Measurement of strong-phase difference between D^0 and $\overline{D^0} \rightarrow K_S^0 K^+ K^-$ decay at BESIII



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Outline

Flavour Physics

CKM Matrix

Measurement of angle γ

BPGGSZ method- model-dependent and model-independent methods

Strong-phase coefficients

BEPCII and **BESIII**

Dalitz plots

Dalitz plot binning

Bin yields, corrections, backgrounds etc.

Extraction of strong-phase parameters

Fitter

Systematic uncertainties

Impact of strong-phase measurements on γ A sneek-peek into LHCb result

Summary

Flavour Physics

....

Standard Model is an QFT that describes the fundamental particles and its interaction (Weak , em and strong)

All test made on SM has been successful. Not full story!

Flavour sector is less well known (more open questions)

Flavour relates the existence of family of quarks and how they couple to each other

Why flavour physics is interesting?

- Why 3 generations of quarks? Why only 3?
- Extreme hierarchy of masses (2.4 to 1.75×10^5) MeV/ c^2
- CP violation explained in SM; but not enough to explain matter –antimatter asymmetry

These mysteries makes flavour physics of SM of great interest



Charm Factories: BESIII, CLEO-c

B- Factories: BELLE, BELLEII, BaBar, LHCb

Mixing between weak eigenstates and flavor eigenstates in three generations.

$$\begin{pmatrix} d'\\ s'\\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub}\\ V_{cd} & V_{cs} & V_{cb}\\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d\\ s\\ b \end{pmatrix} \qquad V_{CKM}V_{CKM}^{\dagger} = I = V_{CKM}^{\dagger}V_{CKM}$$



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3 angles and 1 phase required to write it down

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

(b-d unitary triangle)

$$V_{ud}V_{ub}^{*}$$
 α $V_{td}V_{tb}^{*}$ γ β $V_{cd}V_{cb}^{*}$

$$lpha/\phi_2 = (84.9^{+5.1}_{-4.5})^\circ$$

 $eta/\phi_1 = (22.2 \pm 0.7)^\circ$
 $\gamma/\phi_3 = (71.1^{+4.6}_{-5.3})^\circ$
http://ckmfitter.in2p3.fr



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$$Q$$

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One aim of flavor physics experiments is to measure CKM parameters precisely



Direct measurements

Indirect measurements

Direct measurements

- Measure it using tree level decays
- Theoretical uncertainty $\mathcal{O}(10^{-7})$



Large experimental uncertainties, potential for further improvement in coming years

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Indirect measurements

- Infer the value of *γ* using other sides of triangle, assuming the triangle is closed
- NP effects can play potential for different central value



Uncertainties from LQCD

Direct measurements

- Measure it using tree level decays
- Theoretical uncertainty $\mathcal{O}(10^{-7})$ JHEP 01, 051 (2014)



Large experimental uncertainties, potential for further improvement in coming years

Indirect measurements

- Infer the value of γ using other sides of triangle, assuming the triangle is closed
- NP effects can play potential for different central value



Uncertainties from LQCD

Precise measurement required for meaningful comparison

Measurement of γ/ϕ_3



$$B^{\pm} \rightarrow DK^{\pm}$$
 where $D = D^0$ or $\overline{D^0}$ (PLB **265**, 172 (1991))

Measurement of γ/ϕ_3



Common final states-possibility of interference – access to phase term

$$\Gamma \propto |f(B^- \to DK^-)|^2 = A_B^2 + A_B^2 r_B^2 + 2A_B^2 r_B^2 \cos(\delta_B - \gamma)$$

$$r_B = \left| \frac{f(B^- \to \overline{D^0} K^-)}{f(B^- \to D^0 K^-)} \right|$$

Sensitivity to γ comes from interference

 $B^{\pm} \rightarrow DK^{\pm}$ where $D = D^0$ or $\overline{D^0}$ (PLB 265, 172 (1991))

 $b \rightarrow c \bar{u} s + b \rightarrow u \bar{c} s$



Other related modes with D^* or K^* in final states can also be used

Measurement of γ/ϕ_3



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Sensitivity to γ comes from interference

 $B^{\pm} \rightarrow D\pi^{\pm}$ modes has larger BF but small value of r_B less sensitive to γ

 $B^{\pm} \rightarrow DK^{\pm}$ where $D = D^0$ or $\overline{D^0}$ (PLB 265, 172 (1991))

 $b \rightarrow c \bar{u} s + b \rightarrow u \bar{c} s$



Other related modes with D^* or K^* in final states can also be used

Methods to measure γ from $B^{\pm} \rightarrow DK^{\pm}$

- Gronau, London and Wyler method (GLW): D decay to CP eigenstate, K^-K^+ , $K_S^0\pi^0$... PLB 265, 172 (1991) Atwood, Dunietz and Soni method (ADS): D decay to CF and DCS states, $K^-\pi^+$, $K^-\pi^+\pi^0$ Phys. Rev. Lett. 78, 3357 (1997)
- □ Bondar, Poluektov, Giri, Grossman, Soffer and Zupan method (BPGGSZ): *D* decay to multibody final states, $K_S^0 K^+ K^-$, $K_S^0 \pi^+ \pi^-$, $K_S^0 \pi^+ \pi^- \pi^0$... PRD 68, 054018 (2003)

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BPGGSZ method

Dalitz plot analysis of multibody final states.

Dalitz plot: A scatter plot of decay in terms of Lorentz invariant quantities Phil. Mag. 44, 1068 (1953)

$$d\Gamma = \frac{1}{(2\pi)^3} \frac{1}{32M^3} |\mathcal{A}|^2 dm_{ab}^2 dm_{bc}^2$$

Multibody final states: decays proceeds through various intermediate final states.



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$$\Box D \rightarrow K_S^0 K^+ K^- \text{ coordinates:} \quad m_{\pm}^2 = (P_{K_S^0} \pm P_{K^{\pm}})^2$$



$$\Box D \to K_{S}^{0}K^{+}K^{-} \text{ coordinates:} \quad m_{\pm}^{2} = (P_{K_{S}^{0}} \pm P_{K^{\pm}})^{2}$$
$$\Box f_{B^{-}}(m_{\pm}^{2}, m_{-}^{2}) \propto f_{D}(m_{\pm}^{2}, m_{-}^{2}) + r_{B}e^{i(\delta_{B}-\gamma)}f_{\overline{D}}(m_{\pm}^{2}, m_{-}^{2})$$



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$$\Box \text{ Neglecting the CP violation and second-order effects of charm mixing:} f_{\overline{D}}(m_+^2, m_-^2) \equiv f_D(m_-^2, m_+^2)$$



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$$f_D(m_{+,}^2 m_{-}^2) = |f_D(m_{+,}^2 m_{-}^2)| e^{i\delta_D(m_{+,}^2 m_{-}^2)}$$

 $d\Gamma (B^- \to D(K_S^0 K^+ K^-) K^-) \propto \left(|f_D(m_+^2, m_-^2)|^2 + r_B^2 |f_D(m_-^2, m_+^2)|^2 + 2r_B \Re[f_D(m_+^2, m_-^2) f_D^*(m_-^2, m_+^2) e^{-i(\delta_B - \gamma)}] \right) dm_+^2 dm_-^2 dm_+^2 dm_-^2 dm_+^2 dm_+^$

 $\Delta \delta_D(m_+^2, m_-^2) = \delta_D(m_+^2, m_-^2) - \delta_D(m_-^2, m_+^2)$

(Strong-phase difference)



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$$f_D(m_{+,}^2 m_{-}^2) = |f_D(m_{+,}^2 m_{-}^2)| e^{i\delta_D(m_{+,}^2 m_{-}^2)}$$

 $d\Gamma \left(B^{-} \rightarrow D\left(K_{S}^{0}K^{+}K^{-}\right)K^{-}\right) \propto \left(\mid f_{D}\left(m_{+}^{2}, m_{-}^{2}\right)\mid^{2} + r_{B}^{2}\mid f_{D}\left(m_{-}^{2}, m_{+}^{2}\right)\mid^{2} + 2r_{B} \Re[f_{D}\left(m_{+}^{2}, m_{-}^{2}\right)f_{D}^{*}\left(m_{-}^{2}, m_{+}^{2}\right)e^{-i(\delta_{B}-\gamma)}]\right) \ dm_{+}^{2}dm_{-}^{2}dm_{+}^{2}dm_{-}^{2}dm_{+}$

$$\Delta \delta_D(m_+^2, m_-^2) = \delta_D(m_+^2, m_-^2) - \delta_D(m_-^2, m_+^2)$$

(Strong-phase difference)

Knowledge of D decay dynamics crucial



 $\begin{array}{l} \Box \ D \rightarrow K_S^0 K^+ K^- \ \text{coordinates:} \quad m_{\pm}^2 = (P_{K_S^0} \pm P_{K^{\pm}})^2 \\ \Box \ f_{B^-} \big(m_+^2, m_-^2 \big) \propto f_D \big(m_+^2, m_-^2 \big) + \ r_B e^{i(\delta_B - \gamma)} f_{\overline{D}}(m_+^2, m_-^2) \\ \Box \ \text{Neglecting the CP violation and second-order effects of charm mixing:} \\ f_{\overline{D}}(m_+^2, m_-^2) \equiv f_D \big(m_-^2, m_+^2 \big) \end{array}$

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Knowledge of D decay dynamics crucial

Multibody decays proceeds through various intermediate resonant state hence large Strong-phase variations expected over Dalitz plot.

 $\mathcal{BF}(D^0 \to K^0_S K^+ K^-) = (4.45 \pm 0.19) \times 10^{-3}$ (PDG) PTEP **2020**, 083C01 (2020) ²⁴

Model-dependent measurements

 $\Box f_D(m_+^2, m_-^2)$ from an amplitude model for $D \to K_S^0 K^+ K^-$



Model-dependent measurements

 $\Box f_D(m_+^2, m_-^2)$ from an amplitude model for $D \to K_S^0 K^+ K^-$



Model-independent measurements

Binned Dalitz plot

□ Requires amplitude weighted average values of $\Delta \delta_D(m^2_+, m^2_-)$ in each bins





Model-dependent measurements

Model-independent measurements



Model-dependent measurements

Model-independent measurements



Steps:





Steps:

□ Divide Dalitz plot into 2 \mathcal{N} bins (indexed *i*), symmetrically around $m_+^2 = m_-^2$ line (Note: no assumptions on bin shapes) PRD 68, 054018 (2003) □ Yield of $B^{\pm} \rightarrow D(K_S^0 K^+ K^-) K^{\pm}$ decay in *i*th bin

$$N_i^{\mp} \propto (K_{\pm i} + r_B^2 K_{\mp i} + 2\sqrt{K_i K_{-i}} (x_{B \mp} c_i + y_{B \mp} s_i))$$

- $x_{B\pm} = r_B \cos(\delta_B \pm \gamma) \qquad y_{B\pm} = r_B \sin(\delta_B \pm \gamma) \qquad r_B^2 = x_{B\pm}^2 + y_{B\pm}^2$
- K_i : Flavor-tagged $D^0 \rightarrow K^0_S K^+ K^-$ events



Square binning scheme

Steps:

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 K_i : Flavor-tagged $D^0 \rightarrow K_S^0 K^+ K^-$ events



$$c_{i} = \frac{\int_{i} |f_{D}(m_{+}^{2}, m_{-}^{2})||f_{D}(m_{-}^{2}, m_{+}^{2})| \times \cos\left(\Delta\delta_{D}(m_{+}^{2}, m_{-}^{2})\right) dm_{+}^{2}, dm_{-}^{2}}{\sqrt{F_{i}F_{-i}}}$$

$$s_{i} = \frac{\int_{i} |f_{D}(m_{+}^{2}, m_{-}^{2})| |f_{D}(m_{-}^{2}, m_{+}^{2})| \times \sin\left(\Delta\delta_{D}(m_{+}^{2}, m_{-}^{2})\right) dm_{+}^{2}, dm_{-}^{2}}{\sqrt{F_{i}F_{-i}}}$$

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Needs to be measured first.
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35

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Determination of strong-phase difference from $D^0 \overline{D^0}$ sample

 c_i and s_i can be measured using $e^+e^- \rightarrow \psi(3770) \rightarrow D^0\overline{D^0}$

Data from charm factories collected at \sqrt{s} = 3.773 GeV


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Data from charm factories collected at \sqrt{s} = 3.773 GeV

 $D^{0}\overline{D^{0}} \text{ are in quantum correlated } (C = -1 \text{ state})$ $e^{+}e^{-} \rightarrow \psi(3770) \rightarrow \frac{1}{\sqrt{2}} \left[D^{0}\overline{D^{0}} - \overline{D^{0}}D^{0} \right]$ $e^{+}e^{-} \rightarrow \psi(3770) \rightarrow \frac{1}{\sqrt{2}} \left[D_{CP-}D_{CP+} - D_{CP+}D_{CP-} \right]$ $D_{CP\pm} = \frac{(D^{0} \pm \overline{D^{0}})}{\sqrt{2}}$



Both D has opposite CP to each other: reconstructing one D in CP eigenstates gives the CP of other D

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Tagging





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Tagging

Singletag (ST): Only one **D** meson in reconstructed in an event For eg: $D^0 \rightarrow K_S^0 K^+ K^- vs \overline{D^0} \rightarrow anything$ or $D^0 \rightarrow anything vs \overline{D^0} \rightarrow K_S^0 K^+ K^-$





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Double tag (DT): Both **D** meson reconstructed in an event For eg: $D^0 \rightarrow K_S^0 K^+ K^- vs \overline{D^0} \rightarrow K^+ K^$ or $D^0 \rightarrow K^+ K^- vs \overline{D^0} \rightarrow K_S^0 K^+ K^-$

Flavor identification not possible





 c_i can determined from *CP*-tagged $K_S^0 K^+ K^-$ events (DT $K_S^0 K^+ K^-$ vs *CP*± tag modes)

 F_+ = 1 (0) for pure CP+ (CP-) states

$$\langle M_i^{\pm} \rangle = \frac{S_{\pm}}{S_f} \left(K_i - 2\boldsymbol{c_i}(2F_+ - 1)\sqrt{K_iK_{-i}} + K_{-i} \right) \times \boldsymbol{\epsilon}_{\mathrm{DT},i}$$

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 c_i and s_i can be determined from $K_S^0 K^+ K^-$ vs $K_S^0 h^+ h^-$ ($h = K, \pi$)

$$\left\langle M_{ij} \right\rangle = \frac{N_D^0 \overline{D^0}}{2S_f^2} \left(K_i K_{-j} + K_{-i} K_j - 2\sqrt{K_i K_{-j} K_{-i} K_j} (c_i c_j + s_i s_j) \right) \times \epsilon_{DT,ij}$$

Done!

 $F_+ = 1$ (0) for pure CP+ (CP-) states

For $K_S^0 K^+ K^-$ vs $K_S^0 \pi^+ \pi^- c_j, s_j$ corresponds to strong phase parameters of $D^0 \rightarrow K_S^0 \pi^+ \pi^-$

> PRL 124, 241802 (2020) PRD 124, 241802 (2020)

 c_i can determined from *CP*-tagged $K_S^0 K^+ K^-$ events (DT $K_S^0 K^+ K^-$ vs *CP*± tag modes)

$$\langle M_i^{\pm} \rangle = \frac{S_{\pm}}{S_f} \left(K_i - 2c_i(2F_+ - 1)\sqrt{K_iK_{-i}} + K_{-i} \right) \times \epsilon_{\mathrm{DT},i}$$

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Done!

For $D^0 o K_L^0 K^+ K^-$ we can measure c'_i and s'_i

 c'_i can determined from *CP*-tagged $K^0_L K^+ K^-$ events

$$\left\langle M_i^{\prime \pm} \right\rangle = \frac{S_{\pm}}{S_f} \left(K_i^{\prime} + 2c_i^{\prime} (2F_+ - 1) \sqrt{K_i^{\prime} K_{-i}^{\prime}} + K_{-i}^{\prime} \right) \times \epsilon_{\mathrm{DT},i}$$

 $F_{+} = 1$ (0) for pure CP+ (CP-) states

For $K_S^0 K^+ K^-$ vs $K_S^0 \pi^+ \pi^- c_j, s_j$ corresponds to strong phase parameters of $D^0 \rightarrow K_S^0 \pi^+ \pi^-$

> PRL 124, 241802 (2020) PRD 124, 241802 (2020)

 c'_i and s'_i not required for γ measurements

Required for improving the precision of c_i and s_i

 c'_i and s'_i can be determined from $K^0_L K^+ K^-$ vs $K^0_S h^+ h^ (h = K, \pi)$

$$\left\langle M_{ij}^{\prime}\right\rangle = \frac{N_{D}^{0}\overline{D^{0}}}{2S_{f}^{2}} \left(K_{i}K_{-j}^{\prime} + K_{-i}K_{j}^{\prime} - 2\sqrt{K_{i}K_{-j}^{\prime}K_{-i}K_{j}^{\prime}}(c_{i}c_{j}^{\prime} + s_{i}s_{j}^{\prime})\right) \times \epsilon_{DT,ij}$$

Beijing Spectrometer Experiment (BESIII)



BEPC II

□ Two ring e^+e^- symmetric collider; circumference: 240 m □ Design $\mathcal{L} = 1 \times 10^{33} \text{ cm}^{-2} \text{s}^{-1}$ □ $\sqrt{s} = 2 - 4.6 \text{ GeV}$

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BESIII

□ Hermiticity: 93% of 4π □ MDC: $\frac{\sigma_p}{p} = 0.5\%$ @ 1 GeV □ TOF: $\sigma = 80$ ps (100 ps) in barrel and end-cap regions □ ECL: $\frac{\sigma_E}{E} = 2.5\%$ @ 1 GeV □ SC solenoid: 1T

arXiv:hep-ex/0809.1869

"Second generation charm factory after CLEO-c" (Since 2009)

Analysis strategy

Tag modes

CP and flavor modes will be reconstructed both as ST and DT

□ Signal:
$$K_S^0 K^+ K^-$$
, $K_L^0 K^+ K^-$
□ *CP* even: $K^+ K^-$, $\pi^+ \pi^-$, $K_S^0 \pi^0 \pi^0$, $\pi^+ \pi^- \pi^0$, $K_L^0 \pi^0$, $K_L^0 \omega$, $K_L^0 \eta$, $K_L^0 \eta'$
□ *CP* odd: $K_S^0 \pi^0$, $K_S^0 \omega$, $K_S^0 \eta$, $K_S^0 \eta'$
□ Flavor: $K^- \pi^+$, $K^- \pi^+ \pi^0$, $K^- e^+ v_e$
□ Mixed *CP*: $K_S^0 \pi^+ \pi^-$, $K_L^0 \pi^+ \pi^-$

Reconstruction of final state particles

$$K_S^0 \to \pi^+\pi^-, \omega \to \pi^+\pi^-\pi^0, \pi^0 \to \gamma\gamma, \eta \to \gamma\gamma, \eta' \to \pi^+\pi^-\eta$$

Particles are combined to reconstruct final states

- □ Fully reconstructed tags: Full kinematic reconstruction possible,
- Partially reconstructed tags: Particles having missing particle in final states; no full reconstruction, inference using missing energy and momentum

Branching fraction (1.43 ± 0.06) % (PDG) PTEP **2020**, 083C01 (2020) $F_{+} = 0.973 \pm 0.017$

PLB **747**, 9 (2015)





Data yields

ST Yield

d Total PDF = Signal MC shape ⊗ Gaussian + Argus PLB **241**, 278 (1990)

	$N_{\rm ST}$ by integrating signa				
Mode	ST				
	$N_{ m ST}$	$\epsilon_{ m ST}(\%)$			
Flavor-tags					
$K^{-}\pi^{+}$	524307 ± 742	63.31 ± 0.06			
$K^-\pi^+\pi^0$	995683 ± 1117	31.70 ± 0.03			
$K^- e^+ \nu_e$	752387 ± 12795				
CP-even tags					
K^+K^-	53481 ± 247	61.02 ± 0.11			
$\pi^+\pi^-$	19339 ± 163	64.52 ± 0.11			
$K^0_{ m S}\pi^0\pi^0$	19882 ± 233	14.86 ± 0.08			
$\pi^+\pi^-\pi^0$	99981 ± 618	37.65 ± 0.11			
$K^0_{ m L}\pi^0$	209445 ± 14796				
$K^0_{ m L}\eta(\gamma\gamma)$	40009 ± 2543				
$K^0_{ m L}\omega$	207376 ± 11498				
$K_{ m L}^0\eta'(\pi^+\pi^-\eta)$	33683 ± 1909				
CP-odd tags					
$K^0_{ m S}\pi^0$	65072 ± 281	36.92 ± 0.11			
$K^0_{ m S}\eta(\gamma\gamma)$	9524 ± 134	32.94 ± 0.11			
$K^0_{ m S}\omega$	19262 ± 157	12.14 ± 0.07			
$K^0_{ m S}\eta'(\pi^+\pi^-\eta)$	3301 ± 62	12.46 ± 0.07			



For partially reconstructed tags

$$N_{ST} = 2 \times N_{D^0 \overline{D^0}} \times \mathcal{BF}$$

Note: $\mathcal{BF}(K_L^0 X) = \mathcal{BF}(K_S^0 X)$ for $(X = \eta, \eta', \omega)$ (PDG) PTEP **2020**, 083C01 (2020)

Difference expected to be around 10% PLB **349**, 363 (1995) 1.88



For partially reconstructed tag sideband subtraction done on $M_{\rm miss}^2$ or $U_{\rm miss}$ distribution

$$Y_{S} = \frac{(N_{S} - N_{S}^{P}) - \delta(N_{L} - N_{L}^{P}) - \gamma(N_{H} - N_{H}^{P})}{1 - \delta\alpha - \gamma\beta}$$

Full DT yields are not required bin-by-bin yields are only required Sideband-subtraction done on each bin of the Dalitz plot



Dalitz Plot

 $K_S^0 K^+ K^-$ vs. *CP*-even tags



Dalitz Plot



Equal- $\Delta \delta_D$ binning scheme



Equal- $\Delta \delta_D$ binning scheme



Equal- $\Delta \delta_D$ binning scheme



Equal- $\Delta \delta_D$ binning scheme



Equal- $\Delta \delta_D$ binning scheme

Dalitz plot divided into bins satisfying condition

$$2\pi(i-3/_2)/\mathcal{N} \leq \Delta \delta_D(m_+^2,m_-^2) < 2\pi(i-1/_2)/\mathcal{N}$$
 for $i = 1, 2, 3 \dots \mathcal{N}$

Gain in statistical sensitivity by 30 % for Equal- $\Delta \delta_D$ compared to rectangular bins !

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- Modified optimal binning etc.

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Compare the Q^2 values EPJC 47, 347 (2006)

$$\mathcal{Q}^{2} = \frac{\sum_{i} \left[\left(\frac{1}{\sqrt{\Gamma_{i}}} \frac{d\Gamma_{i}}{dx} \right)^{2} + \left(\frac{1}{\sqrt{\Gamma_{i}}} \frac{d\Gamma_{i}}{dy} \right)^{2} \right]}{\int \left[\left(\frac{1}{\sqrt{|A_{B^{-}}|^{2}}} \frac{d|A_{B^{-}}|^{2}}{dx} \right)^{2} + \left(\frac{1}{\sqrt{|A_{B^{-}}|^{2}}} \frac{d|A_{B^{-}}|^{2}}{dy} \right)^{2} \right] dm_{+}^{2} dm_{-}^{2}}$$

 $\mathcal{N} = 2 \ (0.94^{+0.16}_{-0.06})$ $\mathcal{N} = 3 \ (0.87^{+0.14}_{-0.06})$ $\mathcal{N} = 4 \ (0.94^{+0.21}_{-0.06})$ for the binning used in this analysis

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Effects of Model based binning on γ

- No model-dependent systematics
- Wrong model effects sensitivity but no bias on result

Bin yields

Tag	Bin 1	Bin 2
$K_{\rm S}^0 K^+ K^-$ vs. $K^+ K^-$	6	36
$K_{\rm S}^{0}K^+K^-$ vs. $\pi^+\pi^-$	3	7
$K_{\rm S}^{0}K^{+}K^{-}$ vs. $K_{\rm S}^{0}\pi^{0}\pi^{0}$	2	5
$K_{\rm S}^0 K^+ K^-$ vs. $\pi^+ \pi^- \pi^0$	14	37
$K_{\rm S}^0 K^+ K^-$ vs. $K_{\rm L}^0 \pi^0$	25	67
$K_{ m S}^0 K^+ K^-$ vs. $K_{ m L}^0 \omega$	11	32
$K_{\rm S}^0 K^+ K^-$ vs. $K_{\rm L}^0 \eta(\gamma \gamma)$	5	14
$K_{\rm S}^0 K^+ K^-$ vs. $K_{\rm L}^0 \eta' (\pi^+ \pi^- \eta)$	3	4
$K_{\rm S}^{0}K^{+}K^{-}$ vs. $K_{\rm S}^{0}\pi^{0}$	30	7
$K^0_{ m S}K^+K^-$ vs. $K^0_{ m S}\omega$	15	0
$K^0_{\rm S}K^+K^-$ vs. $K^0_{\rm S}\eta(\gamma\gamma)$	7	2
$K_{\rm S}^0 K^+ K^-$ vs. $K_{\rm S}^0 \eta' (\pi^+ \pi^- \eta)$	2	0
$K_{\rm L}^0 K^+ K^-$ vs. $K^+ K^-$	95	17
$K_{\rm L}^0 K^+ K^-$ vs. $\pi^+ \pi^-$	27	4
$K_{\rm L}^0 K^+ K^-$ vs. $K_{\rm S}^0 \pi^0 \pi^0$	36	9
$K_{\rm L}^0 K^+ K^-$ vs. $\pi^+ \pi^- \pi^0$	197	57
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Migration correction

Finite detector resolution leads to migration of events

Migration matrix **U** constructed using signal MC sample

$$\mathbf{U}_{ij} = \frac{m_{ji}}{\sum_{k=-\mathcal{N}, k\neq 0}^{\mathcal{N}} m_{jk}}$$

 m_{ji} : Events generated in bin j and reconstructed in bin i

Vector of corrected yield ${\bf N},$ related to uncorrected yield ${\bf N}^{\rm rec}$ by

 $N = U^{-1}N^{rec}$

i	$U_{i,1}$	$U_{i,2}$	$U_{i,3}$	$U_{i,-1}$	$U_{i,-2}$	$U_{i,-3}$
1	0.968	0.020	0.001	0.011	0.000	0.000
2	0.036	0.967	0.001	0.000	0.001	0.003
3	0.007	0.001	0.992	0.000	0.000	0.000
-1	0.010	0.000	0.000	0.972	0.018	0.000
$^{-2}$	0.000	0.000	0.000	0.032	0.967	0.001
-3	0.000	0.000	0.000	0.006	0.006	63 0.988

Background Analysis



Background Analysis



Extraction of c_i and s_i

The uncorrected yields in related bins are combined according to the symmetry relations

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 $c_i^{(\prime)}$, $s_i^{(\prime)}$ obtained by minimizing the negative log likelihood expression

$$-2 \ln \mathcal{L} = -2 \sum_{i} \ln P \left(N_{i}^{\pm}, \langle N_{i}^{\pm} \rangle \right)_{K_{\mathrm{S}}^{0}K^{+}K^{-}, CP} -2 \sum_{i} \ln P \left(N_{i}^{\prime\pm}, \langle N_{i}^{\prime\pm} \rangle \right)_{K_{\mathrm{L}}^{0}K^{+}K^{-}, CP} -2 \sum_{i,j} \ln P \left(N_{ij}, \langle N_{ij} \rangle \right)_{K_{\mathrm{S}}^{0}K^{+}K^{-}, K_{\mathrm{S}}^{0}K^{+}K^{-}} -2 \sum_{i,j} \ln P \left(N_{ij}^{\prime}, \langle N_{ij}^{\prime} \rangle \right)_{K_{\mathrm{S}}^{0}K^{+}K^{-}, K_{\mathrm{L}}^{0}K^{+}K^{-}} -2 \sum_{i,j} \ln P \left(N_{ij}, \langle N_{ij} \rangle \right)_{K_{\mathrm{S}}^{0}K^{+}K^{-}, K_{\mathrm{L}}^{0}\pi^{+}\pi^{-}} -2 \sum_{i,j} \ln P \left(N_{ij}^{\prime}, \langle N_{ij}^{\prime} \rangle \right)_{K_{\mathrm{S}}^{0}K^{+}K^{-}, K_{\mathrm{L}}^{0}\pi^{+}\pi^{-}} -2 \sum_{i,j} \ln P \left(N_{ij}^{\prime}, \langle N_{ij}^{\prime} \rangle \right)_{K_{\mathrm{L}}^{0}K^{+}K^{-}, K_{\mathrm{S}}^{0}\pi^{+}\pi^{-}} +\chi^{2}.$$

 $\langle N \rangle$ is expected migration corrected yield

N = M + B $\langle N \rangle = \langle M \rangle + \langle B \rangle$

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$$\chi^{2} = \sum_{i} \left(\frac{c_{i}' - c_{i} - \Delta c_{i}}{\sigma_{\Delta c_{i}}} \right)^{2} + \sum_{i} \left(\frac{s_{i}' - s_{i} - \Delta s_{i}}{\sigma_{\Delta s_{i}}} \right)^{2}$$
$$\Delta c_{i} = c_{i,\text{BaBar}}' - c_{i,\text{BaBar}}$$

Systematic uncertainties

General strategy: smearing the input quantity by a Gaussian with in the measured uncertainty to produce a new value of input quantity.

Repeat fit 1000 times and build a distribution of $c_i^{(\prime)}$ and $s_i^{(\prime)}$ values; width of distribution gives systematic uncertainty

Accounting for correlation: vector of correlated variable

 $X = \mu + AZ$

A is Cholesky decomposition of covariance matrix, Z is vector of unit Gaussian

Systematic	c_1	<i>c</i> ₂	s_1	<i>s</i> ₂	c'_1	c'_2	s'_1	s'_2
ST yield	0.002	0.003	0.000	0.000	0.002	0.001	0.000	0.000
$K_i^{(')}$ statistics	0.000	0.003	0.004	0.005	0.001	0.002	0.004	0.004
$K^0 \pi^+ \pi^- (c_i^{(\prime)}, s_i^{(\prime)})$	0.001	0.002	0.037	0.075	0.001	0.002	0.037	0.076
$K^0 \pi^+ \pi^- (K_i^{(\prime)})$	0.000	0.001	0.007	0.033	0.000	0.000	0.006	0.033
$N_{D\bar{D}}$	0.000	0.000	0.001	0.001	0.000	0.000	0.001	0.001
MC statistics	0.001	0.002	0.000	0.000	0.001	0.003	0.000	0.000
Background	0.002	0.003	0.011	0.022	0.002	0.003	0.011	0.022
DCS correction	0.000	0.000	0.000	0.001	0.000	0.000	0.000	0.001
Stat	0.034	0.040	0.144	0.198	0.035	0.034	0.144	0.198
Syst total	0.003	0.007	0.039	0.085	0.003	0.006	0.039	0.086
Total	0.034	0.041	0.149	0.215	0.035	0.035	0.149	0.216



Systematic uncertainty for $\mathcal{N}=$ 2 bins 69

Fit results



Impact on measurement of γ

Simulate $B^{\pm} \rightarrow D(K_S^0 K^+ K^-) K^{\pm}$ decay in each bins

Smear values of c_i and s_i within their uncertainties accounting for correlation

10000 ToyMC sample generated

$$N_i^{\mp} = \frac{a_B}{a_D} (K_{\pm i} + r_B^2 K_{\mp i} + 2\sqrt{K_i K_{-i}} (x_{B\mp} c_i + y_{B\mp} s_i))$$

 a_B set to a larger value

Values of r_B , δ_B and γ set to http://ckmfitter.in2p3.fr latest world averages



Summary

Explained the model-independent BPGGSZ method for measuring CKM angle γ Role of charm factories in the model independent measurement of γ explained $\Box c_i$ and s_i measured with a better precision to date

Generation Measurements are statistically dominated no irreducible systematic uncertainty Uncertainty on γ due to c_i and s_i are 2.4°, 1.3°, 1.3° for $\mathcal{N} = 2, 3, 4$ Equal- $\Delta \delta_D$ bins

Phys. Rev. D 102, 052008 (2020) [arXiv: hep-ex/2007.07959]

Measurements of γ using $B^{\pm} \rightarrow D(K_S^0 h^+ h^-) K^{\pm}$ at LHCb (2020) JHEP 02, 169 (2021) Uncertainty due to c_i and $s_i \sim 1^\circ \rightarrow 50\%$ improvement with BESIII value 9 $fb^{-1} P\bar{P}$ data @ 7,8, 13 TeV

1900 $K_S^0 K^+ K^-$ events

$$\mathcal{N}=2$$


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 $\mathcal{N}=2$

Similar analysis using Belle+BelleII data under progress!



Other areas:

The value of c_i and s_i required for model-independent determination of charm mixing parameters and for searching CP violation in $D \rightarrow K_S^0 K^+ K^-$ decays

Phys.Rev.Lett 122, 231802 (2019) [arXiv:hep-ex/1903.03074]

New $D \rightarrow K_S^0 K^+ K^-$ model at BESIII

arXiv:hep-ex/2006.02800

Other activities at BESIII

- Involved in the similar measurements of c_i and s_i for $D \to K_S^0 \pi^+ \pi^-$ with groups at Oxford University
- Quantum-correlated studies of $D \rightarrow K_S^0 \pi^+ \pi^- \pi^0$ decays (undergraduate project of Mr. Pratyush Anand, Master student at ETH Zurich)

TOF Calibration work at BESIII

Plastic based scintillator was replaced by MRPC in 2015 ETOF has 72 modules, each modules has 8 MRPC strips Need to calibrate each strip



BACKUP

Fit results: First uncertainty is statistical and second uncertainty is systematic

\mathcal{N}	i	c_i	s_i	c_i'	s_i'
2	1	$0.704 \pm 0.034 \pm 0.003$	$-0.038 \pm 0.144 \pm 0.039$	$0.730 \pm 0.035 \pm 0.003$	$-0.028 \pm 0.144 \pm 0.039$
	2	$-0.760 \pm 0.040 \pm 0.007$	$0.590 \pm 0.198 \pm 0.085$	$-0.785 \pm 0.034 \pm 0.006$	$0.669 \pm 0.198 \pm 0.086$
3	1	$0.724 \pm 0.035 \pm 0.003$	$-0.037 \pm 0.174 \pm 0.049$	$0.751 \pm 0.036 \pm 0.003$	$-0.037 \pm 0.174 \pm 0.049$
	2	$-0.576 \pm 0.050 \pm 0.009$	$0.616 \pm 0.146 \pm 0.047$	$-0.512\pm0.050\pm0.009$	$0.640 \pm 0.146 \pm 0.047$
	3	$-0.174 \pm 0.173 \pm 0.040$	$-0.669 \pm 0.370 \pm 0.119$	$-0.382\pm0.145\pm0.040$	$0.045 \pm 0.384 \pm 0.116$
4	1	$0.783 \pm 0.034 \pm 0.003$	$-0.242\pm0.173\pm0.051$	$0.802 \pm 0.034 \pm 0.003$	$-0.239 \pm 0.174 \pm 0.051$
	2	$-0.053\pm0.106\pm0.017$	$0.306 \pm 0.294 \pm 0.125$	$0.070 \pm 0.106 \pm 0.017$	$0.286 \pm 0.294 \pm 0.124$
	3	$-0.654 \pm 0.057 \pm 0.011$	$0.659 \pm 0.210 \pm 0.059$	$-0.630 \pm 0.056 \pm 0.011$	$0.709 \pm 0.210 \pm 0.059$
	4	$0.090 \pm 0.208 \pm 0.041$	$-0.713 \pm 0.387 \pm 0.195$	$-0.290 \pm 0.201 \pm 0.036$	$0.122 \pm 0.422 \pm 0.206$

Average of CLEO and BESIII results

CLEO and BESIII results are compatible: Weighted average will be equally good

Just add a term to fit

$$\chi^2_{\text{avg}} = (\mathbf{P} - \mathbf{P}^{\text{CLEO}})^T V^{-1} (\mathbf{P} - \mathbf{P}^{\text{CLEO}})$$

Average of CLEO and BESIII results, uncertainties statistical and systematic added in quadratures

\mathcal{N}	i	c_i	s_i	c_i'	s_i'
2	1	0.713 ± 0.032	0.107 ± 0.132	0.737 ± 0.032	0.116 ± 0.132
	2	-0.758 ± 0.037	0.394 ± 0.173	-0.782 ± 0.033	0.473 ± 0.174
3	1	0.738 ± 0.030	0.112 ± 0.102	0.765 ± 0.030	0.111 ± 0.102
	2	-0.573 ± 0.044	0.550 ± 0.113	-0.503 ± 0.044	0.574 ± 0.113
	3	-0.129 ± 0.155	-0.619 ± 0.317	-0.412 ± 0.138	0.089 ± 0.327
4	1	0.796 ± 0.030	-0.082 ± 0.173	0.817 ± 0.030	-0.080 ± 0.173
	2	-0.018 ± 0.099	0.393 ± 0.262	0.105 ± 0.098	0.375 ± 0.261
	3	-0.691 ± 0.048	0.551 ± 0.200	-0.657 ± 0.048	0.601 ± 0.200
	4	0.183 ± 0.182	-0.646 ± 0.415	-0.321 ± 0.185	0.218 ± 0.438

Detector described using cylindrical coordinates with IP as the origin

Charged Particle

- Polar angle: |cos (θ)| < 0.93
- Radial distance: $|V_{xy}| < 1$ cm
- Longitudinal distance: |V_Z| < 10 cm

Neutral showers

- 0 < time < 700 ns
- Energy: > 0.025 GeV
 (> 0.050 GeV) for barrel (end-cap)
- Distance from MDC track exit $> 10\sigma$

Particle identification

•
$$\pi^{\pm}$$
: $\mathcal{L}_{\pi} > \mathcal{L}_{K}$

•
$$K^{\pm}$$
: $\mathcal{L}_K > \mathcal{L}_{\pi}$
• e^{\pm} : $\frac{\mathcal{L}_e}{\mathcal{L}_{\pi} + \mathcal{L}_K + \mathcal{L}_e} > 0.8$

 $\pi^0
ightarrow \gamma \gamma (\eta
ightarrow \gamma \gamma)$ reconstruction

- Kinematic fit: $\chi^2 < 20$
- 0.110 < M(γγ) < 0.165 GeV
 (0.480 < M(γγ) < 0.580 GeV)
- No. of γ in end-cap > 1

K_S^0 reconstruction

- Kinematic fit: $\chi^2 < 100$
- 0.487 < M(γγ) < 0.511 GeV
- Flight significance: $L/\sigma_L > 2$
- No track quality requirement

 $\omega \rightarrow \pi^+ \pi^- \pi^0$ and

$$\eta' \rightarrow \pi^+\pi^-\eta$$

- 0.760 < M($\pi^+\pi^-\pi^0$) < 0.805 (GeV)
- 0.938 < M(π⁺π⁻η) < 0.978
 GeV

Particles are combined to reconstruct final states

□ Fully reconstructed tags: Full kinematic reconstruction possible,

Partially reconstructed tags: Particles having missing particle in final states; no full reconstruction, inference using missing energy and momentum

Cabbibo-Kobayashi-Maskawa (CKM) matrix

Mixing between weak eigenstates and flavor eigenstates in three generations.

$$\begin{pmatrix} d'\\ s'\\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub}\\ V_{cd} & V_{cs} & V_{cb}\\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d\\ s\\ b \end{pmatrix}$$

$$V_{CKM}V_{CKM}^{\dagger} = I = V_{CKM}^{\dagger}V_{CKM}$$

Can be parameterized using 3 angles and 1 phase (rephasing quark fields)

Various parameterizations (Chau and Keung parameterization)

$$\begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}s_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

 $C_{ij}(S_{ij}) = \cos \theta_{ij}(\sin \theta_{ij})$ δ is the irreducible phase

Wolfenstein parameterization

$$V_{\rm CKM} \approx \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

$$\begin{aligned} s_{12} &\equiv \lambda & \lambda = 0.22 \\ s_{23} &\equiv A\lambda^2 \\ s_{13}e^{-i\delta} &\equiv A\lambda^3(\rho - i\eta) \end{aligned}$$