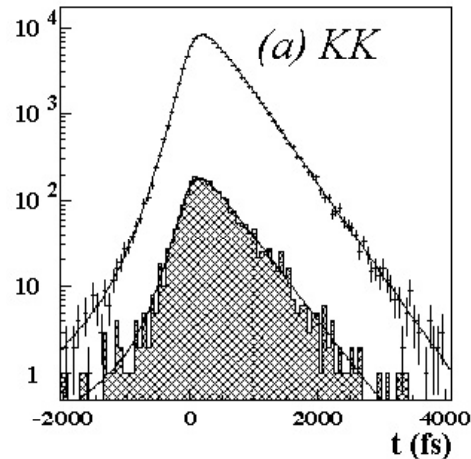
$$|D_{1,2}\rangle = p|D^0\rangle \pm q|\bar{D}^0\rangle$$



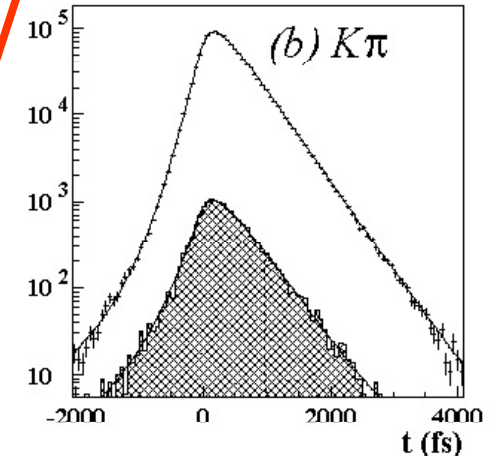
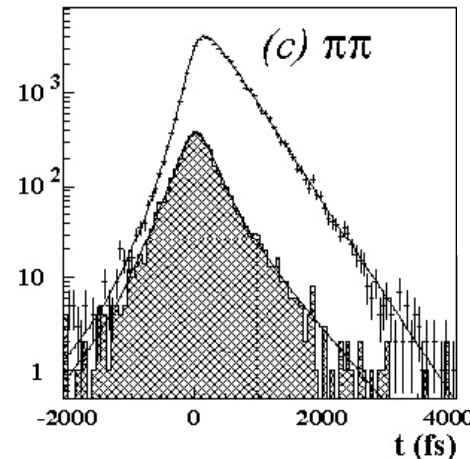


# $D^0$ mixing in $K^+K^-$ , $\pi^+\pi^-$

## Decay time distributions for $KK$ , $\pi\pi$ , $K\pi$



+



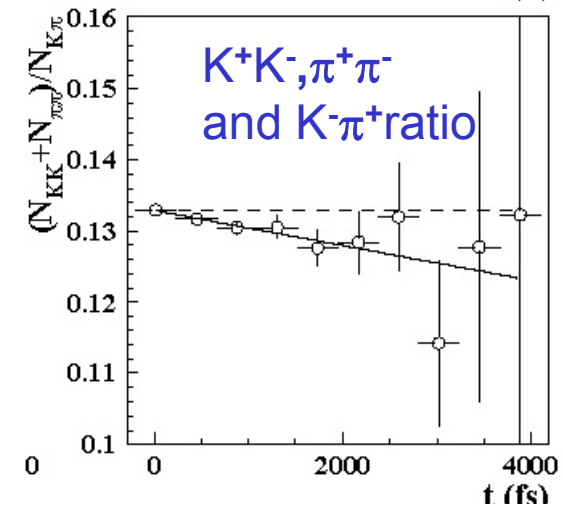
Difference of lifetimes  
visually observable  
in the ratio of the distributions  $\rightarrow$

Real fit:

$$y_{CP} = (1.31 \pm 0.32 \pm 0.25) \%$$

evidence for  $D^0$  mixing  
(regardless of possible CPV)

$\rightarrow y_{CP}$  is on the high side of SM expectations





# $D^0$ mixing in $K_S \pi^+ \pi^-$

## time-dependent Dalitz plot analysis

different decays identified through Dalitz plot analysis

CF:  $D^0 \rightarrow K^{*-} \pi^+$

DCS:  $D^0 \rightarrow K^{*+} \pi^-$

CP:  $D^0 \rightarrow \rho^0 K_S$

## time-dependence:

$$\mathcal{M}(m_-^2, m_+^2, t) \equiv \langle K_S \pi^+ \pi^- | D^0(t) \rangle =$$

$$= \frac{1}{2} \mathcal{A}(m_-^2, m_+^2) [e^{-i\lambda_1 t} + e^{-i\lambda_2 t}] + \frac{1}{2} \frac{q}{p} \bar{\mathcal{A}}(m_-^2, m_+^2) [e^{-i\lambda_1 t} - e^{-i\lambda_2 t}]$$

$\langle f | D^0 \rangle$

$\langle f | \bar{D}^0 \rangle$

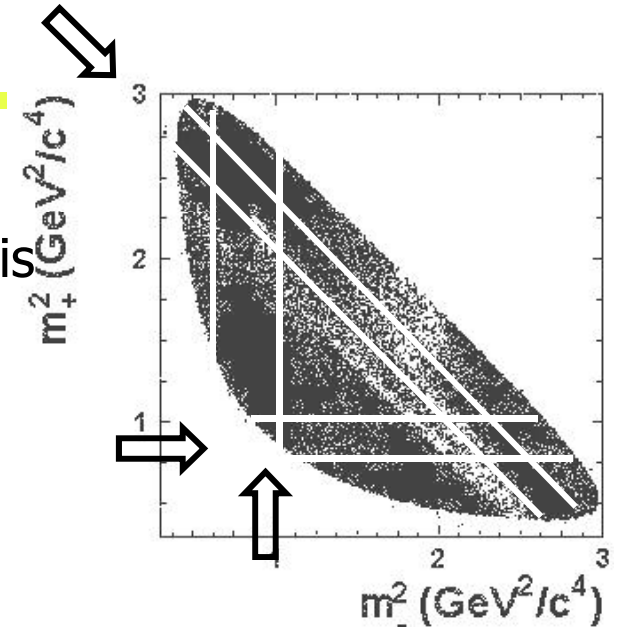
analogous for  $\bar{\mathcal{M}} = \langle f | \bar{D}^0(t) \rangle$

$m_{\pm}^2 = m^2(K_S \pi^{\pm})$ : Dalitz variables

$$\lambda_{1,2} = m_{1,2} - i\Gamma_{1,2}/2 = f(x, y)$$

Rate: terms with  $\cos(\mathbf{x}\Gamma t) \exp(-\Gamma t)$ ,  $\sin(\mathbf{x}\Gamma t) \exp(-\Gamma t)$ ,

- exp $(-(1+\mathbf{y}) \Gamma t) \rightarrow$  sensitive to x and y





# $D^0$ mixing in $K_S \pi^+ \pi^-$

Signal

$$N_{\text{sig}} = (534.4 \pm 0.8) \times 10^3$$
$$P \approx 95\%$$

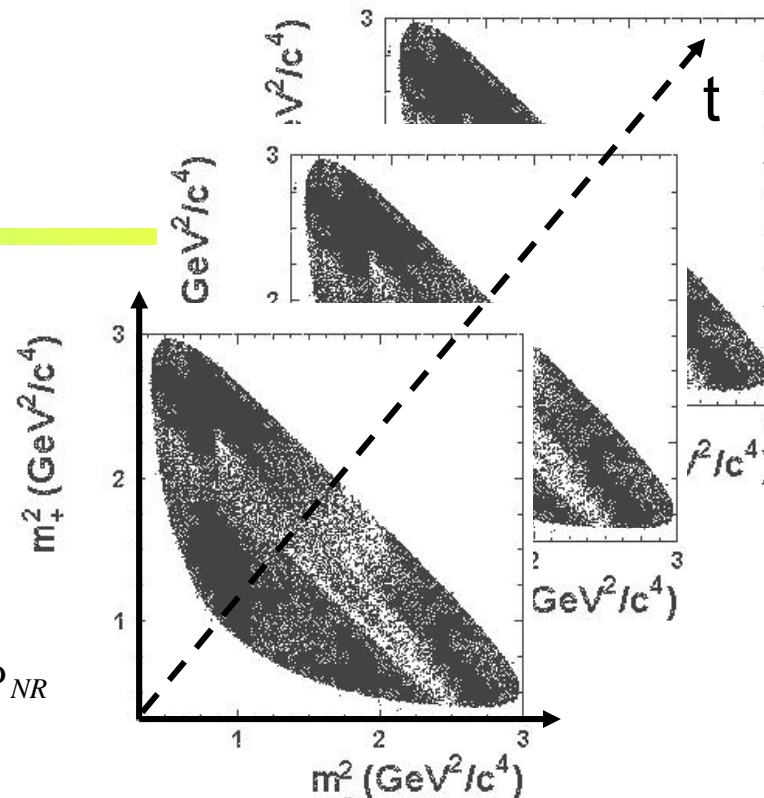
Dalitz model

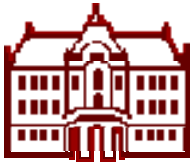
$$\mathcal{A}(m_-^2, m_+^2) = \sum a_r e^{i\Phi_r} B(m_-^2, m_+^2) + a_{NR} e^{i\Phi_{NR}}$$

18 resonant BW terms + non-resonant contribution

Fit  $|\mathcal{M}(m_-^2, m_+^2, t)|^2$  to the data distribution  $\Rightarrow x, y$

$$x = (0.80 \pm 0.29 \pm {}^{0.13}_{0.16})\%$$
$$y = (0.33 \pm 0.24 \pm {}^{0.10}_{0.14})\%$$





# Search for CP violation in the $D^0$ system

Relevant CKM elements of the 2x2 submatrix:

$$\begin{pmatrix} 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda + \frac{1}{2}A^2\lambda^5[1 - 2(\rho + i\eta)] & 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4(1 + 4A^2) & A\lambda^2 \\ A\lambda^3[1 - (1 - \frac{1}{2}\lambda^2)(\rho + i\eta)] & -A\lambda^2 + \frac{1}{2}A\lambda^4[1 - 2(\rho + i\eta)] & 1 - \frac{1}{2}A^2\lambda^4 \end{pmatrix}$$

$$\text{phase: } \sim \frac{2\eta A^2 \lambda^5}{\lambda} \sim O(10^{-3})$$

CPV in  $D^0$  very small,  $\leq 10^{-3}$ ;  
parameterization:

$$\frac{q}{p} \neq 1; \quad \frac{q}{p} \equiv \left(1 + \frac{A_M}{2}\right) e^{i\varphi}; \quad A_M, \varphi \neq 0$$

$D^0 \rightarrow K^+\pi^-, K^+K^- / \pi^+\pi^-, K_S \pi^+\pi^-$

t evolution depends also on CPV parameters

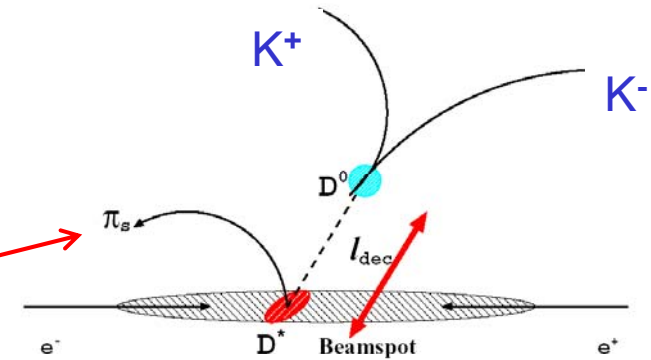
- x, y at upper limit of SM expectation → search for CPV
- at current level of sensitivity: positive signal clear indication of NP



# Search for CP violation

CPV in  $D^0 \rightarrow K^+K^- / \pi^+\pi^-$

Tag the D meson flavour (D or D-bar) by  $D^*$  charge (=charge of the 'slow' pion),  
 $D^{*+} \rightarrow \pi^+D^0$



$$y_{CP} \equiv \frac{\tau(K^- \pi^+)}{\tau(K^- K^+)} - 1 = y \cos \varphi - \frac{1}{2} A_M x \sin \varphi$$

$$A_\Gamma = \frac{\tau(\bar{D}^0 \rightarrow K^- K^+) - \tau(D^0 \rightarrow K^- K^+)}{\tau(\bar{D}^0 \rightarrow K^- K^+) + \tau(D^0 \rightarrow K^- K^+)} = \frac{1}{2} A_M y \cos \varphi - x \sin \varphi$$

$$A_\Gamma = (0.01 \pm 0.30 \pm 0.15) \%$$

indirect CPV





# Search for CP violation - continued

## CPV in $D^0 \rightarrow K_S \pi^+ \pi^-$

95% C.L. contours for  $(x, y)$ :

- CPV allowed: dash-dotted: statistical, dashed: statistical and systematic

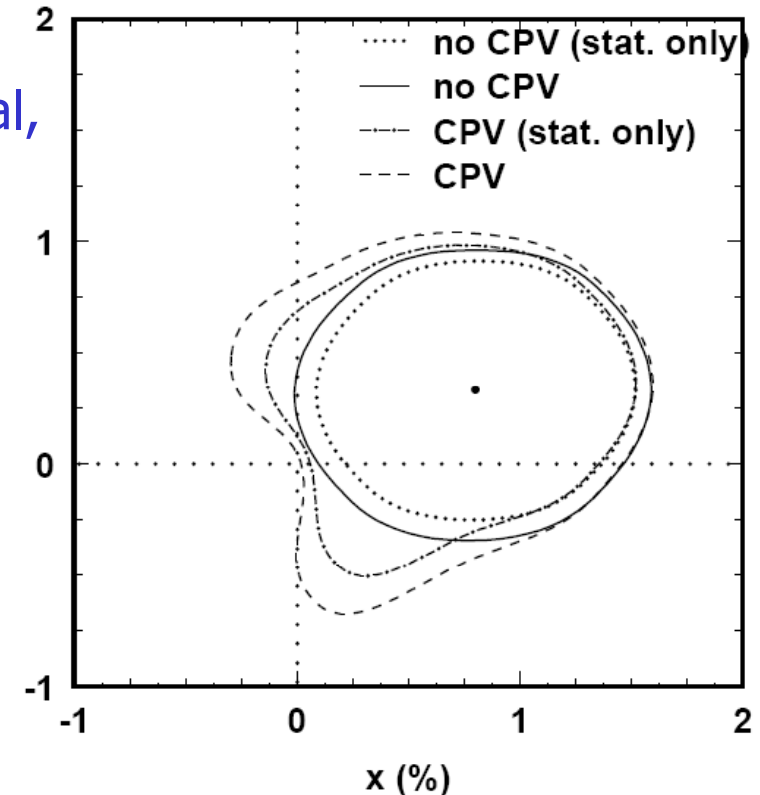
(No CPV assumed: dotted and solid)

Dalitz plot fit separately for  $D^0$  and  $\bar{D}^0$ :

- Fit parameters consistent for both samples  $\rightarrow$  no direct CPV

- Parameters  $|q/p|$  and  $\phi = \arg(q/p)$  consistent with CP conservation

Fit assuming no direct CPV  $\rightarrow$   
Parameters of CPV in mixing and interf. in mixing and decay:



$$|q/p| = 0.95 \pm_{0.20}^{0.22}$$
$$\phi = \arg(q/p) = (-2 \pm_{11}^{10})^\circ$$



# D<sup>0</sup> mixing

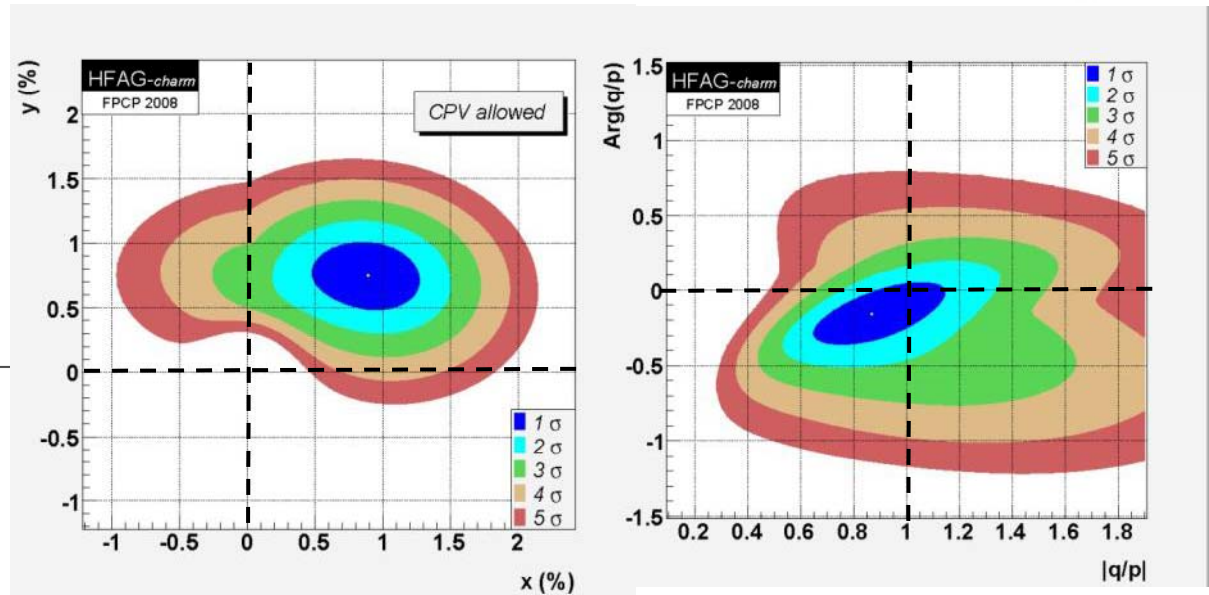
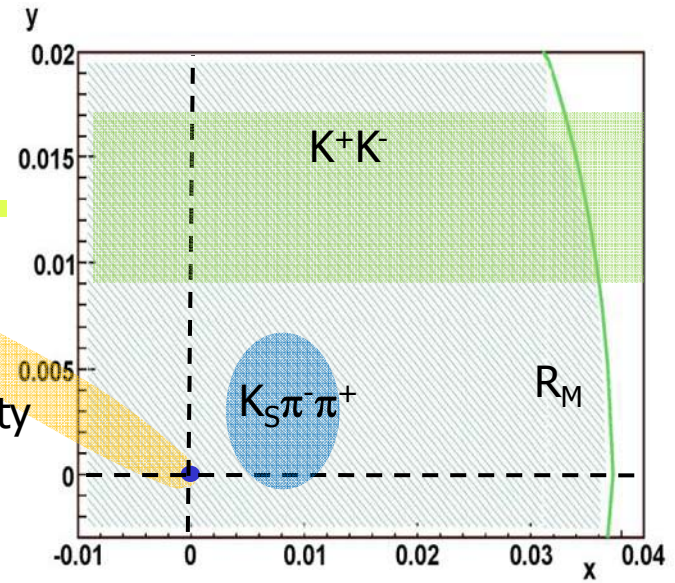
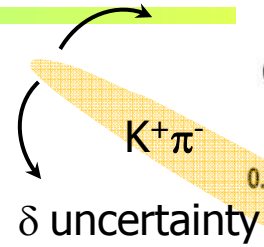
Average of results

—  $\chi^2$  fit including correlations among measured quantities

Parameter	CPV-allowed	
$x$ (%)	$0.89^{+0.26}_{-0.27}$	
$y$ (%)	$0.75^{+0.17}_{-0.18}$	
$\delta$ (°)	$21.9^{+11.3}_{-12.4}$	
$R_D$ (%)	$0.3348 \pm 0.0086$	
CPV	$A_D$ (%)	$-2.0 \pm 2.4$
	$ q/p $	$0.87^{+0.18}_{-0.15}$
	$\phi$ (°)	$-9.1^{+8.1}_{-7.8}$
	$\delta_{K\pi\pi}$ (°)	$33.0^{+25.9}_{-26.6}$

$(x,y) \neq (0,0)$ :  $6.7\sigma$ ;  
 CP even state heavier and shorter lived;  
 no CPV within  $1\sigma$

$\chi^2/n.d.f. = 23.5/18$



$x(D^0) \approx 0.01$ ;  $x(K^0) \approx 1$ ;  $x(B_d) \approx 0.8$ ;  $x(B_s) \approx 25$



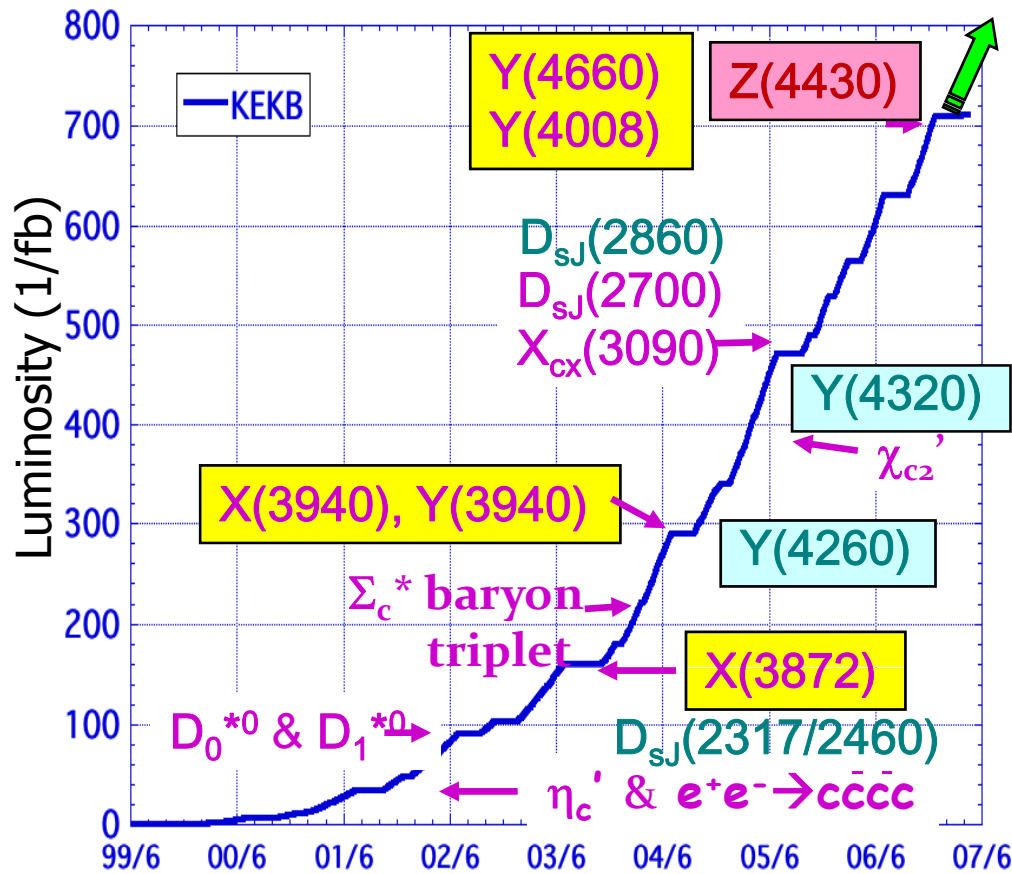
## B factories: a success story

- Measurements of CKM matrix elements and angles of the unitarity triangle
- Observation of direct CP violation in B decays
- Measurements of rare decay modes (e.g.,  $B \rightarrow \tau \nu$ ,  $D \tau \nu$ ) by fully reconstructing the other B meson
- Observation of D mixing
- CP violation in  $b \rightarrow s$  transitions: probe for new sources if CPV
- Forward-backward asymmetry ( $A_{FB}$ ) in  $b \rightarrow s l^+ l^-$  has become a powerful tool to search for physics beyond SM.
- Observation of new hadrons

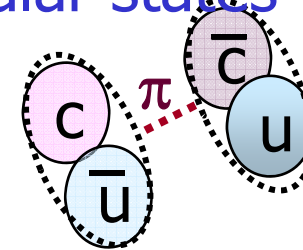


# New hadrons at B-factories

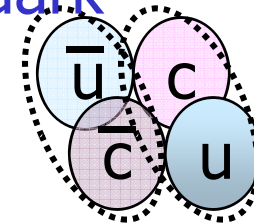
Discoveries of many new hadrons at B-factories have shed light on new class of hadrons beyond the ordinary mesons.



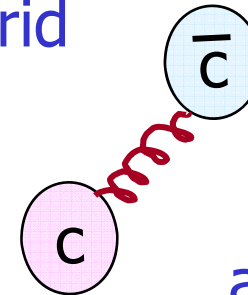
Molecular states



Tetra-quark



Hybrid



and more...



# Back-up slides

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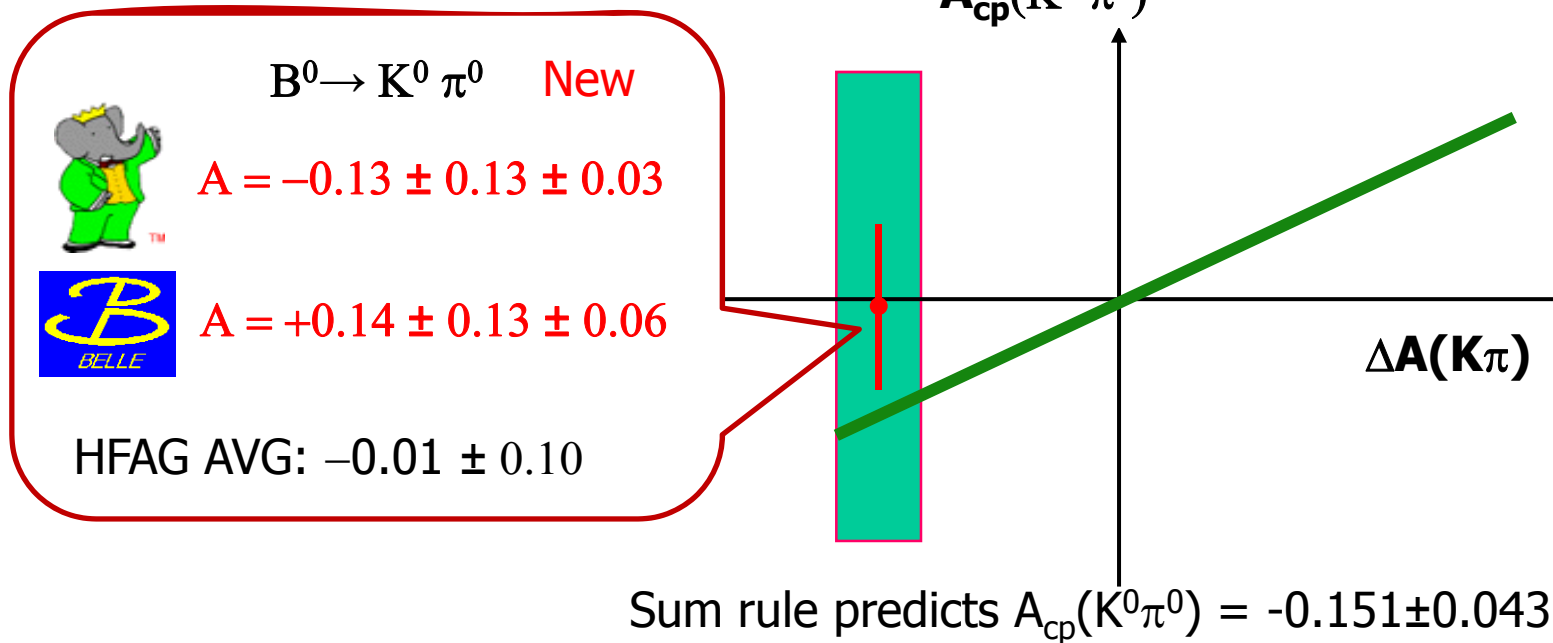
# Model-indep. check of NP

M. Gronau, PLB 627, 82 (2005);

- $A_{cp}(K\pi)$  sum rule

D. Atwood & A. Soni, Phys. Rev. D 58, 036005(1998).

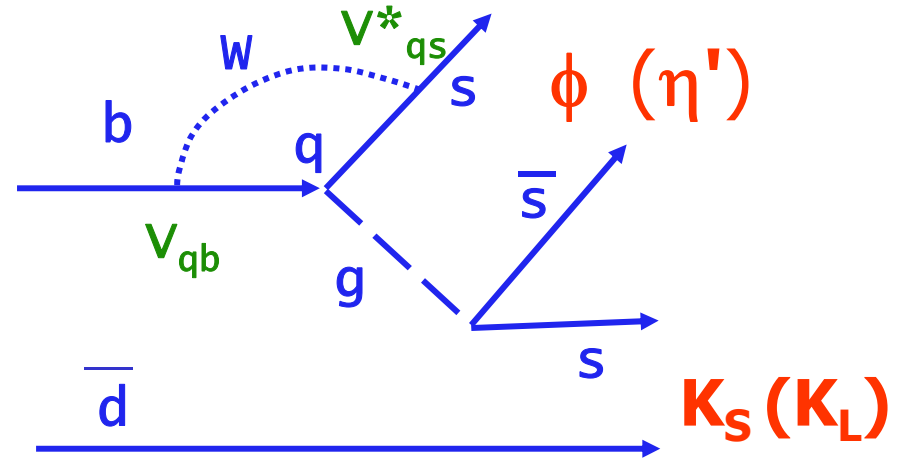
$$A_{CP}(K^+\pi^-) + A_{CP}(K^0\pi^+) \frac{\mathcal{B}(K^0\pi^+) \tau_0}{\mathcal{B}(K^+\pi^-) \tau_+} = A_{CP}(K^+\pi^0) \frac{2\mathcal{B}(K^+\pi^0) \tau_0}{\mathcal{B}(K^+\pi^-) \tau_+} + A_{CP}(K^0\pi^0) \frac{2\mathcal{B}(K^0\pi^0)}{\mathcal{B}(K^+\pi^-)}$$





# b → sss decays

Pure penguin diagrams



$$A(s\bar{s}s) = V_{cb}V_{cs}^* (P_s^c - P_s^t) + V_{ub}V_{us}^* (P_s^u - P_s^t).$$

$$V_{cb}V_{cs}^* = A\lambda^2$$

$$V_{ub}V_{us}^* = A\lambda^4(\rho - i\eta)$$

First term dominates →

$\lambda$  same as for  $J/\psi K_S$

$$\lambda_{\phi K_S} = \eta_{\phi K_S} \left( \frac{V_{tb}^* V_{td}}{V_{tb} V_{td}^*} \right) \left( \frac{V_{cd}^* V_{cb}}{V_{cd} V_{cb}^*} \right)$$

$$\text{Im}(\lambda_{\phi K_S}) = \sin 2\phi_1 = \sin 2\beta$$