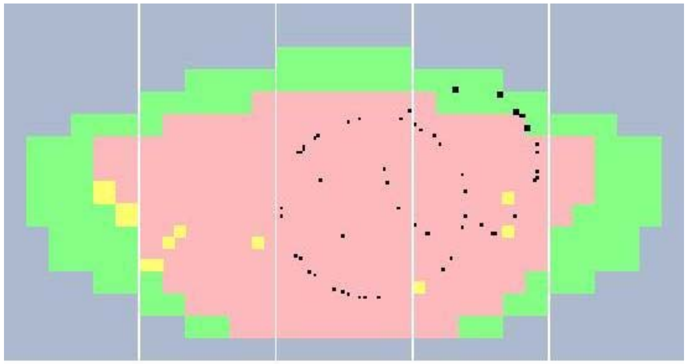
EDIT-2020: Excellence in Detectors and Instrumentation Technologies
17-28 February 2020 | DESY, Hamburg



Particle Identification

Peter Križan

University of Ljubljana and J. Stefan Institute

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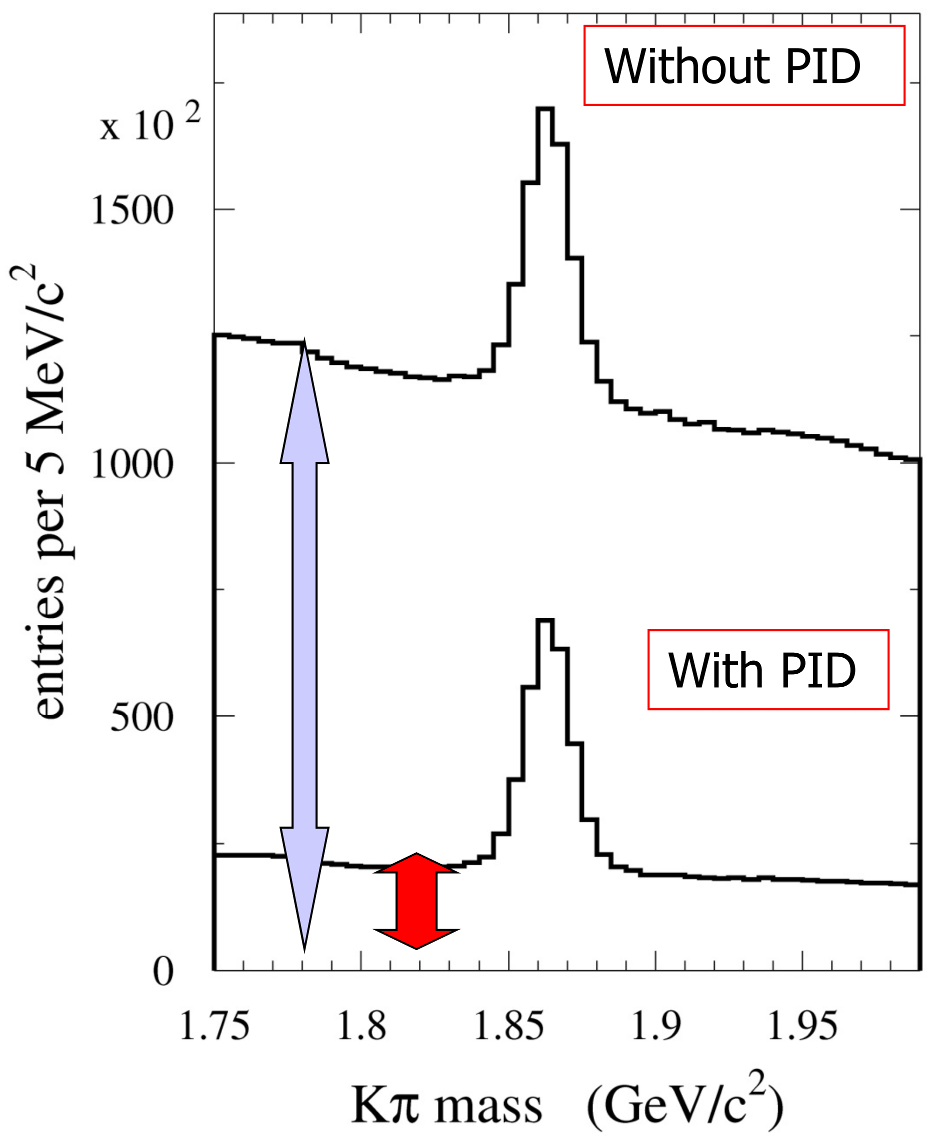
Why particle ID?

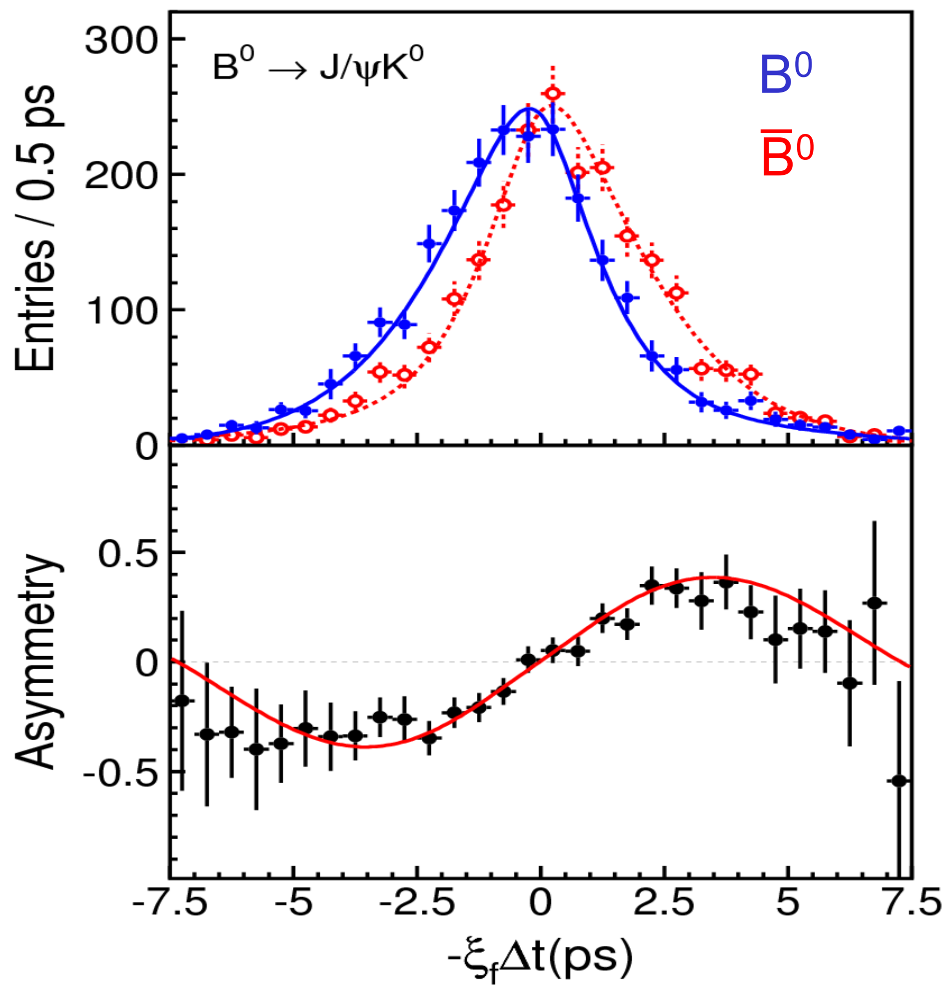


Example 1: BaBar (B factory)

Particle identification reduces the fraction of wrong Kπ combinations (combinatorial background) by ~5x

Why particle ID?

Particle identification at B factories (Belle and BaBar): was essential for the observation of CP violation in the B meson system.

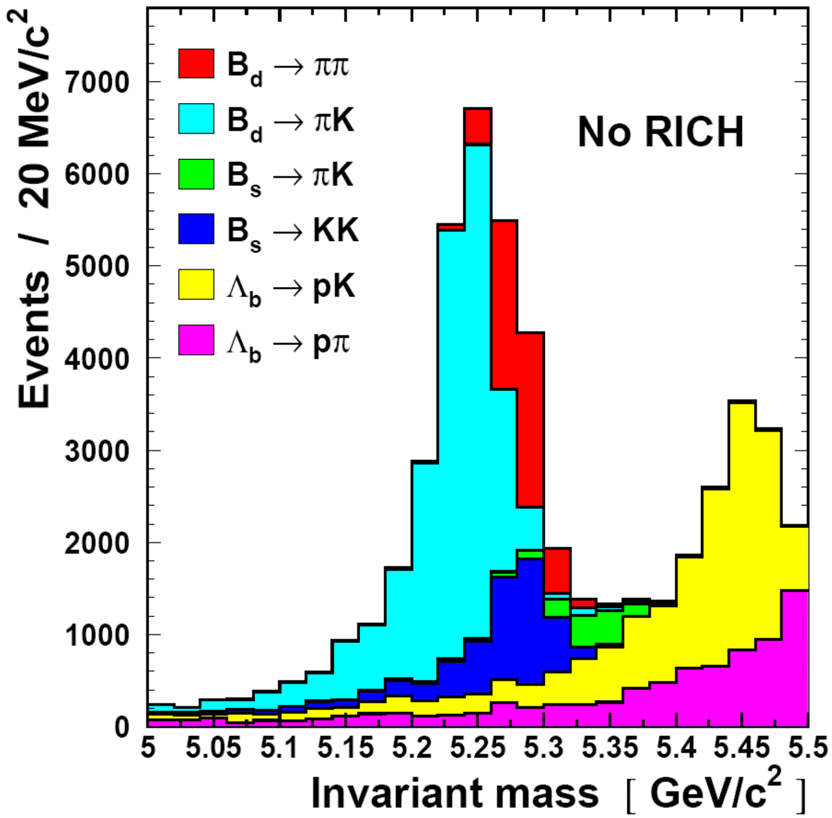


B^0 and its anti-particle decay differently to the same final state $J/\psi\ K^0$

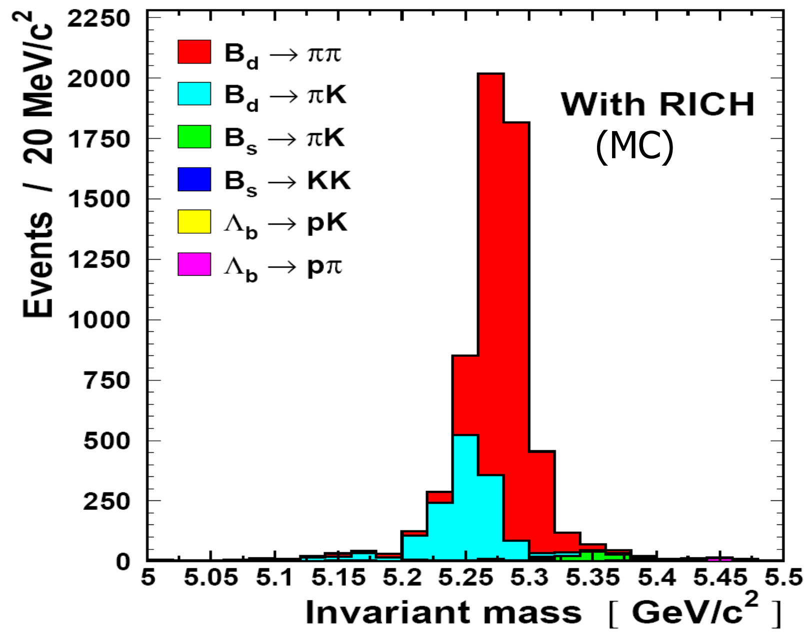
Flavour of the B: from decay products of the other B: charge of the kaon, electron, muon

→particle ID is compulsory

Why particle ID?

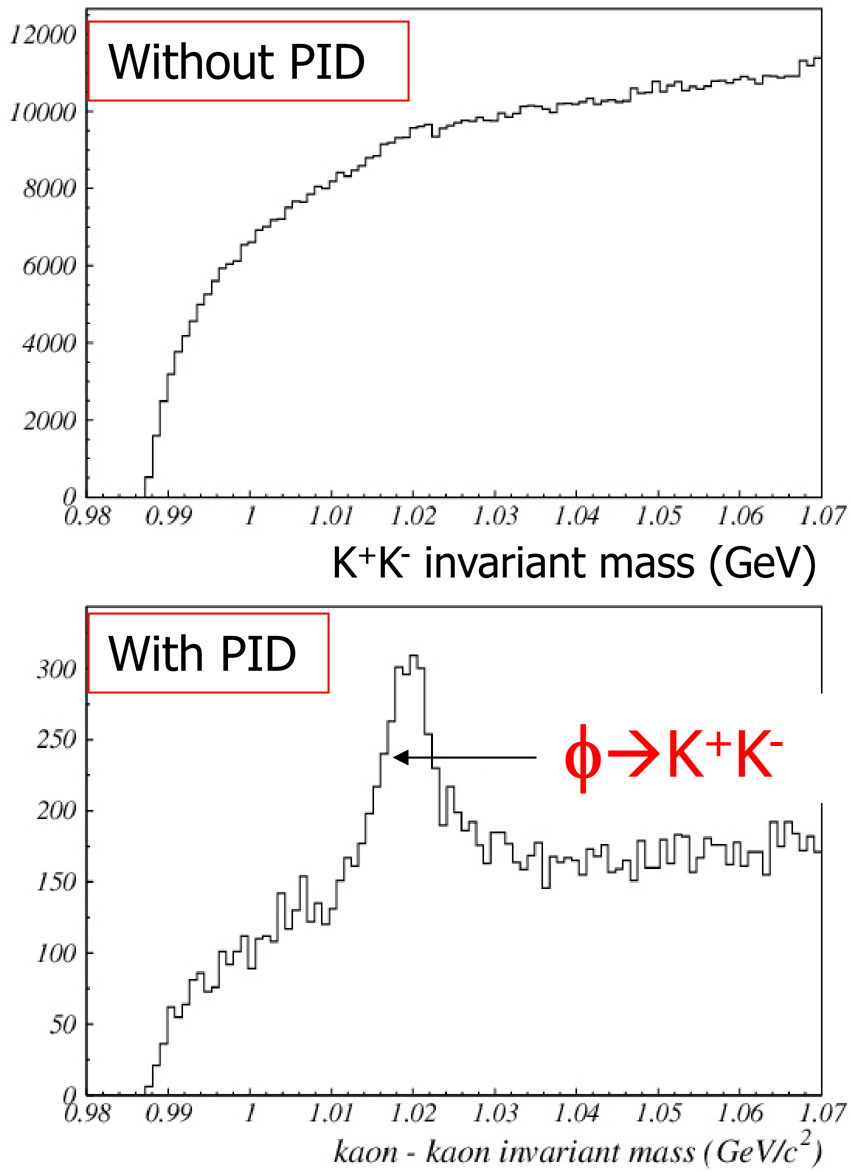


Example 2: LHCb



Need to distinguish $B_d \to \pi\pi$ from other similar topology 2-body decays and to distinguish B from anti-B using K tag.

Why particle ID?



Example 3: HERA-B

K^+K^- invariant mass.

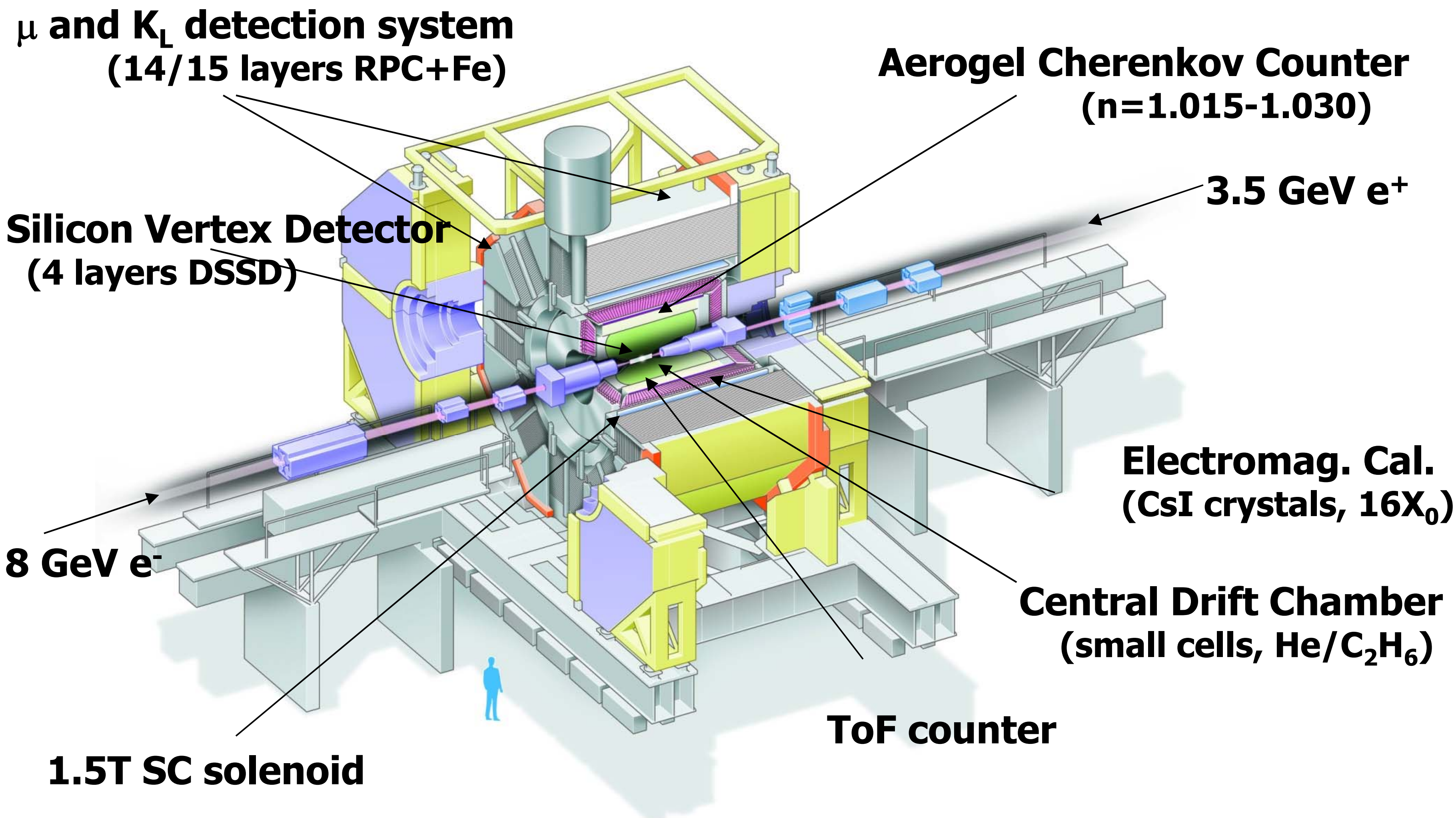
The inclusive $\phi \rightarrow K^+K^-$ decay only becomes visible after particle identification is taken into account.

Why particle ID?

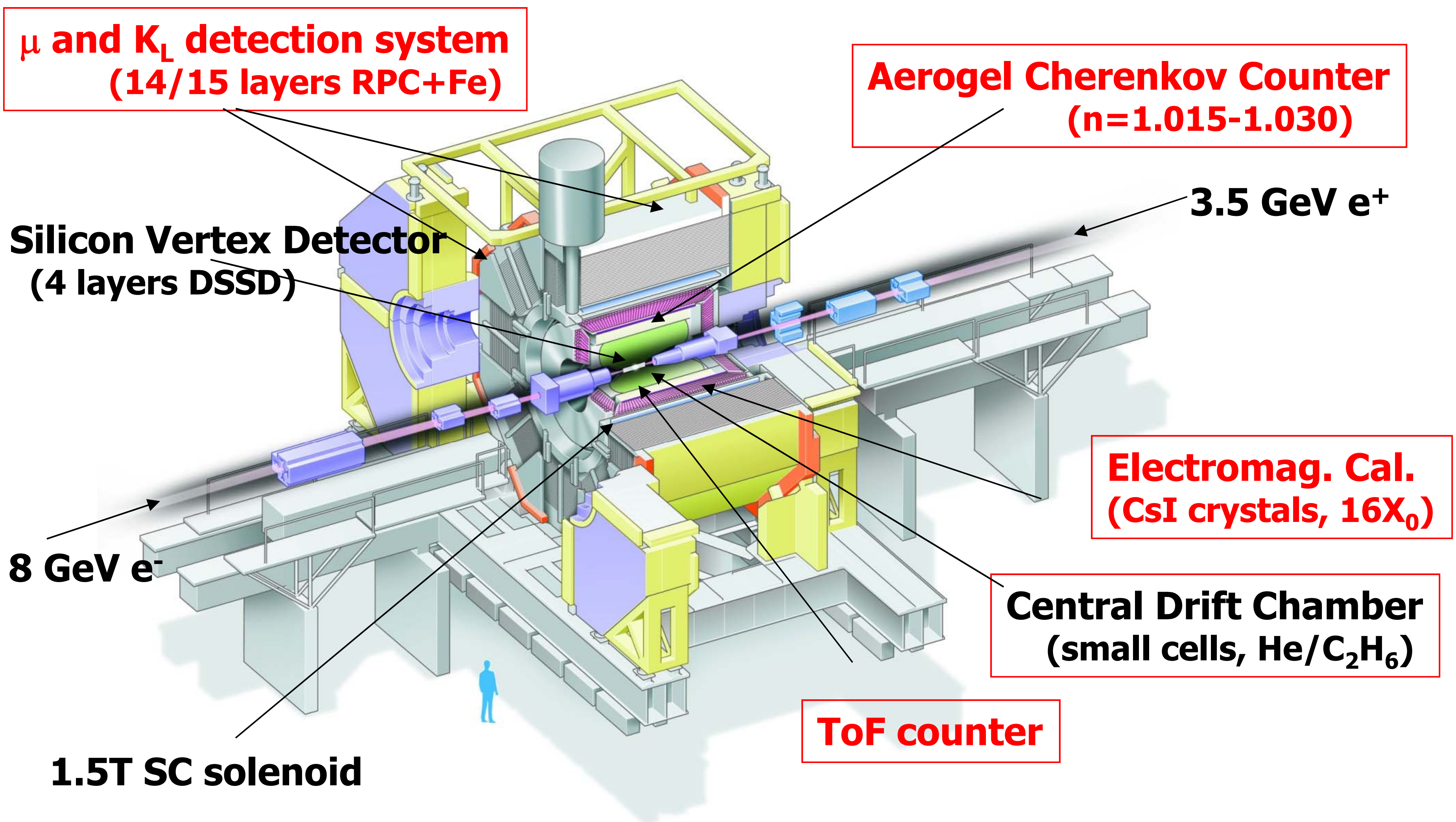
PID is also needed in:

- General purpose LHC experiments: final states with electrons and muons
- Searches for exotic states of matter (quark-gluon plasma)
- Spectroscopy and searches for exotic hadronic states
- Studies of fragmentation functions

Example: Belle

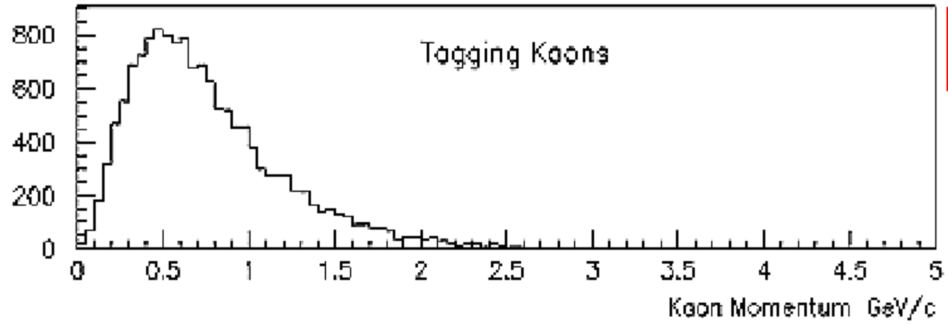


Particle identification systems in Belle



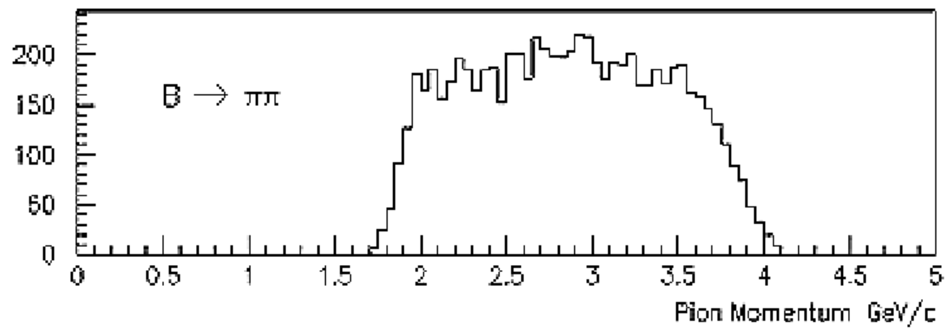
Particle identification methods depend on the requirements (physics channel, kinematics)

Example: B factory, pion/kaon separation



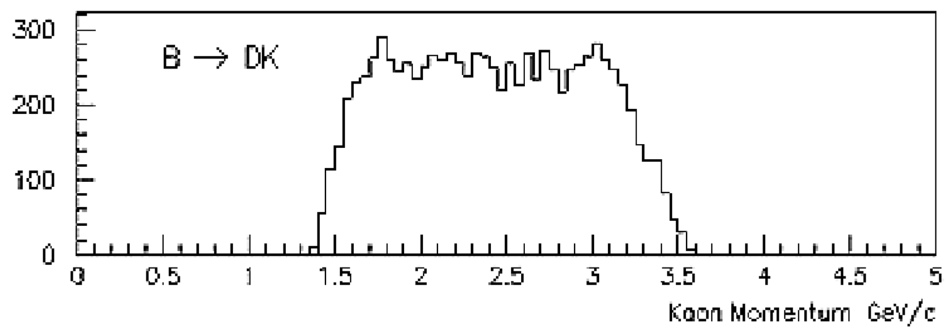
Tagging Kaons

Relatively soft, ms dominated for tracking



$B \to \pi\pi$

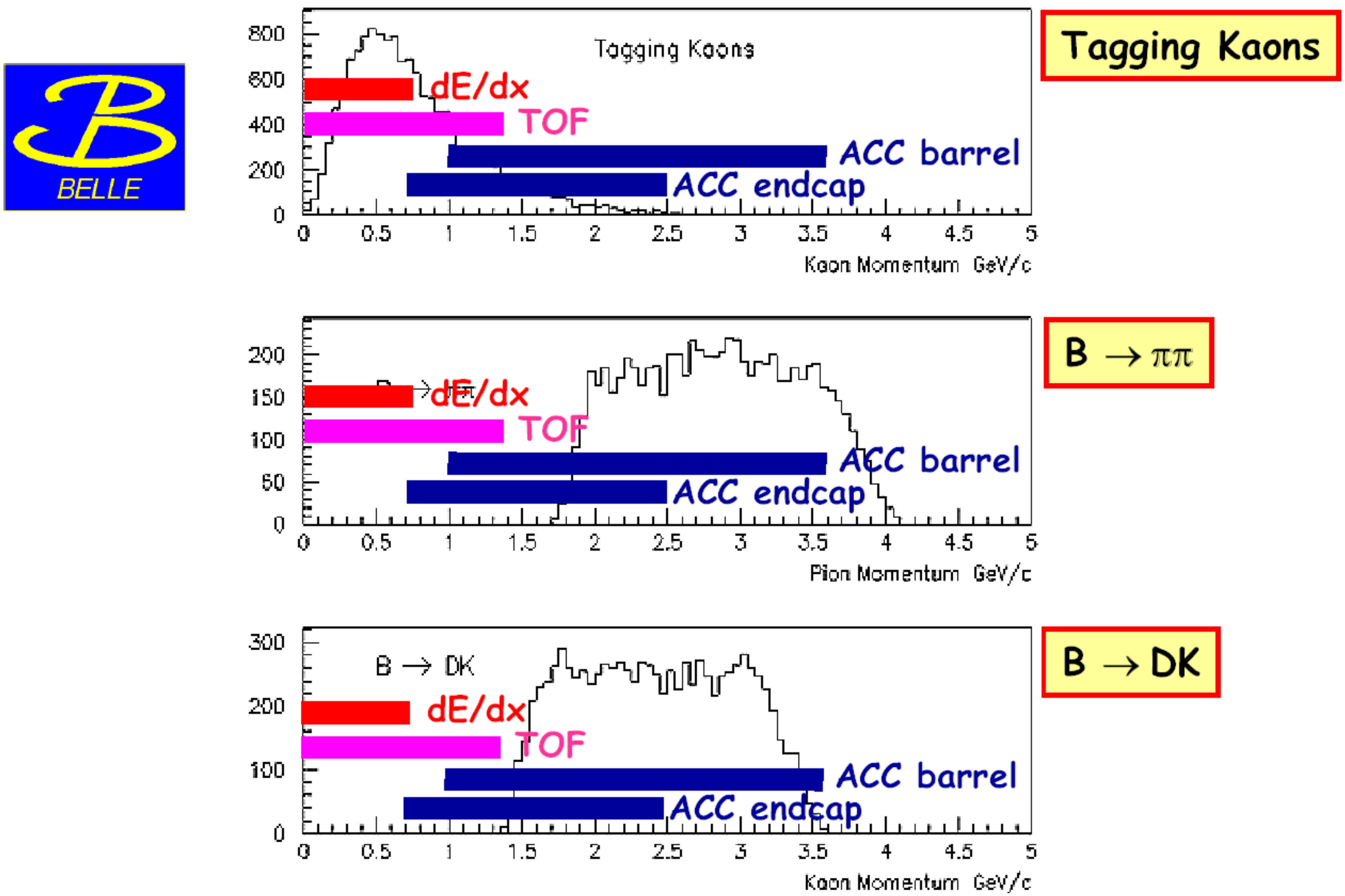
Requires dedicated PID



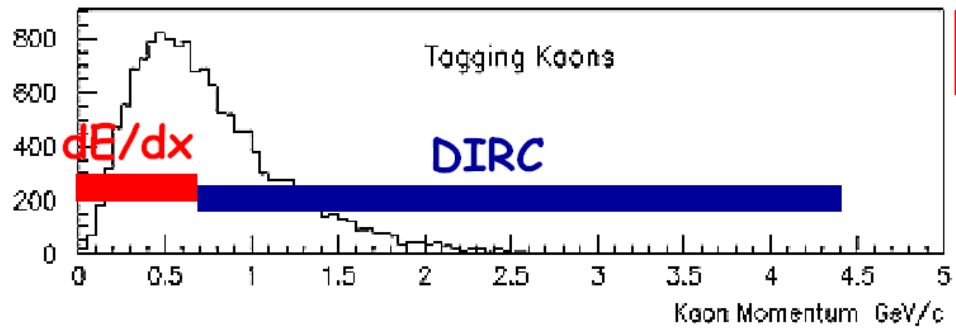
$B \to DK$

Requires dedicated PID

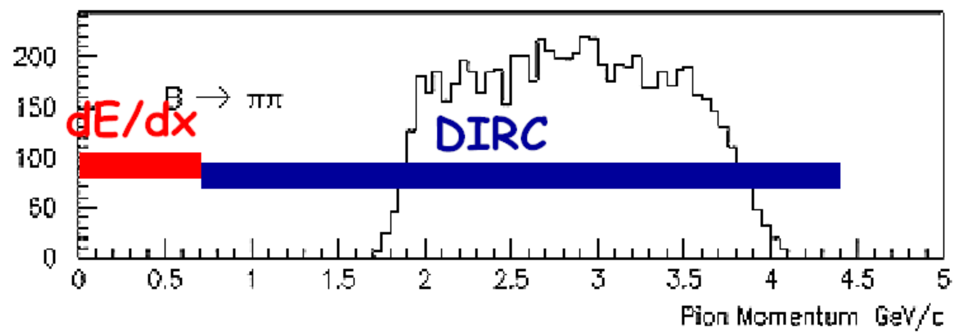
PID coverage of kaon/pion spectra in Belle



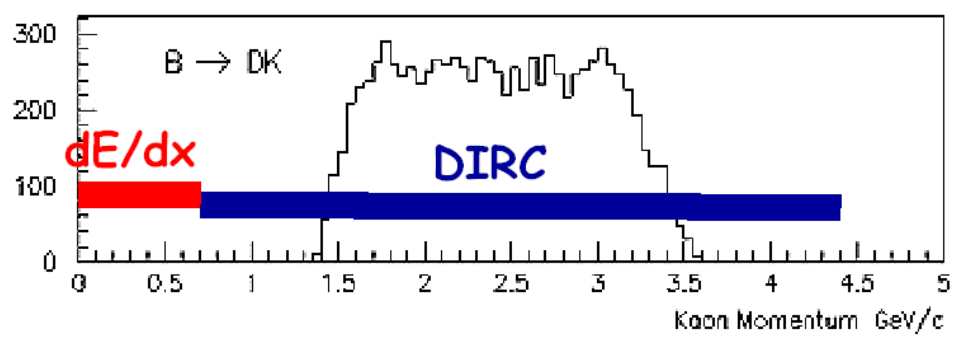
PID coverage of kaon/pion spectra in BaBar



Tagging Kaons



$B \to \pi\pi$



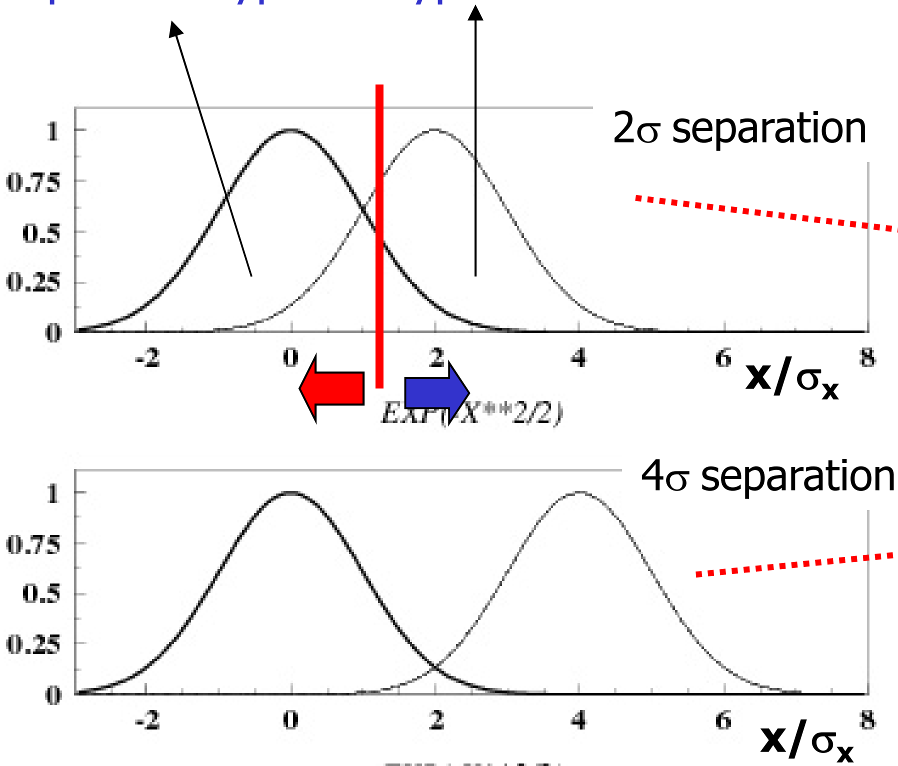
$B \to DK$

Efficiency and purity in particle identification

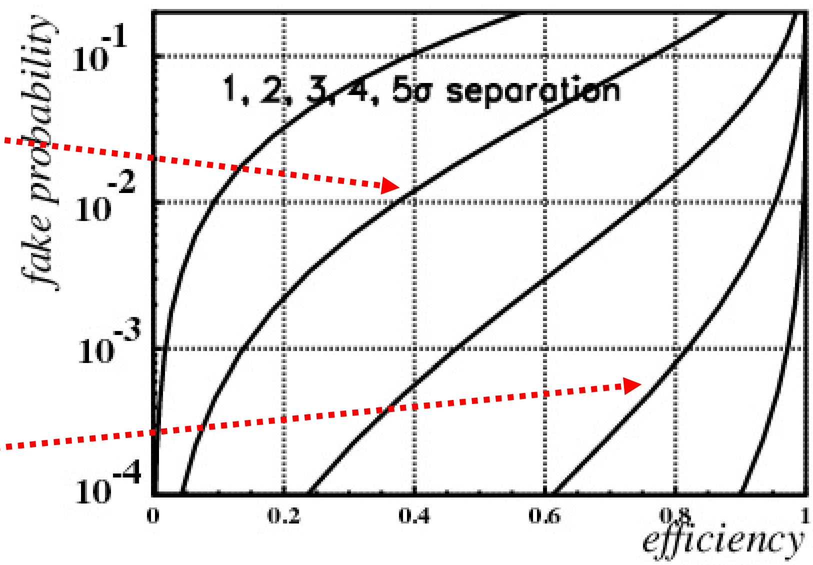
Efficiency and purity are tightly coupled!

Two examples:

particle type 1 type 2



efficiency vs fake probability (for Gaussian distributions)



some discriminating variable **x**, scaled to the resolution σ_x

Identification of charged particles

Particles (e, μ, π, K, p) in the final state are identified by their mass or by the way they interact.

Determination of mass: from the relation between momentum and velocity, $p=\gamma mv$ (p is known - radius of curvature in the magnetic field)

→Measure velocity by:

- time of flight
- ionisation losses dE/dx
- Cherenkov photon angle (and/or yield)
- transition radiation

Mainly used for the identification of hadrons.

Identification through interaction: electrons and muons

- muon systems
- calorimeters

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Identification through interaction: electrons and muons

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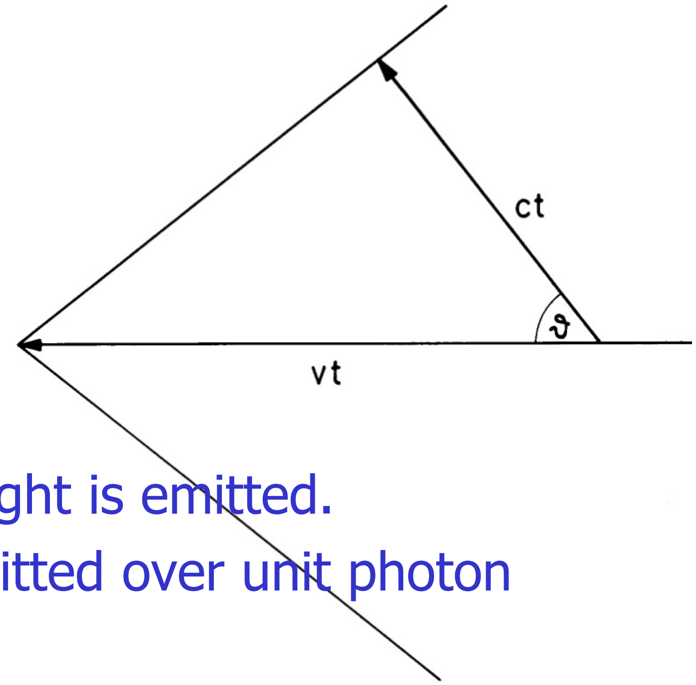
Cherenkov radiation

A charged track with velocity v=βc exceeding the speed of light c/n in a medium with refractive index n emits polarized light at a characteristic (Cherenkov) angle,

cosθ = c/nv = 1/βn



Two cases:

→ $\beta < \beta_t$ = 1/n: below threshold no Cherenkov light is emitted.

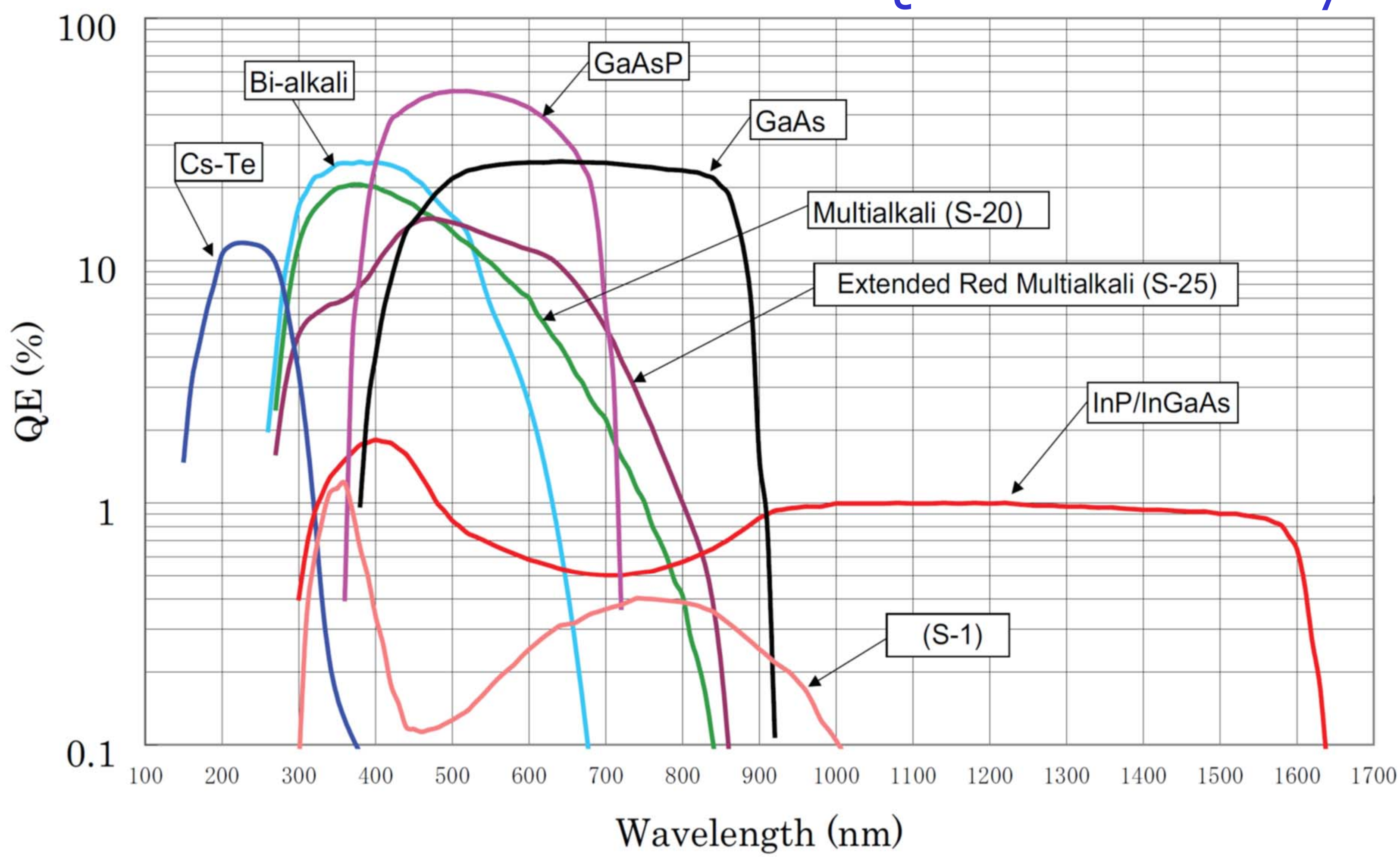
→ $\beta > \beta_t$: the number of Cherenkov photons emitted over unit photon energy E=hν in a radiator of length ***L***:

$$\frac{dN}{dE} = \frac{\alpha}{\hbar c} L \sin^2\theta = 370(cm)^{-1}(eV)^{-1} L \sin^2\theta$$

→Few detected photons

Photon detection efficiency

Quantum efficiency



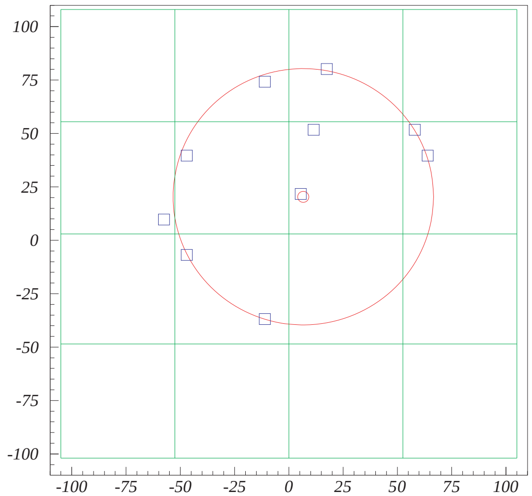
Number of detected photons

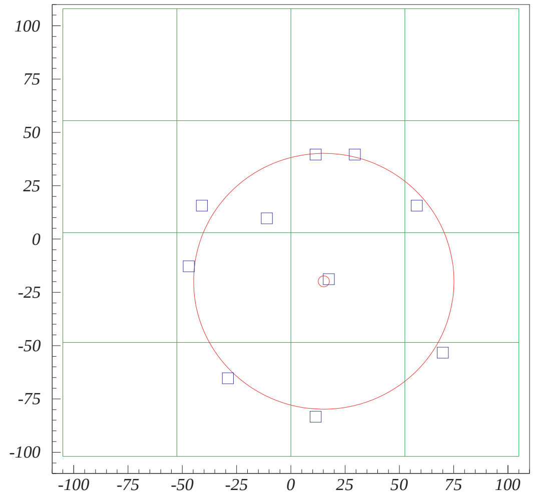
Example: in 1m of air (n=1.00027) a track with β=1 emits N=41 photons in the spectral range of visible light (ΔE ~ 2 eV).

If Čerenkov photons were detected with an average detection efficiency of ε=0.1 over this interval, N=4 photons would be measured.

Few photons detected

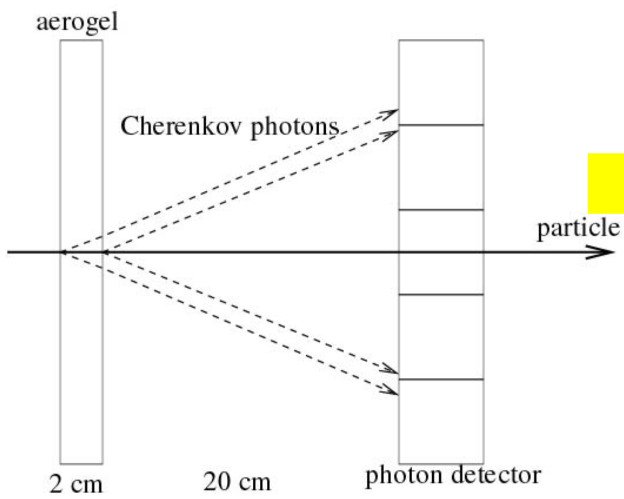
→Important to have a low noise detector

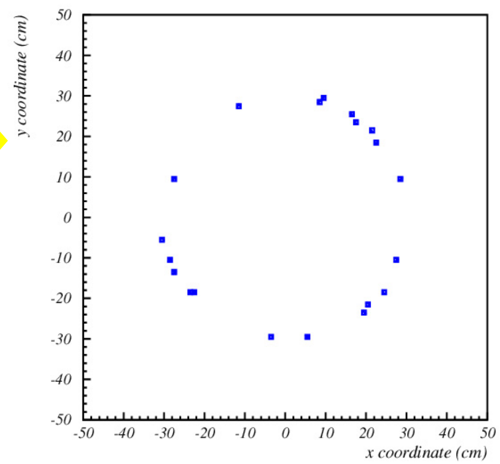




Measuring the Cherenkov angle

Particles above threshold: measure θ



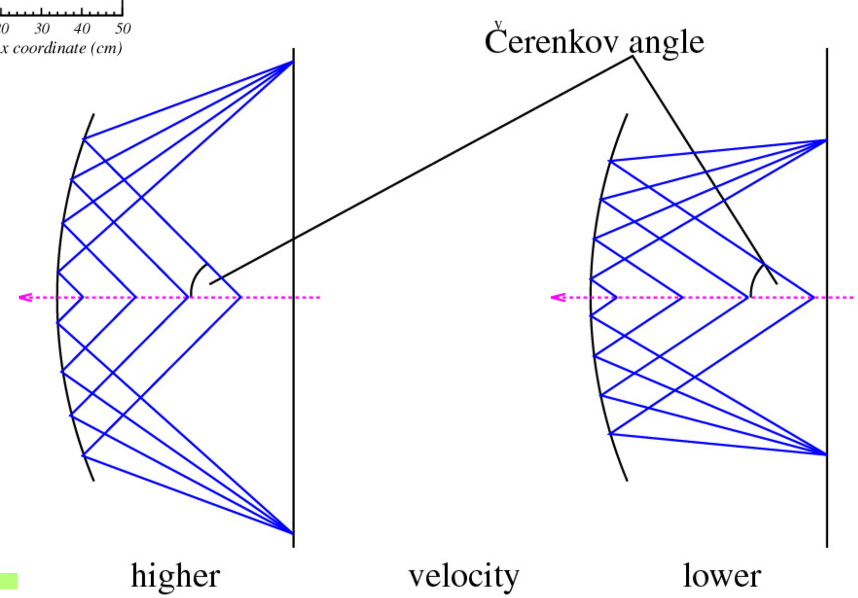


Idea: transform the direction into a coordinate → ring on the detection plane

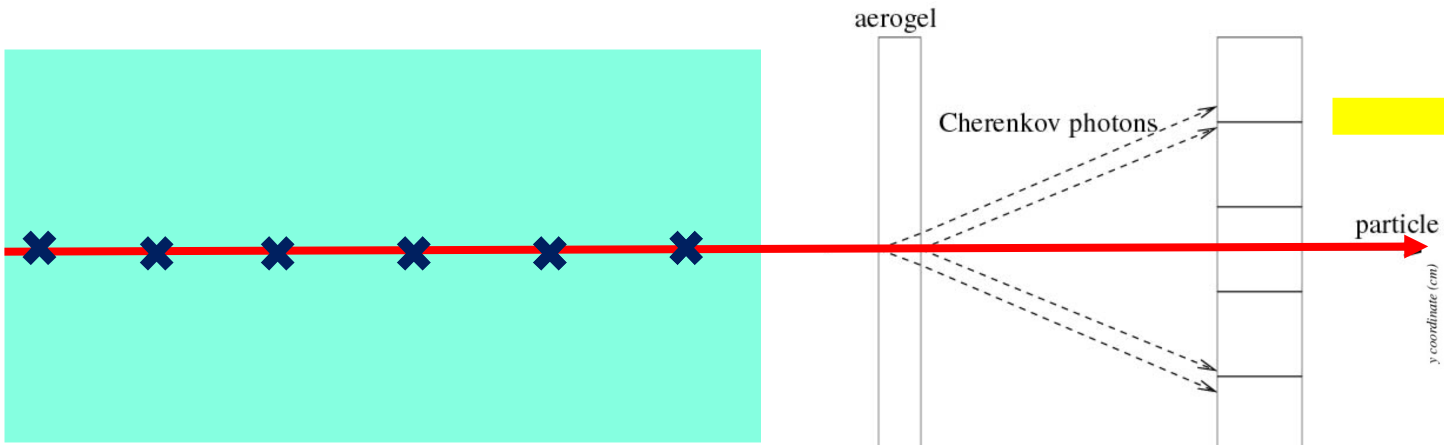
→ Ring Imaging Cherenkov (RICH) counter

Proximity focusing RICH

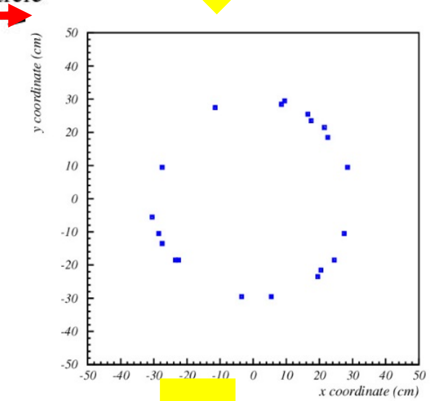
RICH with a focusing mirror

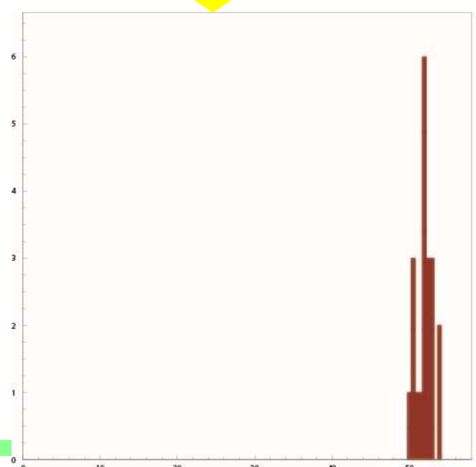


Measuring the Cherenkov angle



Tracking system (measures the direction of the particle)



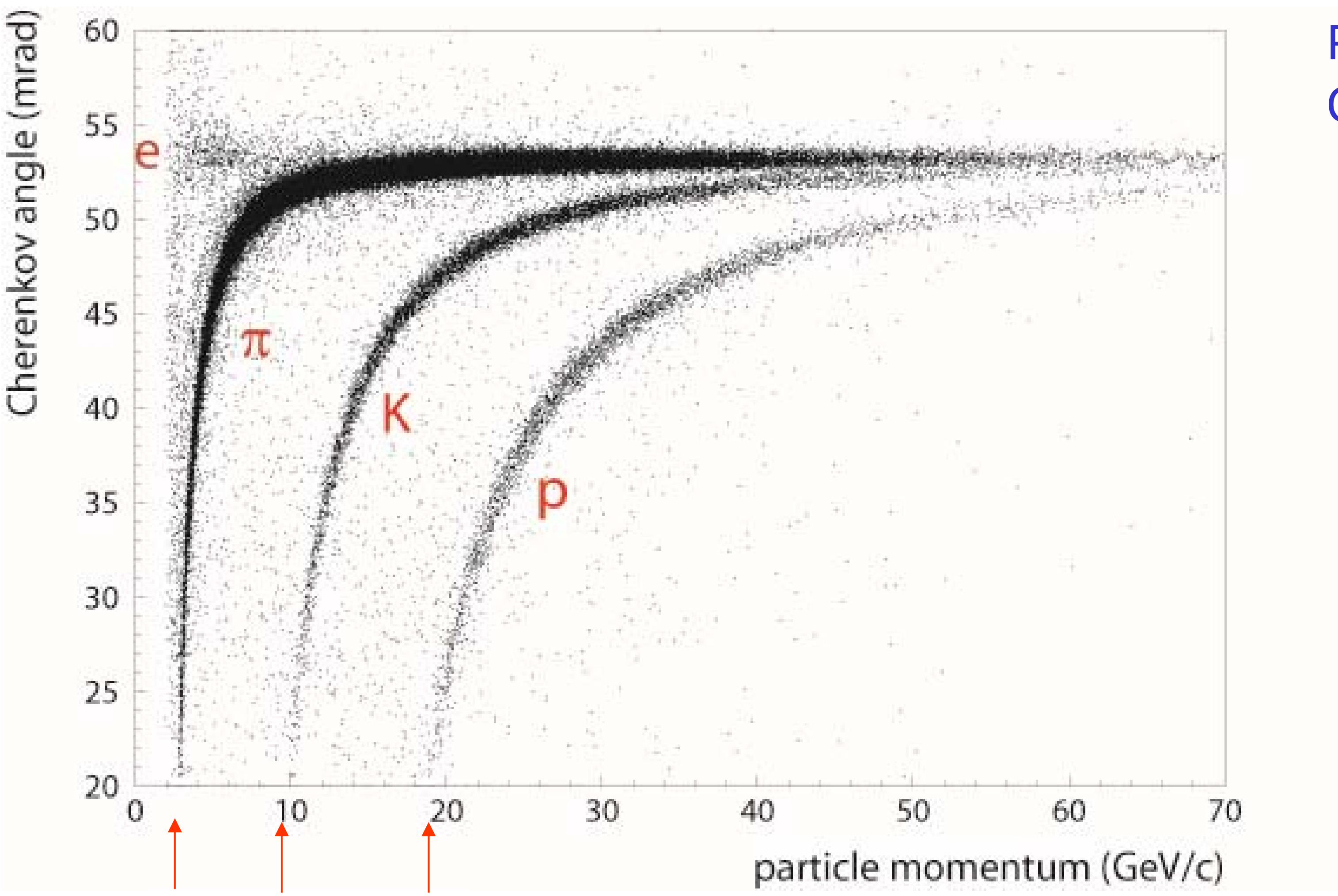


Cherekov angle distribution (mradian)

Tracking system tells us where the particle hit the radiator, and at which angle.

Use this information to calculate the Cherenkov angle for each individual detected photon

Measuring Cherenkov angle

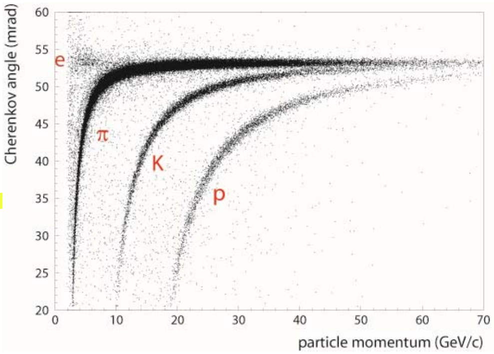


π K p

thresholds

Radiator:
C_4F_{10} gas

Resolution of a RICH counter



Determined by:

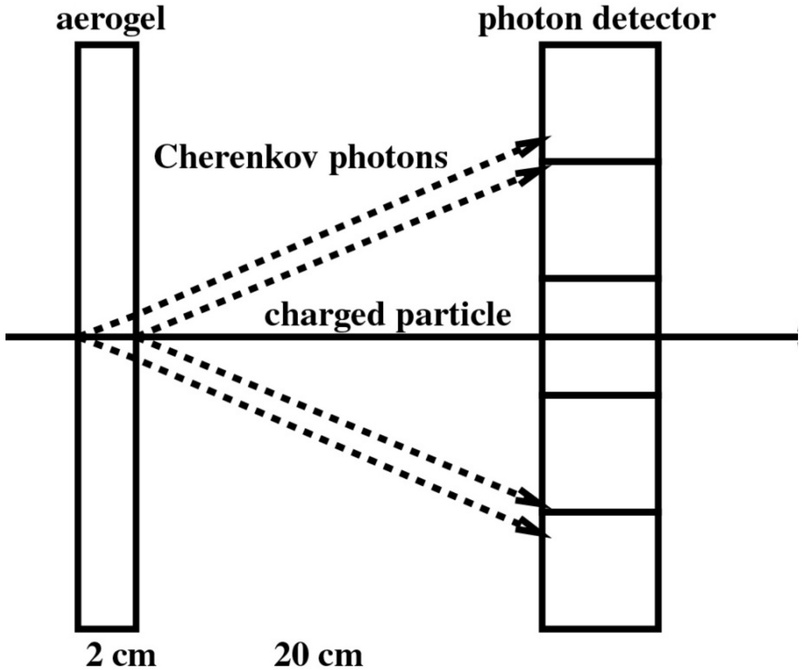
•Photon impact point resolution (~photon detector granularity)

•Emission point uncertainty (not in a focusing RICH)

•Dispersion: $1/\beta = n(\lambda) \cos\theta$

•Errors of the optical system

•Uncertainty in track parameters



Resolution per track:

$$\sigma_{track} = \frac{\sigma_0}{\sqrt{N_{pe}}}$$

single photon resolution

# of detected photons

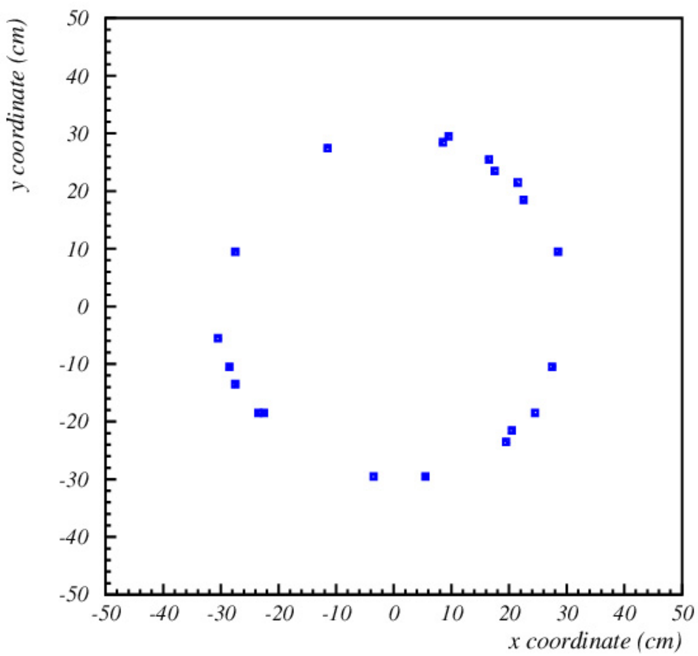
(in the case of low background)

Photon detection in RICH counters

RICH counter: measure photon impact point on the photon detector surface

→ detection of single photons with

- sufficient spatial resolution
- high efficiency and good signal-to-noise ratio (few photons!)
- over a large area (square meters)



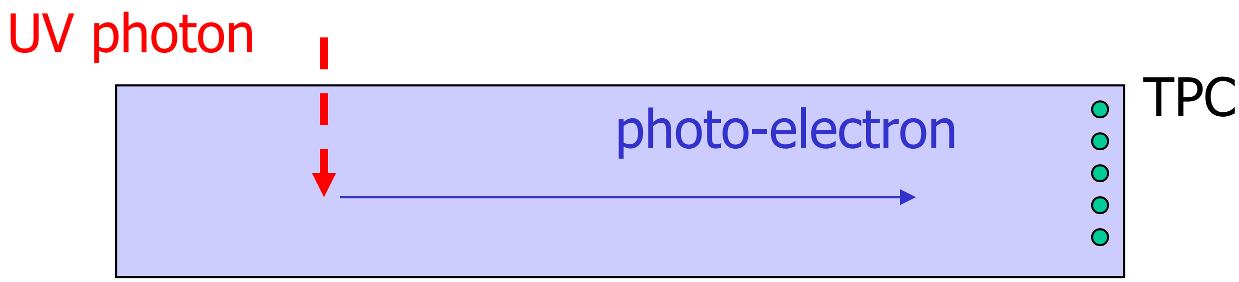
Special requirements:

- Operation in magnetic field
- High rate capability
- Very high spatial resolution
- Excellent timing (time-of-arrival information)

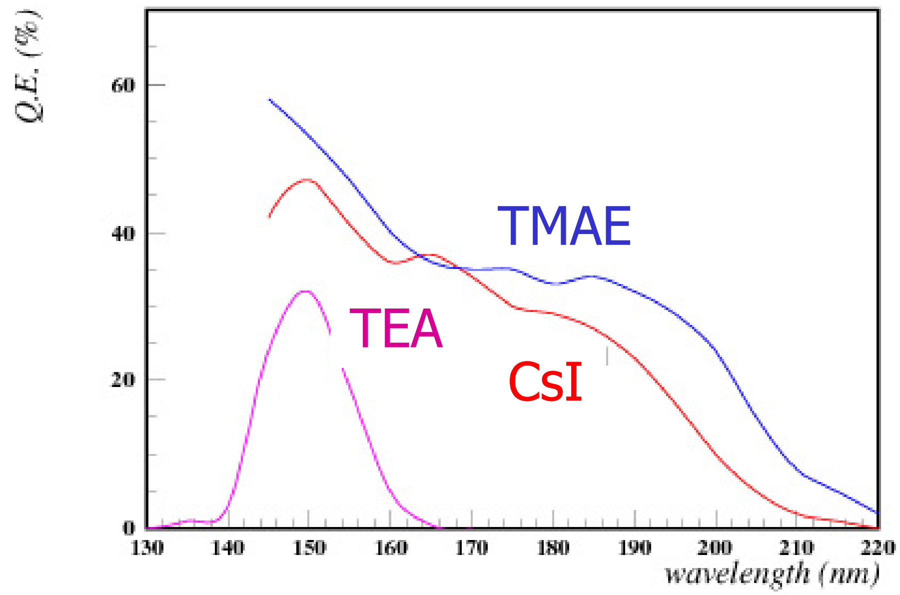
Photon detector is the most crucial element of a RICH counter

First generation of RICH counters

DELPHI, SLD, OMEGA RICH counters: all employed wire chamber based photon detectors (UV photon → photo-electron → detection of a single electron in a TPC)

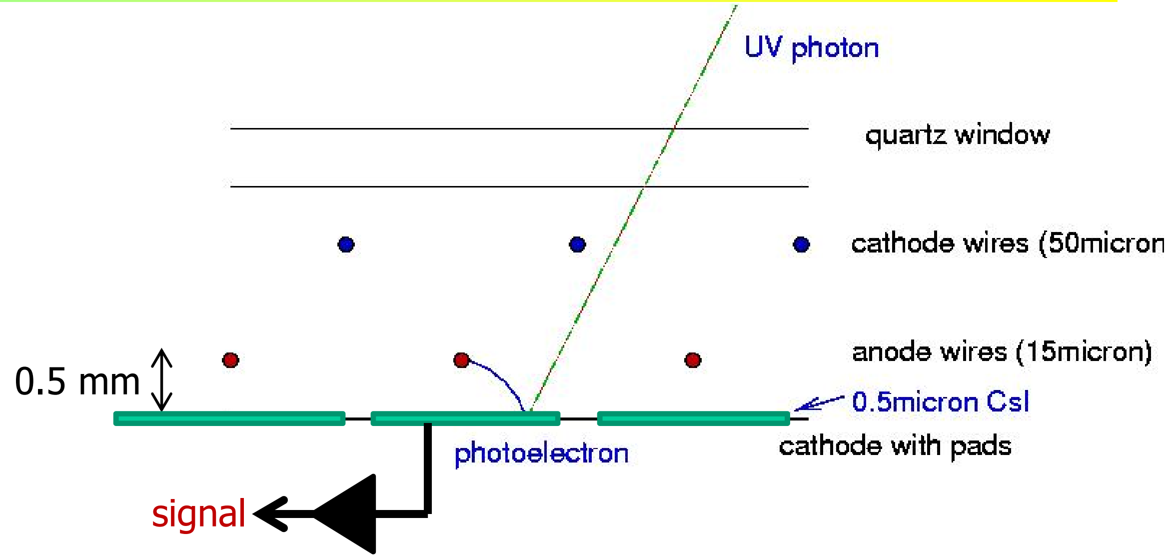


Photosensitive component:
TMAE added to the gas mixture



Fast RICH counters with wire chambers

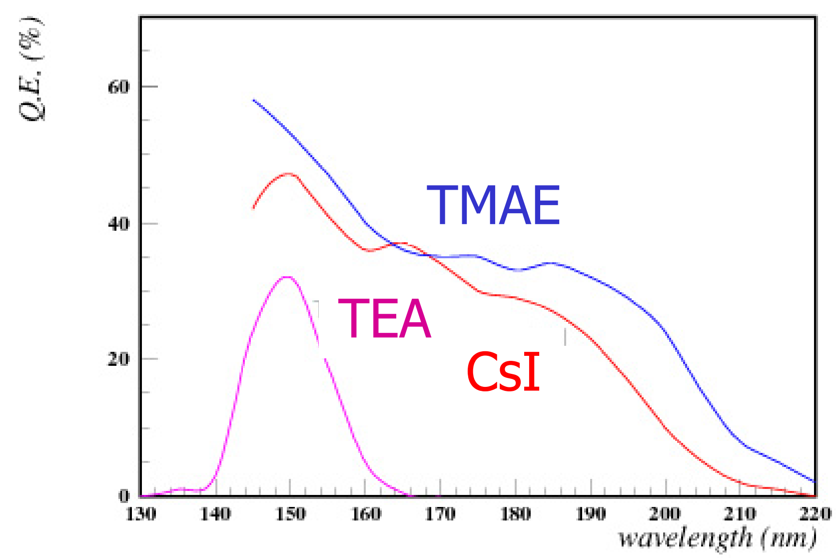
Multiwire chamber with
cathode pad read-out:
→ short drift distances,
fast detector



Photosensitive component:

•in the gas mixture (TEA):
CLEOIII RICH

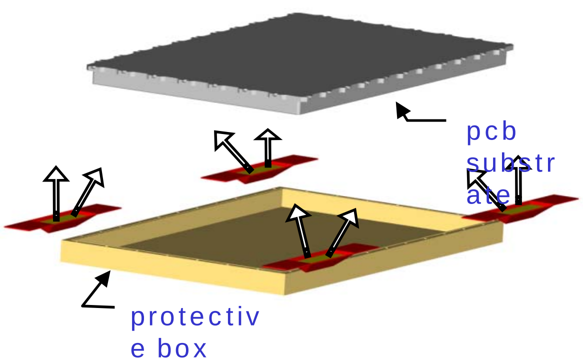
•or a layer on one of the cathodes (CsI on the printed circuit cathode with pads) HADES, COMPASS, ALICE

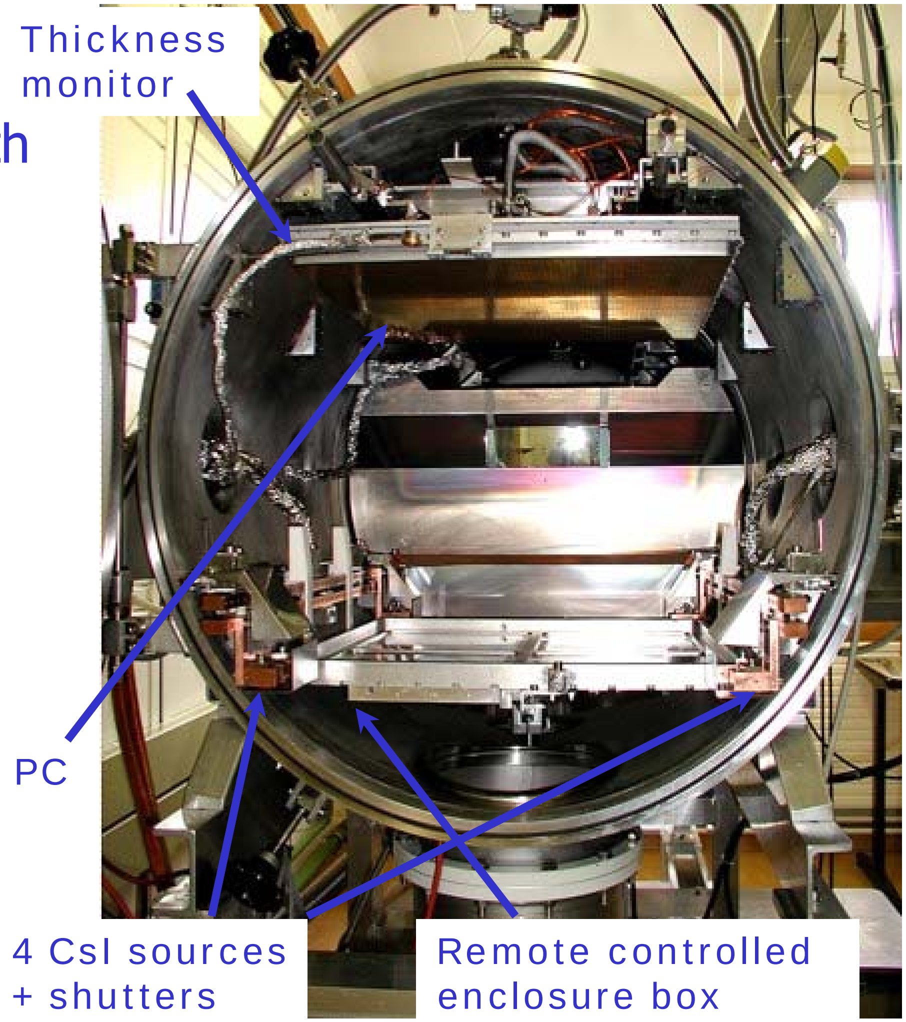


Works in high magnetic field!

CERN CsI deposition plant

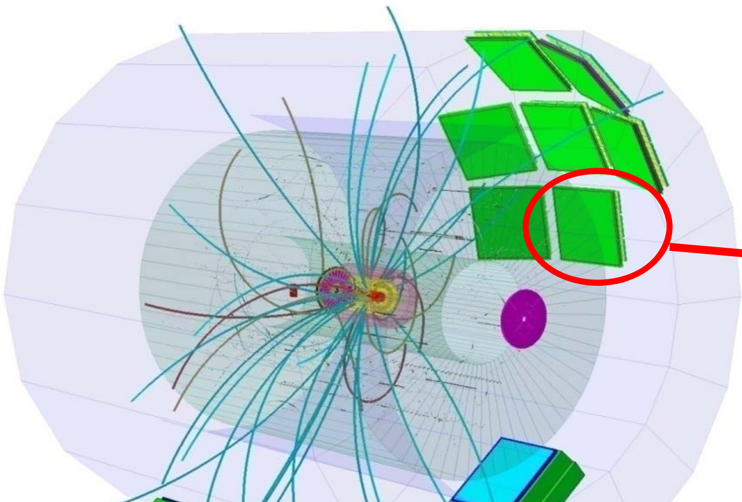
Photocathode produced with a well defined, several step procedure, with CsI vaccum deposition and subsequent heat conditioning

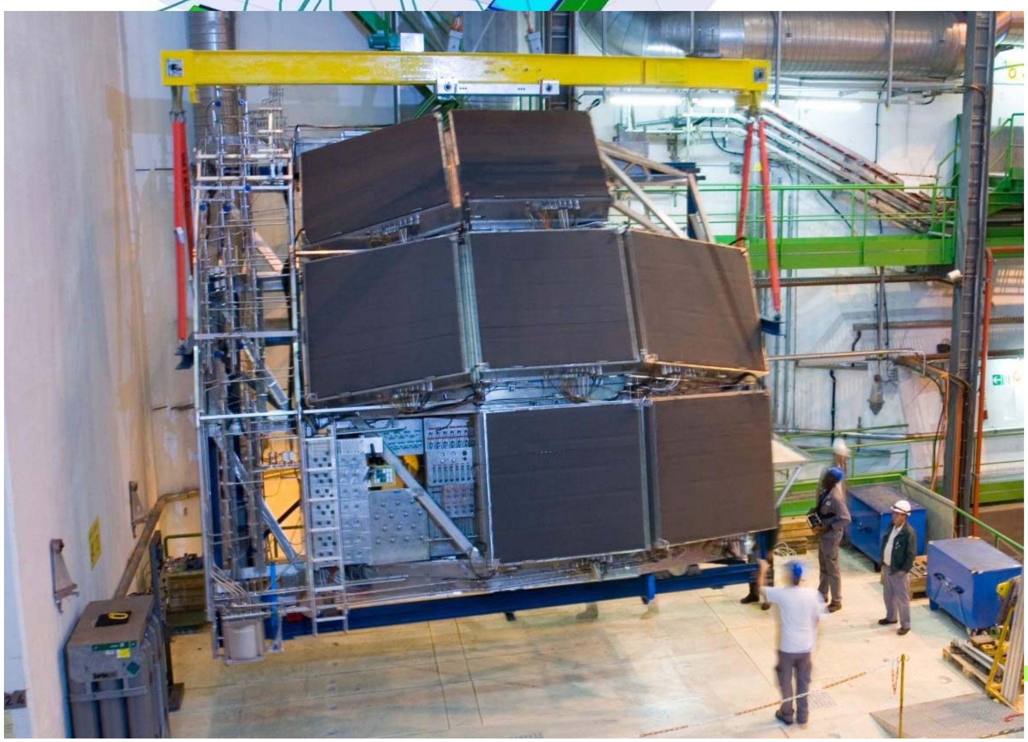




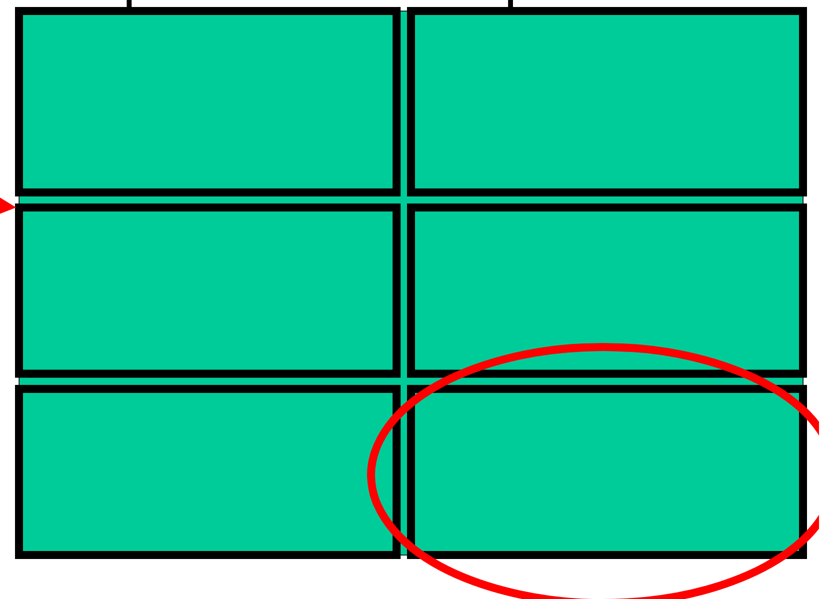
ALICE RICH = HMPID

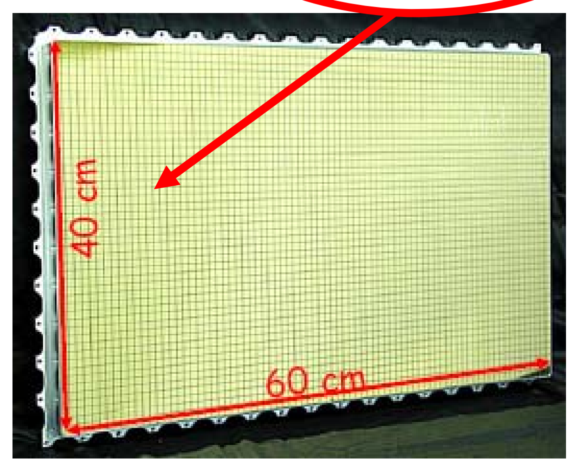
The largest scale (11 m^2) application of CsI photo-cathodes in HEP!





Six photo-cathodes per module





CsI photo-cathode is segmented in 0.8x0.84 cm pads

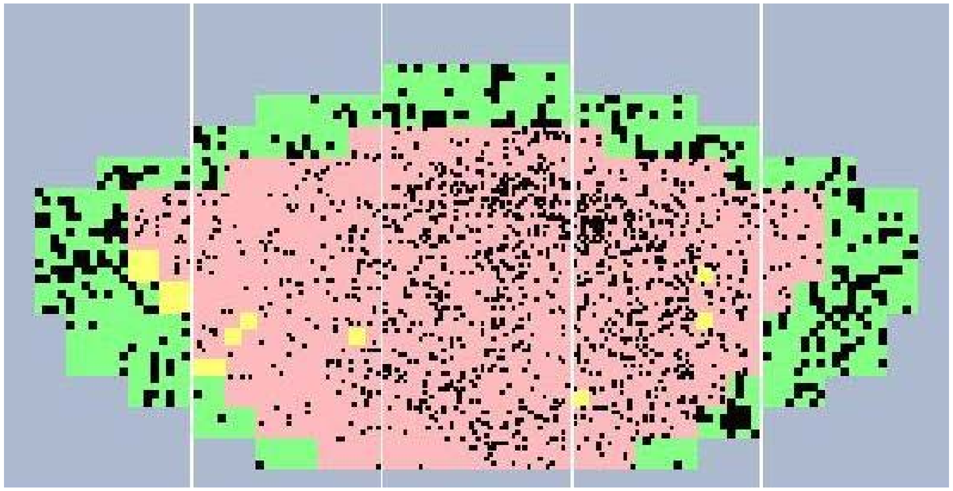
Cherenkov counters with vacuum based photodetectors

Many applications: operation at high rates over extended running periods (years) → wire chamber based photon detectors were found to be unsuitable (problems in high rate operation, ageing, only UV photons, difficult handling in 4π spectrometers)

→ Need vacuum based photon detectors (e.g. PMTs)

Good spacial resolution (pads with ~5 mm size)

→ Solution: multianode PMTs (MaPMTs)



Multianode PMTs

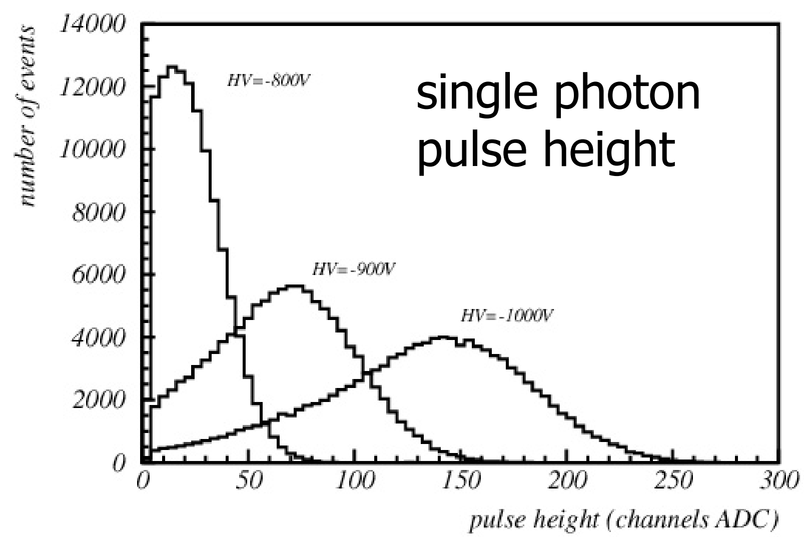
Multianode PMTs (MaPMTs) with metal foil dynodes and 2x2, 4x4 or 8x8 anodes Hamamatsu R5900 (and follow up types 7600, 8500)

→Excellent single photon pulse height spectrum

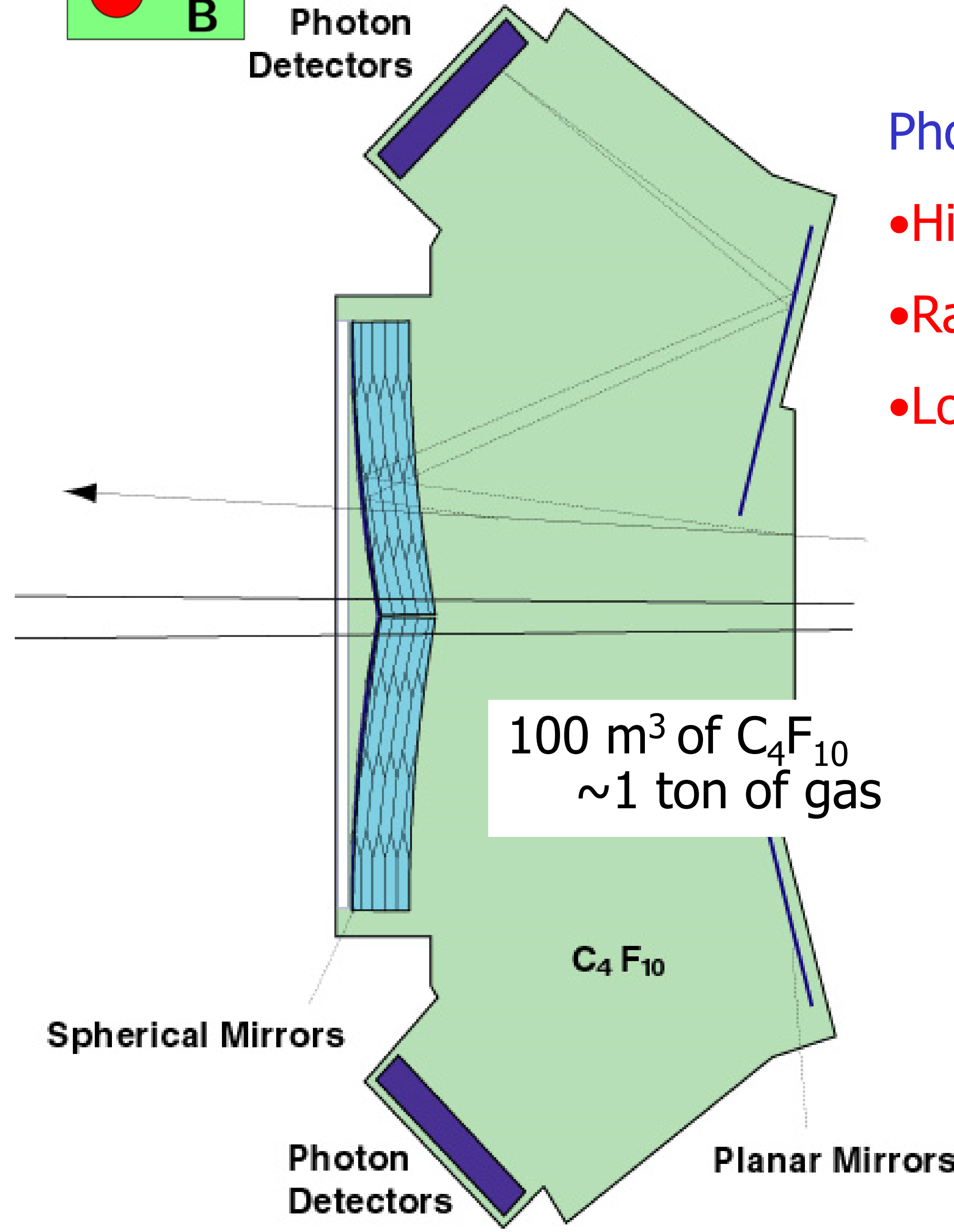
→Low noise (few Hz/ch)

→Low cross-talk (<1%)

→ NIM A394 (1997) 27



HERA-B RICH



100 m^3 of C_4F_{10}
~1 ton of gas

Photon detector requirements:

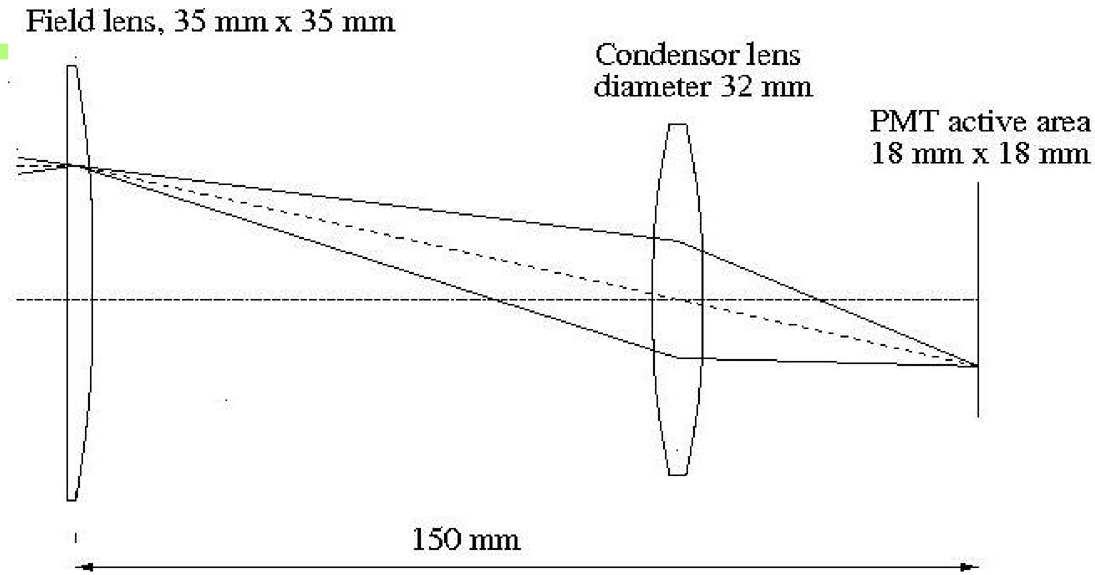
- High QE over ~3m^2
- Rates ~1MHz
- Long term stability

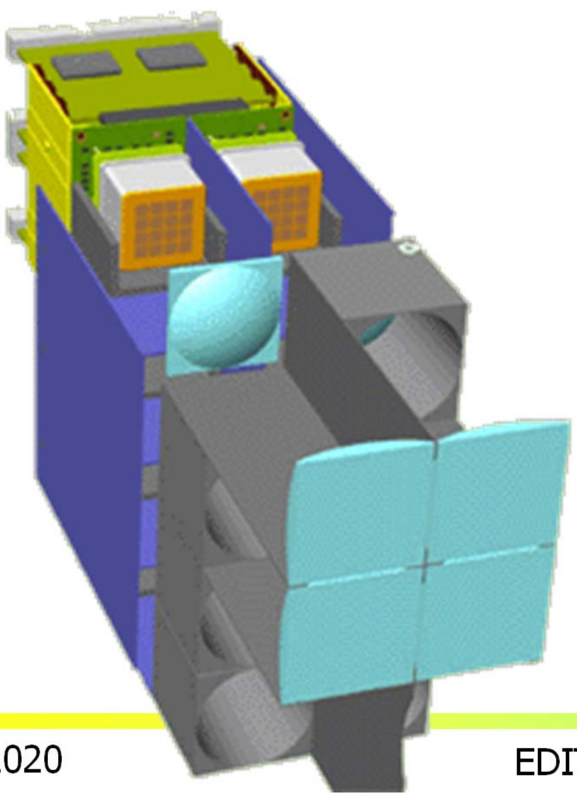
HERA-B RICH photon detector

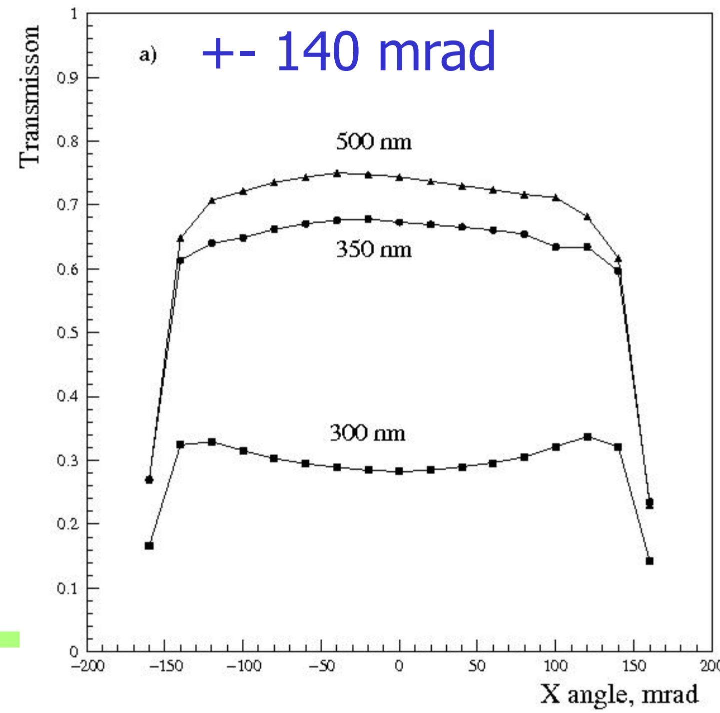
Light collection system (imaging!) to:

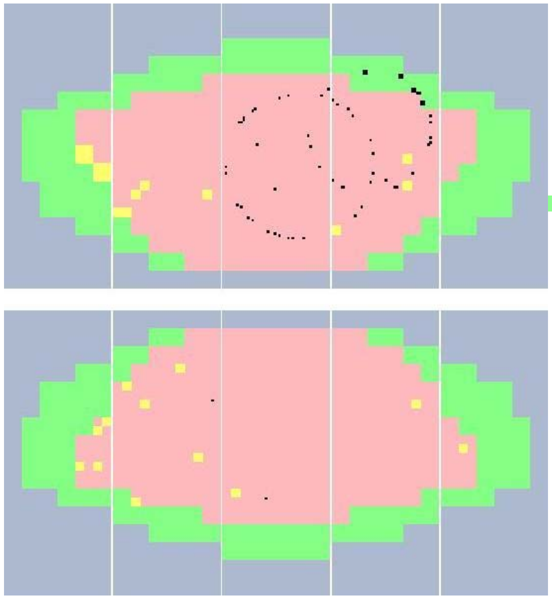
-Eliminate dead areas

-Adapt the pad size





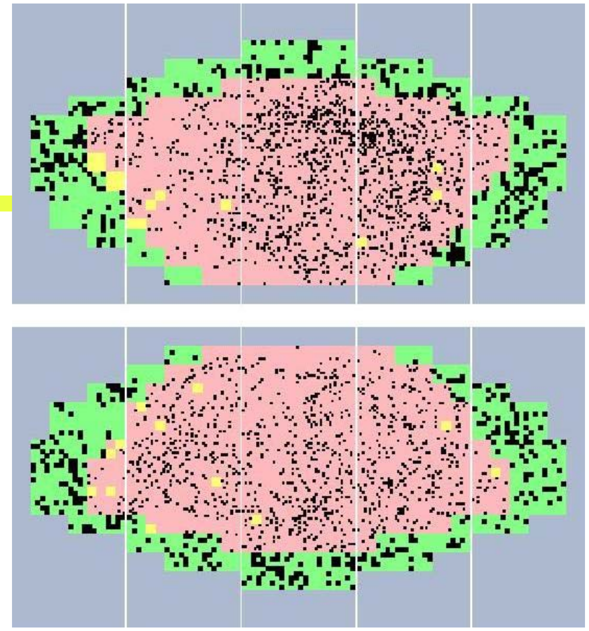




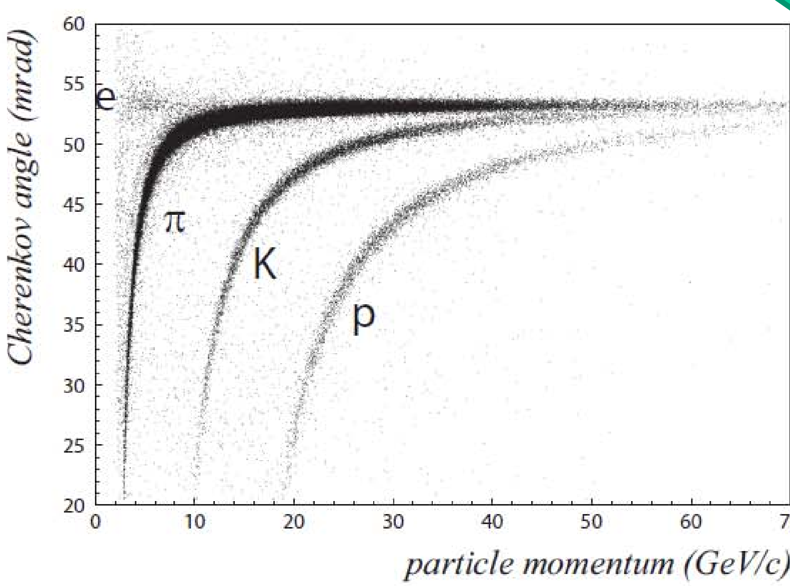
HERA-B RICH

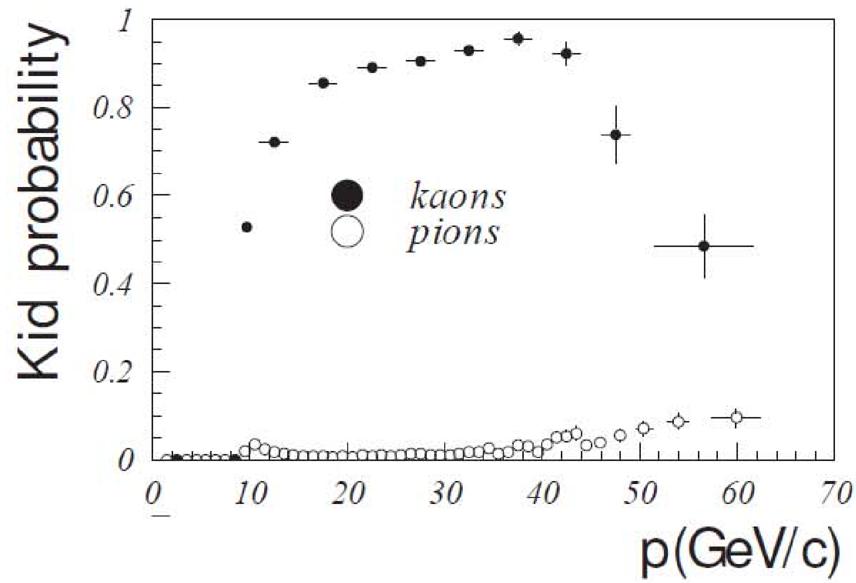
← Little noise, ~30 photons per ring

Typical event →



Worked very well!

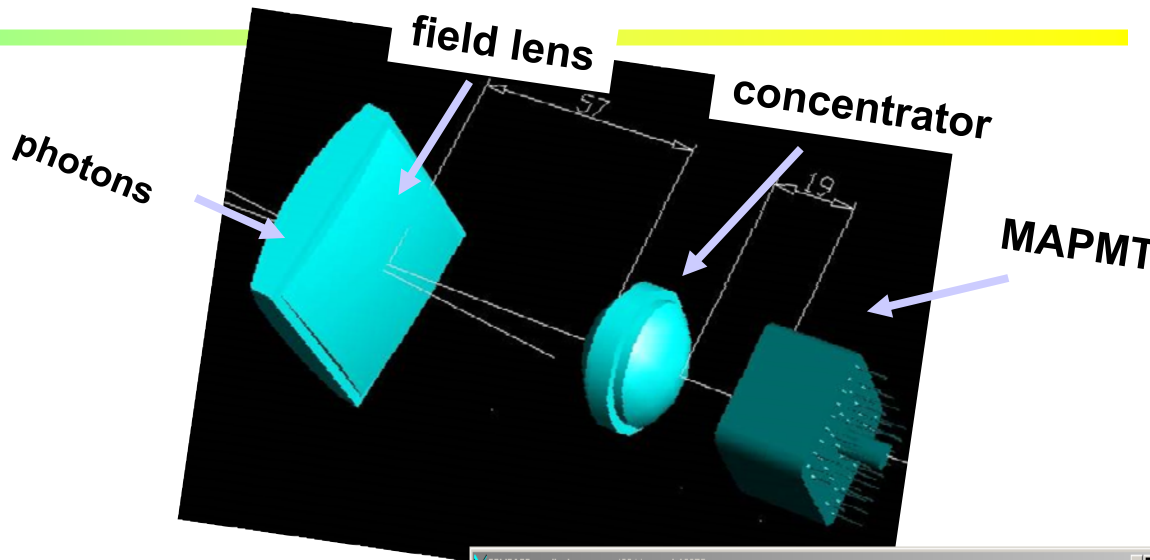




Kaon efficiency and pion fake probability

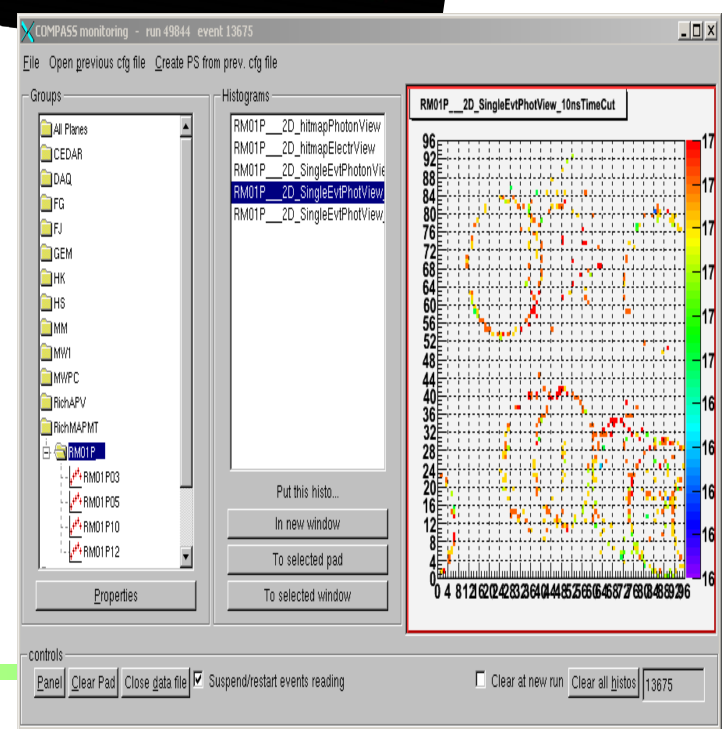
Photon detector for the COMPASS RICH-1

Upgraded COMPASS RICH-1:
similar concept as in the
HERA-B RICH

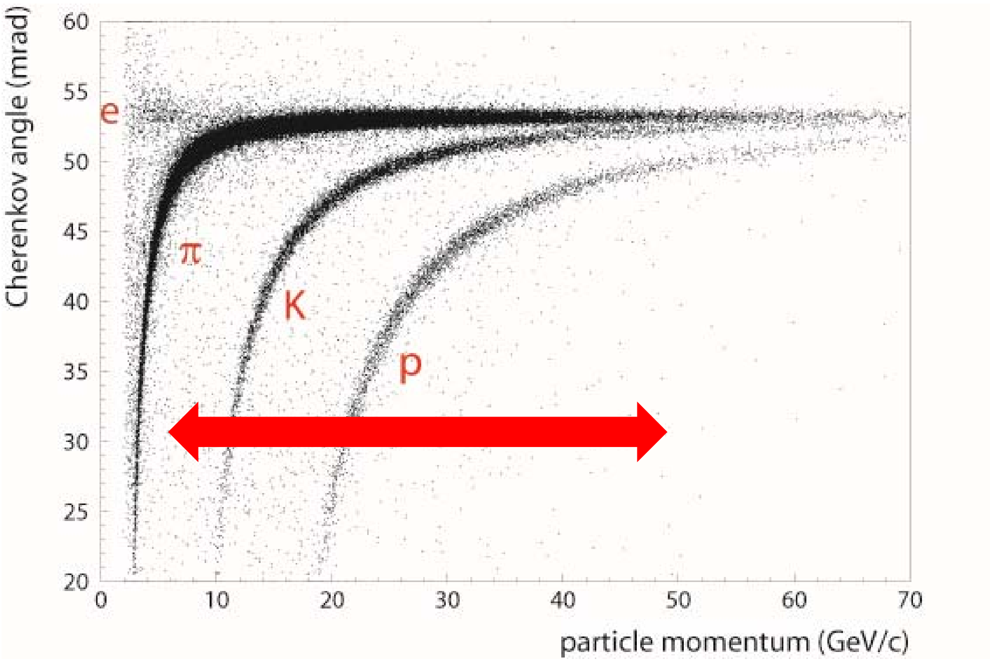


New features:

- UV extended PMTs & lenses (down to 200 nm) → more photons
- surface ratio = (telescope entrance surface) / (photocathode surface) = 7
- fast electronics with <120 ps time resolution



Kinematic range of a RICH counter



Example: kinematic range for kaon/pion separation

Kinematic range for separation of two particle types:

•Lower limit p_{min}: sufficiently above lighter particle threshold

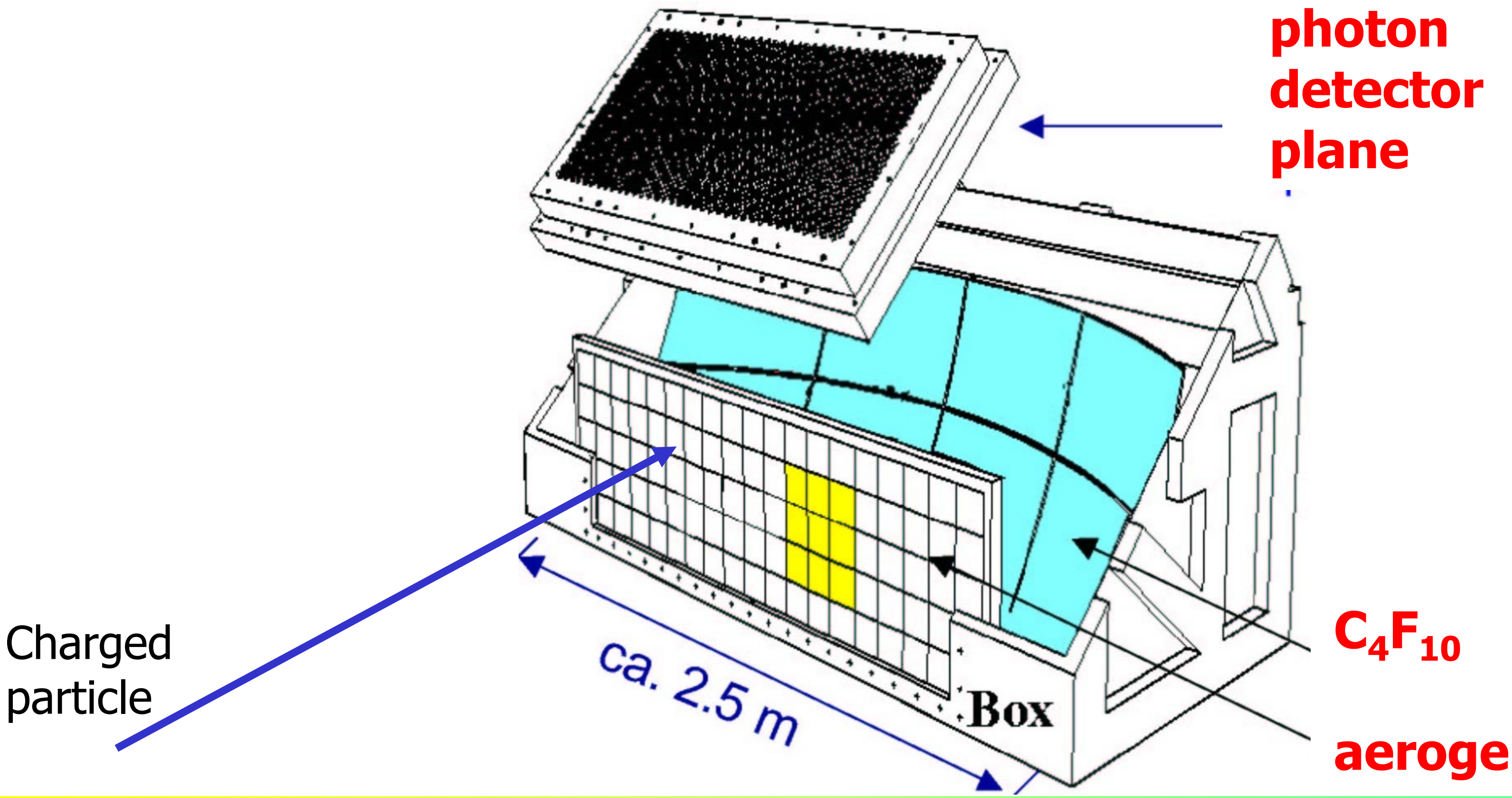
•Upper limit p_{max}: given by Cherenkov angle resolution – overlap of the two bands

Rule of thumb: $p_{max} / p_{min} < 10$

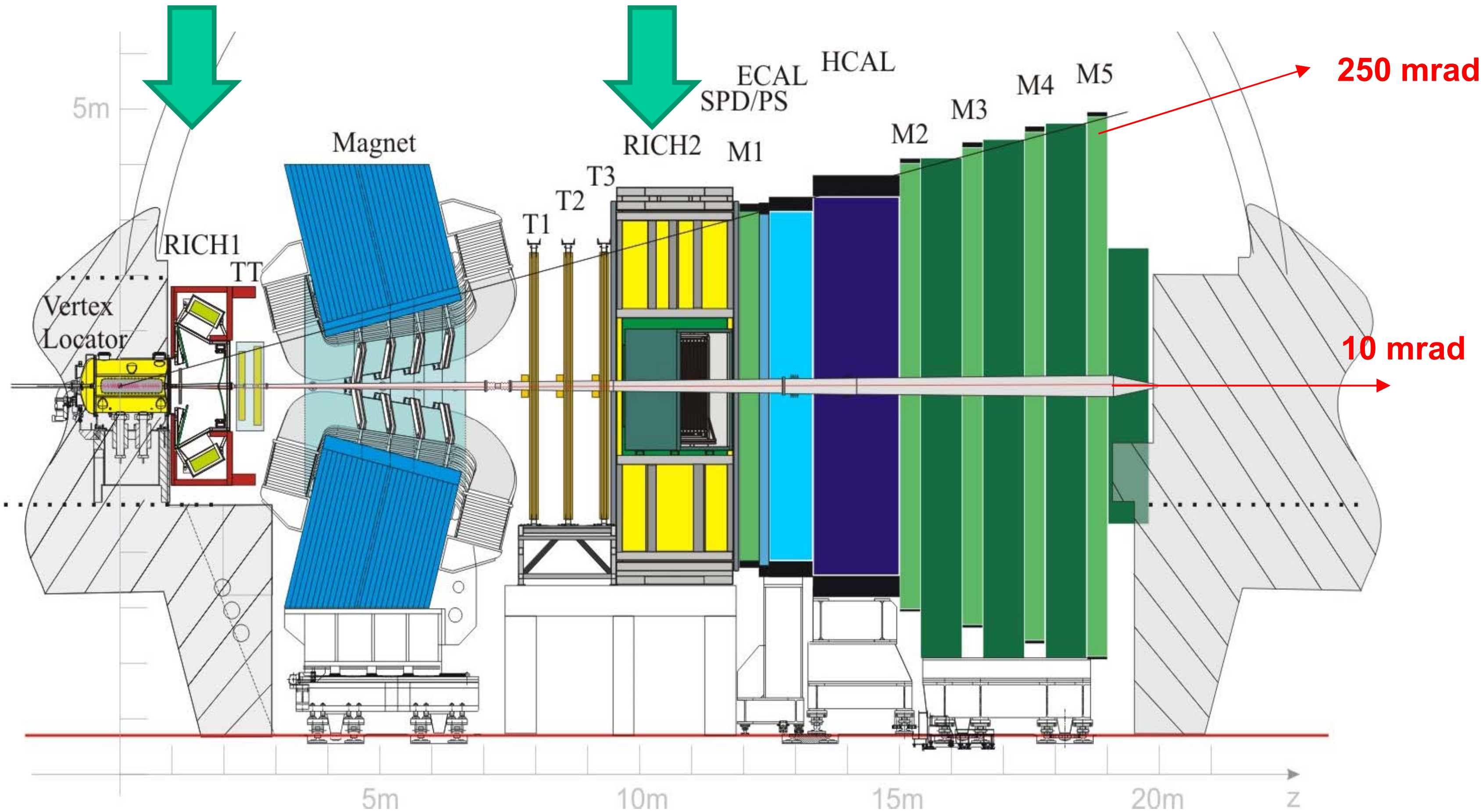
RICHes with several radiators

Extending the kinematic range → need more than one radiator

- DELPHI at LEP, SLD at SLC (liquid +gas)
- HERMES at HERA (aerogel+gas)



The LHCb RICH counters



Vertex reconstruction:
VELO

Trigger:
Muon Chambers
Calorimeters
Tracker

PID:
RICHes
Calorimeters
Muon Chambers

Kinematics:
Magnet
Tracker
Calorimeters

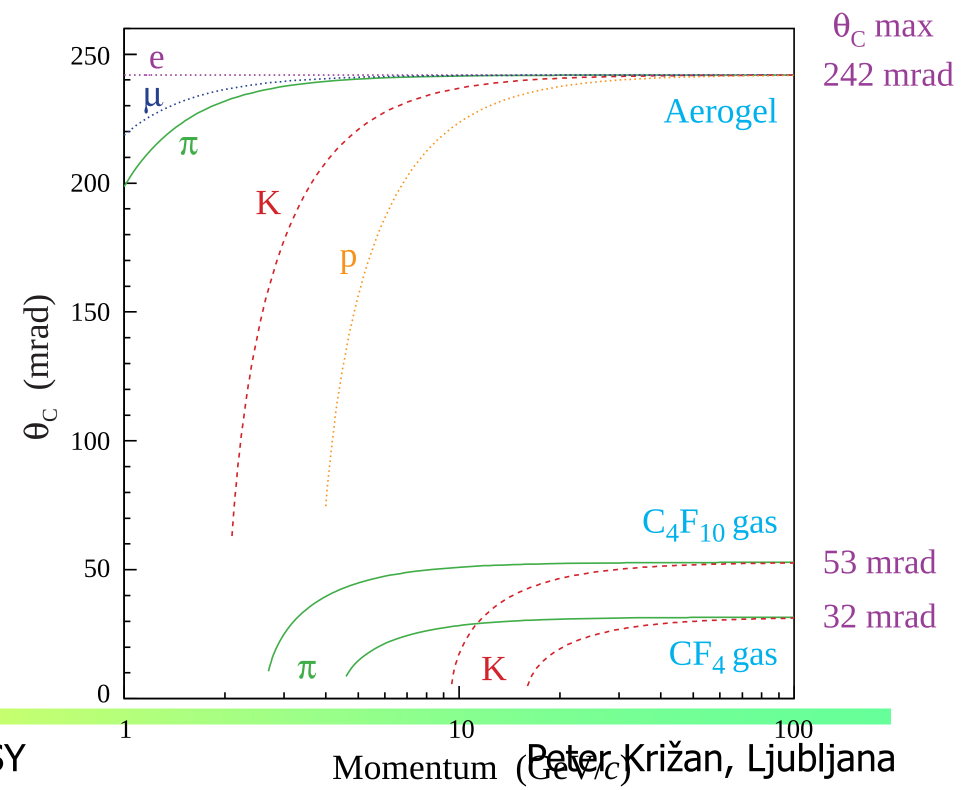
LHCb RICHes

Need:

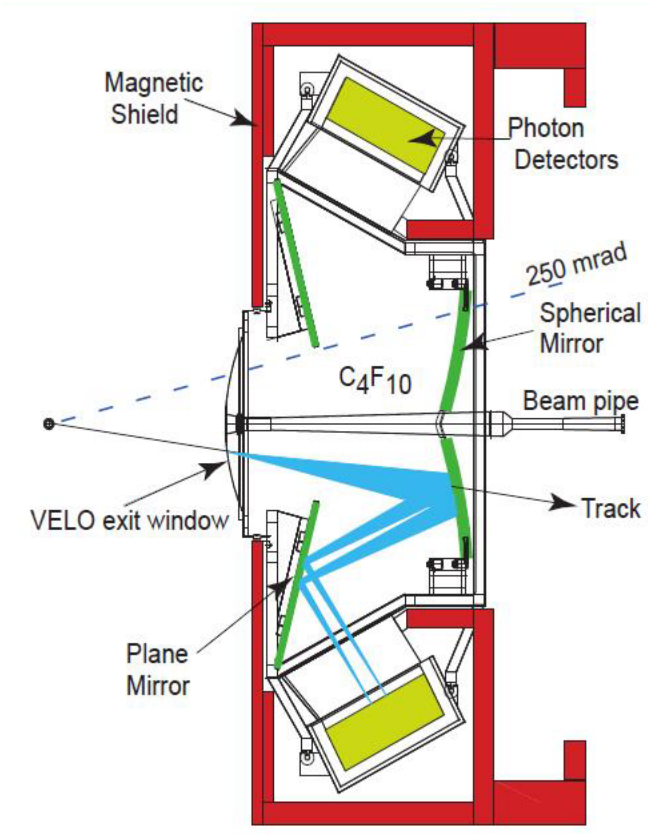
- Particle identification for momentum range ~2-100 GeV/c
- Granularity 2.5x2.5mm^2
- Large area (2.8m^2) with high active area fraction
- Fast compared to the 25ns bunch crossing time
- Have to operate in a small B field

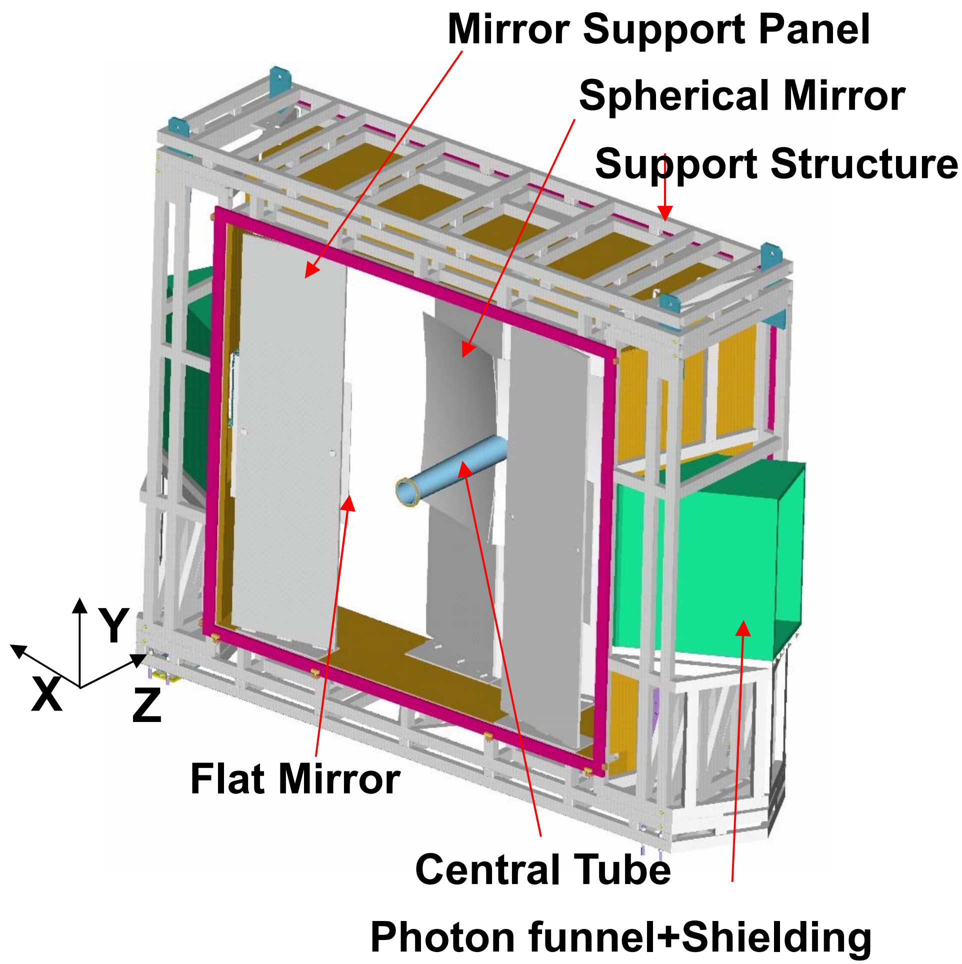
→3 radiators

- Aerogel
- C_4F_{10} gas
- CF_4 gas



LHCb RICHes

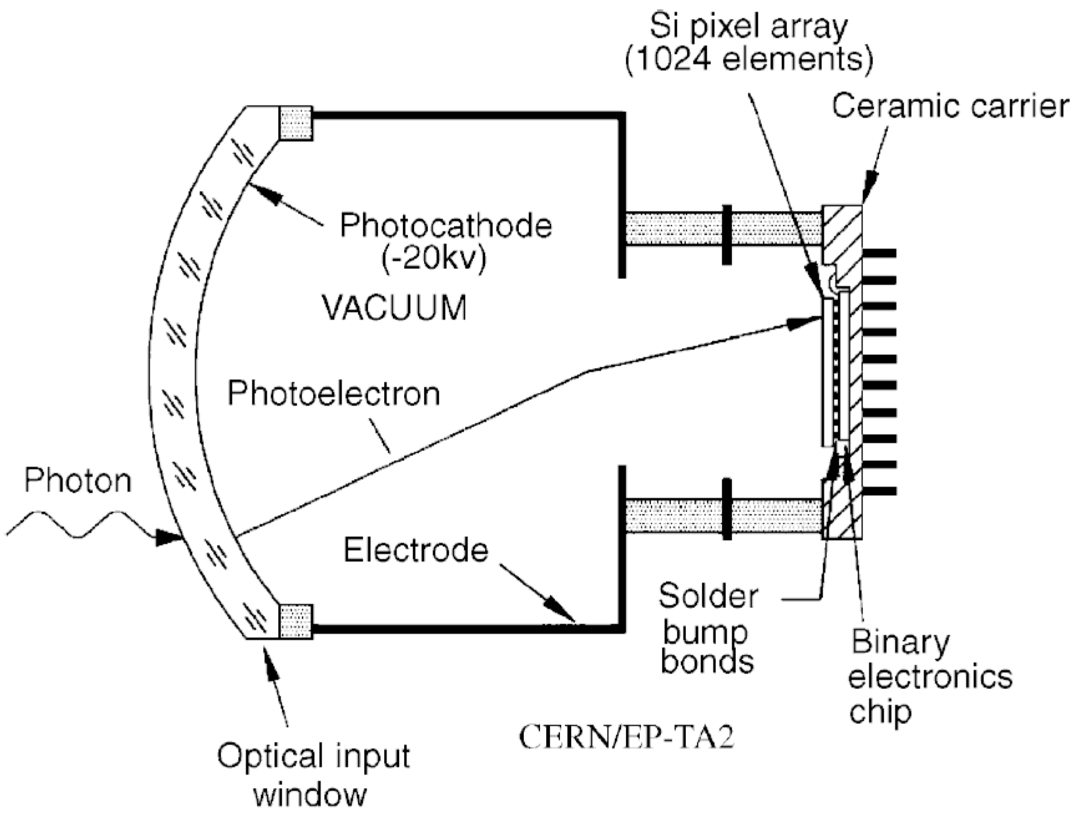

Magnetic Shield
Photon Detectors
250 mrad
Spherical Mirror
C_4F_{10}
Beam pipe
Track
VELO exit window
Plane Mirror


Mirror Support Panel
Spherical Mirror
Support Structure
Y
X
Z
Flat Mirror
Central Tube
Photon funnel+Shielding

LHCb RICHes

Photon detector: hybrid PMT (R+D with DEP) with 5x demagnification (electrostatic focusing).

Hybrid PMT: accelerate photoelectrons in electric field (~20kV), detect them in a pixelated silicon detector.



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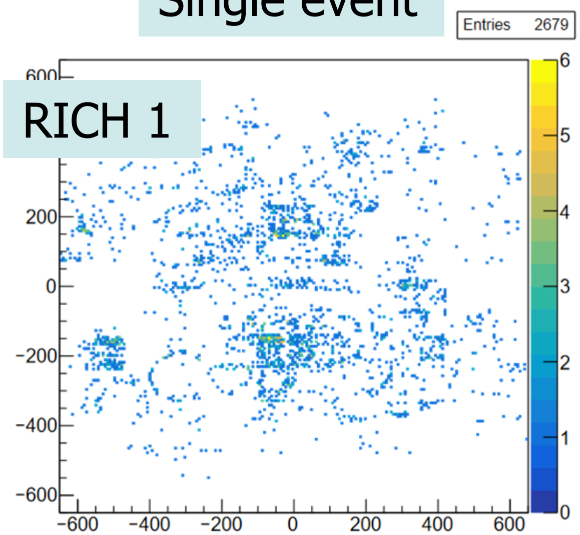
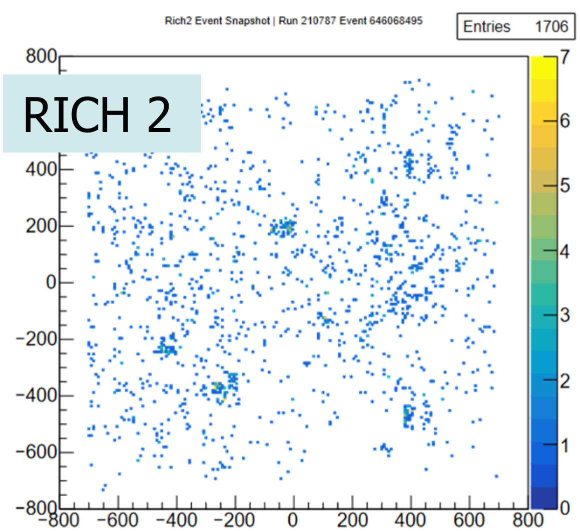
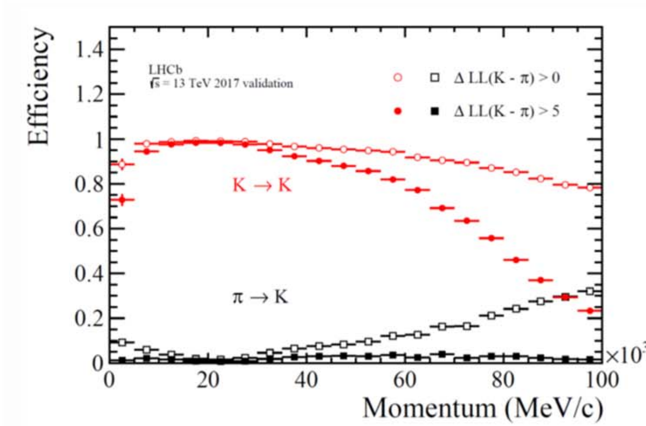
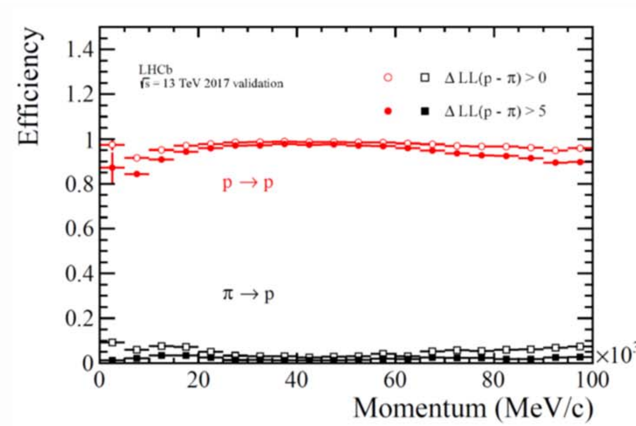
Performance of LHCb RICHes

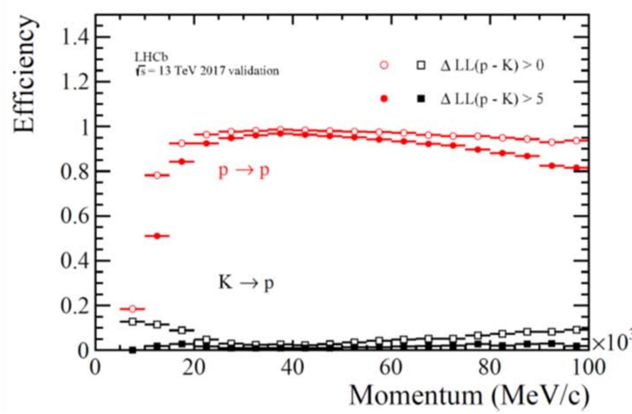
Single event



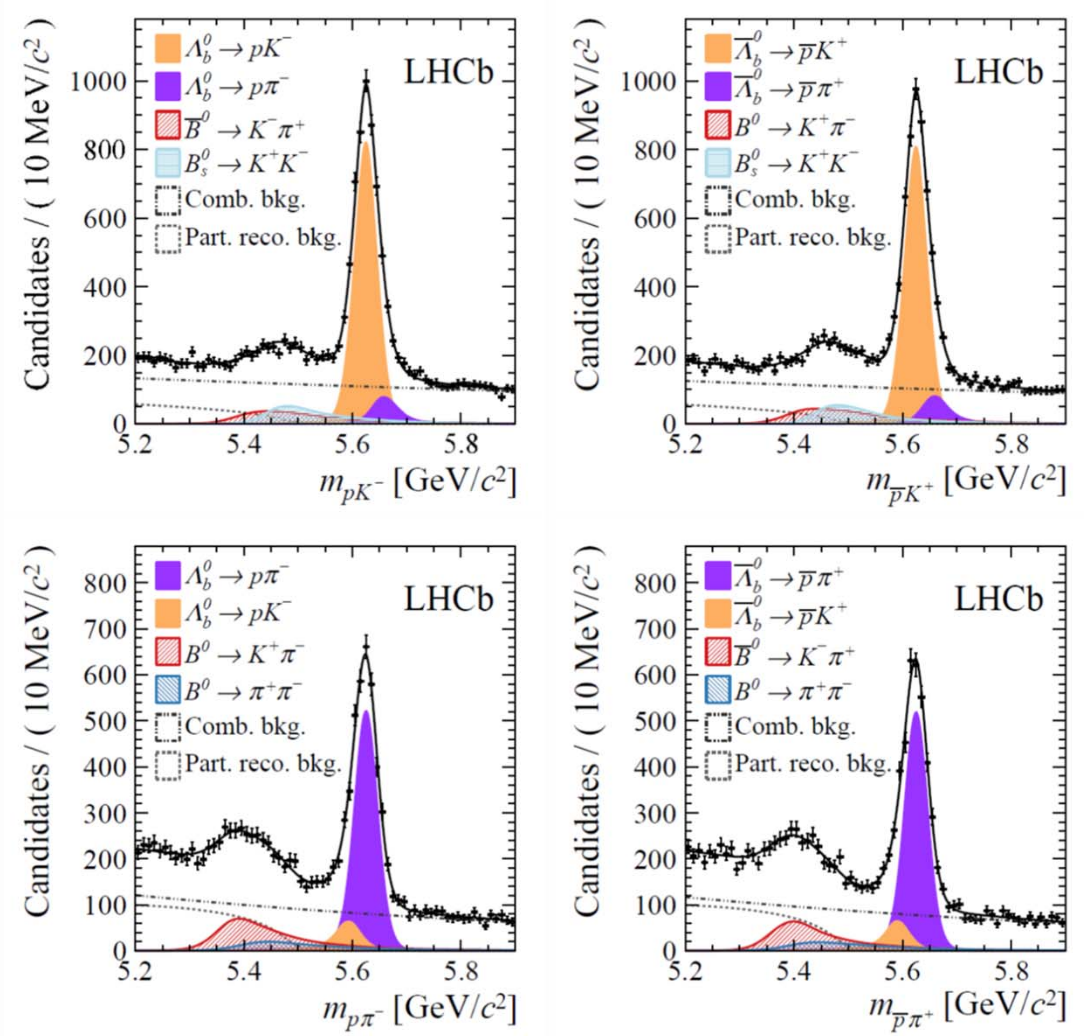






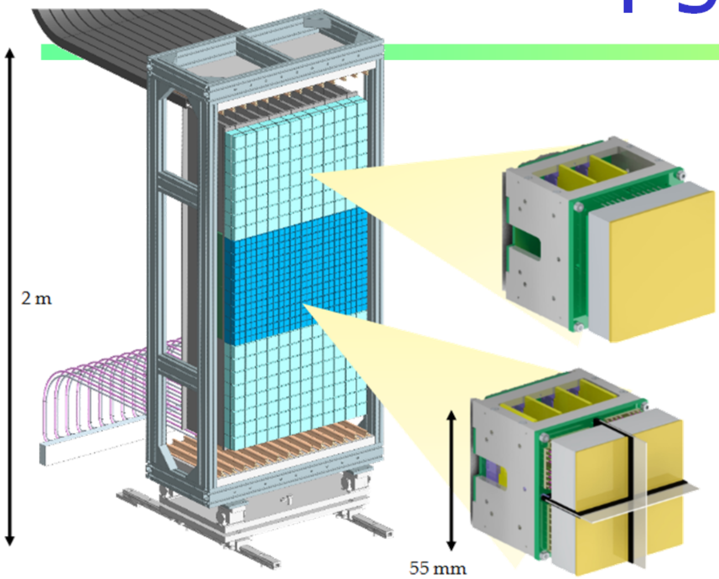


LHCb RICHes: performance

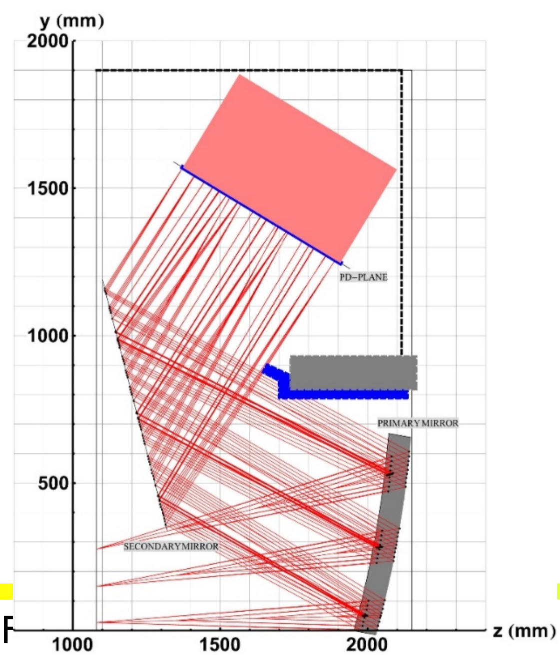


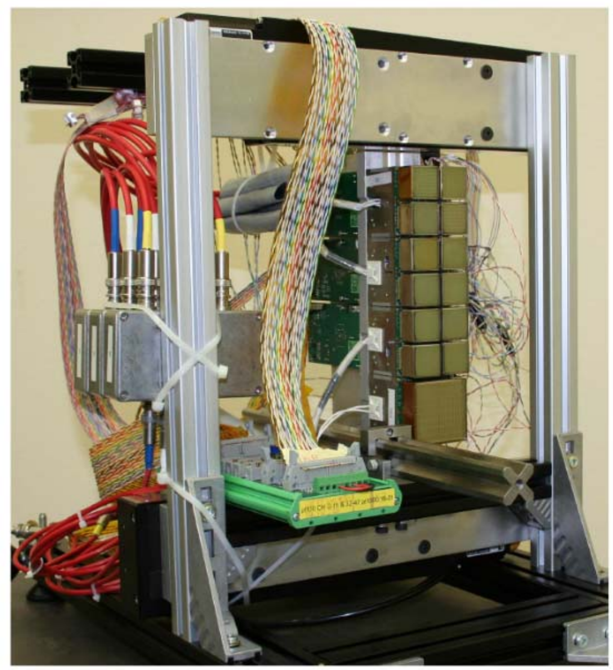
"Search for CP violation in $\Lambda_b^0 \to pK^-$ and $\Lambda_b^0 \to p\pi^-$ decays"
[LHCb-PAPER-2018-025]

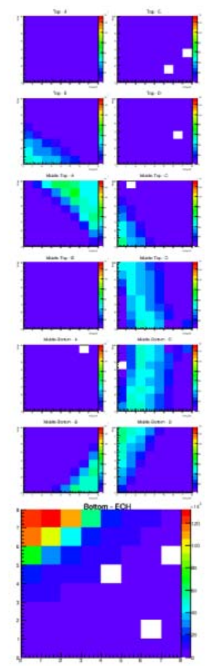
LHCb Upgrade (under way)



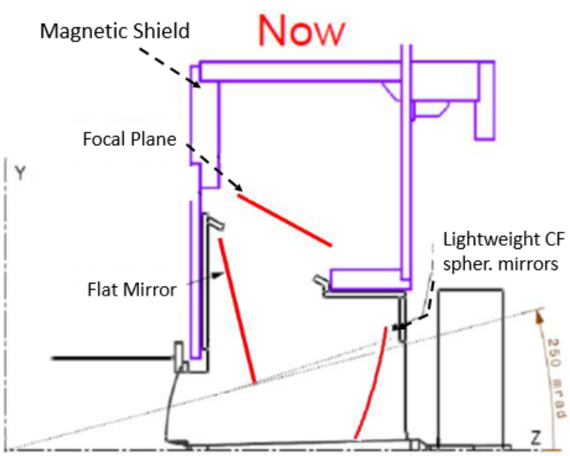
- New photon detectors: MaPMTs Hamamatsu R13743 (H12700) and R13742 (R11265)
- New electronics working at 40 MHz readout rate
- New optics layout for RICH 1

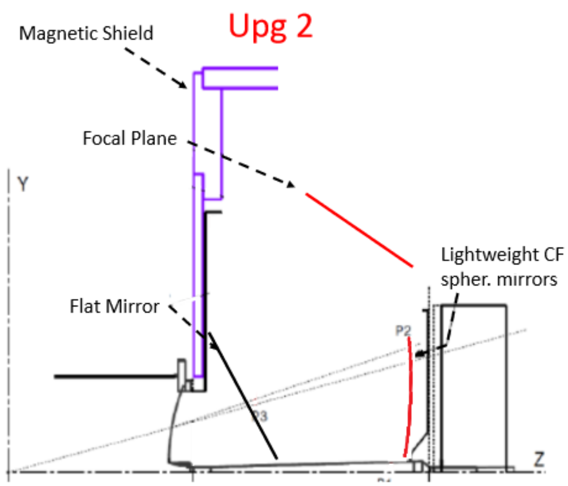






Future LHCb Upgrade



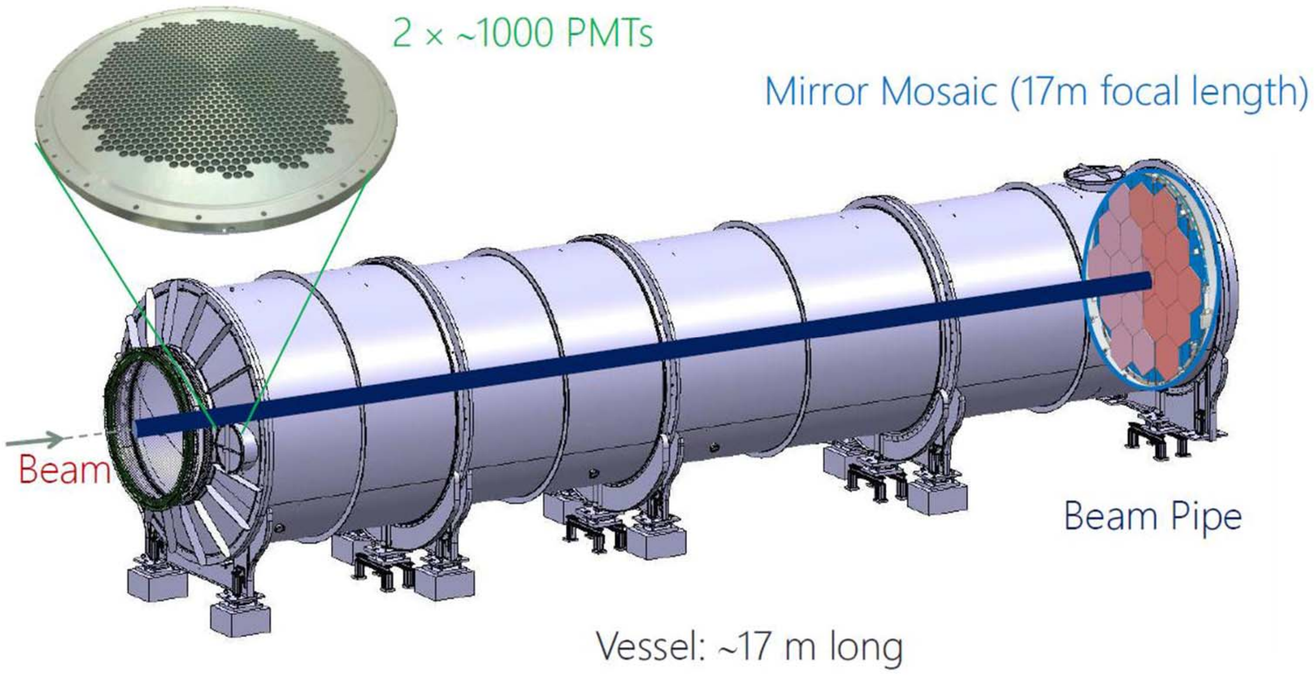


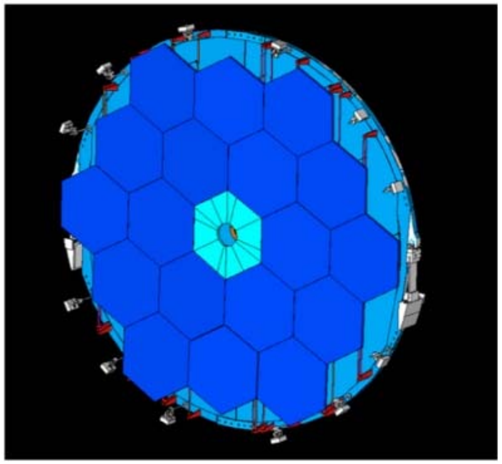
- Provide PID at p-p luminosity of 10^{34} in the forward region
- Incremental improvements in:
  - Improve Cherenkov angle resolution
    - More photons in the green → lower chromatic error
  - Reduced event complexity with timing
  - Enhanced number of photons

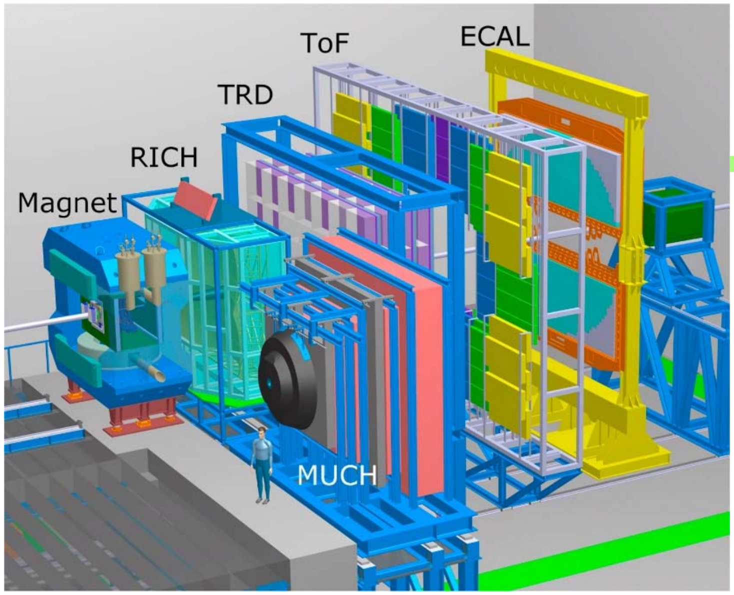
Radiator	C_4F_{10}			CF_4	
Detector Version	RICH 1 Current (HPD)	RICH 1 UPG1	**RICH 1 UPG2**	RICH 2 UPG1	**RICH 2 UPG2**
Average Photoelectron Yield	30	40	**60–30**	22	**30**
Single Photon Errors (mrad)					
Chromatic	0.84	0.58	**0.24–0.12**	0.31	**0.1**
Pixel	0.9	0.44	**0.15**	0.20	**0.07**
Emission Point	0.8	0.37	**0.1**	0.27	**0.05**
Overall	1.47	0.82	**0.3–0.2**	0.46	**0.13**

NA62 RICH

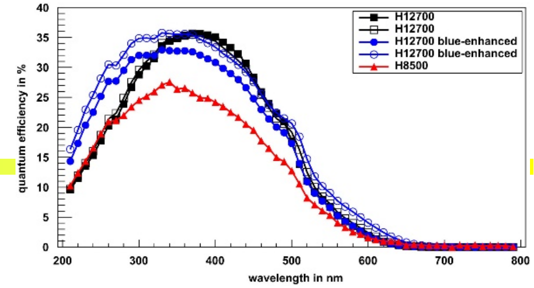
- Momentum range 15-35 GeV/c
- 17m long, 200m^3 cylindrical vacuum proof tank with Neon radiator
- Photon detectors: 2000 PMTs (16mm, 8mm active, with Winstone cone light guides)
- Mirror alignment ~30 μrad
- Single photon resolution: ~140 μrad
- Operational since 2014



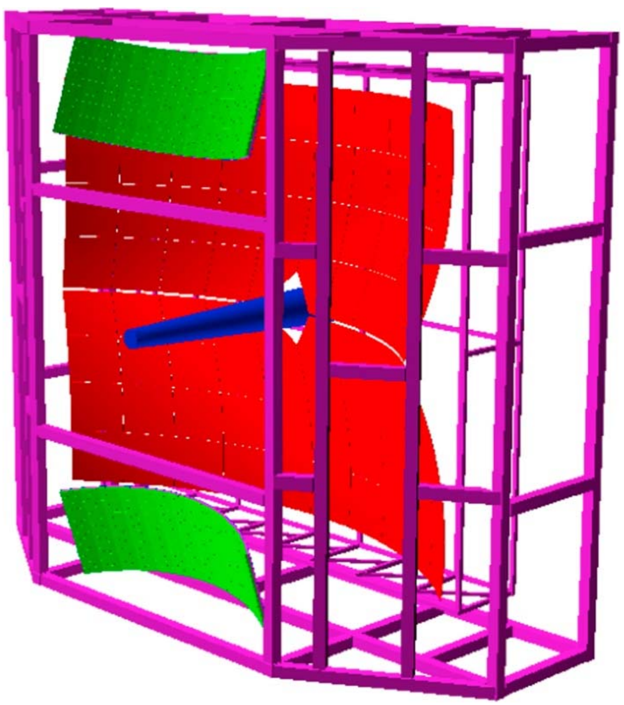


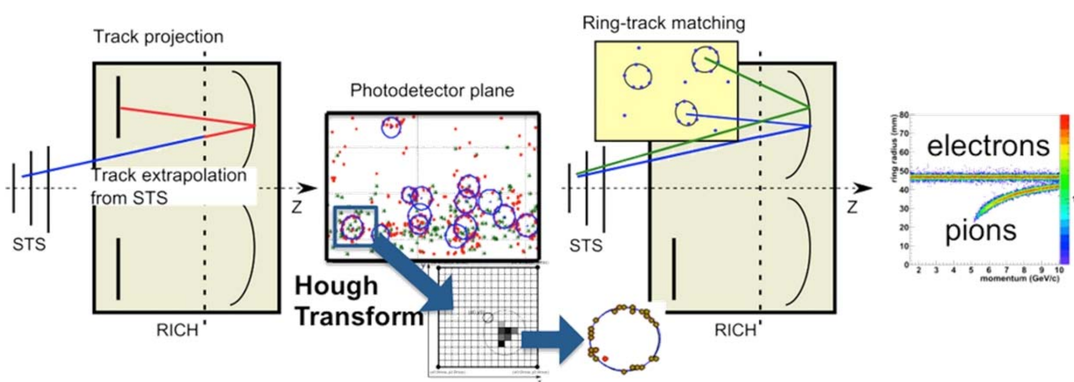


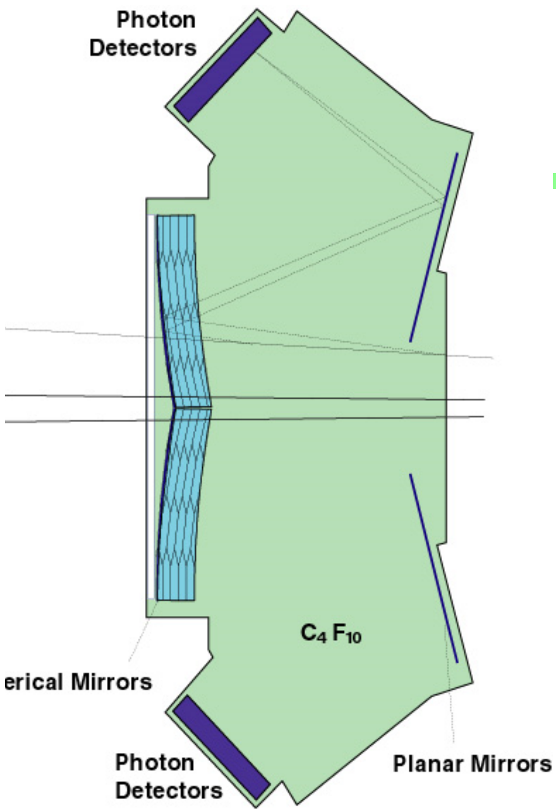
CBM



- RICH with CO_2 radiator
- MaPMTs: Hamamatsu H12700
- Cylindrical photon detection surface
- Extensive testing of MaPMTs for radiation damage
- Up to 1000 tracks per event
- Momentum up to 8 GeV/c
- Pion suppression factor ~5000 (with TRD)

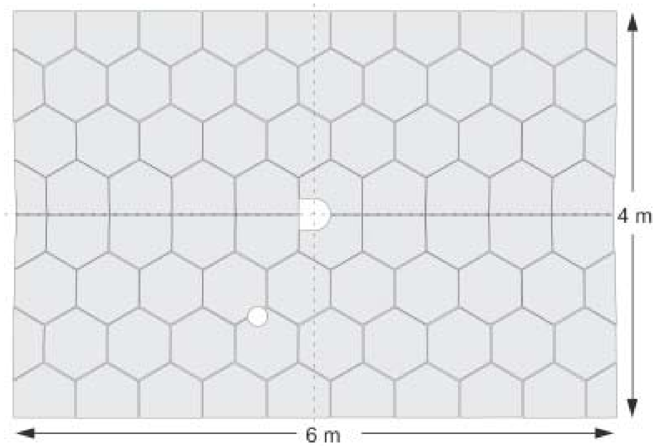






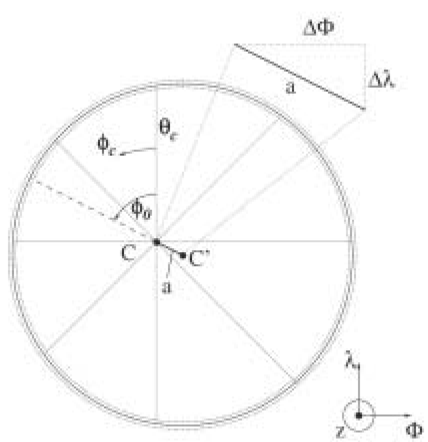
Mirror alignment

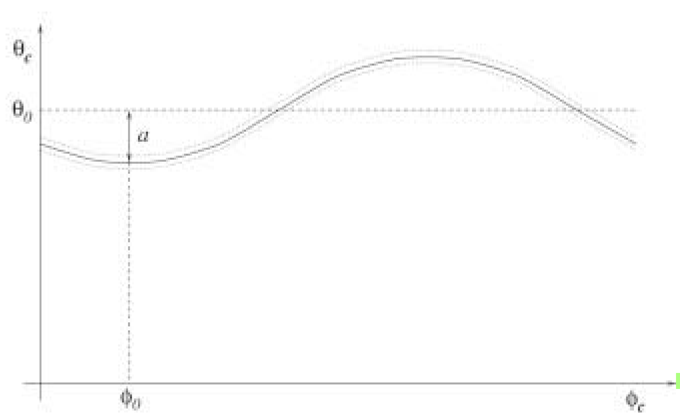
Gas radiator RICHes: large mirrors → tens of segments → need relative alignment



- Spherical mirror: 80 hexagonal segments
- Planar mirrors: 2x18 rectangular segments

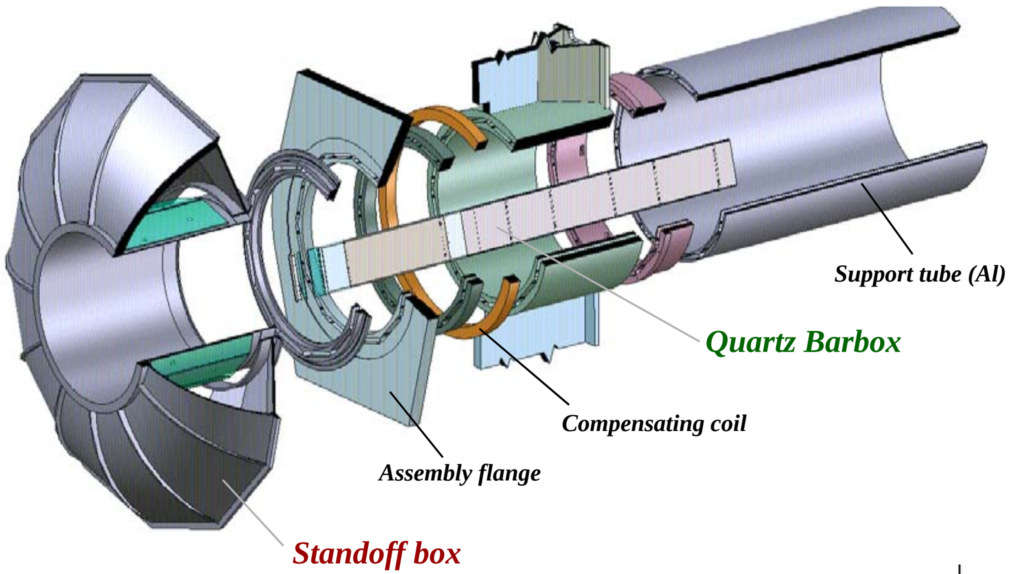
Aligning pairs of spherical and planar segments.

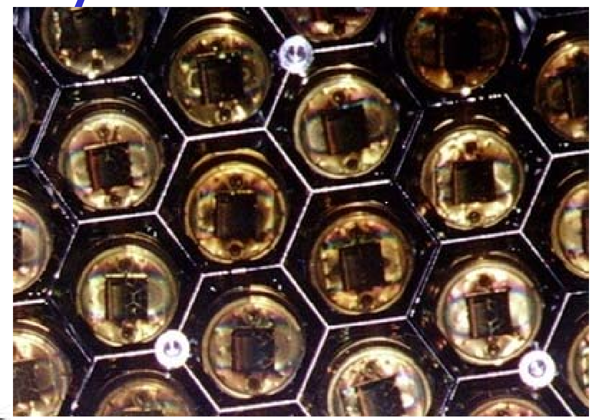


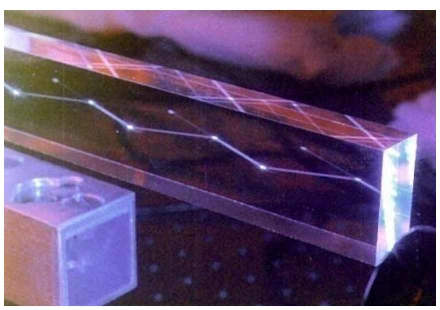


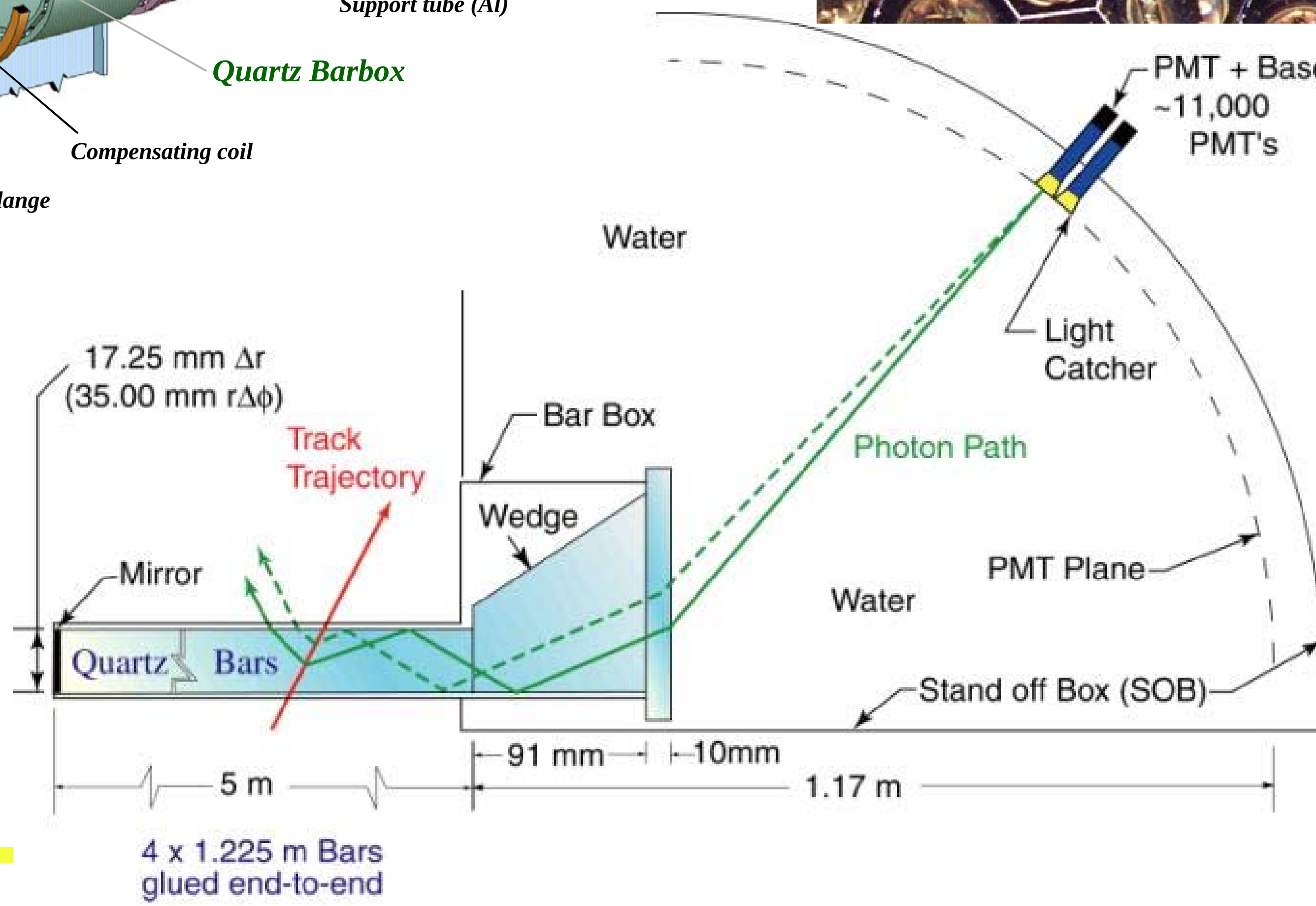
Misalignment: Cherenkov angle depends on the azimuthal angle around the track

DIRC (@BaBar) - detector of internally reflected Cherenkov light



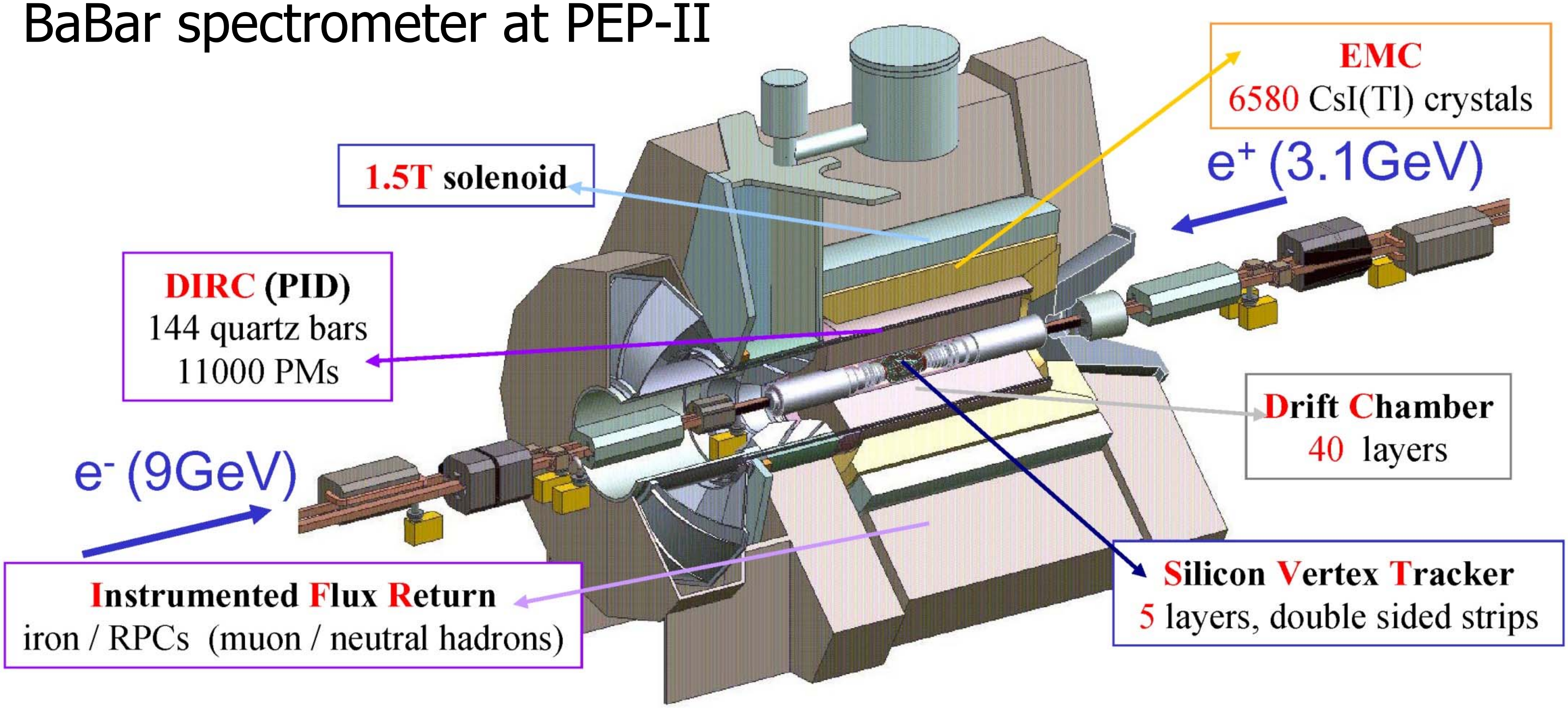




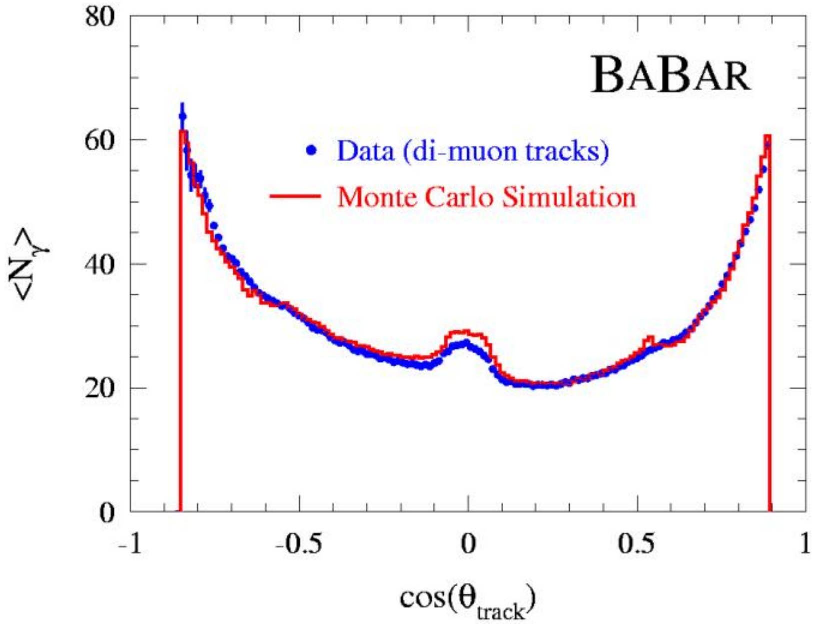


DIRC - detector of internally reflected Cherenkov light

BaBar spectrometer at PEP-II

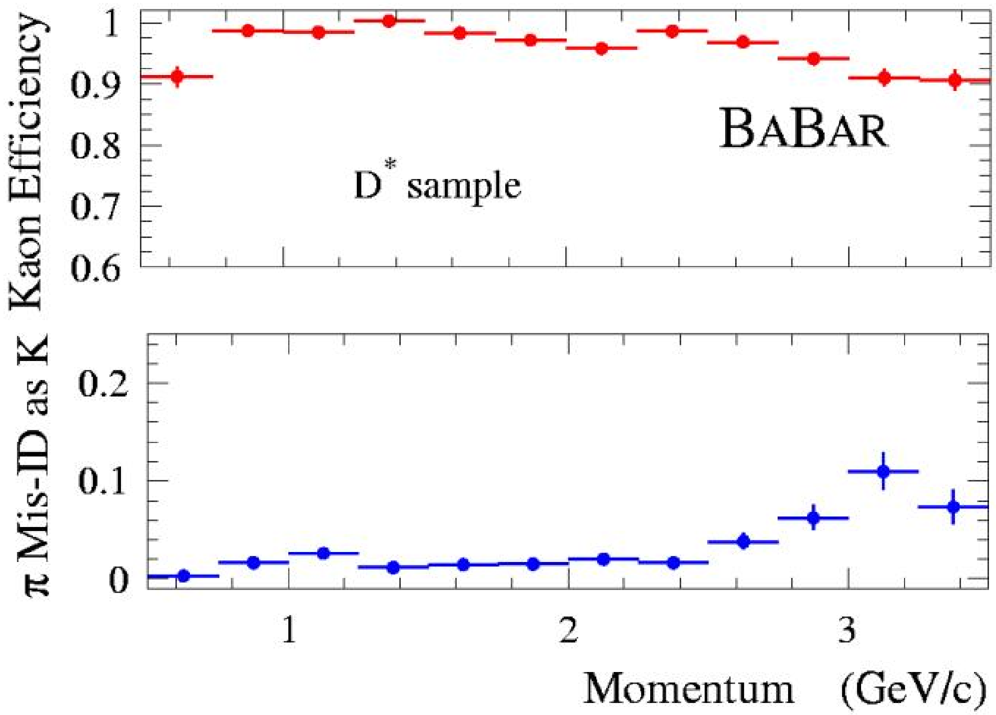


DIRC performance



← Lots of photons!

Excellent π/K separation



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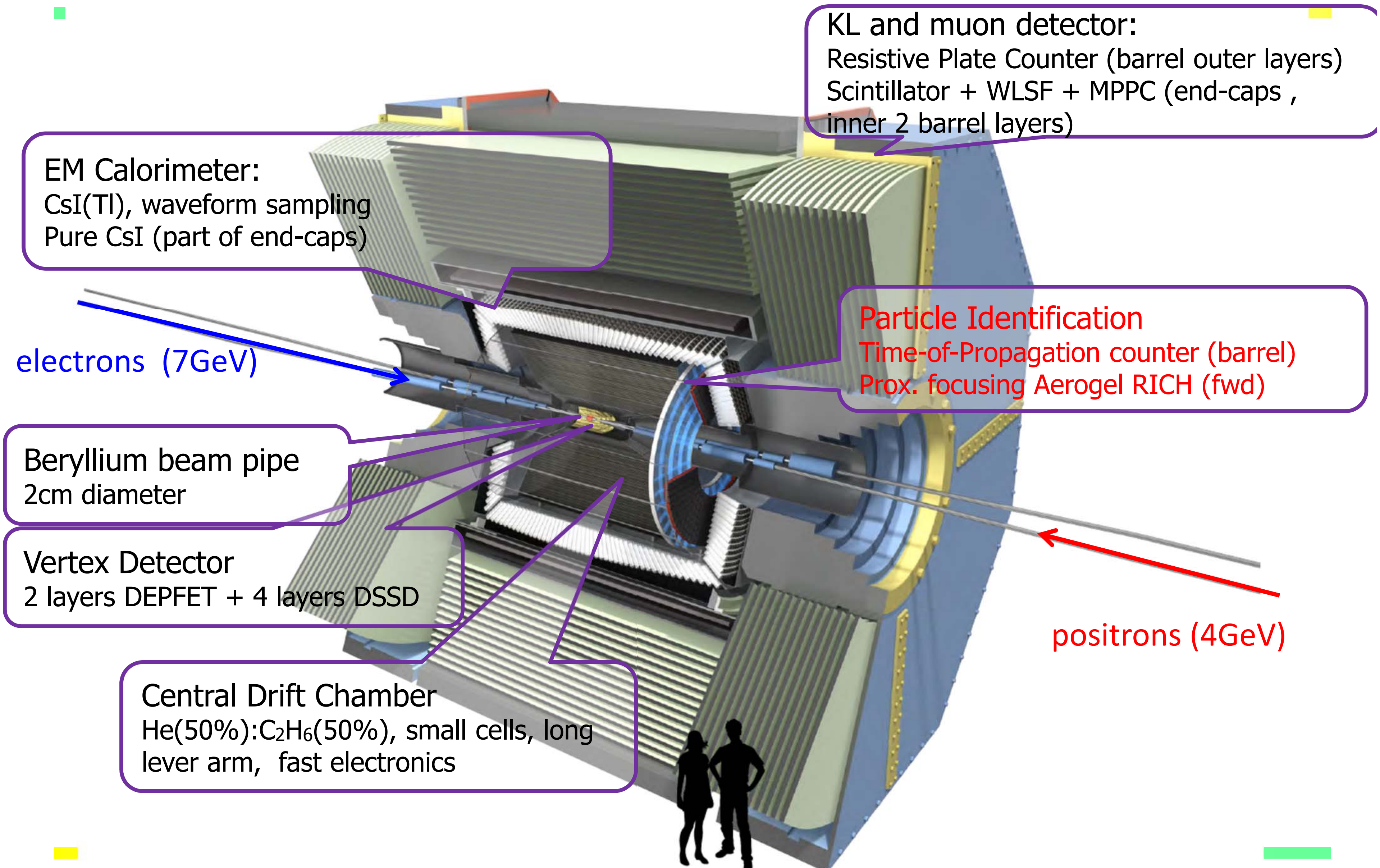
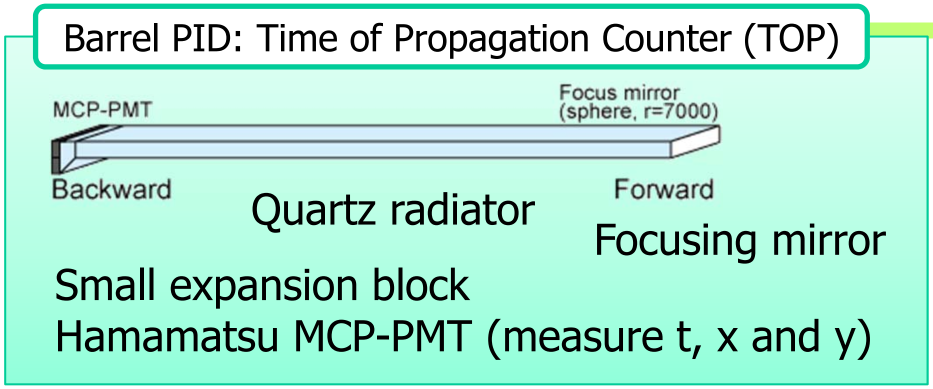
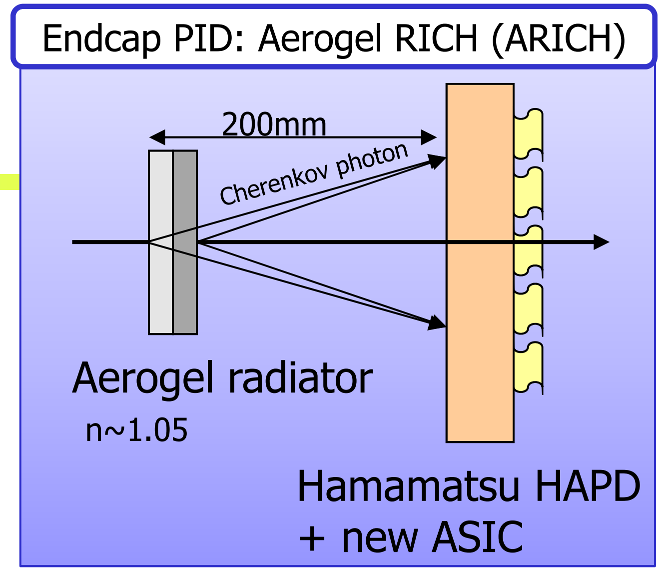
EDIT

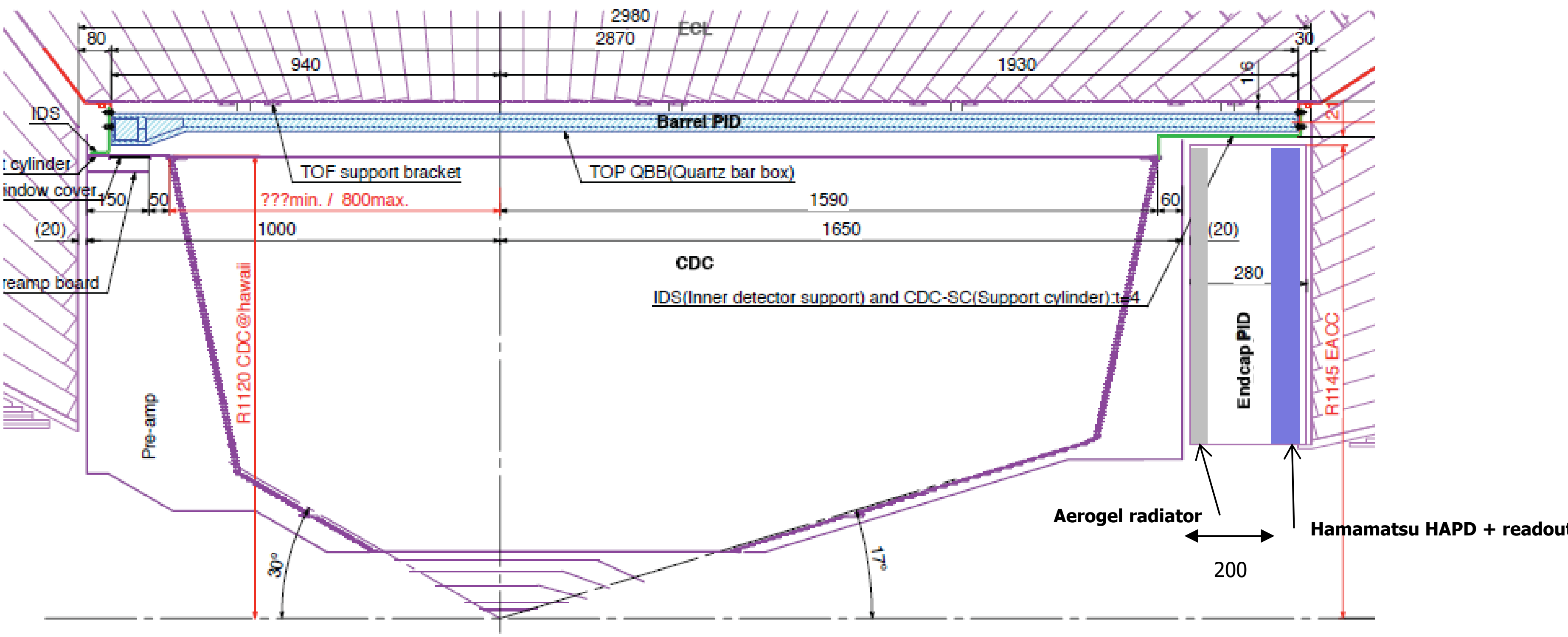
Hadron PID in the Belle II Detector



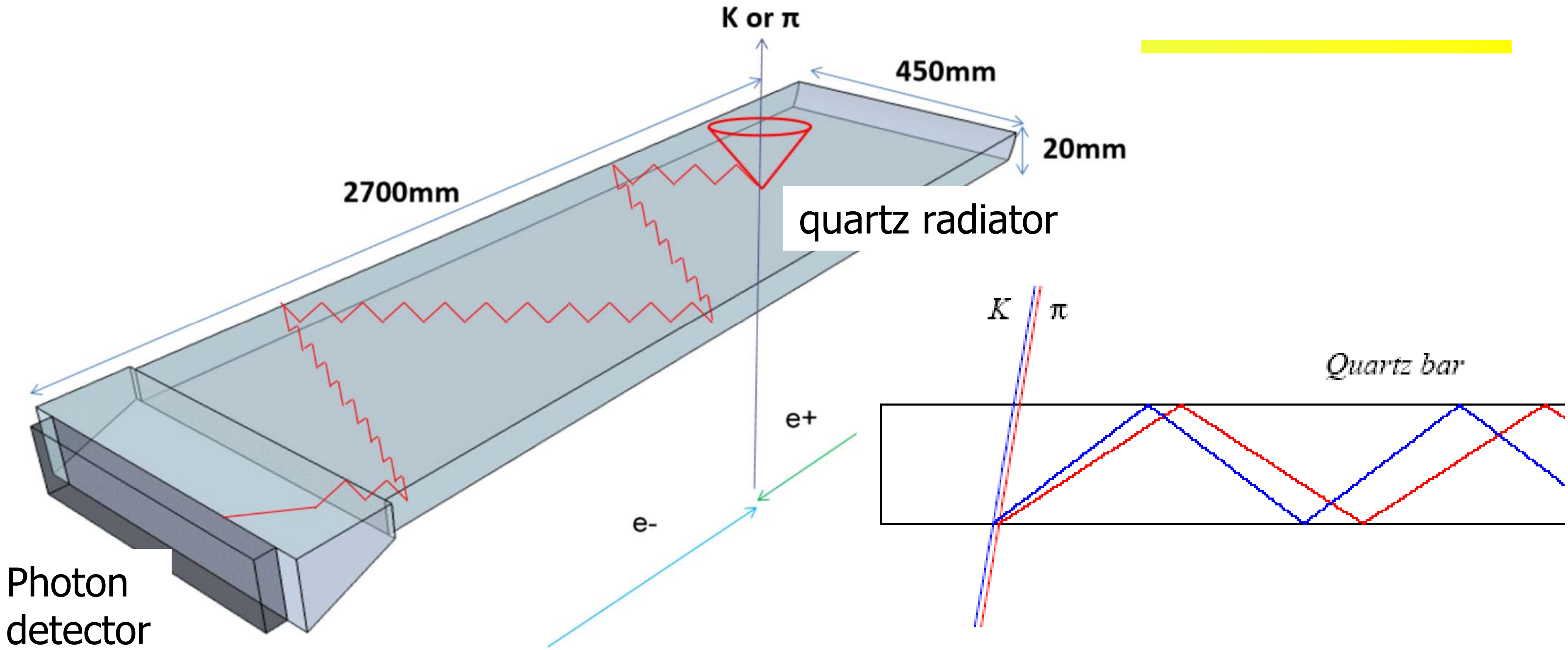
Belle II Cherenkov detectors





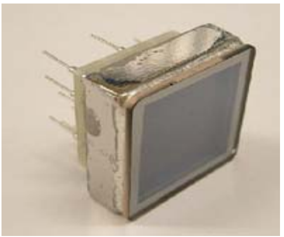


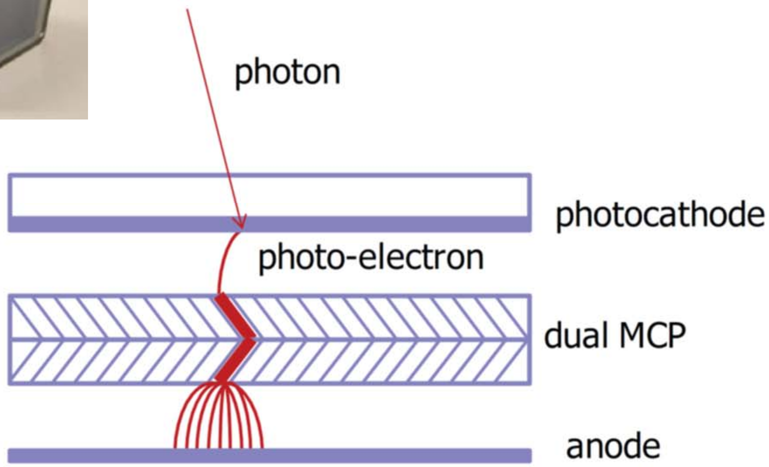
Belle II Barrel PID: Time of propagation (TOP) counter

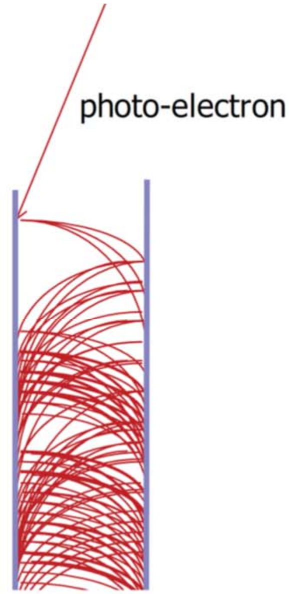


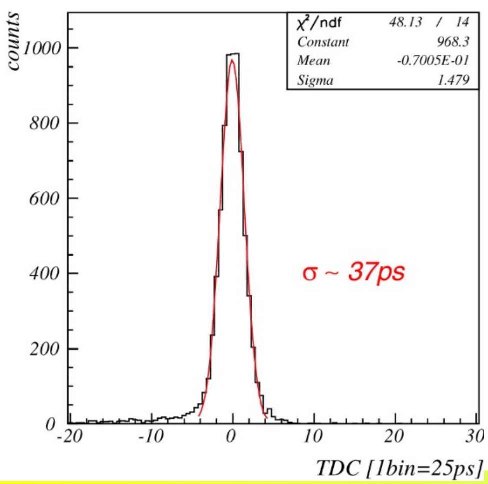
- Similar to the DIRC, Cherenkov ring imaging with precise time measurement.
- Reconstruct Cherenkov angle from two hit coordinates and the time of propagation of the photon
  - Quartz radiator (2cm thick)
  - Photon detector (MCP-PMT)
    - Excellent time resolution ~ 40 ps
    - Single photon sensitivity at 1.5 T

MCP PMTs for a very fast timing





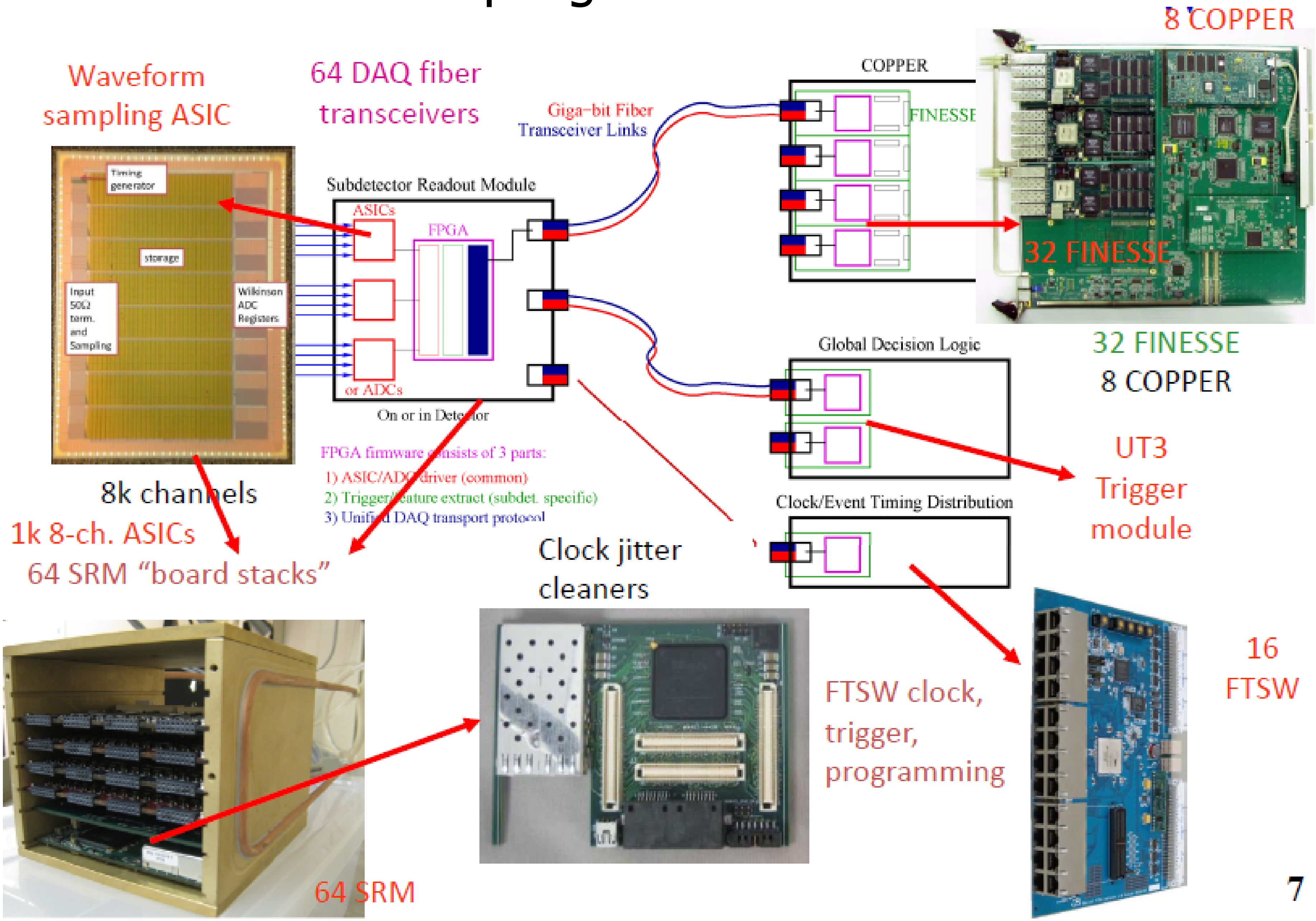


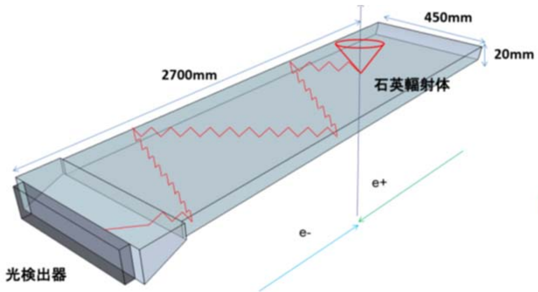


Micro-channel plate PMTs:

Single photon resolution: typically 20ps – 40ps

TOP Waveform sampling readout

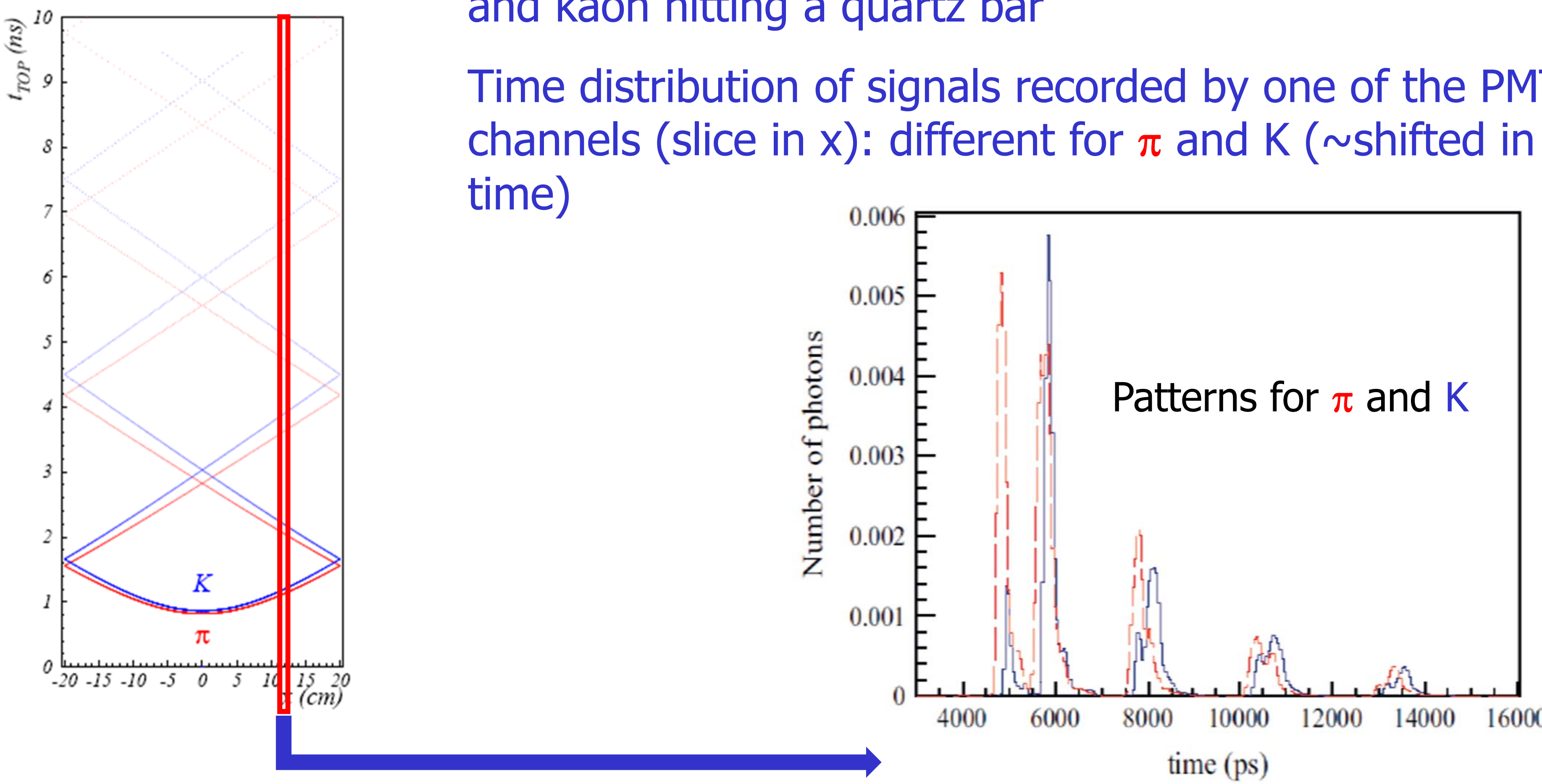




TOP image reconstruction

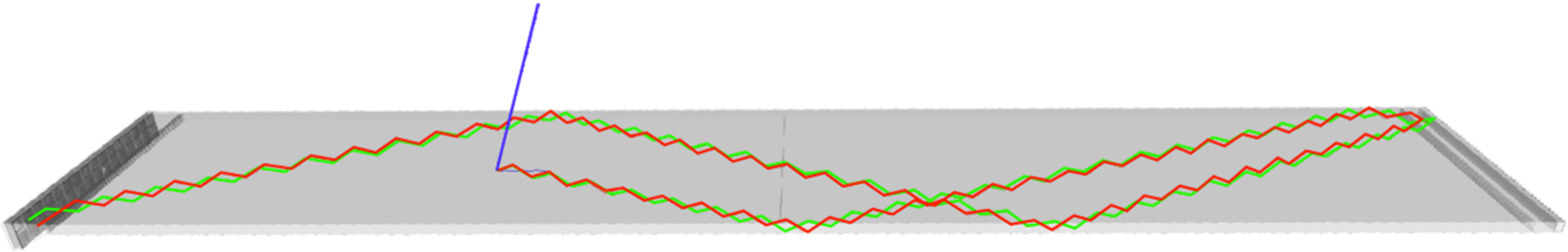
Pattern in the coordinate-time space ('ring') of a pion and kaon hitting a quartz bar

Time distribution of signals recorded by one of the PMT channels (slice in x): different for π and K (~shifted in time)

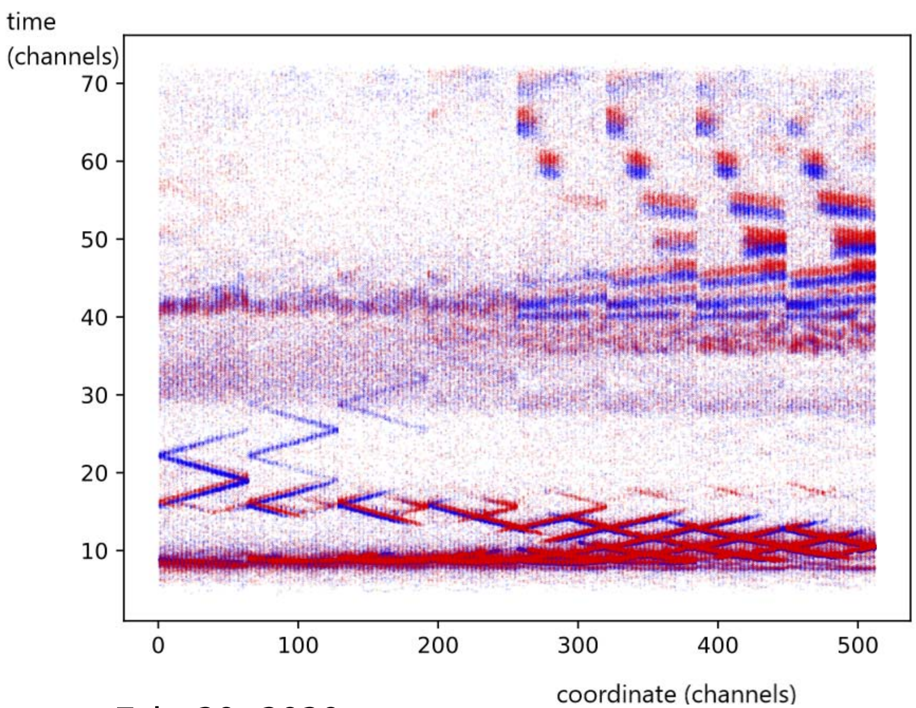


The name of the game: analytic expressions for the 2D likelihood functions
→M. Starič et al., NIMA A595 (2008) 252-255

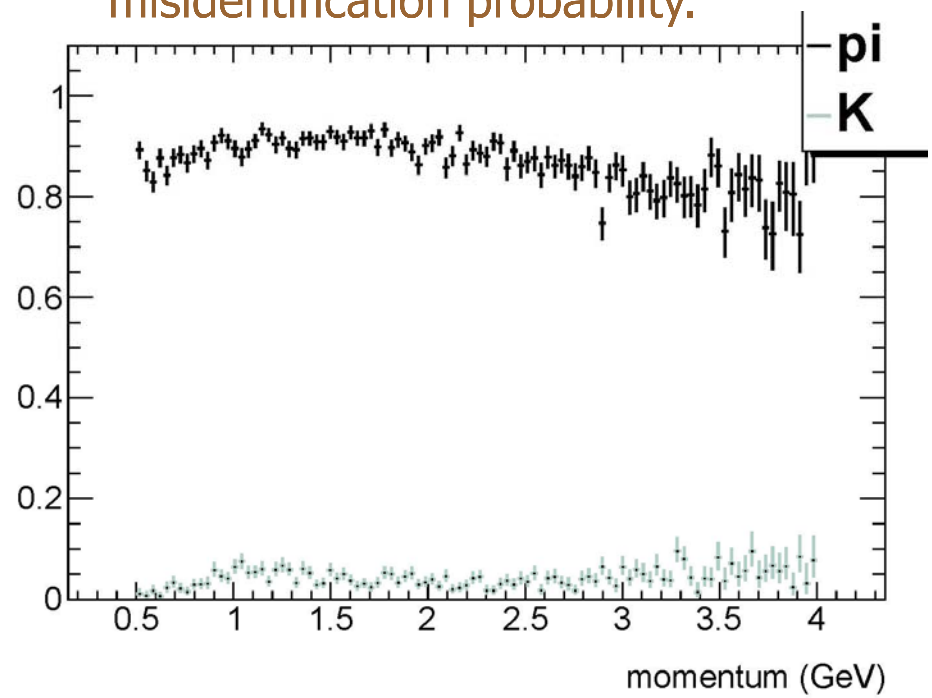
Separation of kaons and pions



Pions vs kaons in TOP:
different patterns in the time vs PMT impact point coordinate

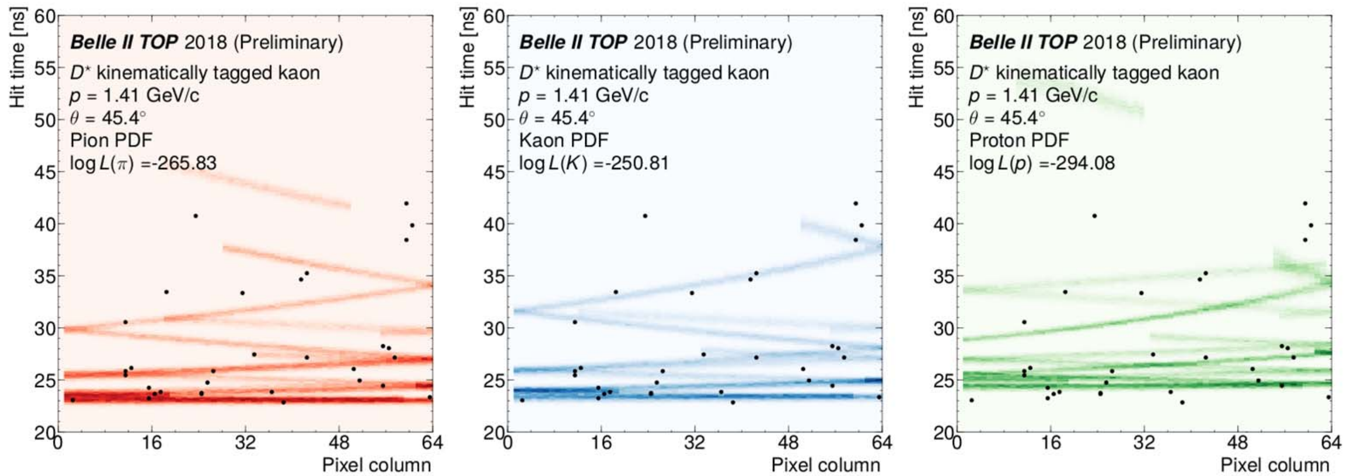


Pions vs kaons:
Expected PID efficiency and misidentification probability.

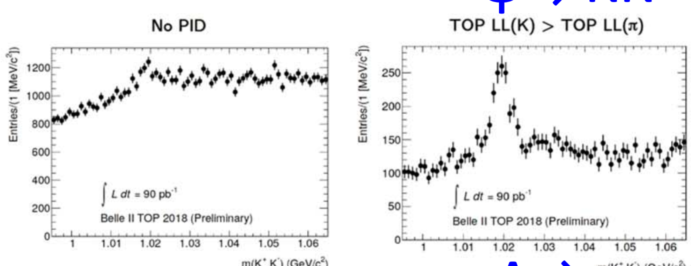


TOP first events

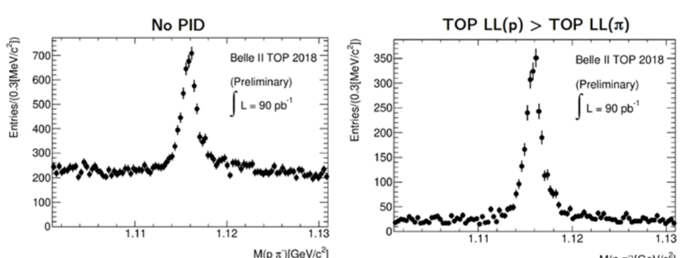
The early data demonstrated that the TOP principle is working



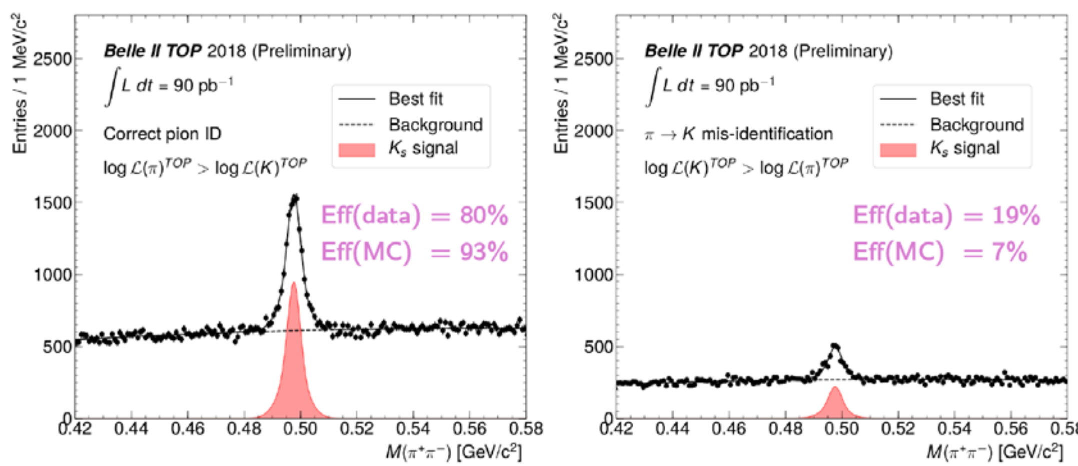
$\phi \to K^+K^-$ with both the tracks in the TOP acceptance $\phi \to KK$



$\Lambda \to p\pi$ with the **proton candidate** in the TOP acceptance $\Lambda \to p\pi$

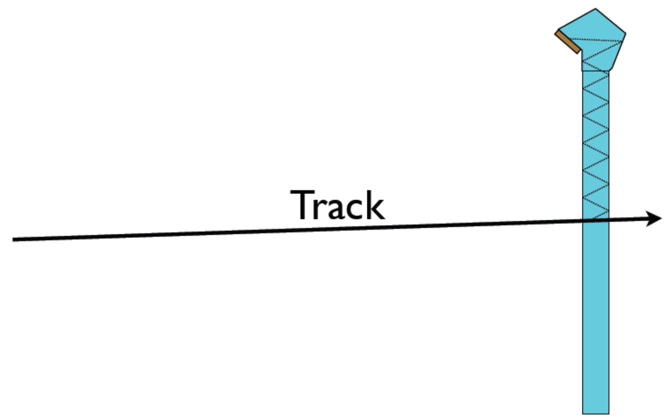


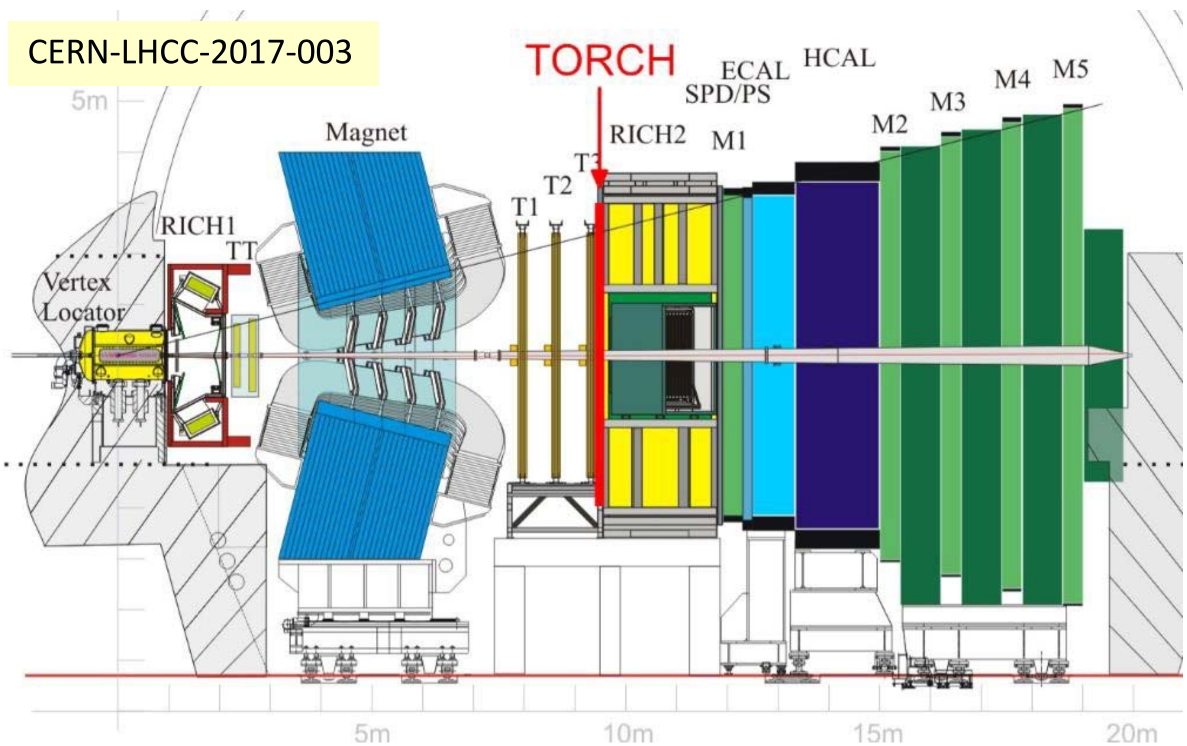
$K_s \to \pi\pi$



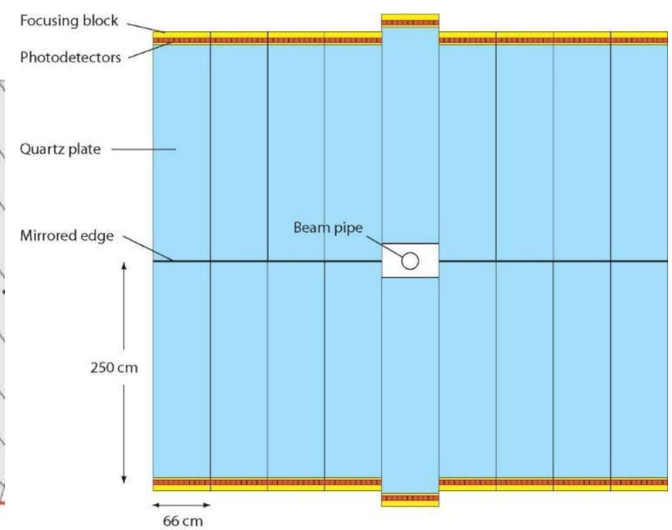
LHCb PID upgrade: TORCH

A special type of Time-of-Propagation counter for the LHCb upgrade

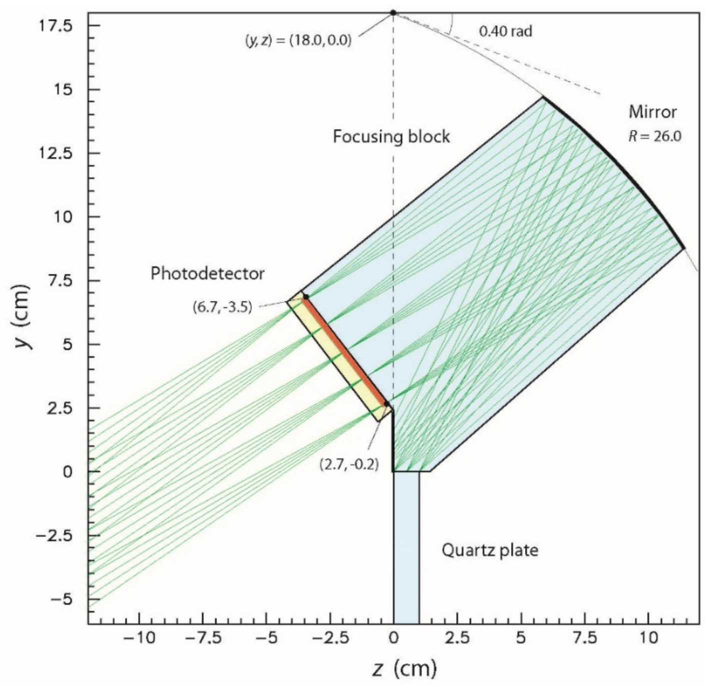




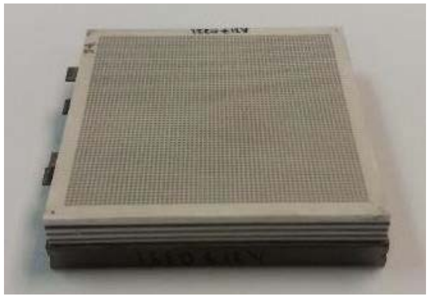
- TORCH area 5 x 6 m^2
- 18 module system
- 11 MCPs per module



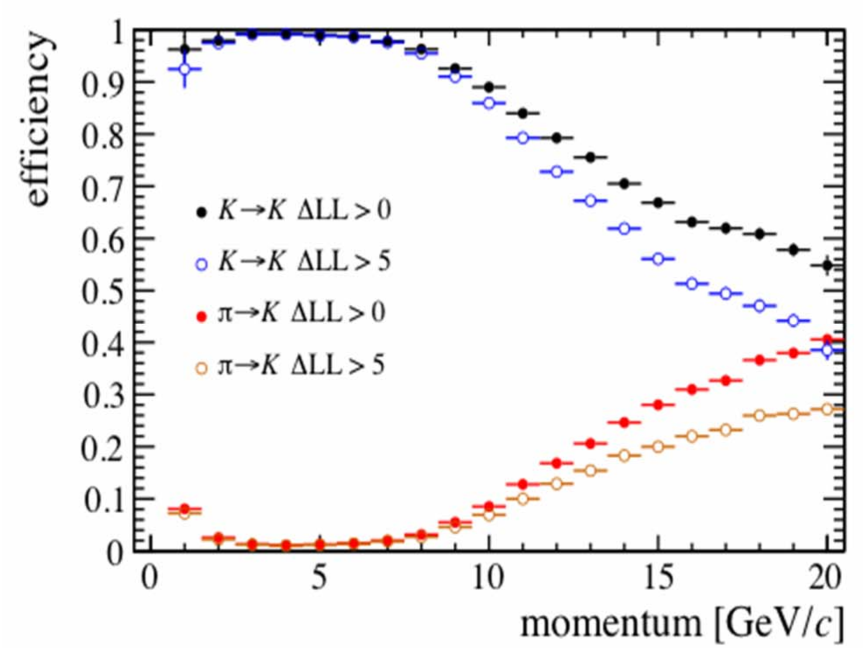
LHCb PID upgrade: TORCH

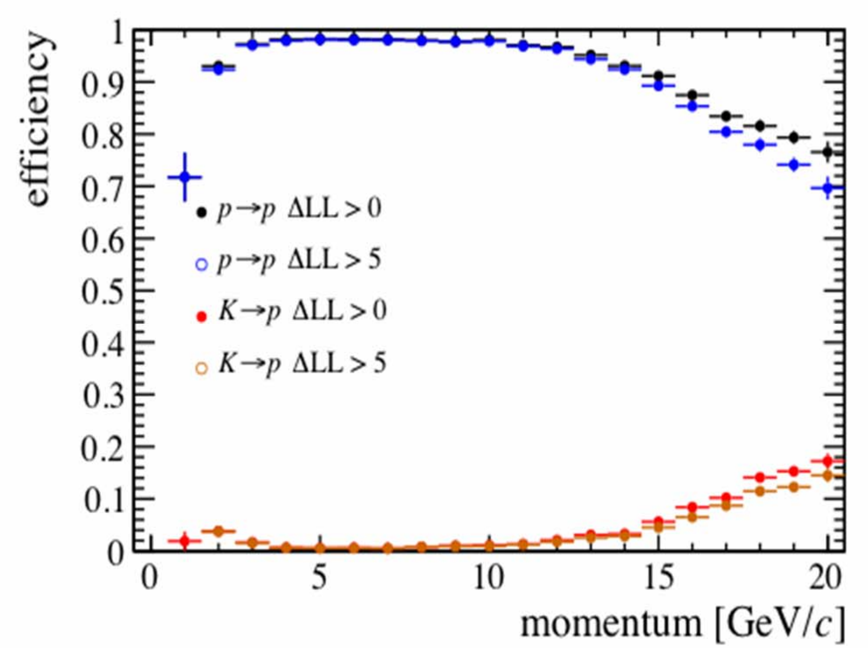


Focusing block with light sensors (MCP PMTs from Photek)

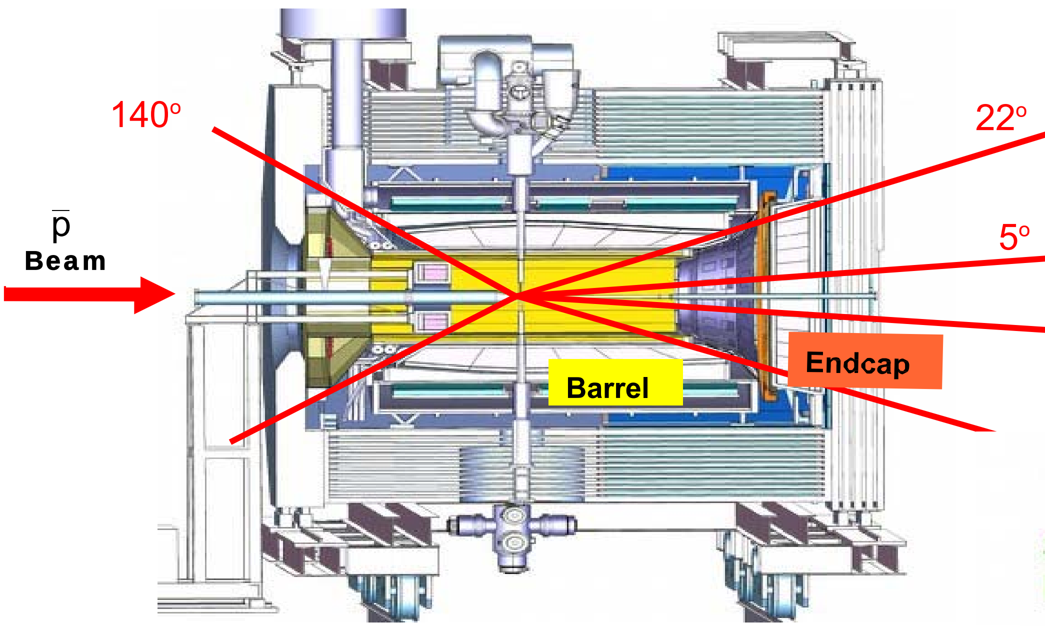


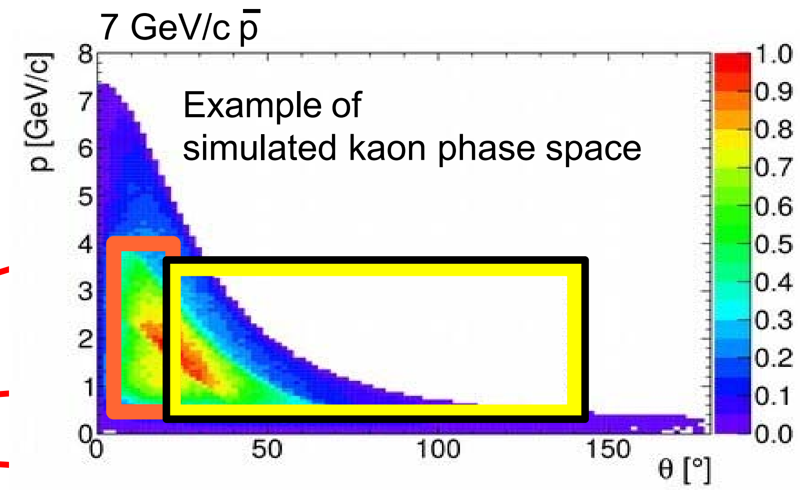
Expected performance

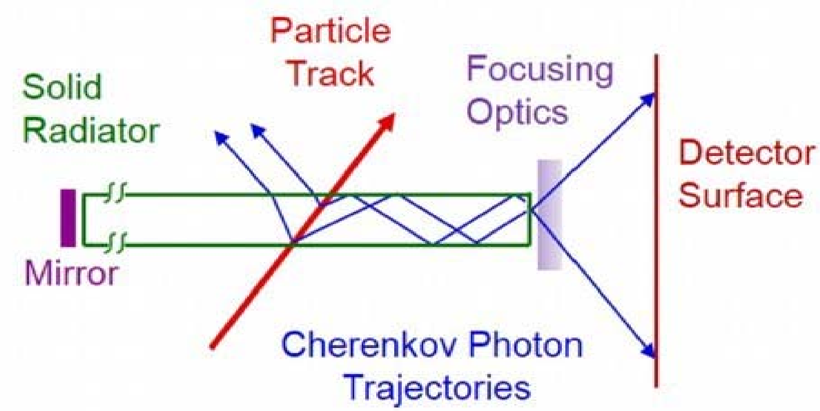




PANDA DIRC counters







Magnitude of photon angles
in radiator preserved

Barrel DIRC

Goal: 3 s.d. π/K
separation up to
3.5 GeV/c

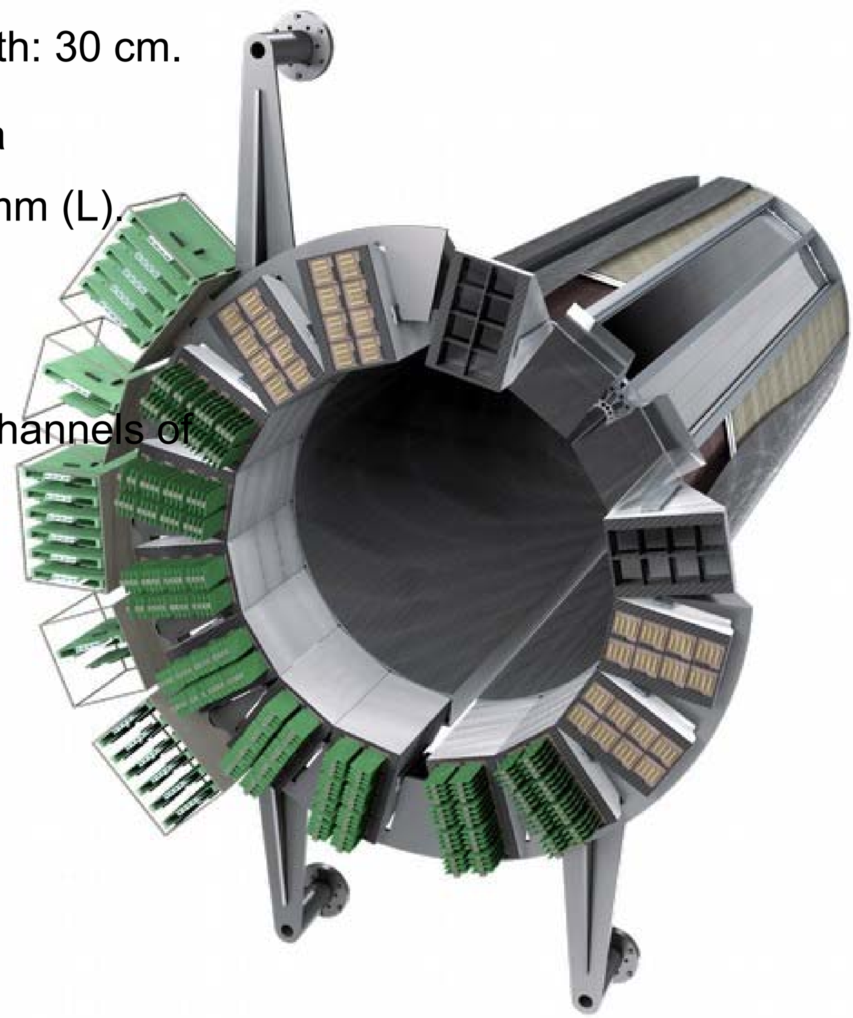
Endcap Disc DIRC

Goal: 3 s.d. π/K
separation up to
4 GeV/c

PANDA Barrel DIRC

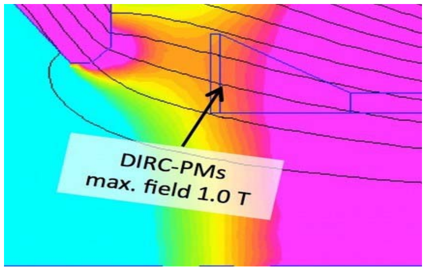
Design: based on BABAR DIRC and SuperB FDIRC with key improvements

- Barrel radius ~48 cm; expansion volume depth: 30 cm.
- 48 narrow radiator bars, synthetic fused silica
  17 mm (T) x 53 mm (W) x 2400 mm (L).



- Compact photon detector:
  30 cm fused silica expansion volume 8192 channels of MCP-PMTs
  in ~1T B field
- Focusing optics: spherical lens system
- Fast photon detection:
  fast TDC plus TOT electronics,
  → 100-200 ps timing

Photon detector



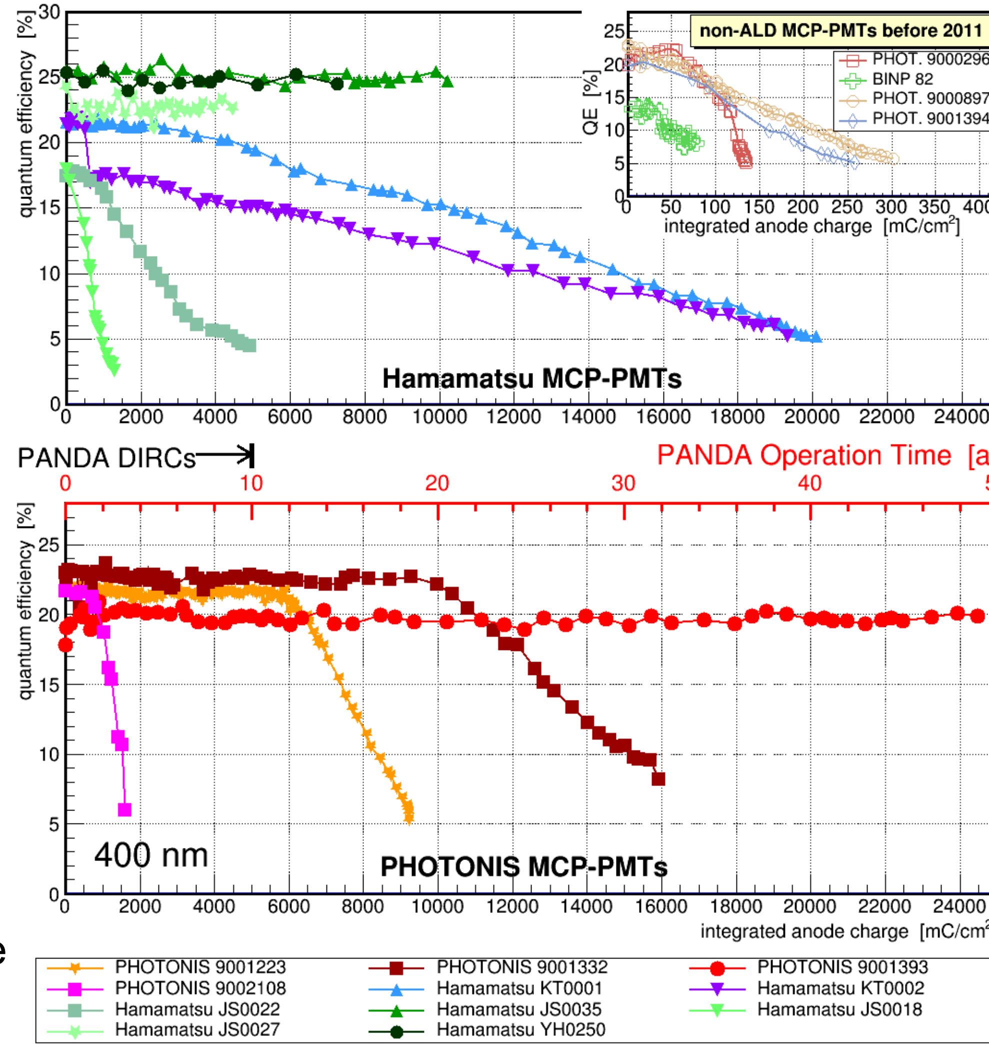
Requirements:

- **few mm** spatial resolution
- **~100 ps** timing resolution

Bar-box:

8 MCP-PMT, 512 pixels (total 8 k readout channels) with **pixel size 6 x 6 mm²** work in **1T magnetic field** survive **10 years** of PANDA (aging)

Most sensors with **ALD coated** MCPs have **lifetime > 5 C/cm²**



Panda Disc DIRC

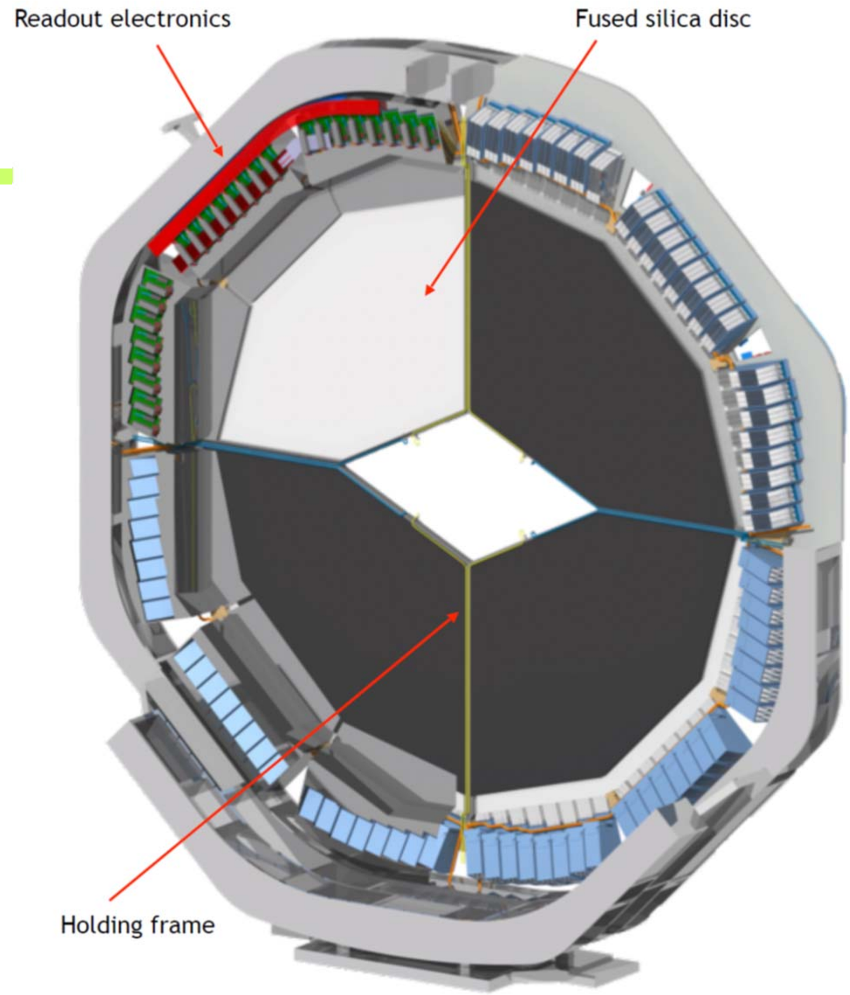
Radiator: fused silica 20 mm thick, *R* =1m

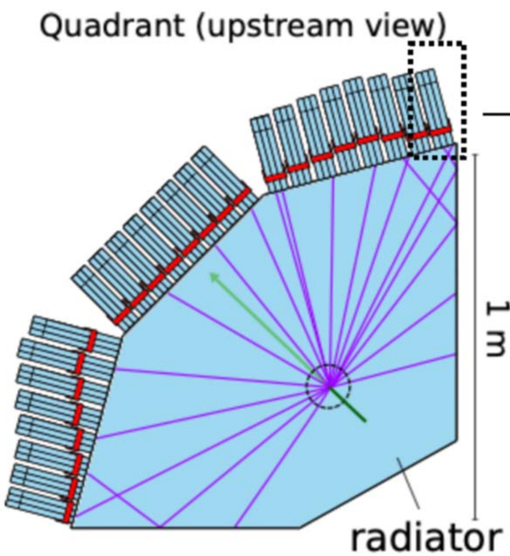
π/K separation up to 4 GeV/c

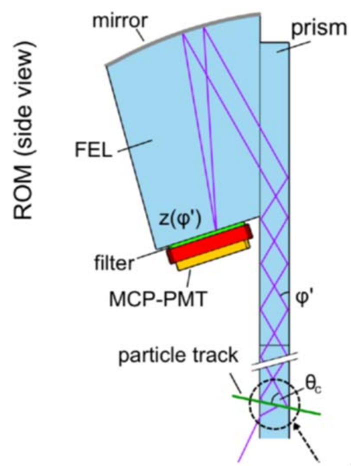
Focusing optics

Photon detector in ~1T field:

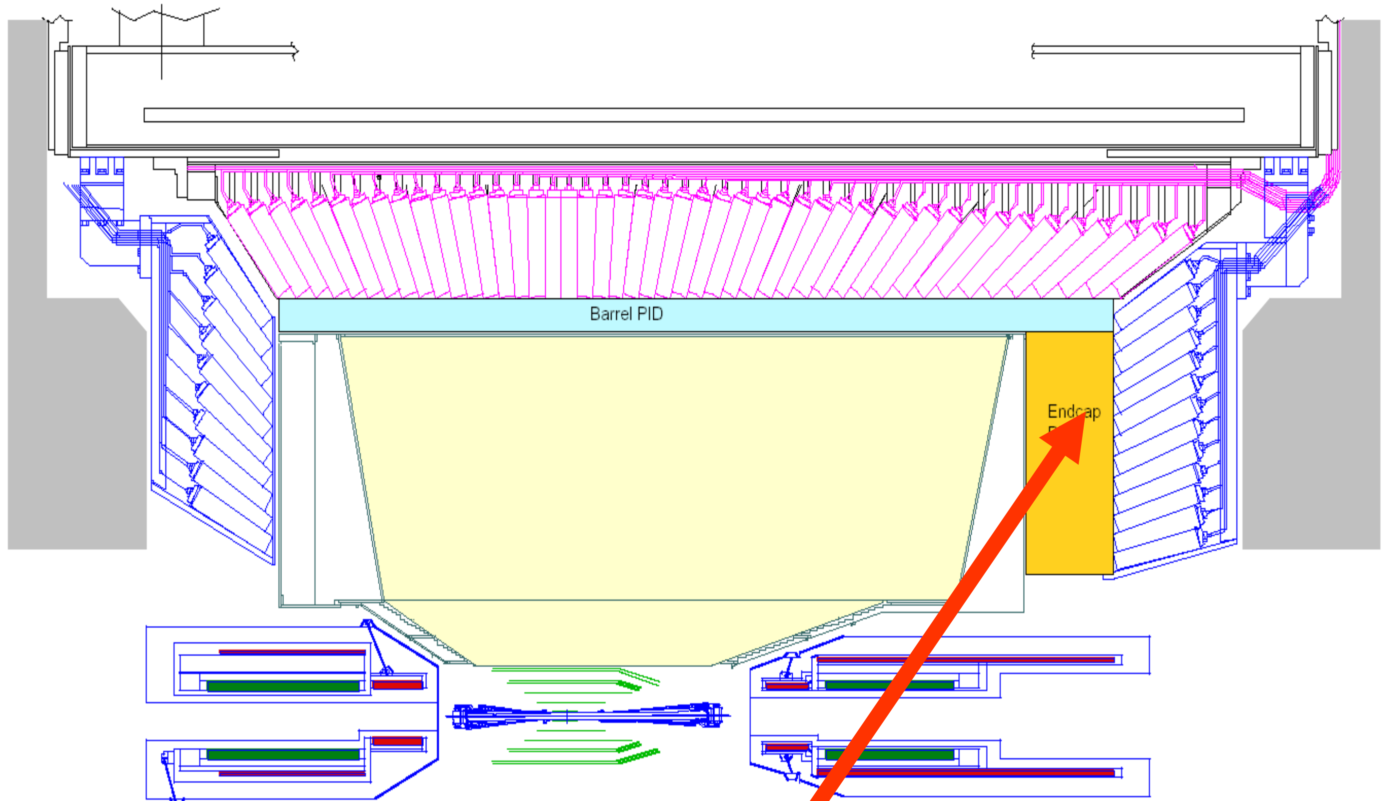
- 96 MCP PMTs with a highly segmented anode, TofPET2 readout ASIC







Belle II PID system

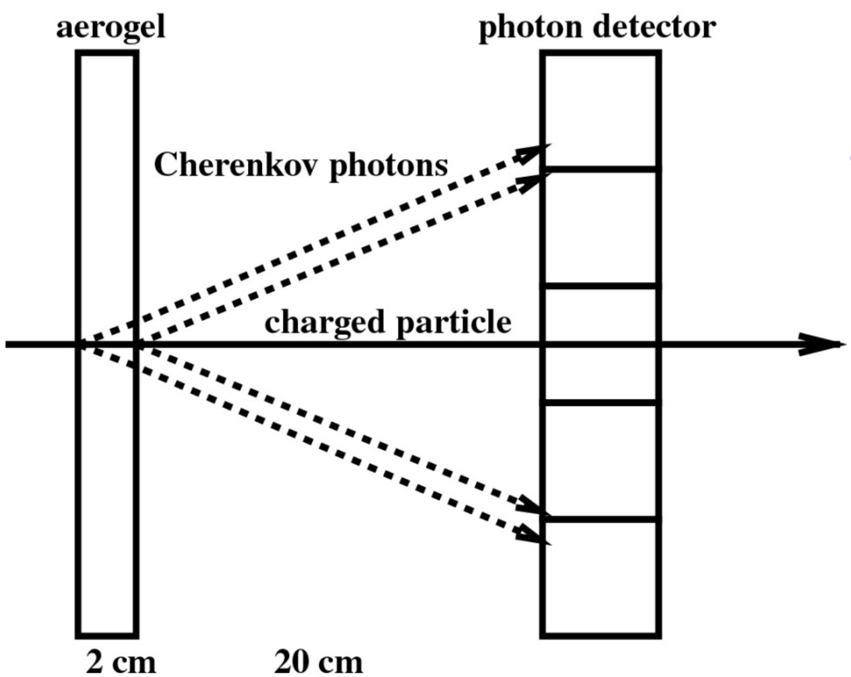


Two novel particle ID devices, both RICHes:

Barrel: Time-of-propagation counter (TOP) counter

Endcap: proximity focusing RICH

Endcap: Proximity focusing RICH



K/π separation at 4 GeV/c:
$\theta_c(\pi)$ ~ 308 mrad (n = 1.05)
$\theta_c(\pi)$– $\theta_c(K)$ ~ 23 mrad

For single photons: $\delta\theta_c$(meas.)=σ_0 ~ 14 mrad,

typical value for a 20mm thick radiator and 6mm PMT pad size

Per track:

$$\sigma_{track} = \frac{\sigma_0}{\sqrt{N_{pe}}}$$

Separation: $[\theta_c(\pi)-\theta_c(K)]/\sigma_{track}$

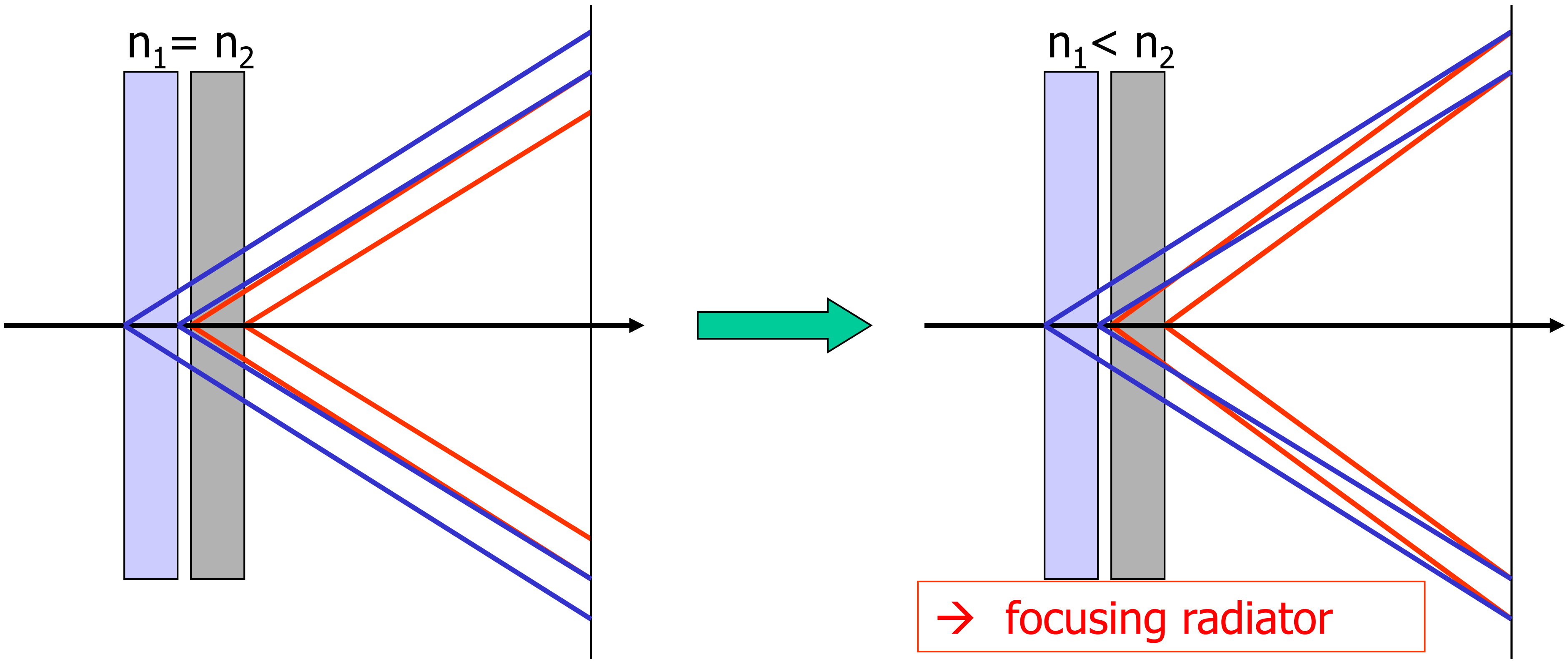
→ 5σ separation with N_{pe}~10

Radiator with multiple refractive indices

Small number of photons from aerogel → need a thick layer of aerogel.
How to improve the resolution by keeping the same number of photons?

normal

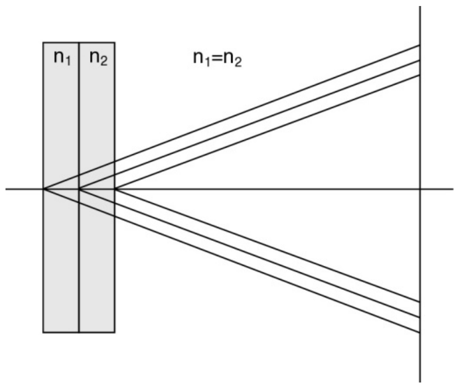
→ stack two tiles with different refractive indices:
"focusing" configuration

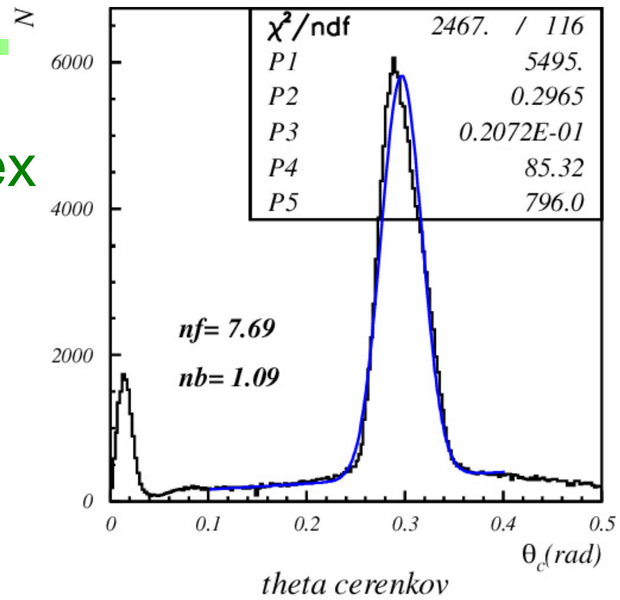


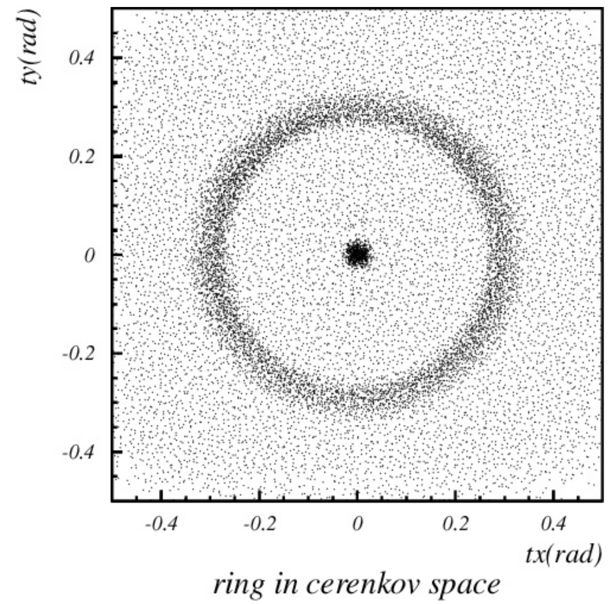
→ focusing radiator

Focusing configuration – data

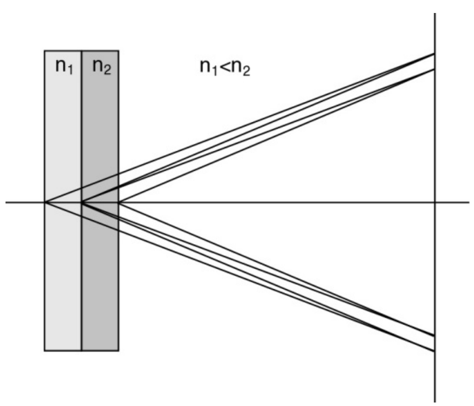
4cm aerogel single index

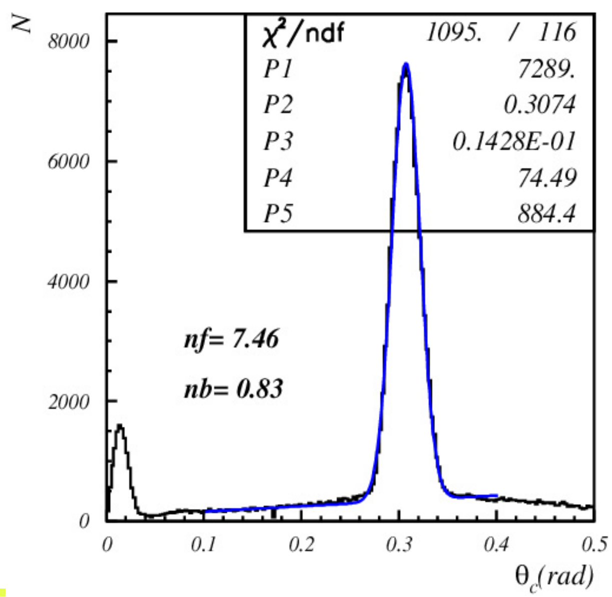


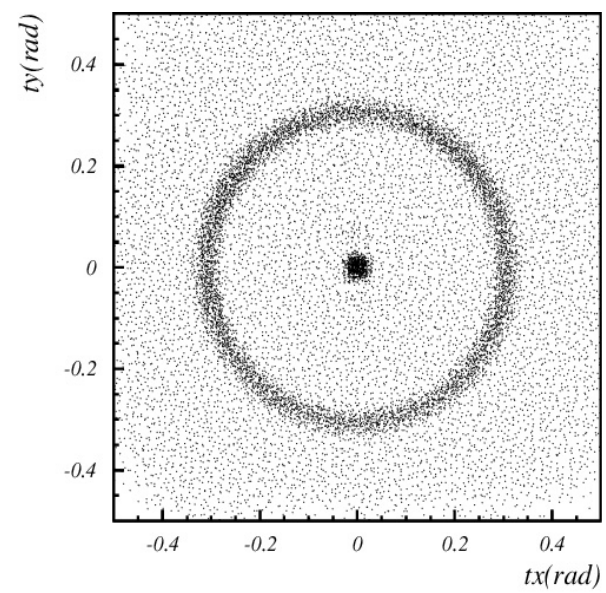




2+2cm aerogel





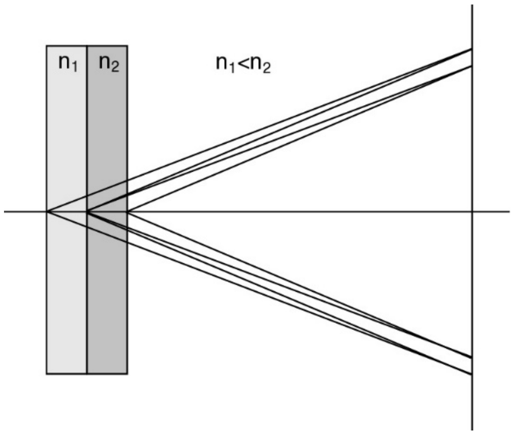


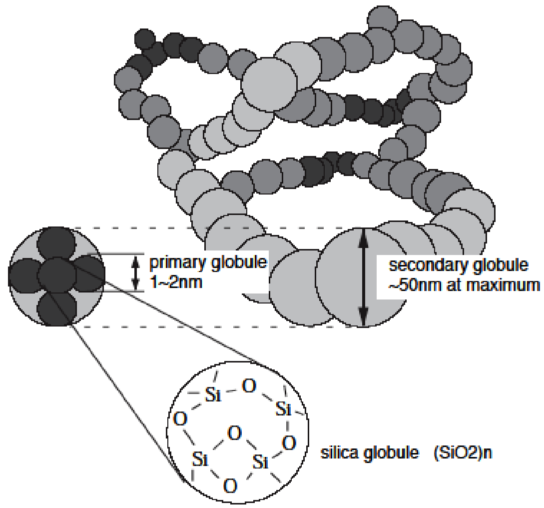
→NIM A548 (2005) 383, NIMA 565 (2006) 457

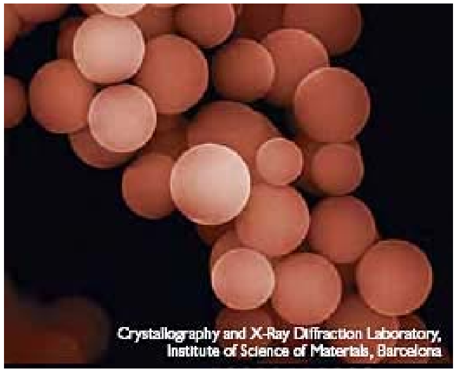
4x4 array of flat pannel MAPMTs

Radiator with multiple refractive indices 2

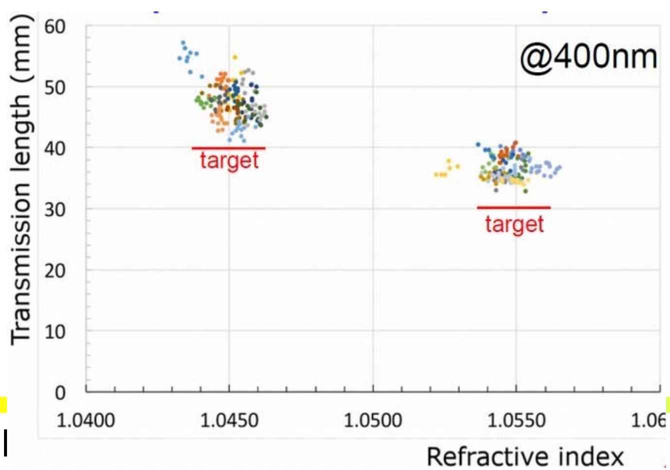
Such a configuration is only possible with aerogel (a form of Si_xO_y) – material with a tunable refractive index between 1.01 and 1.07.

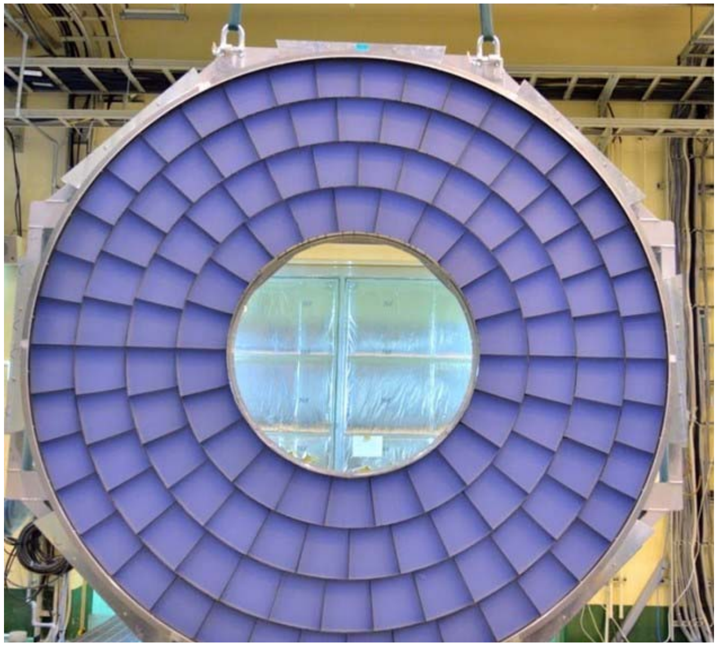






Requires aerogel with high transparency





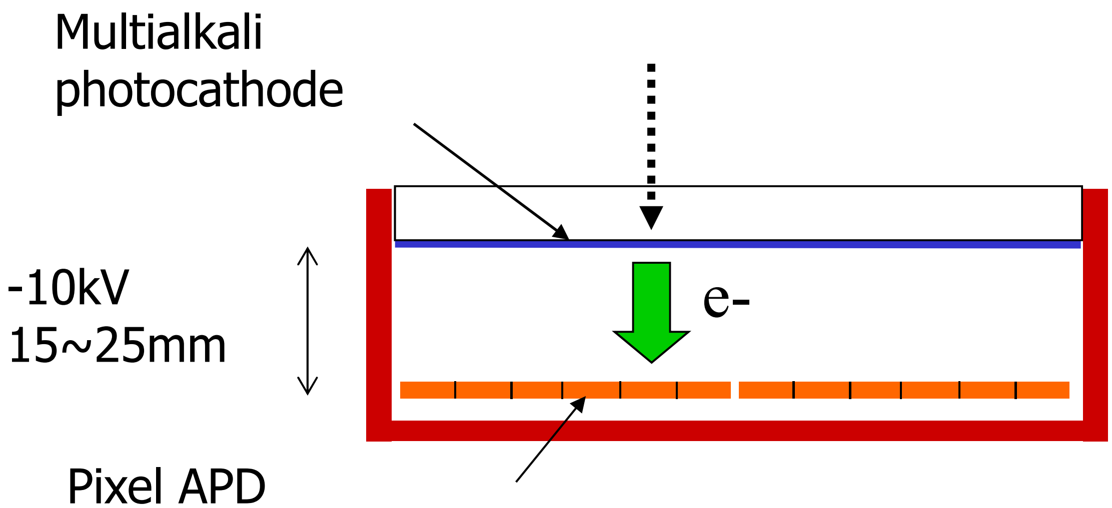
Detector plane covered with 2 x 124 tiles water-jet cut tiles (~ 17x17cm)

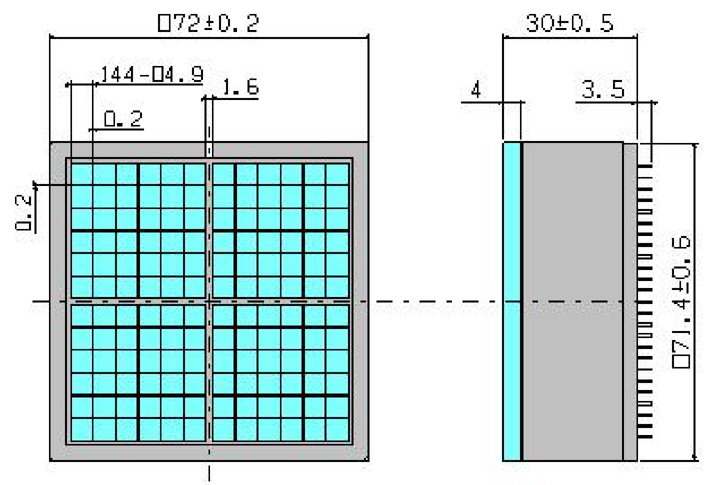
Photon detectors for the aerogel RICH requirements and candidates

Need: Operation in a high magnetic field (1.5 T)
Pad size ~5-6mm

Final choice: large active area HAPD (hybrid avalanche photon detector) of the proximity focusing type

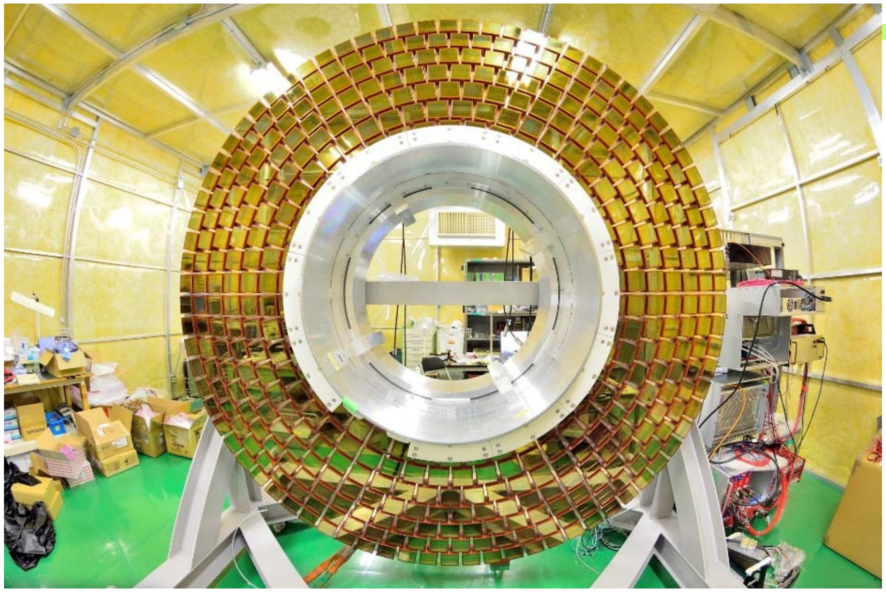
Other candidates: MCP PMT (Photonis 85011), SiPMs

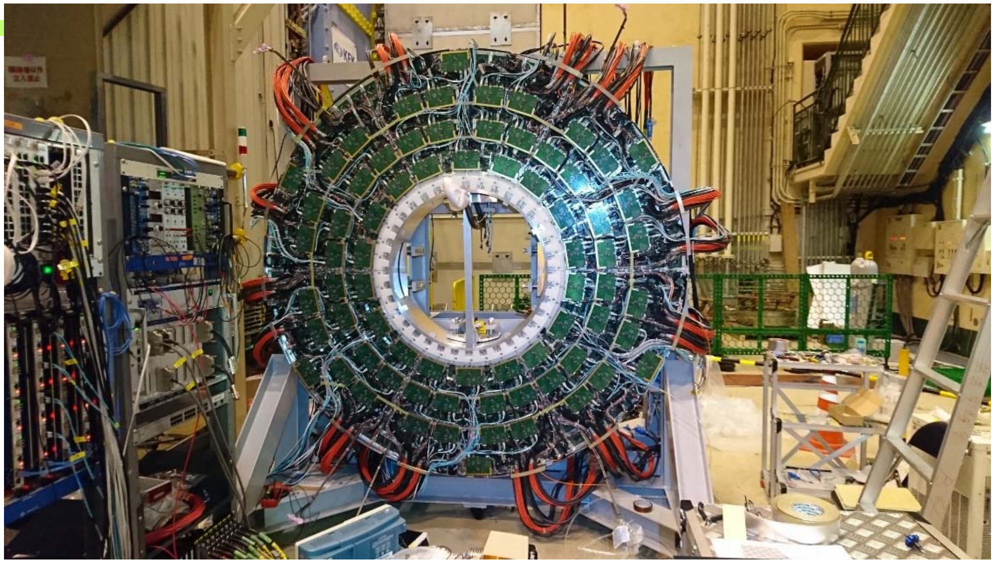


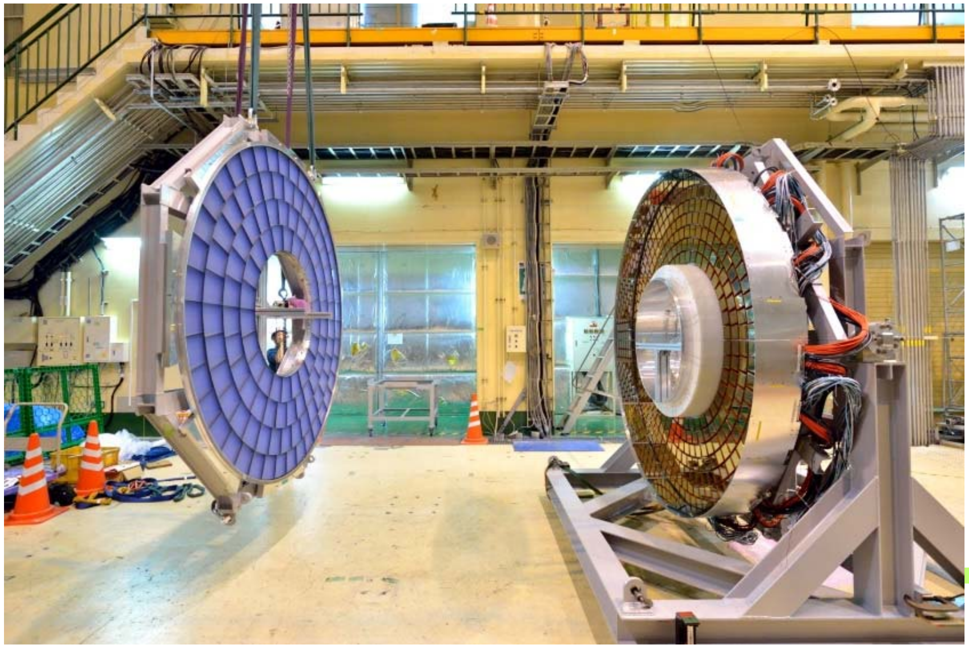


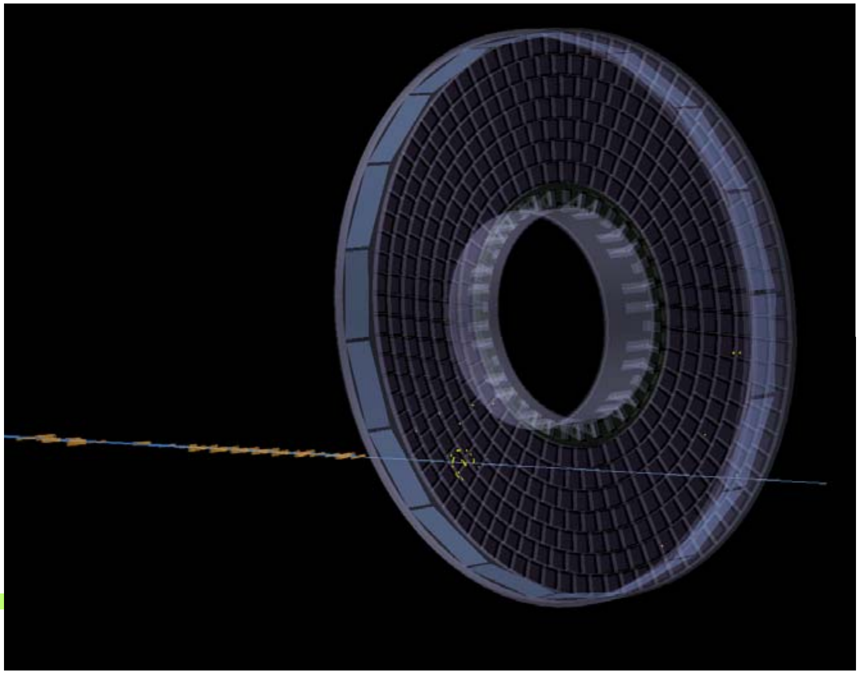
HAPD R&D project in collaboration with HPK.

The big eye of ARICH – and one of the first rings



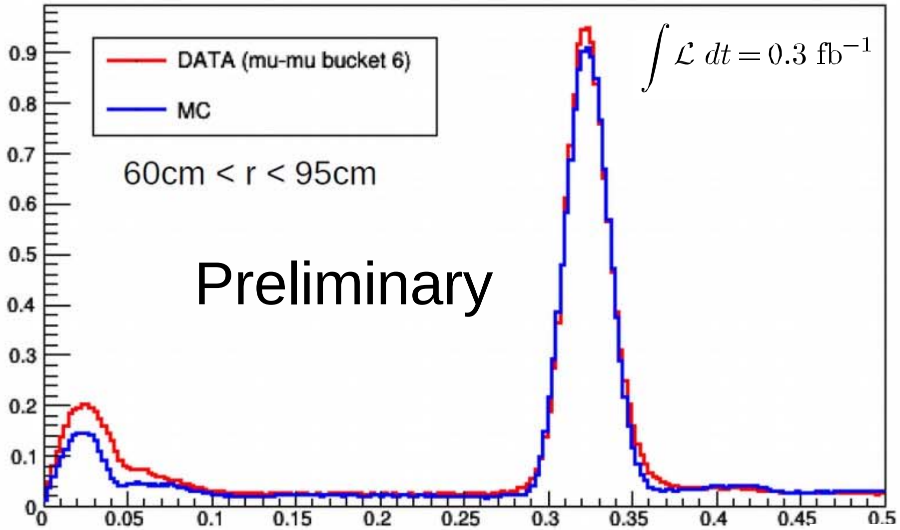






Performance in the early Belle II data

Cherenkov angle distribution in $e^+e^- \to \mu^+\mu^-$



DATA

$N_{sig} = 11.38/\mathrm{track}$

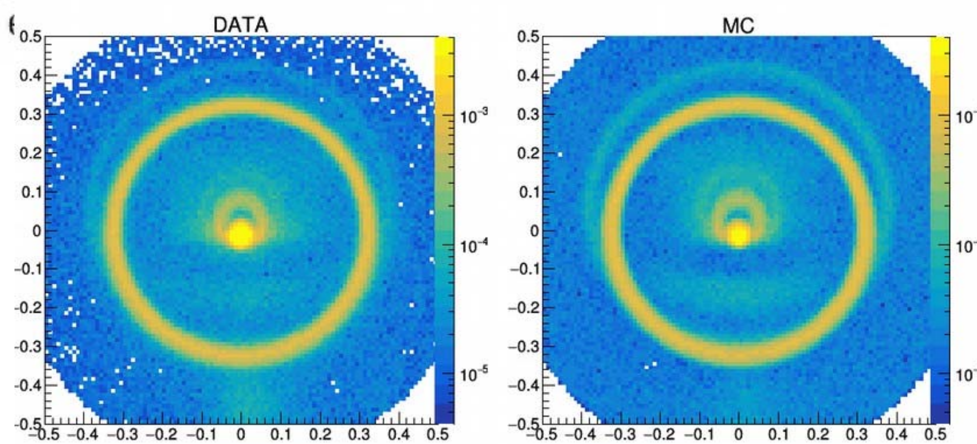
$\sigma_c = 12.7\ \mathrm{mrad}$

MC

$N_{sig} = 11.27/\mathrm{track}$

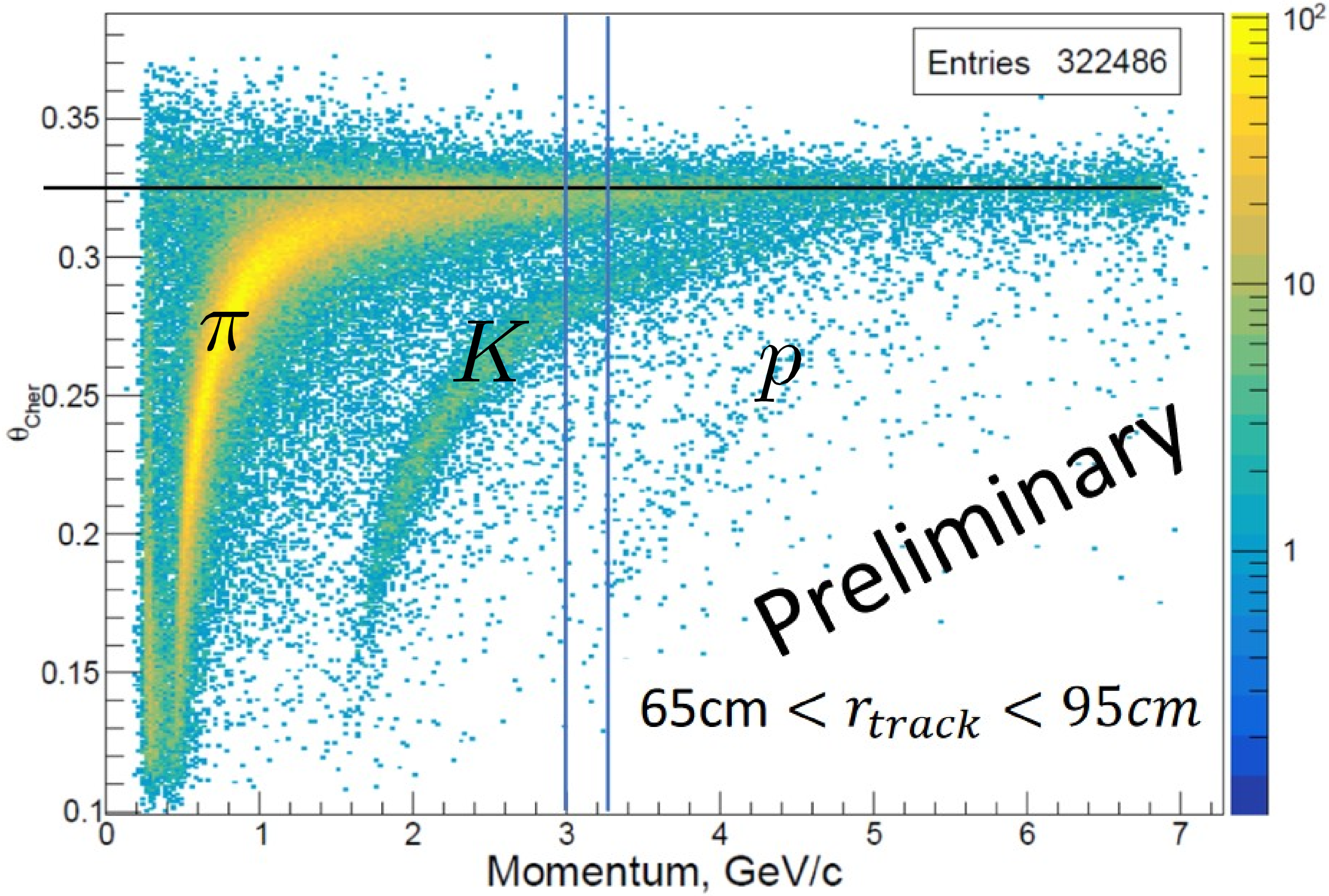
$\sigma_c = 12.75\ \mathrm{mrad}$

Overall a very good DATA/MC agreement !



Chrenkov ring (accumulated)

Cherenkov angle vs momentum in hadronic events



Average Cherenkov angle for tracks from hadronic events

Estimation of π/K separation power using $D^{*\pm}$ decays

- Identify K, π based on track charge in association with the charge of π_{slow}

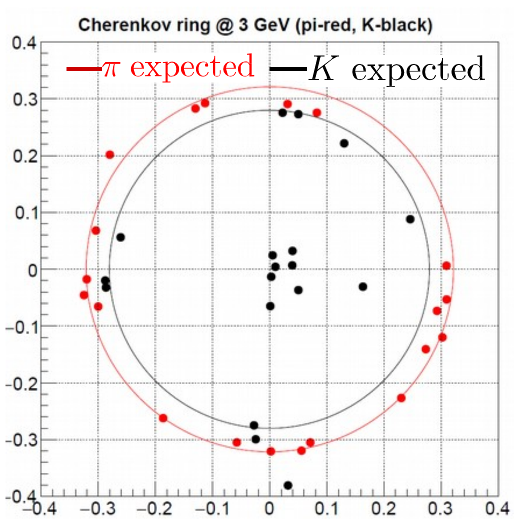
$$D^{*+} \to D^0 \, \pi^+_{\text{slow}} \qquad D^0 \searrow K^- \, \pi^+$$

$$D^{*-} \to \bar{D}^0 \, \pi^-_{\text{slow}} \qquad \bar{D}^0 \searrow K^+ \, \pi^-$$

- Apply selection criteria on

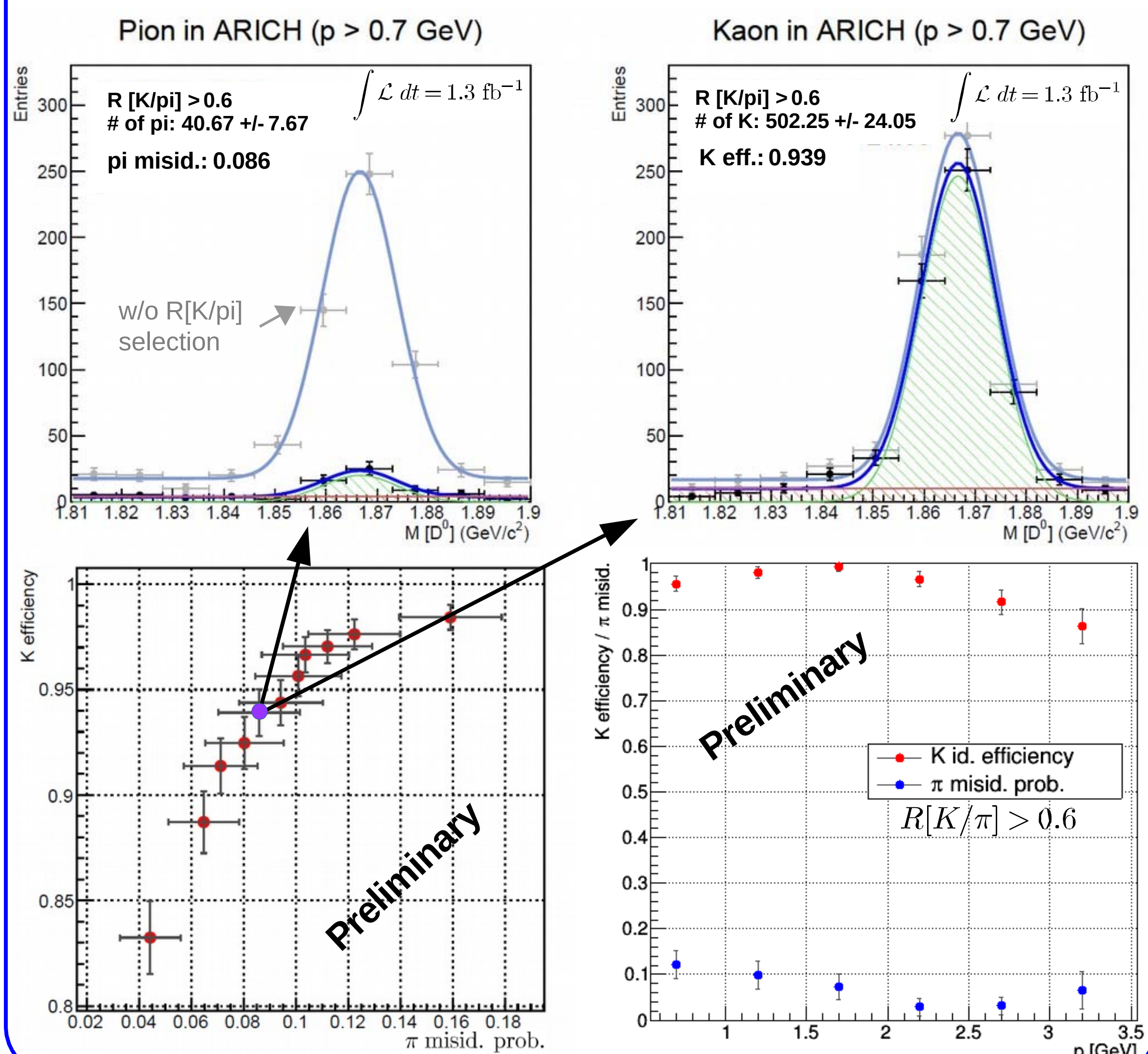
$$R[K/\pi] = \frac{\mathcal{L}_K}{\mathcal{L}_K + \mathcal{L}_\pi}$$

$\mathcal{L}$ - likelihood for given id. hypothesis

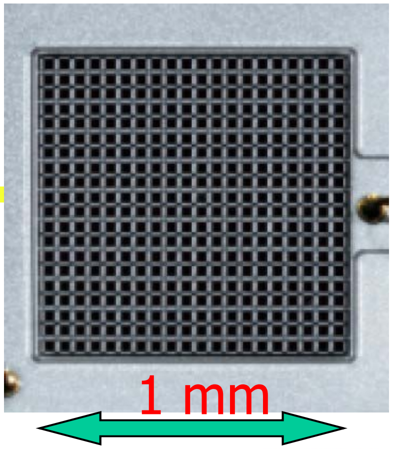


- Only coarse/preliminary calibrations included
  → further improvements expected

D^0 mass peak



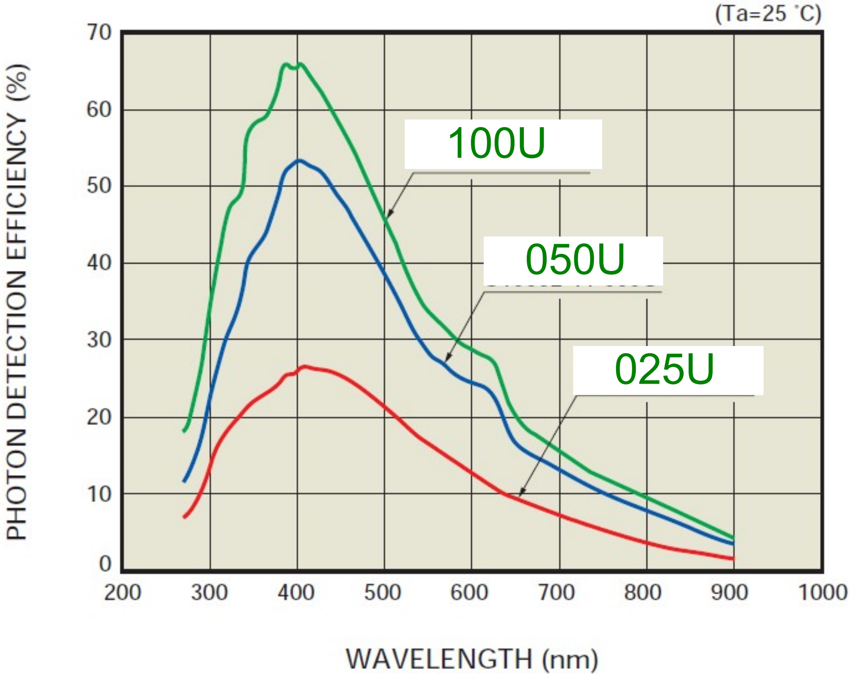
SiPMs as photon detectors?



SiPM is an array of APDs operating in Geiger mode. Characteristics:

- low operation voltage ~ 10-100 V
- gain ~ 10^6
- peak PDE up to 65%(@400nm)

  PDE = QE x ε_{geiger} x ε_{geo} (up to 5x PMT!)
- ε_{geo} – dead space between the cells
- time resolution ~ 100 ps
- works in high magnetic field
- dark counts ~ few 100 kHz/mm^2
- radiation damage (p,n)



Not trivial to use in a RICH where we have to detect single photons!

Dark counts have single photon pulse heights (rate 0.1-1 MHz per mm^2)

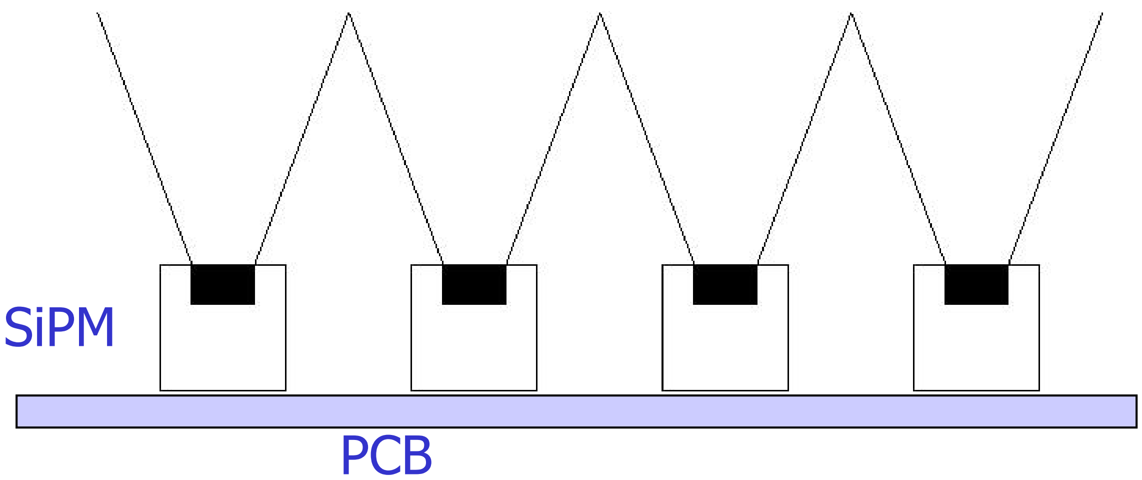
SiPM as photosensor for a RICH counter

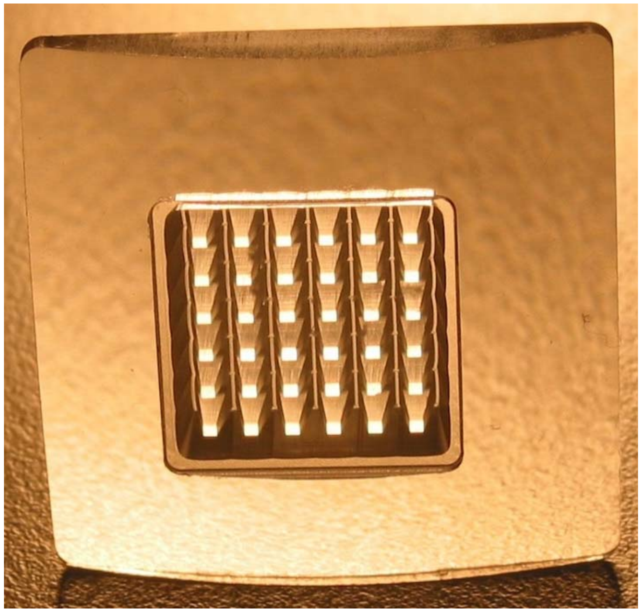
Improve the signal to noise ratio:

•Reduce the noise by a narrow (<10ns) time window (Cherenkov light is prompt!)

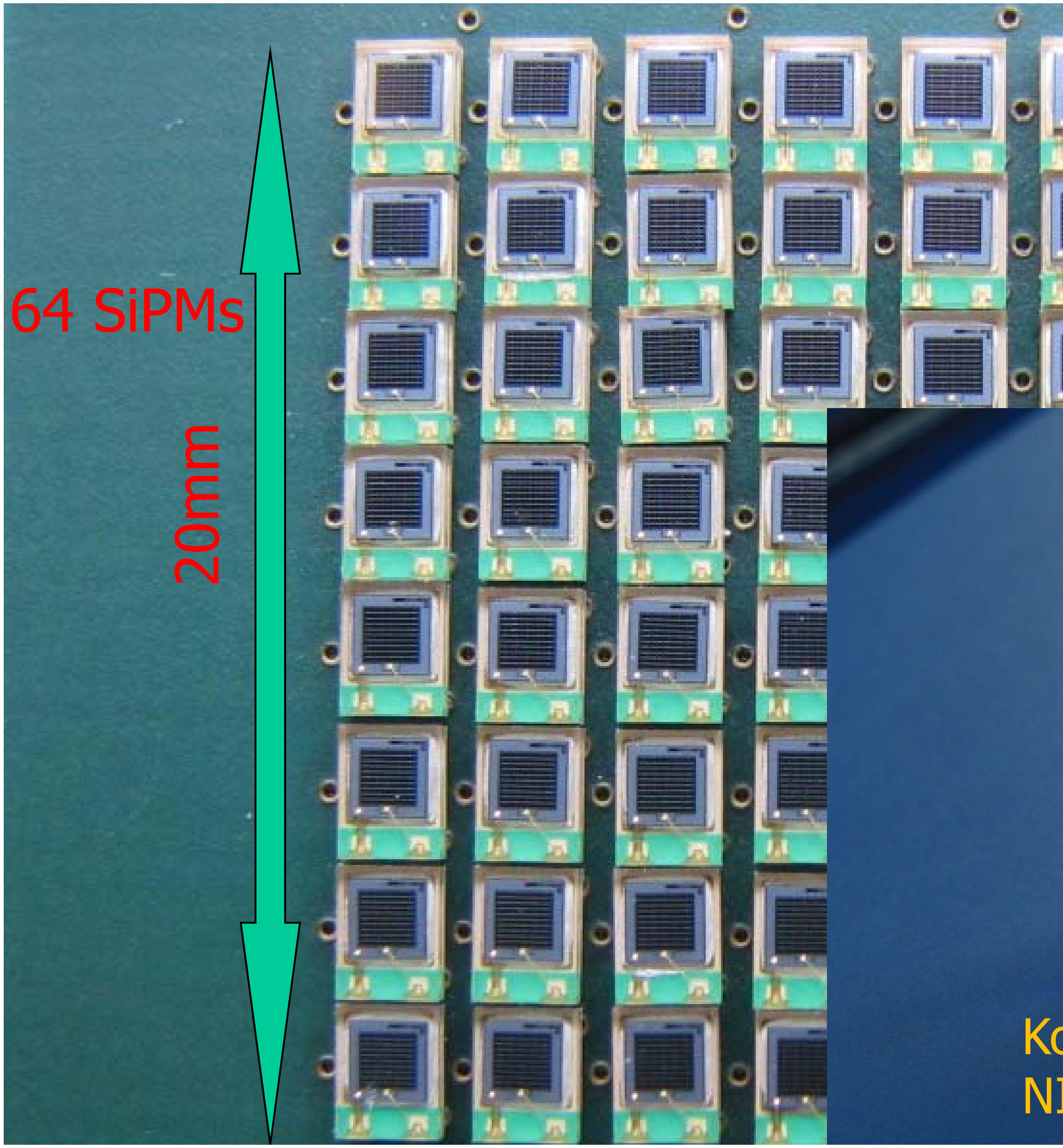
•Increase the number of signal hits per single sensor by using light collectors

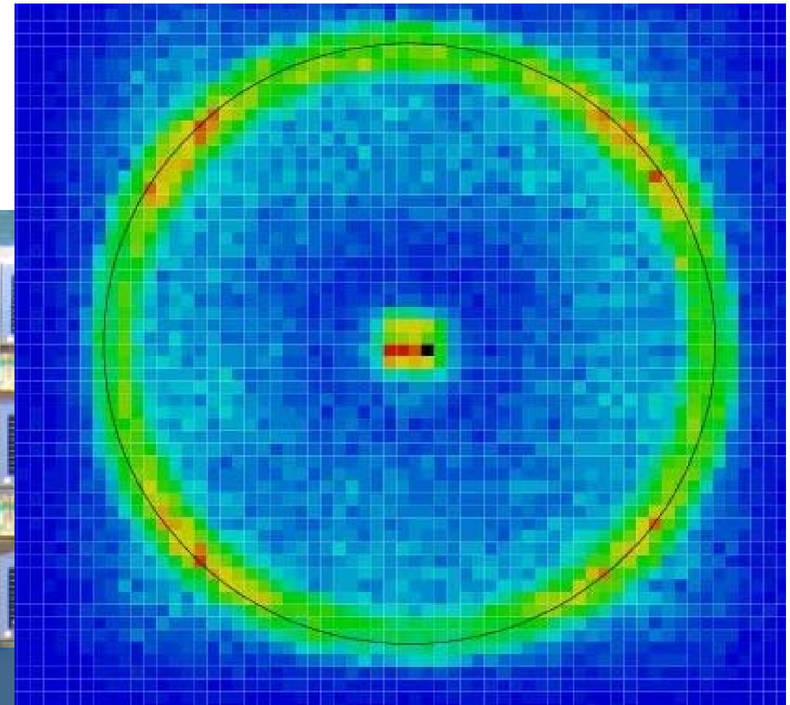
E.g. light collector with reflective walls or plastic light guide

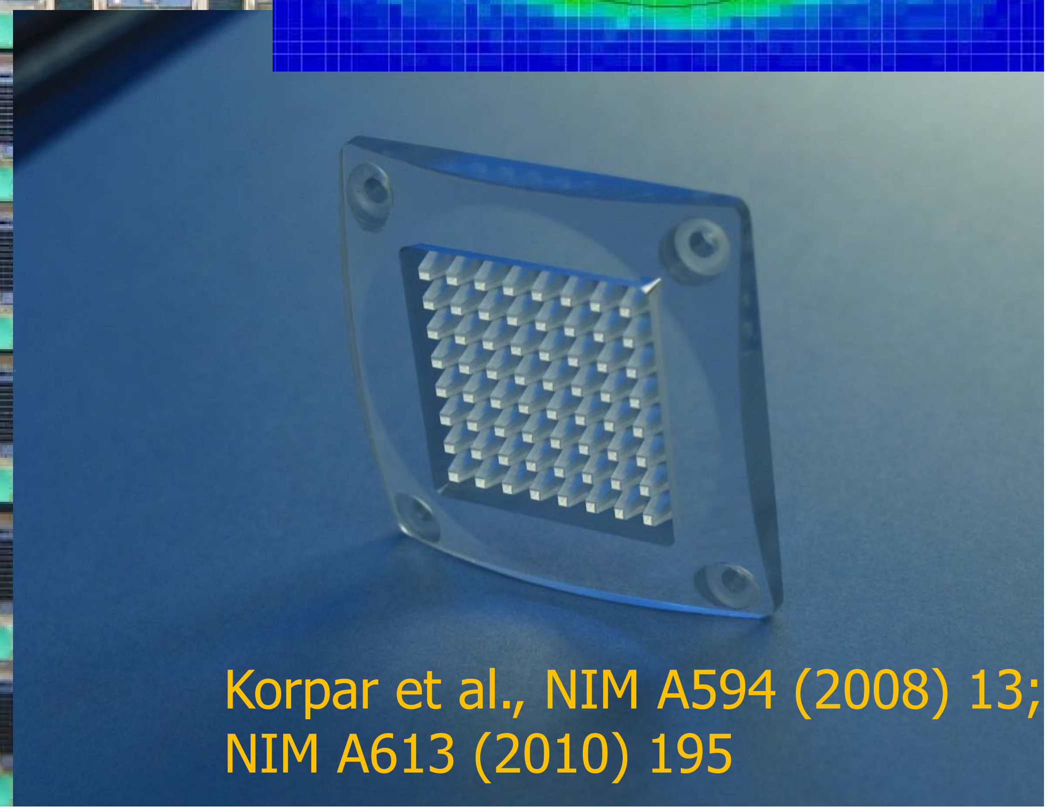




Photon detector with SiPMs and light guides



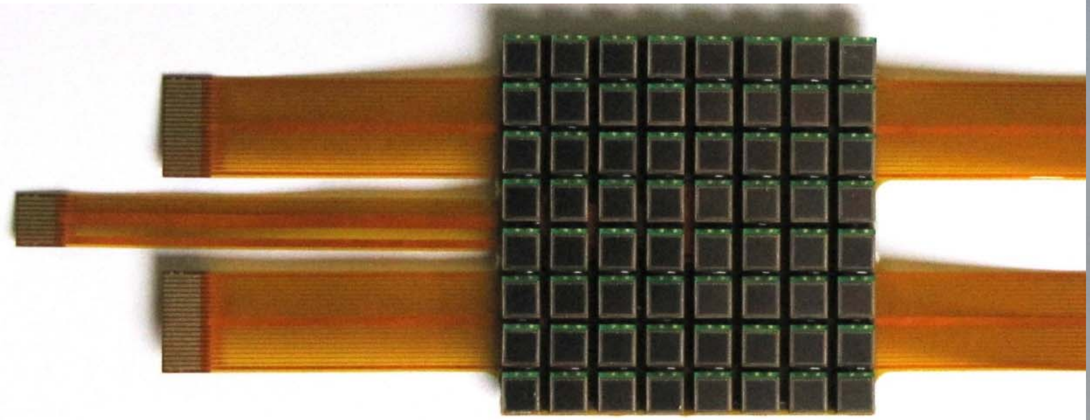




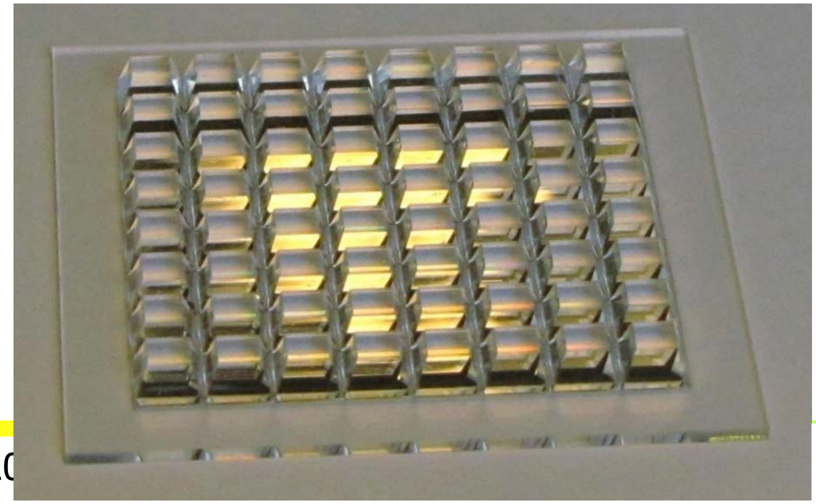
Next step: use arrays of SiPMs

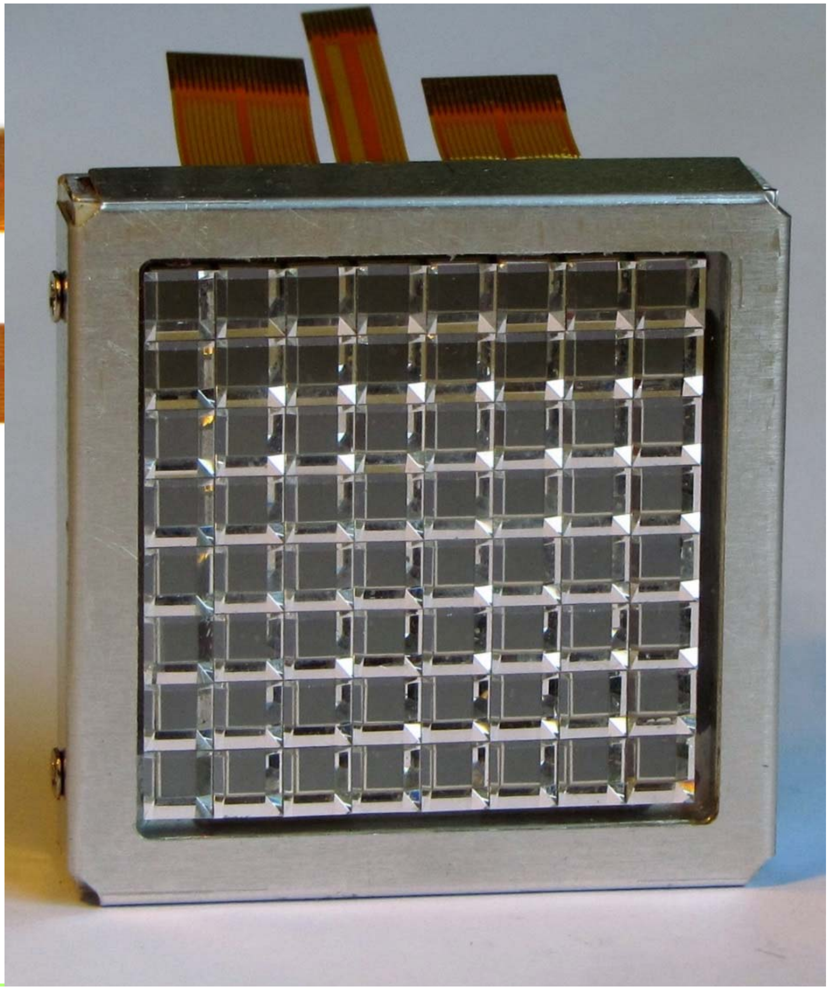
Example: Hamamatsu MPPC S11834-3388DF

- 8x8 SiPM array, with 5x5 mm^2 SiPM channels
- Active area 3x3 mm^2



+ array of quartz light collectors

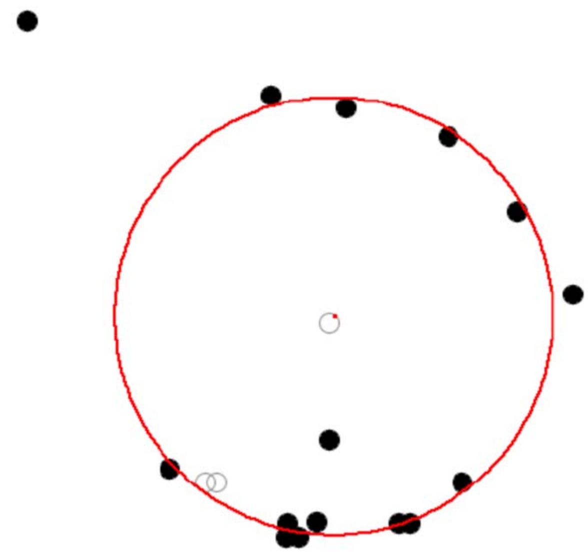




E. Tahirović et al., NIM A787 (2015) 203

Digital SiPM

Digital SiPM (Philips): instead of an analog sum of signals from all cells of a single SiPM, use on board electronics for a digital sum + time stamp

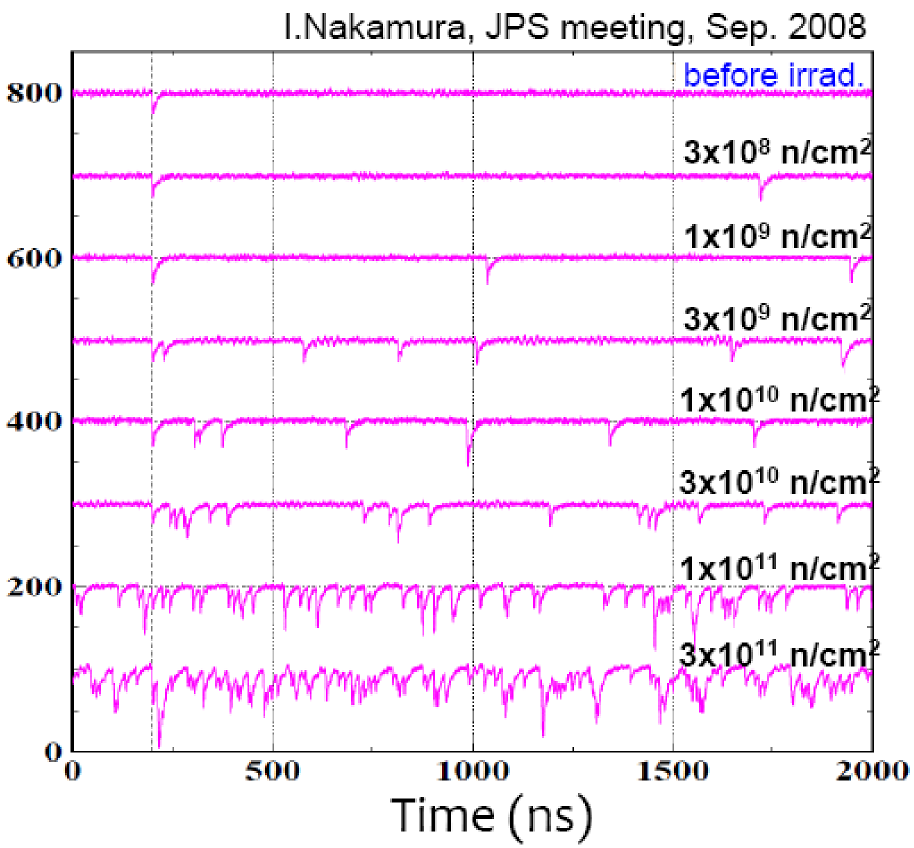


Square matrix 20x20 cm^2

- Sensors: DPC3200-22-44
- 3x3 modules = 6x6 tiles = 24x24 dies = 48x48 pixels in total
- 576 time channels
- 2304 amplitude (position) channels
- 4 levels of FPGA readout: tiles, modules, bus boards, test board

→ A.Y. Barnyakov et al., NIM A732 (2013) 352

SiPMs: Radiation damage



Expected fluence at 50/ab at Belle II:
 2-20 10^{11} n cm^{-2}

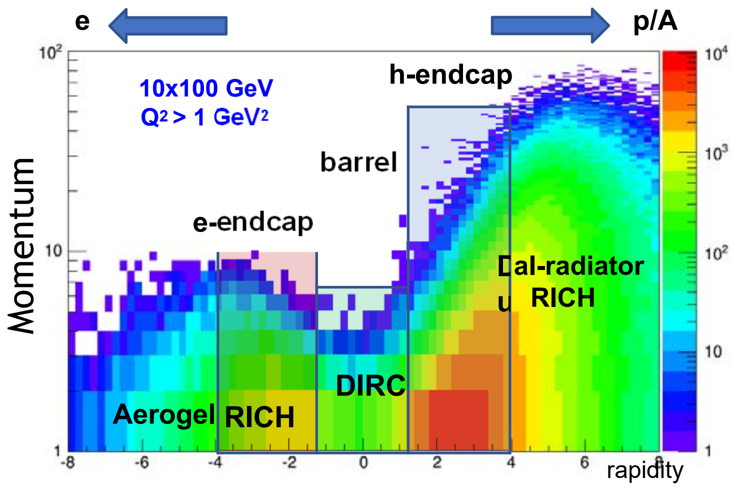
→ Worst than the lowest line

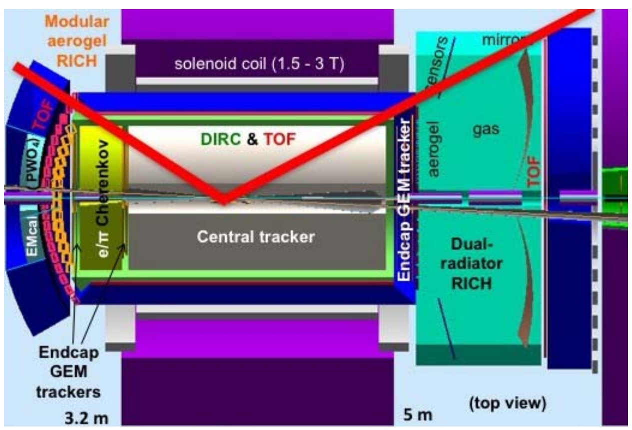
→Need cooling of sensors and wave-form sampling readout electronics

→Annealing?

… and more radiation resistant SiPMs…

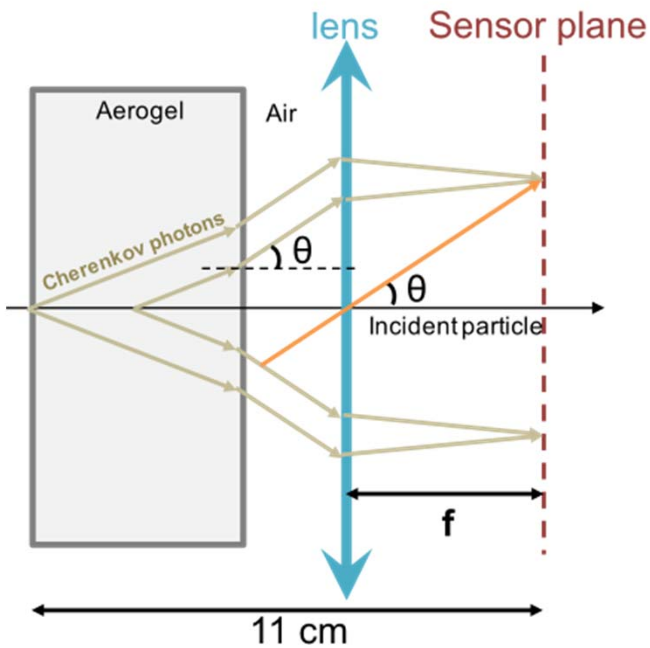
PID Strategies for the Electron Ion Collider

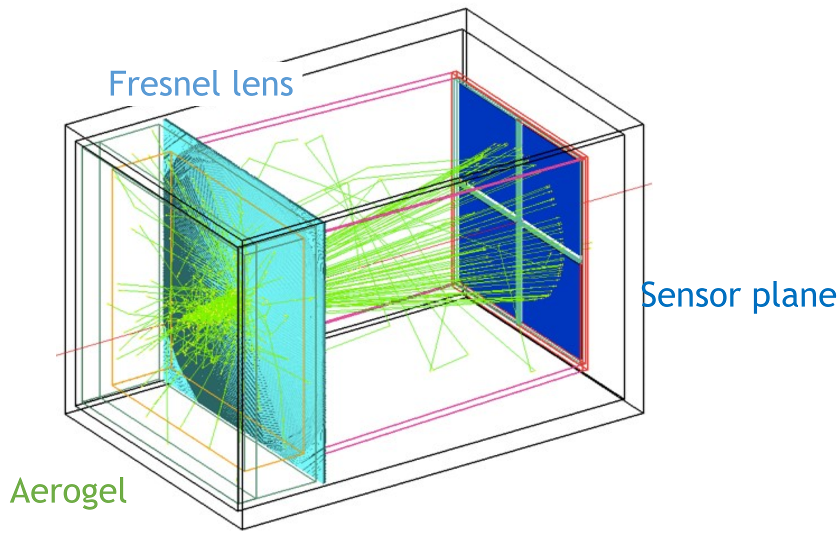




- **h-endcap**: A RICH with two radiators (gas + aerogel) is needed for
  π/K separation up to ~50 GeV/c **dRICH**
- **e-endcap**: A compact aerogel RICH which can be projective
  π/K separation up to ~10 GeV/c **mRICH**
- **barrel**: A high-performance DIRC provides a compact and cost-effective way to cover the area.
  π/K separation up to ~6-7 GeV/c **DIRC**
- **TOF (and/or dE/dx in TPC)**: can cover lower momenta.

EIC mRICH – Working Principle





Geant4 Simulation

Identification of charged particles

Particles (e, μ, π, K, p) in the final state are identified by their mass or by the way they interact.

Determination of mass: from the relation between momentum and velocity, $p=\gamma mv$ (p is known - radius of curvature in magnetic field)

→Measure velocity by:

- time of flight
- ionisation losses dE/dx
- Cherenkov photon angle (and/or yield)
- transition radiation

Mainly used for the identification of hadrons.

Identification through interaction: electrons and muons

- muon systems
- calorimeters

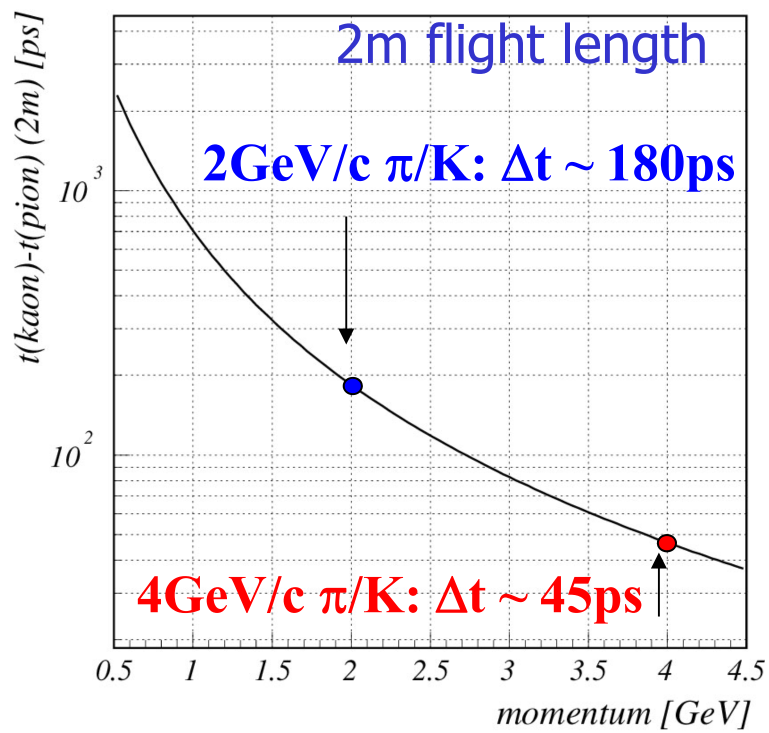
Time-of-Flight (TOF) counters

Measure velocity by measuring the time between
-- the interaction and
-- the passing of the particle through the TOF counter.

Traditionally: plastic scintillator + PMTs

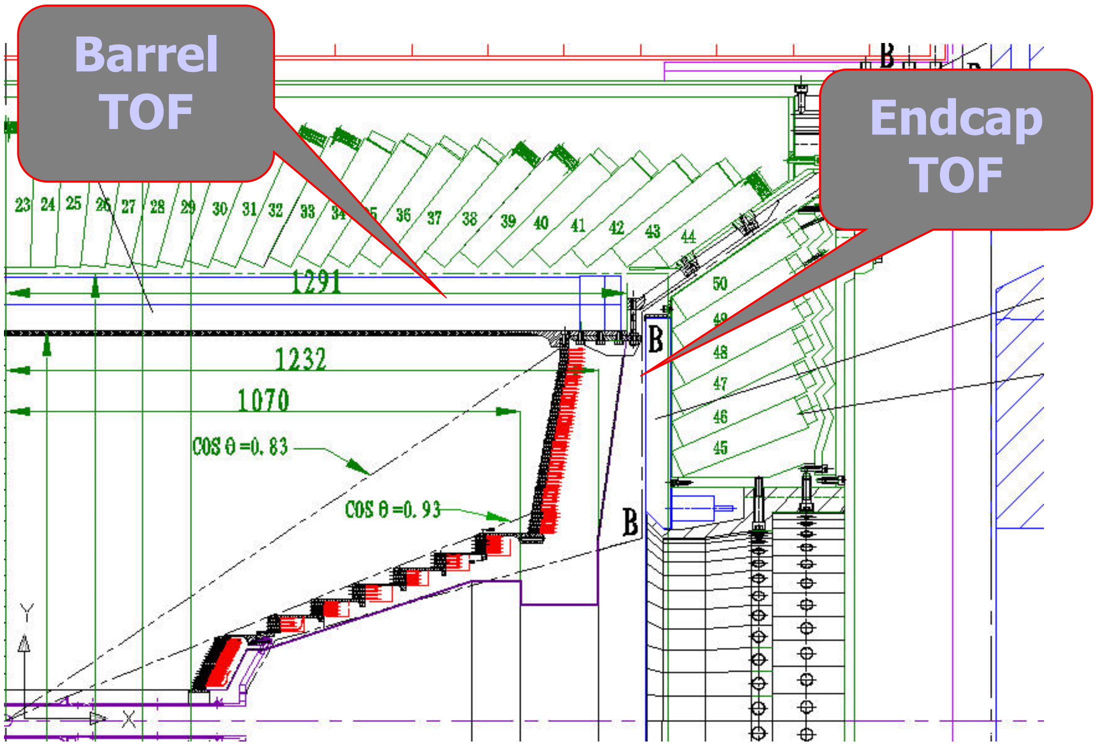
Typical resolution: ~100 ps →π/K separation up to ~1GeV.

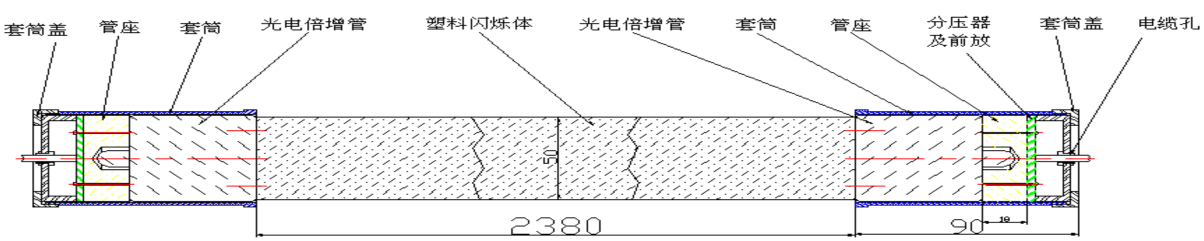
Time difference between π and K:



→ BESSIII

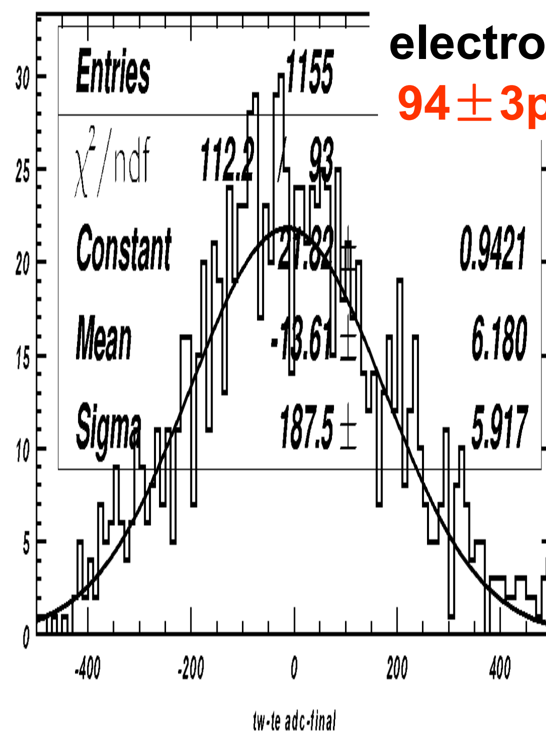
BESIII: Time-Of-Flight counters





TOF module: high quality plastic scintillator: 2.4 m long, 5cm thick, two PMTs with preamplifiers

Beam test results: two TOF modules



electrons
94±3ps

Time-of-Flight (TOF) counters

Measure velocity by measuring the time between the interaction and the passing of the particle through the TOF counter.

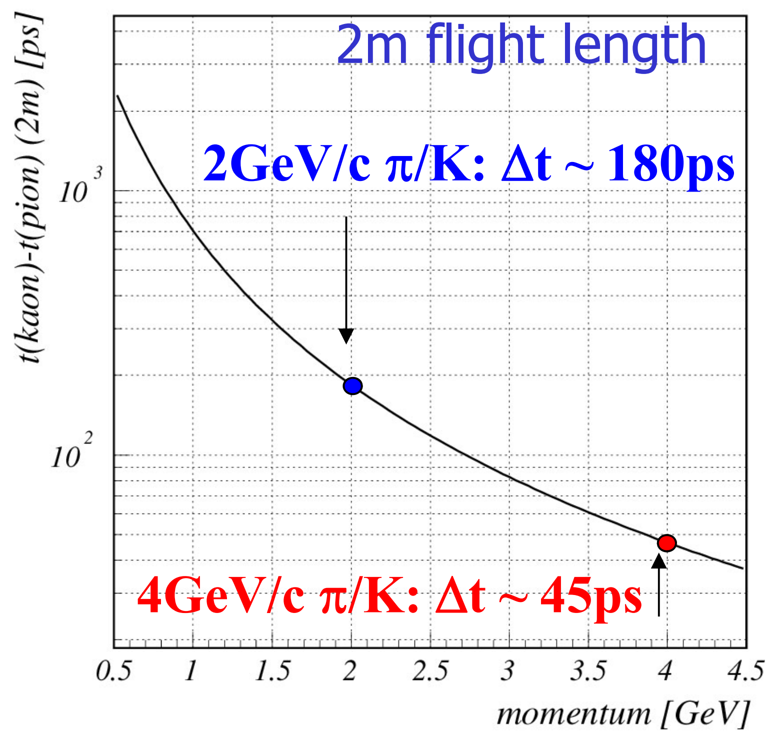
Traditionally: plastic scintillator + PMTs

Typical resolution: ~100 ps →π/K sepration up to ~1GeV.

To go beyond that: need faster detectors:
→use Cherenkov light (prompt) instead of scintillations
→use a fast gas detector (Multi gap RPC)

However: make sure you also know the interaction time very precisely...

Time difference between π and K:



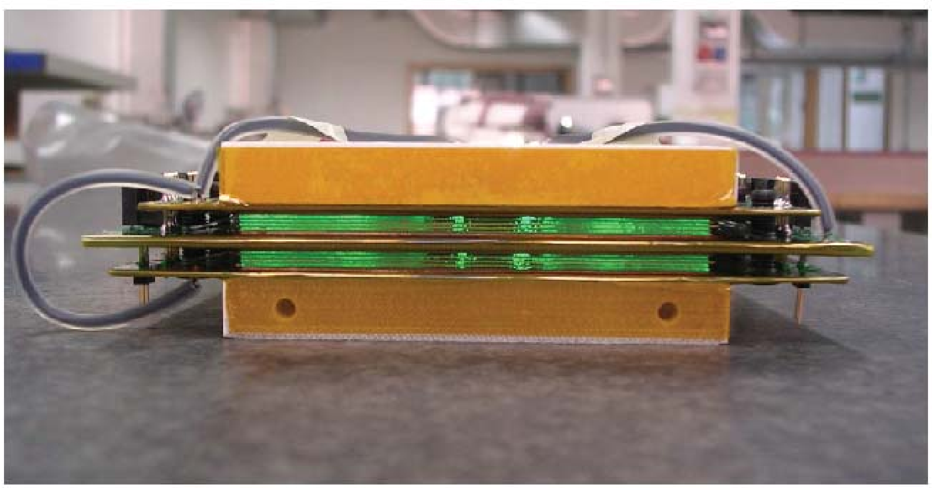
ALICE TOF

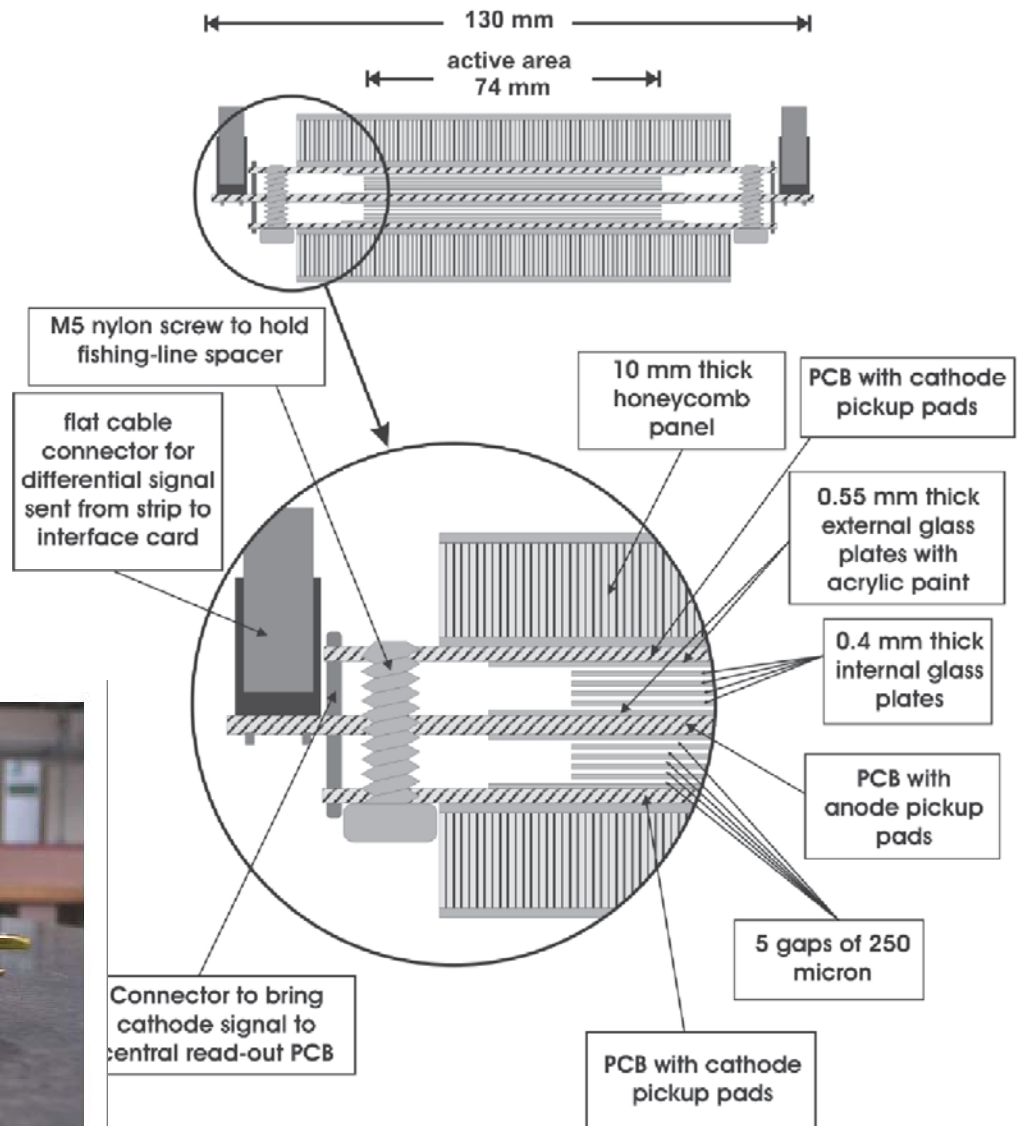
Very fast large area (140m^2) particle detector:

→ MRPC, multi-gap RPC

σ=50ps (incl. read-out)

π/K separation (3σ) up to 2.5 GeV/c at large track densities

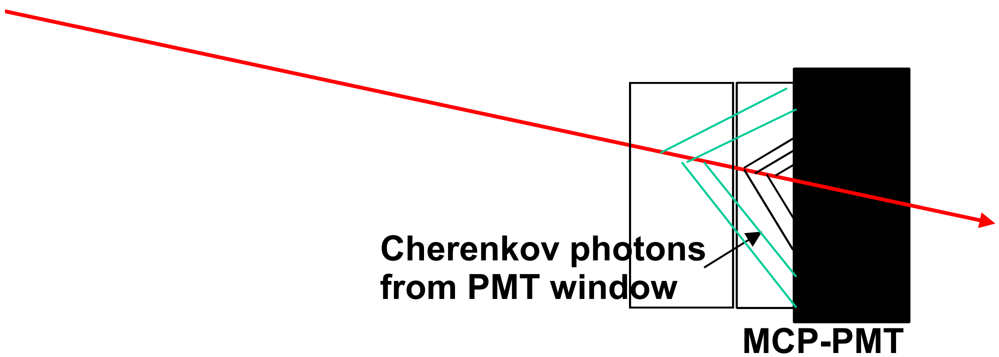


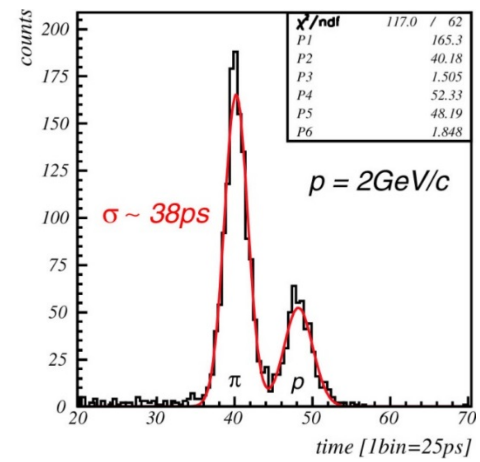


TOF with Cherenkov light

Idea: detect Cherenkov light with a very fast photon detector (MCP PMT).

Cherenkov light is produced in a quartz plate in front of the MCP PMT and in the PMT window.





Proof of principle: beamt test with pions and protons at 2 GeV/c.

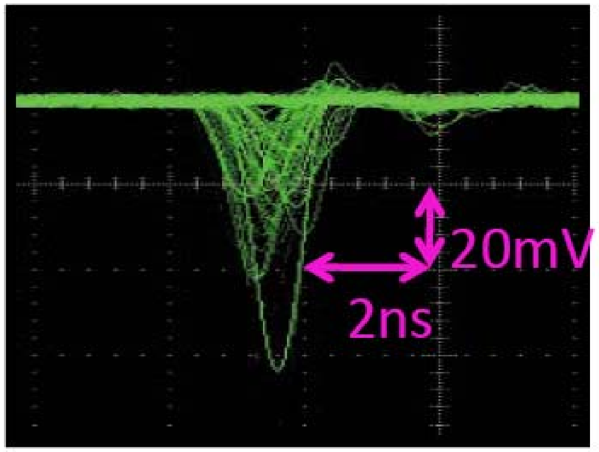
→resolution ~5ps measured

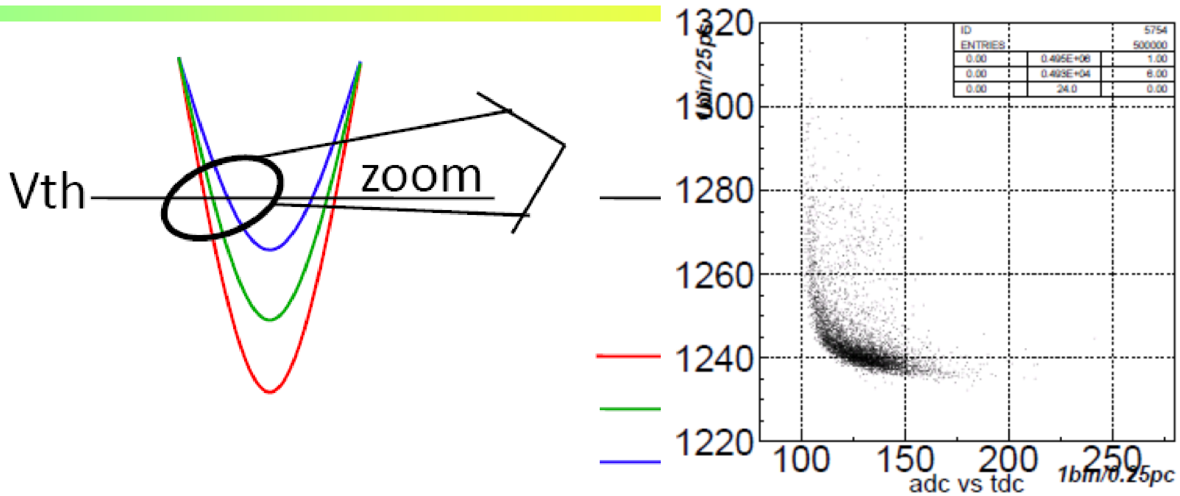
- K. Inami NIMA 560 (2006) 303
- J. Va'vra NIMA 595 (2008) 270

Open issues:

- read-out
- start time

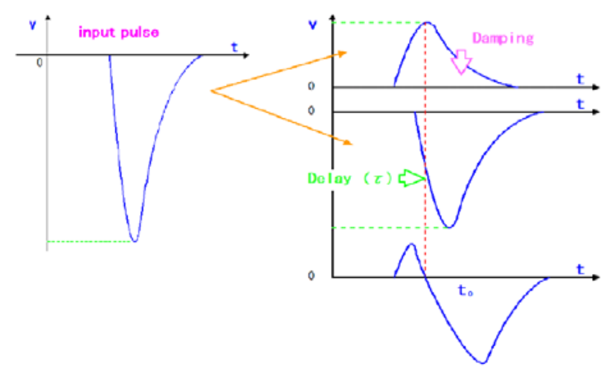
Read out: for precise timing mitigate time walk

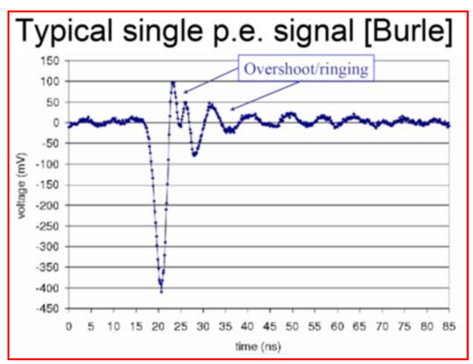




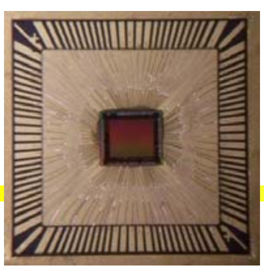
Variation of time determined with a leading edge discriminator: smaller pulses give a delayed signal.→ Has to be corrected!

- Measure both time (TDC) and amplitude (ADC), correct time of arrival by using a ΔT(ADC) correction
- Use constant fraction discriminator (CFD) →
- Wave-form sampling →





e.g. Labrador 3,
G. Varner, U Hawaii



Identification of charged particles

Particles (e, μ, π, K, p) in the final state are identified by their mass or by the way they interact.

Determination of mass: from the relation between momentum and velocity, $p=\gamma mv$ (p is known - radius of curvature in magnetic field)

→Measure velocity by:

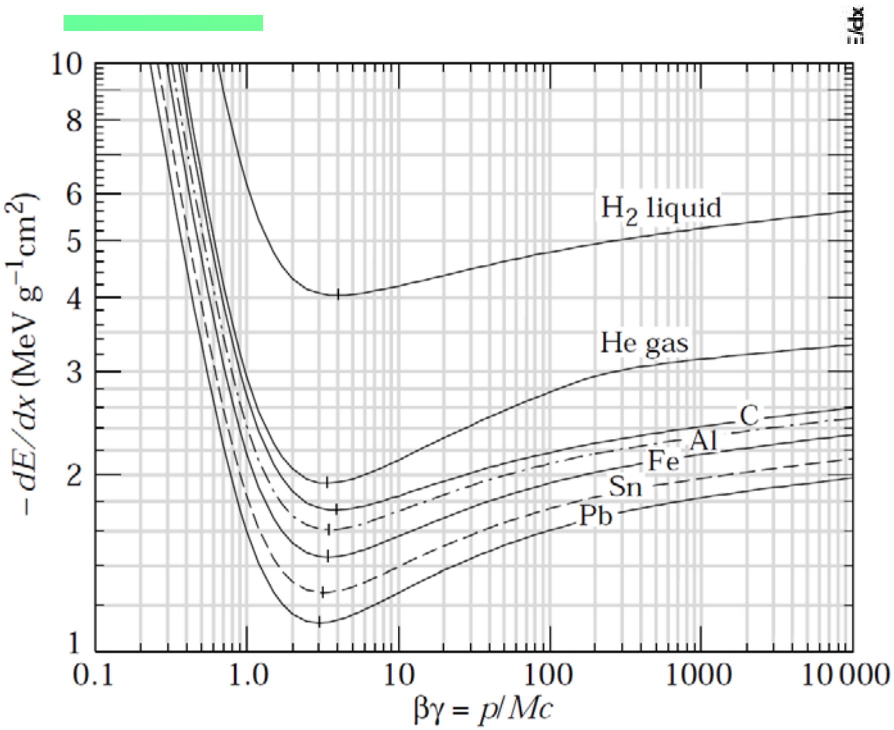
- time of flight
- ionisation losses dE/dx
- Cherenkov photon angle (and/or yield)
- transition radiation

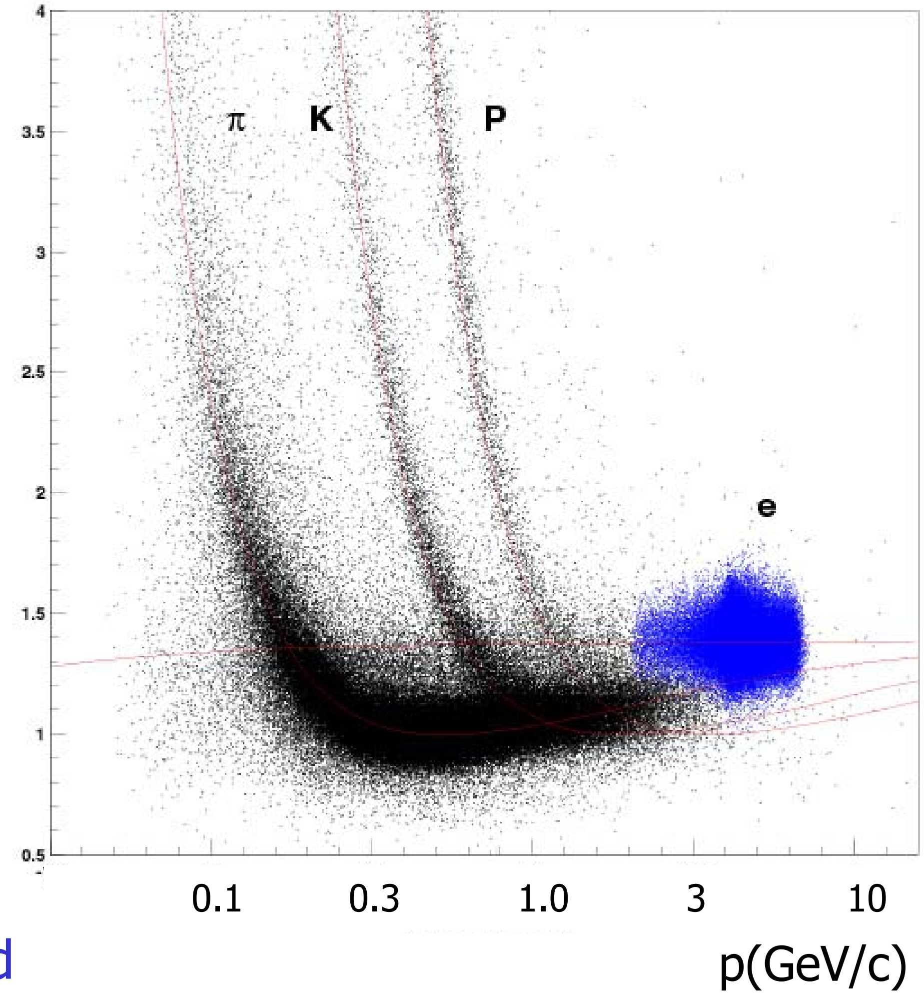
Mainly used for the identification of hadrons.

Identification through interaction: electrons and muons

- muon systems
- calorimeters

Identification with the dE/dx measurement





dE/dx is a function of velocity β

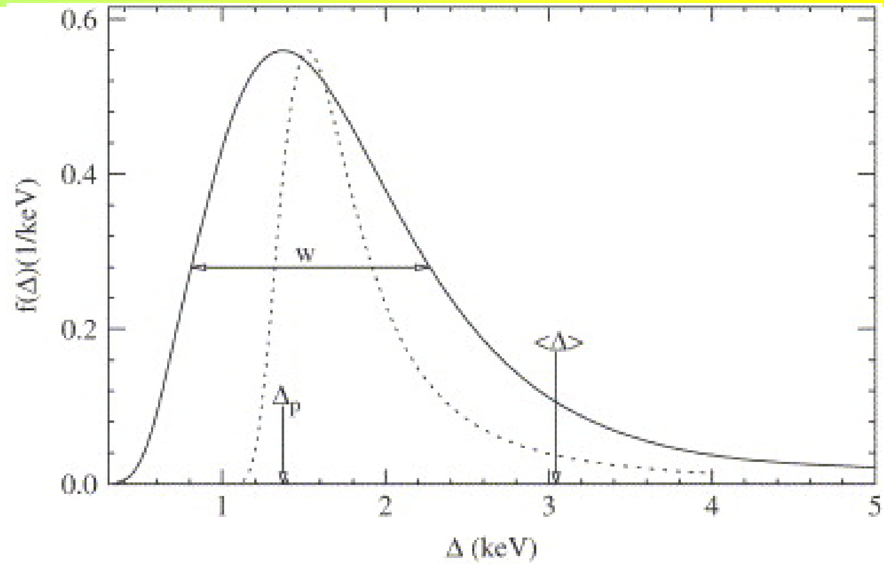
For particles with different mass the Bethe-Bloch curve gets displaced if plotted as a function of p

For good separation: resolution should be ~5%

Measure in each drift chamber layer – use truncated mean

Identification with dE/dx measurement

Problem: long tails (not Gaussian!)



Energy loss distribution for particles with $\beta\gamma$=3.6 traversing 1.2 cm of Ar gas (solid line).

Parameters describing ***f(Δ)*** are

$\Delta_p(x;\beta\gamma)$ **:** the most probable energy loss = the position of the maximum at 1371 eV, and

W : the full-width-at-half-maximum (FWHM) of 1463 eV. The mean energy loss is 3044 eV.

Dotted line: the original Landau function.

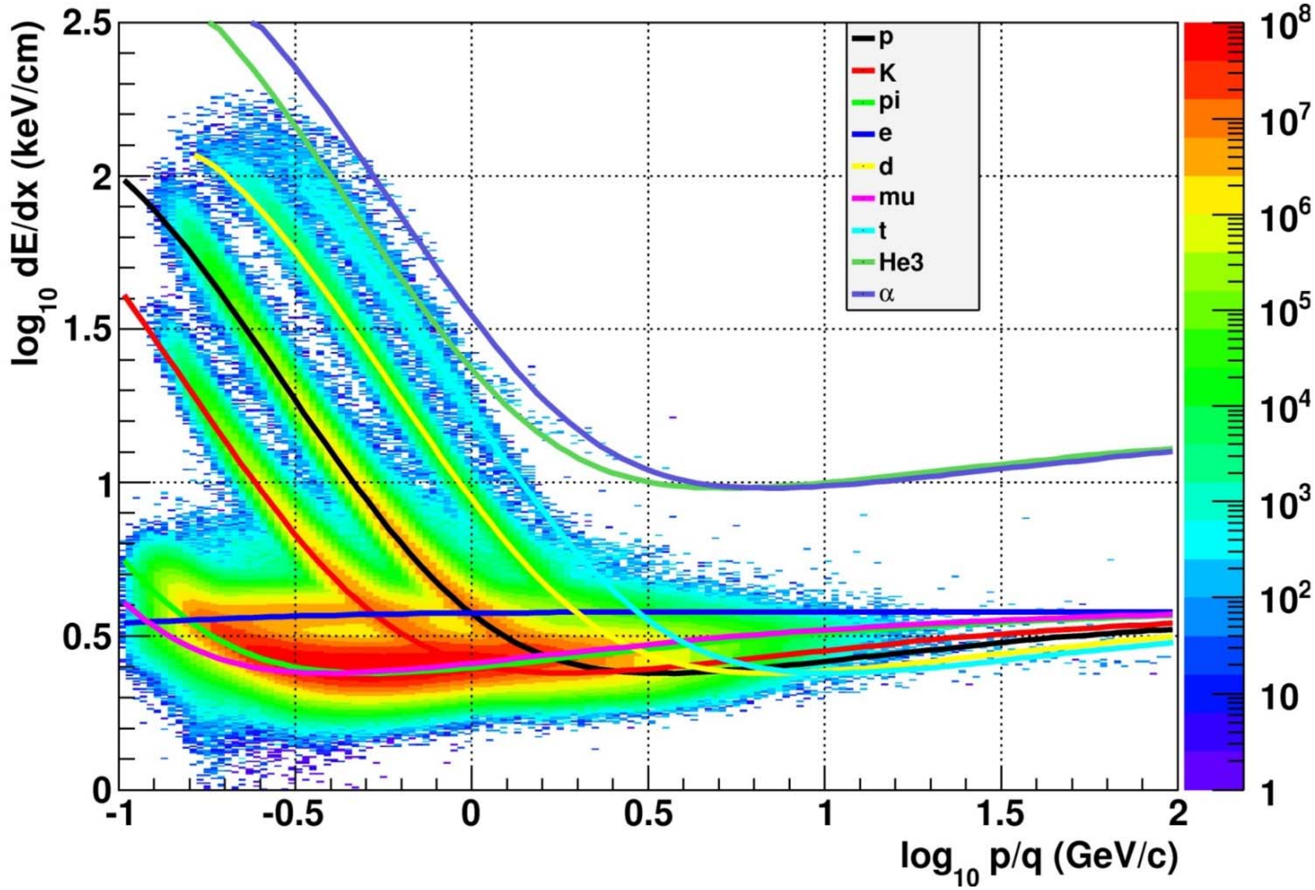
→Many samples along the track (~100 in ALICE TPC), remove the largest ~40% values (reduce the influence of the ling tail) → truncated mean

→ Hans Bichsel: A method to improve tracking and particle identification in TPCs and silicon detectors, NIM A562 (2006) 154

Identification with dE/dx measurement

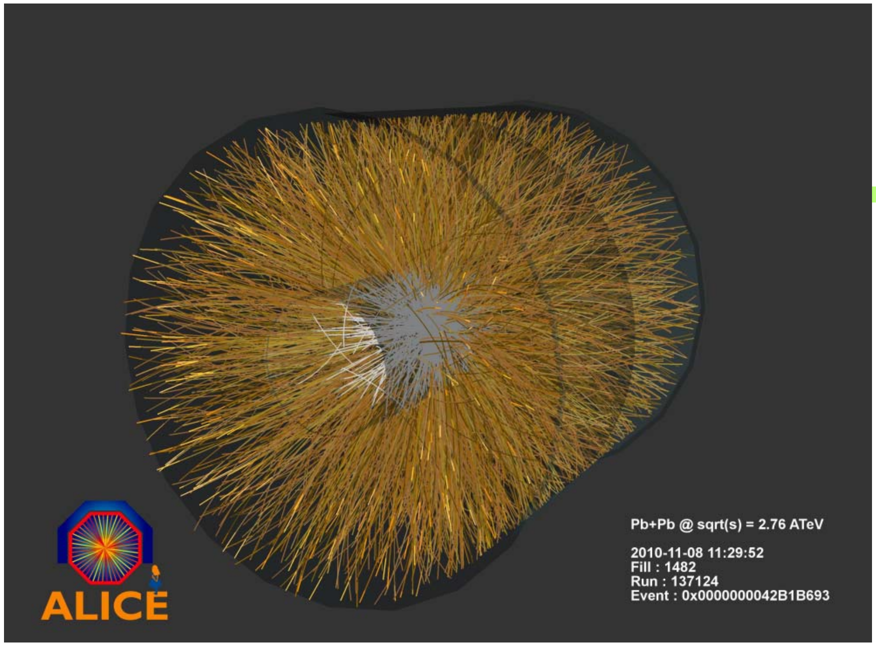
dE/dx performance in the STAR TPC

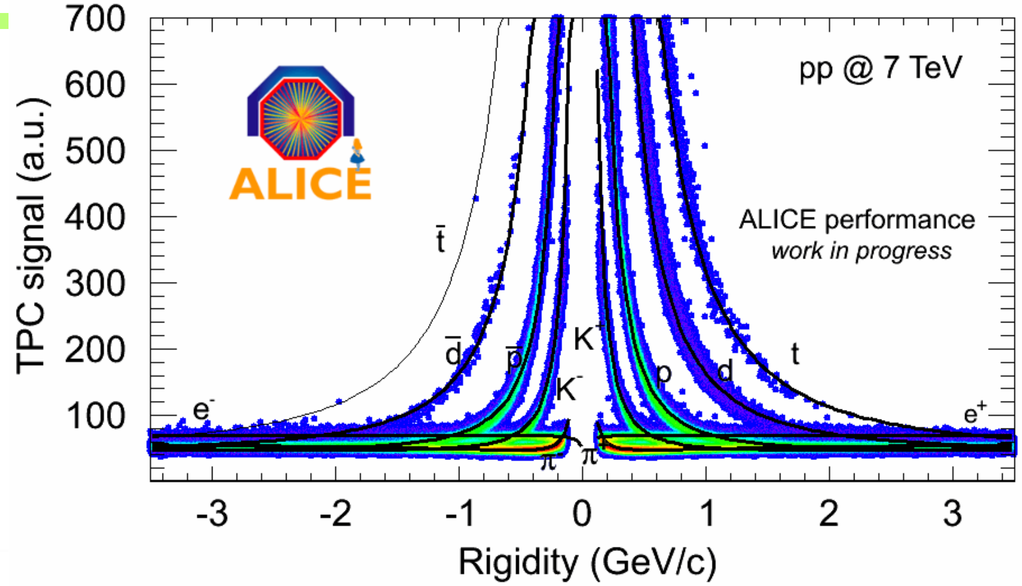
gold-gold
collisions

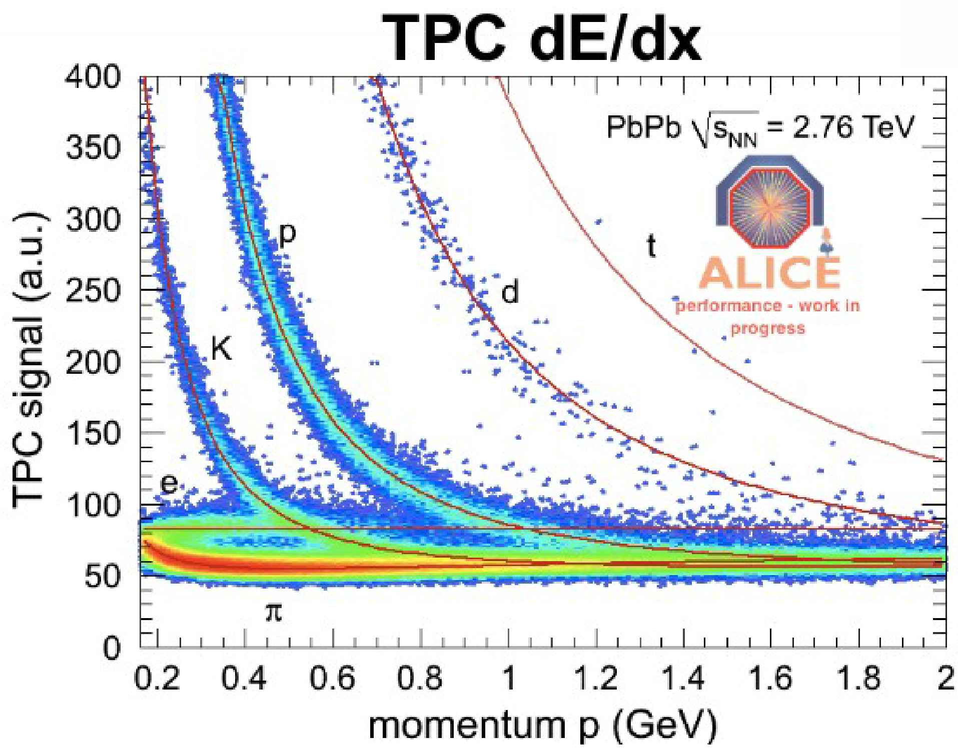


Energy loss in the STAR TPC: truncated mean as a function of momentum. The curves are Bichsel model predictions.

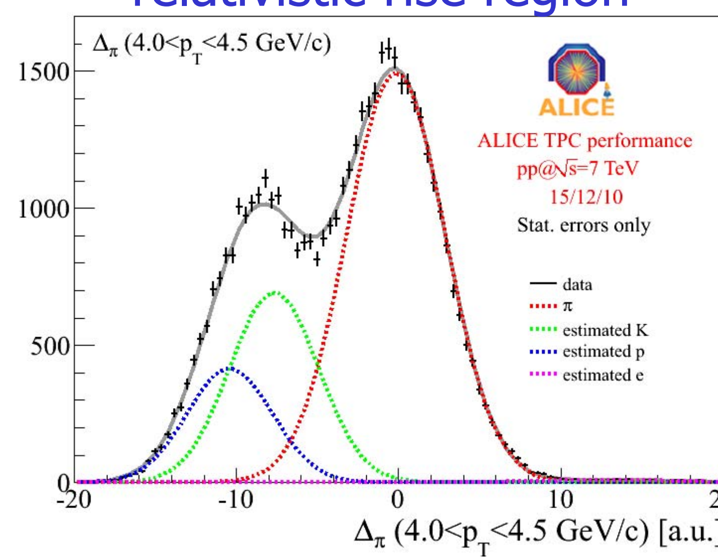
dE/dx in ALICE







relativistic rise region



Y

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→Measure velocity by:

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Mainly used for the identification of hadrons.

Identification through interaction: electrons and muons

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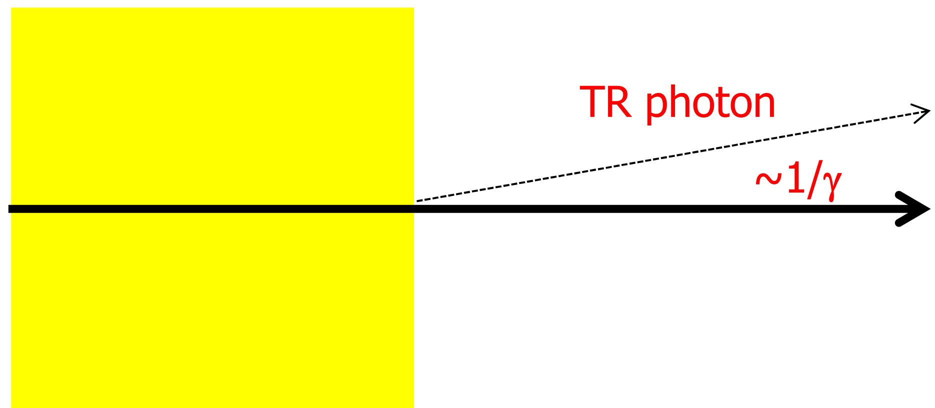
- calorimeters

Transition radiation

E.M. radiation emitted by a charged particle at the boundary of two media with different refractive indices



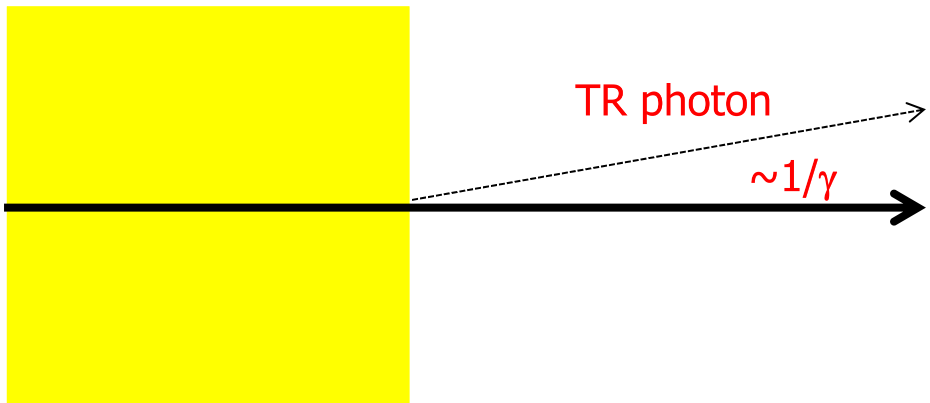
- Accelerated particle emits E.M. radiation
- Transition radiation: particle has a constant velocity, but the phase velocity of the medium changes abruptly at the boundary → radiation

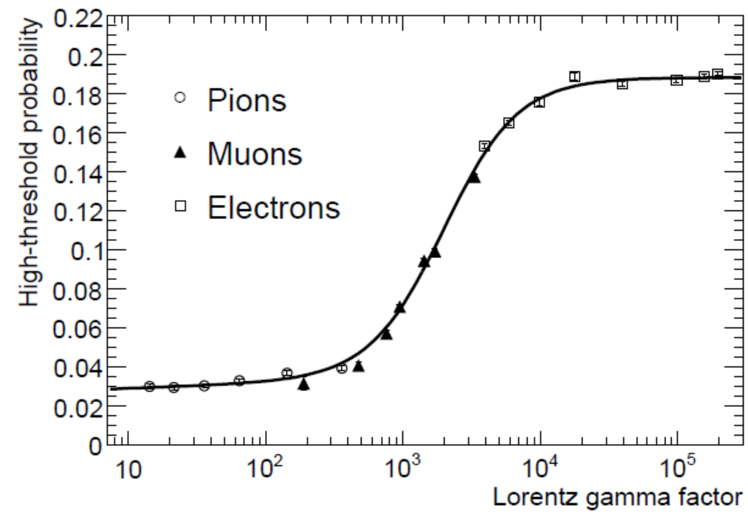
→B. Dolgoshein, NIM A326 (1993) 434-469; J.D. Jackson, Classical Electrodynamics.

→H. Kolanoski, N. Wermes, Teilchendetektoren, Springer.

Transition radiation

E.M. radiation emitted by a charged particle at the boundary of two media with different refractive indices





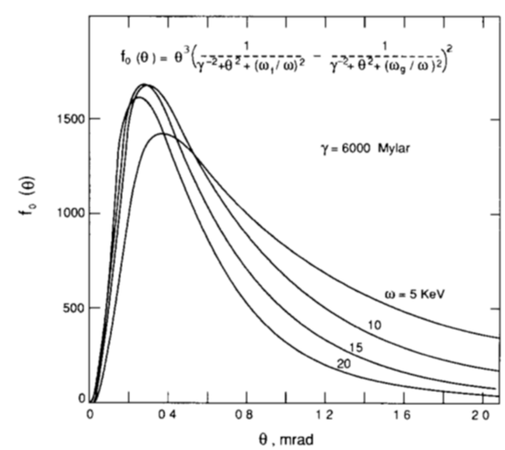
Emission rate depends on γ (Lorentz factor): becomes important at γ~1000

- Electrons at 0.5 GeV
- Pions above 140 GeV

Emission probability per boundary ~α = 1/137

Emission angle ~1/γ

Typical photon energy: ~10 keV → X rays



Transition radiation - detection

Emission probability per boundary $\sim\alpha = 1/137$

→ Need many boundaries

- Stacks of thin foils or
- Porous materials – foam with many boundaries of individual 'bubbles'

Typical photon energy: ~10 keV → X rays

→ Need a wire chamber with a high Z gas (Xe) in the gas mixture

(=large cross section for photoeffect of X rays)

Emission angle $\sim 1/\gamma$

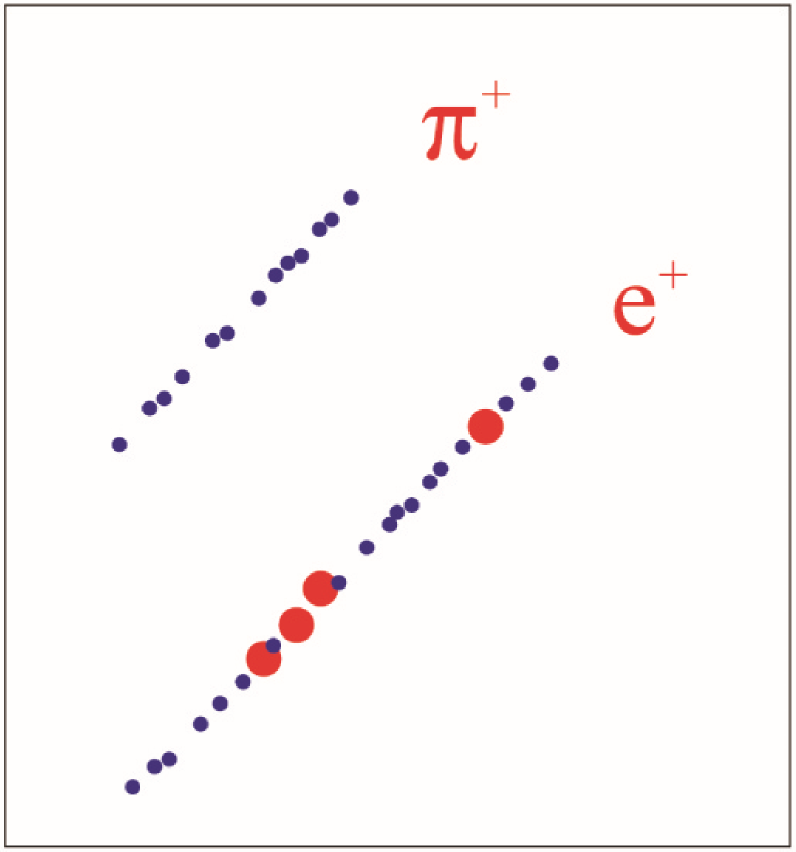
→ Hits from TR photons along the charged particle direction

- Separation of X ray hits (high energy deposit on one place) against ionisation losses (spread out along the track)
- Two thresholds: lower for ionisation losses, higher for X ray detection

Transition radiation - detection

→ Hits from TR photons along the charged particle direction

- Separation of X ray hits (high energy deposit on one place) against ionisation losses (spread out along the track)
- Two thresholds: lower to remove noise, higher to separate X ray conversions from ionisation by charged particles



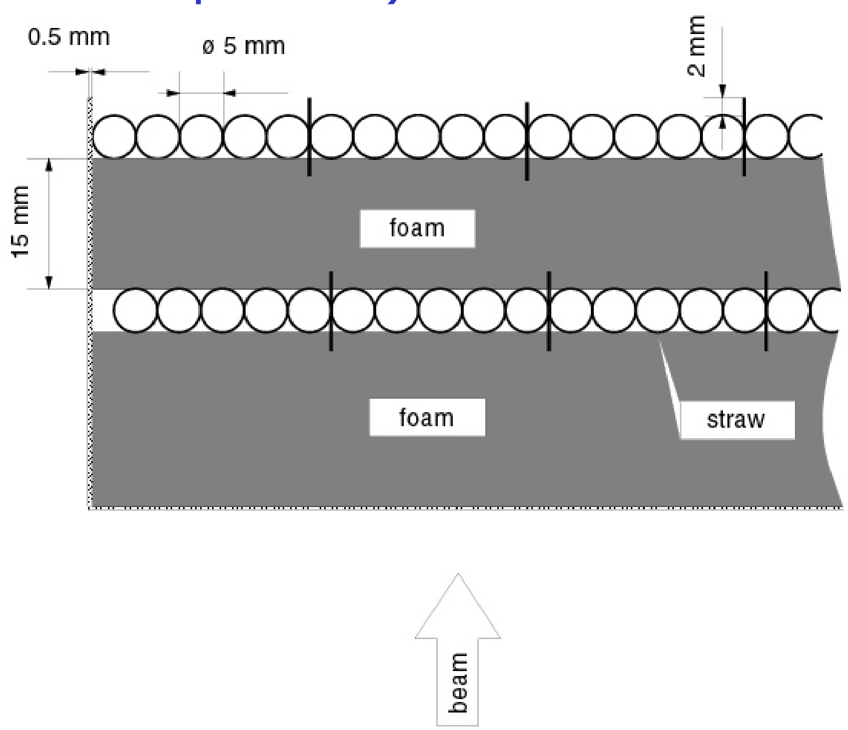
- Small circles: between low and high threshold (ionisation)
- Big circles: high threshold (X ray detection)

(pion below the TR threshold, e above the TR threshold)

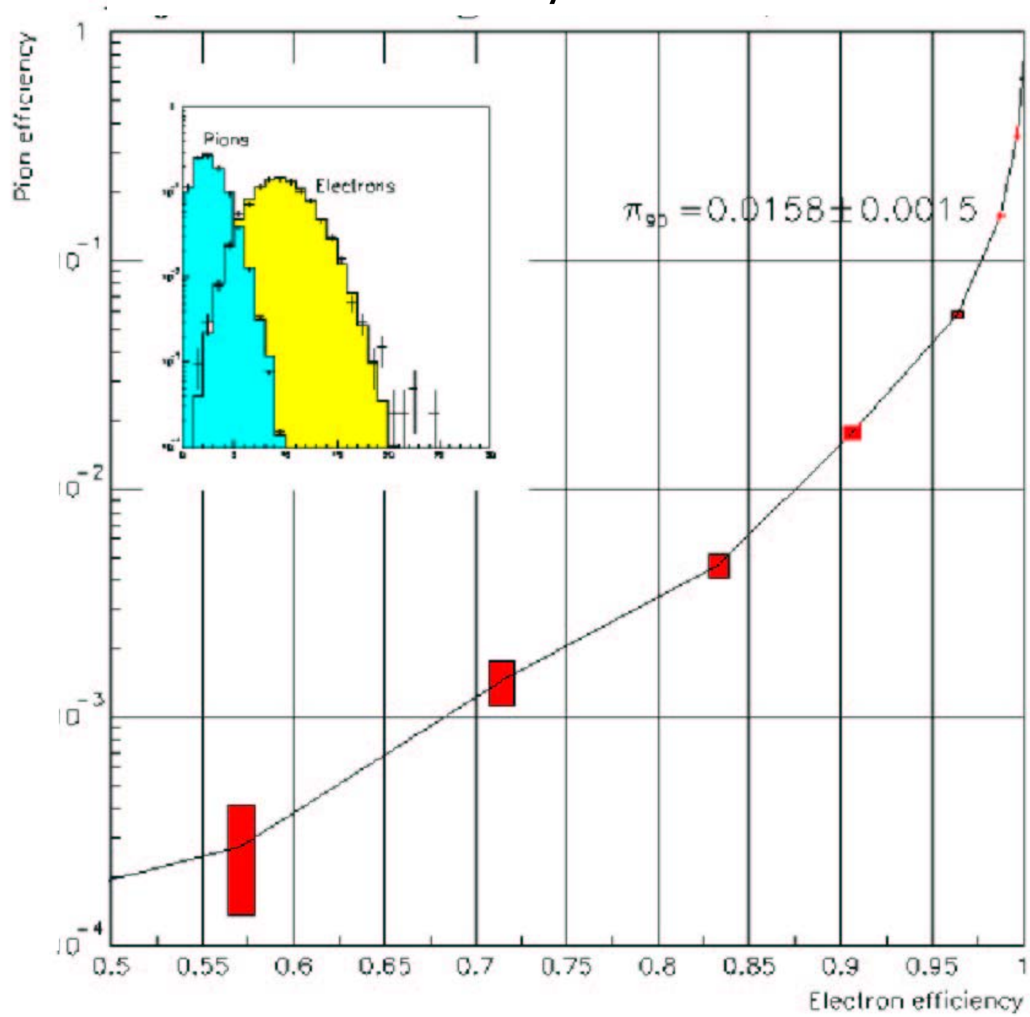
Transition radiation detectors

Example:

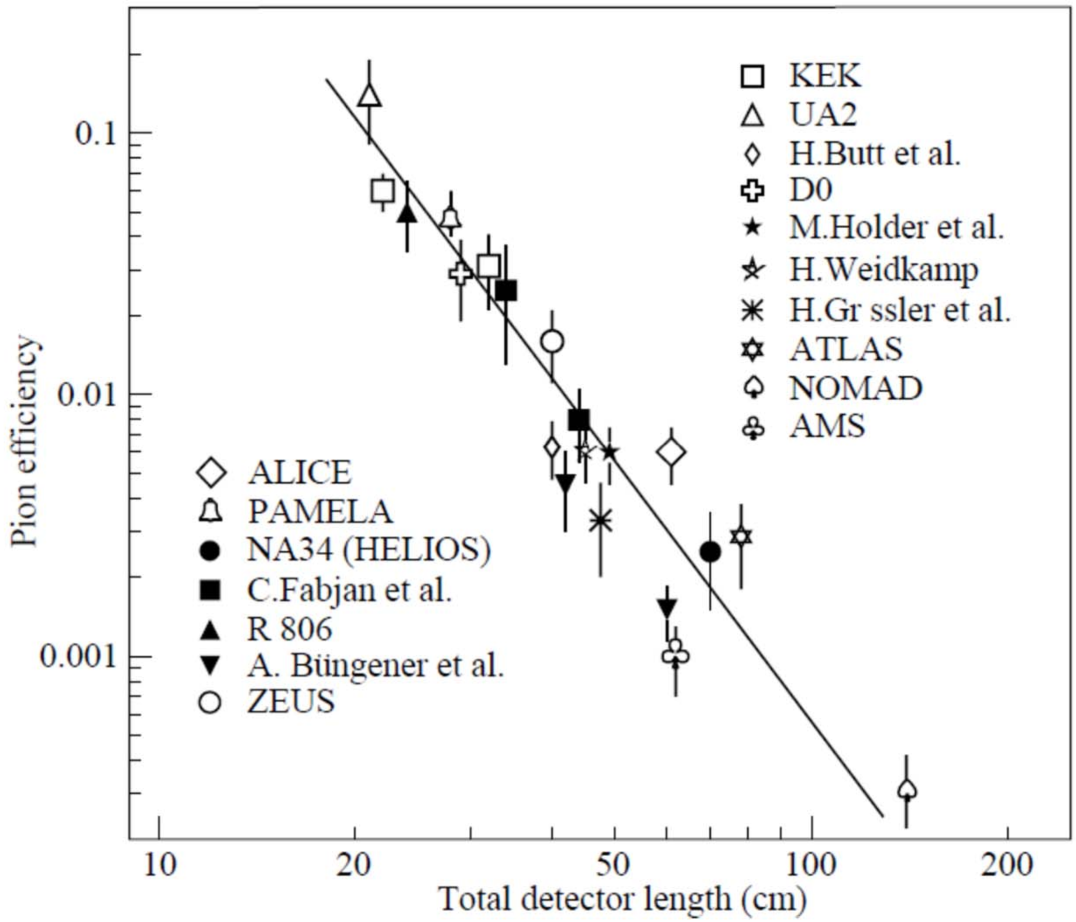
Radiator: organic foam between the detector tubes (straws made of capton foil)



Performance: pion efficiency (fake prob.) vs electron efficiency

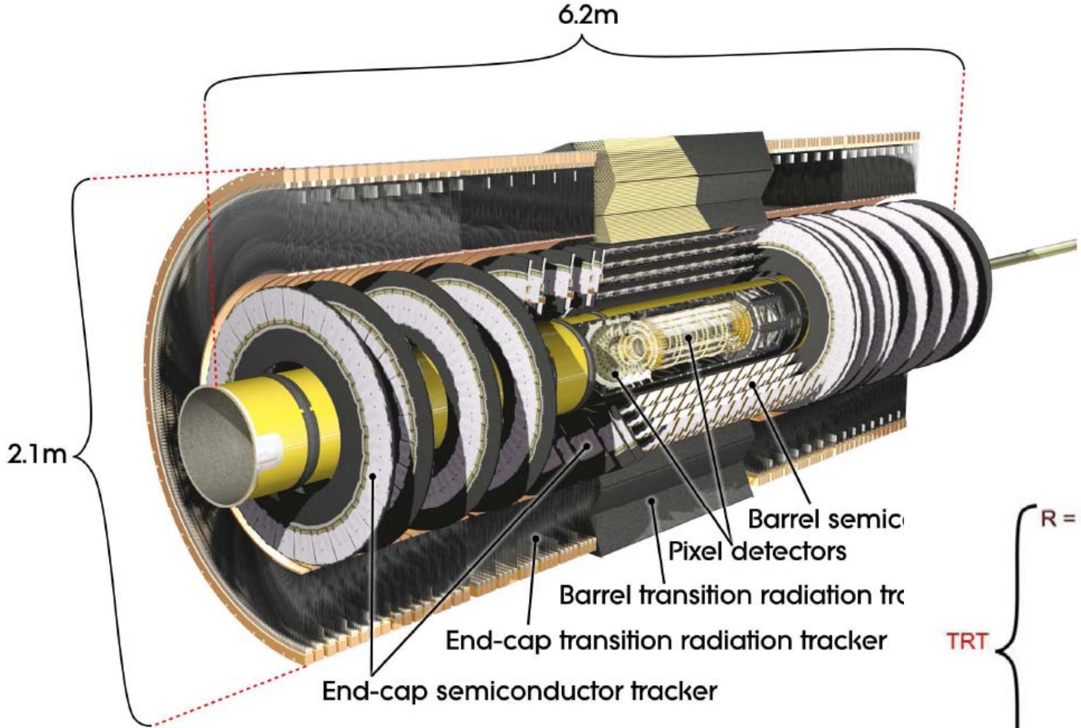


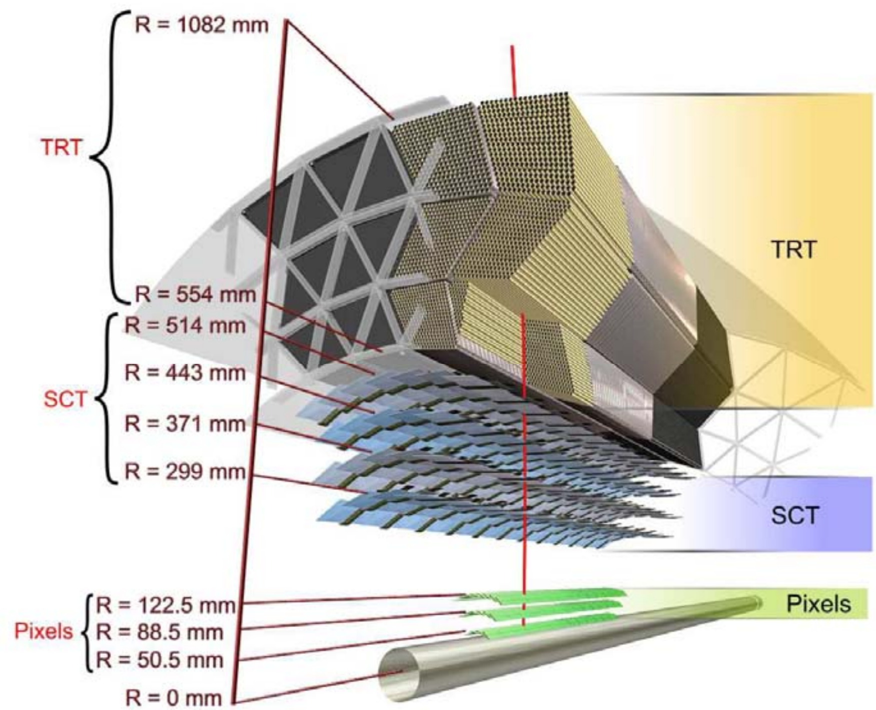
Transition radiation detectors - peformance



Performance: pion efficiency (fake prob.) vs detector length

Transition radiation detector in ATLAS: combination of a tracker and a transition radiation detector



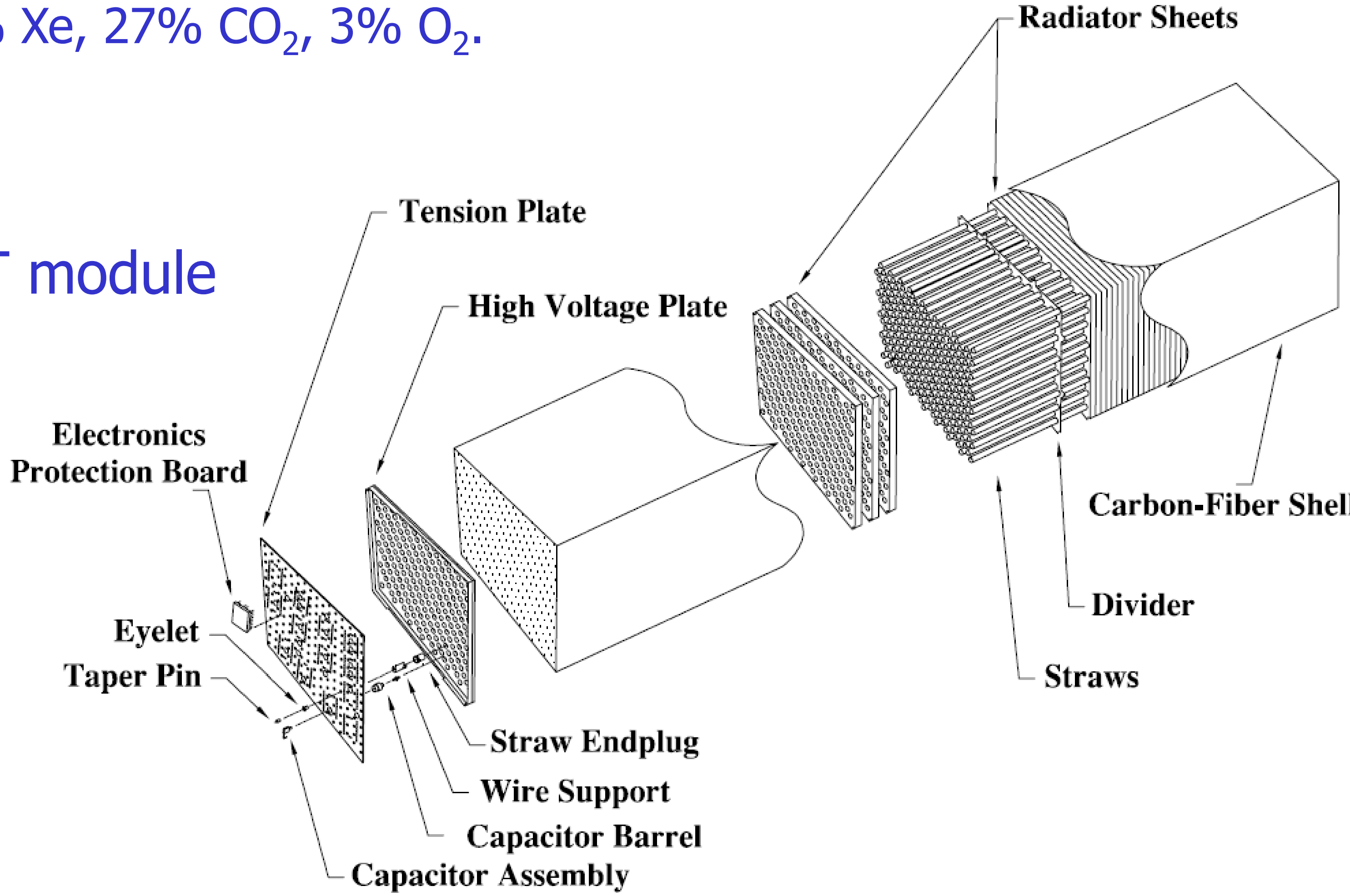


ATLAS TRT

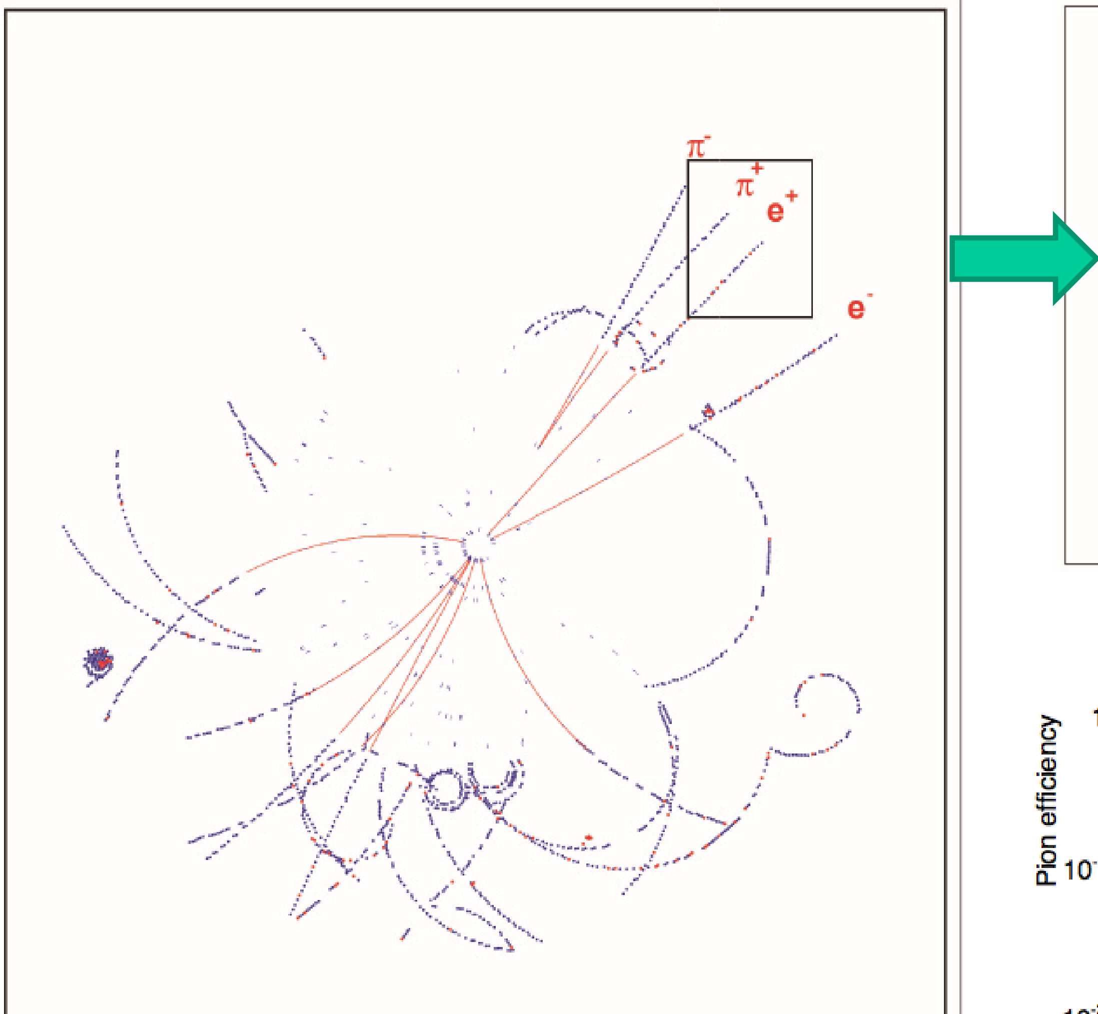
Radiator: 3mm thick layers made of polypropylene-polyethylene fibers with ~19 micron diameter, density: 0.06 g/cm^3

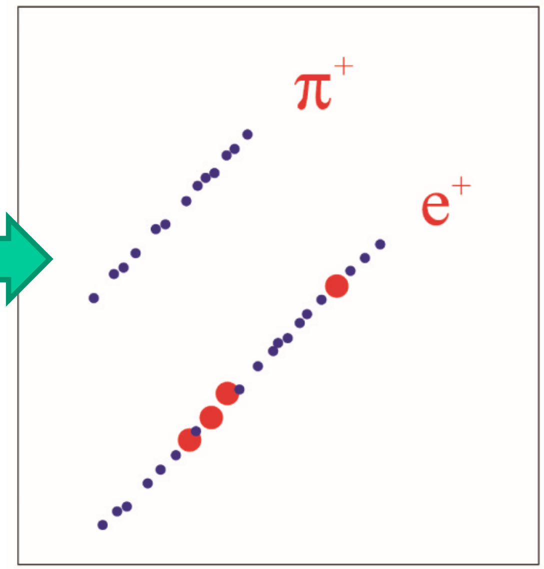
Straw tubes: 4mm diameter with 31 micron diameter anode wires, gas: 70% Xe, 27% CO_2, 3% O_2.

TRT module

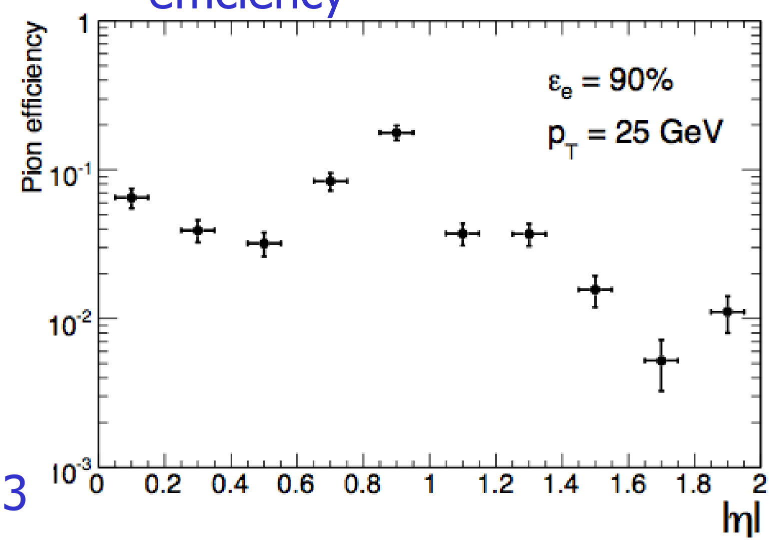


TRT: pion-electron separation





π fake probability at 90% e efficiency



→ JINST 3 (2008) S08003

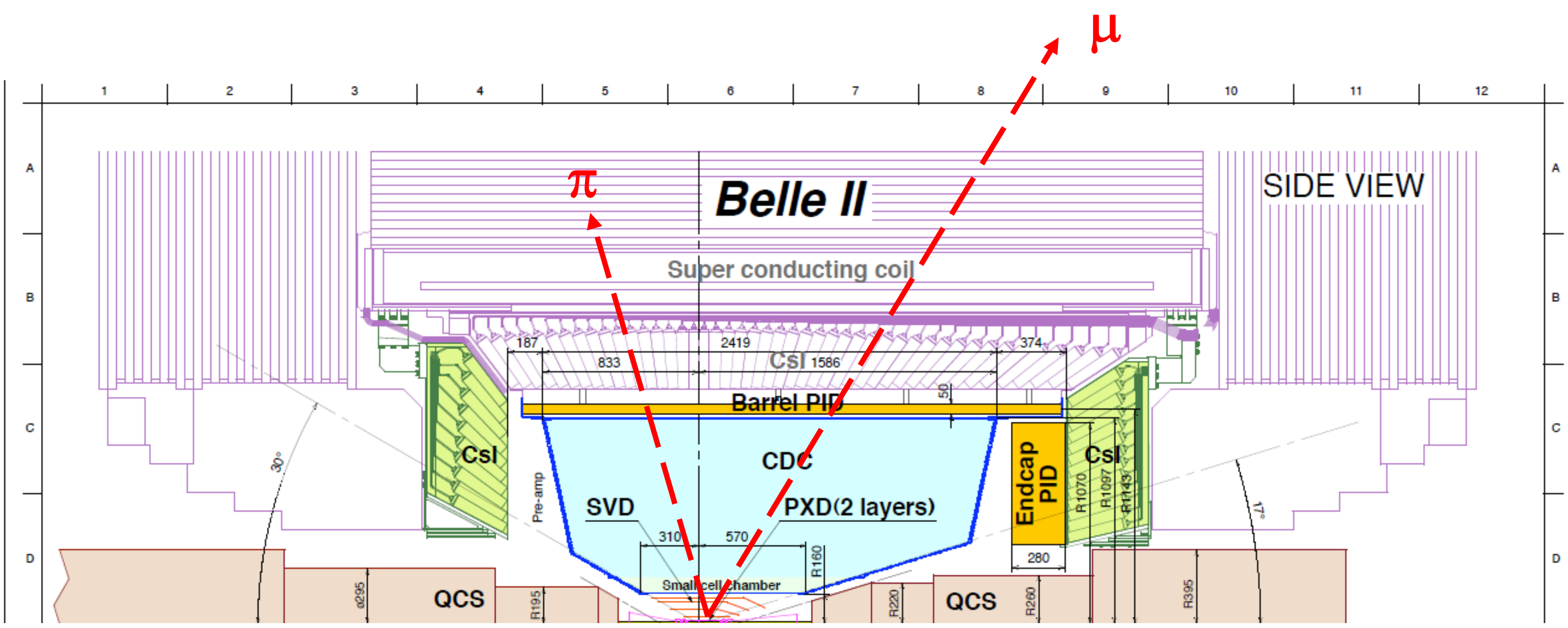
Identification of muons

Separate muons from hadrons (pions and kaons):

Exploit the fact that muons interact only electromag., while hadrons interact strongly

→ need a few interaction lengths to stop hadrons (interaction lengths = about 10x radiation length in iron, 20x in CsI).

→ **A particle is identified as a muon if it penetrates the material.**



Example: muon detection at B factories

Separate muons from hadrons (pions and kaons):

Need a few interaction lengths to stop hadrons (interaction length = about 10x radiation length in iron, 20x in CsI).

Some numbers: 0.8 interaction length (CsI) + 3.9 interaction lengths (iron)

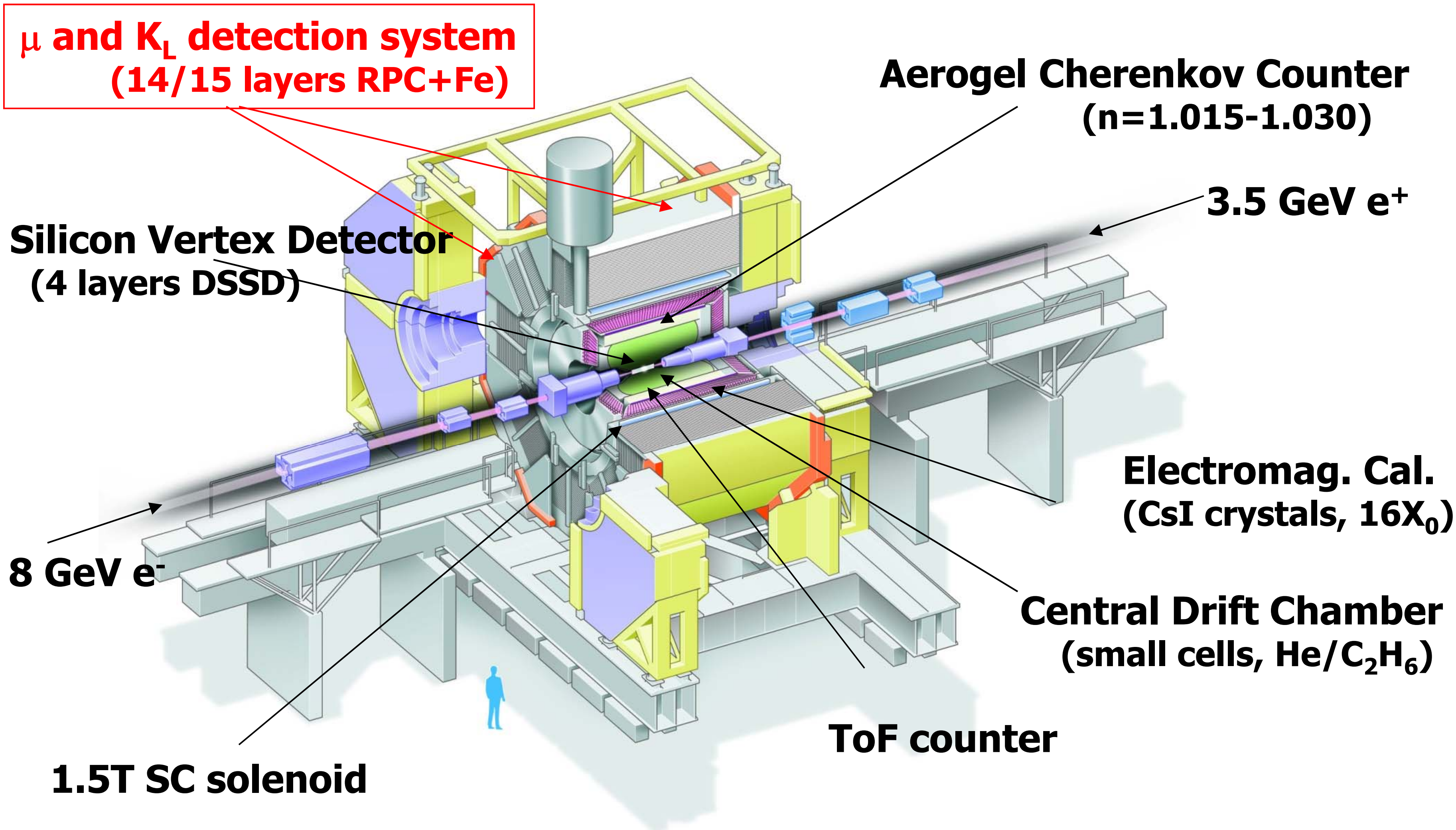
Interaction length: CsI 167 g/cm^2, iron 132 g/cm^2,

$(dE/dx)_{min}$: CsI 1.24 MeV/(g/cm^2), iron 1.45 MeV/(g/cm^2)

→ ΔE_{min} = (0.15+0.75) GeV = 0.90 GeV → reliable identification of muons possible around ~1 GeV

Detect K_L interaction (cluster): again need a few interaction lengths – the same system can be used for both – bonus!

Example: Muon and K_L detection at Belle



Muon and K_L detector

Up to 21 layers of resistive-plate chambers (RPCs) between iron plates of the flux return

Bakelite RPCs at BABAR

Glass RPCs at Belle

(glass was better choice because of ageing effects)

Scintillator strips + RPCs at Belle II

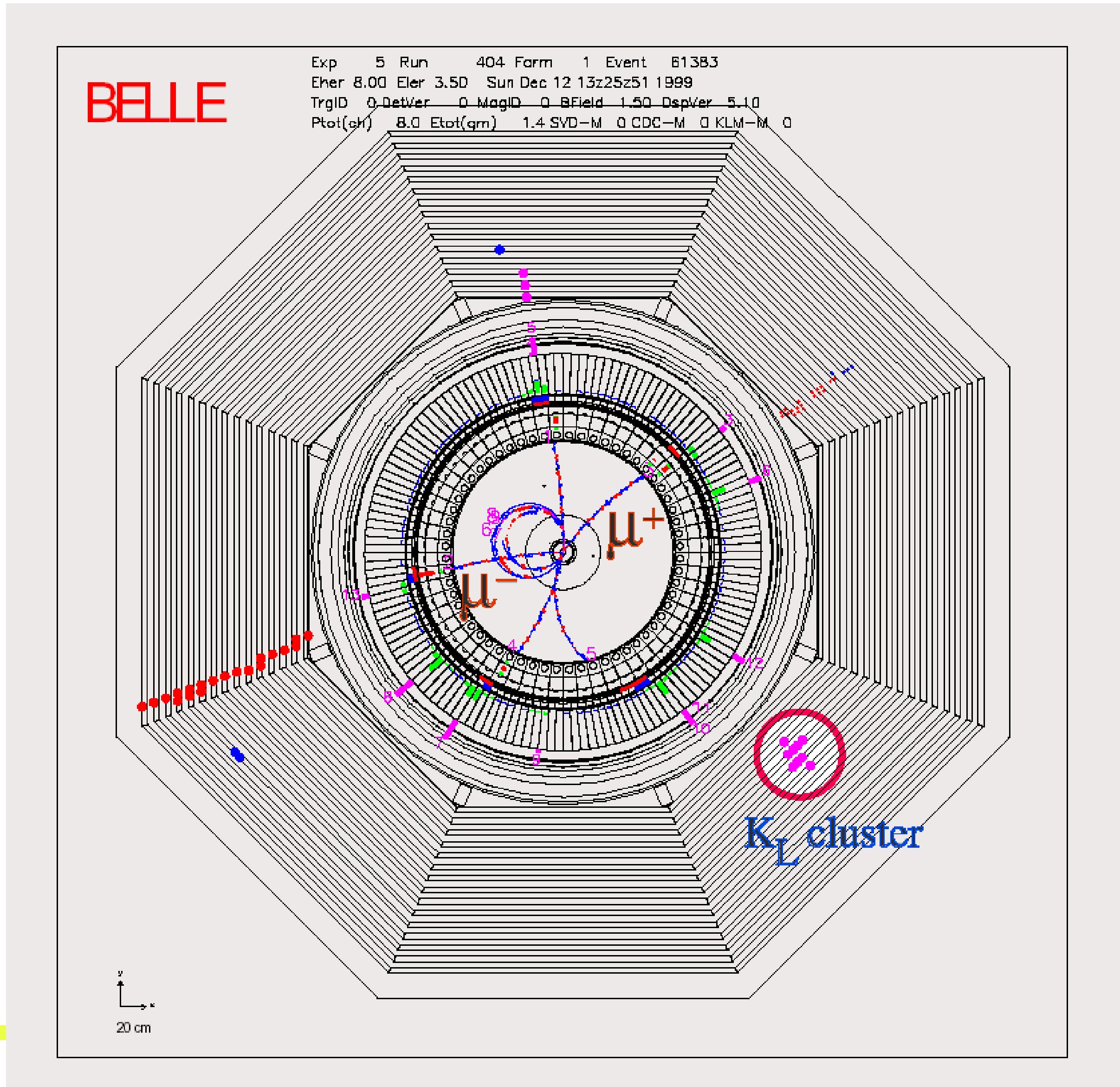
Muon and K_L detector

Example:

event with

- two muons and a
- K_L

and a pion that only partly penetrated



Muon identification performance

Muon identification: efficient for p>800 MeV/c

efficiency

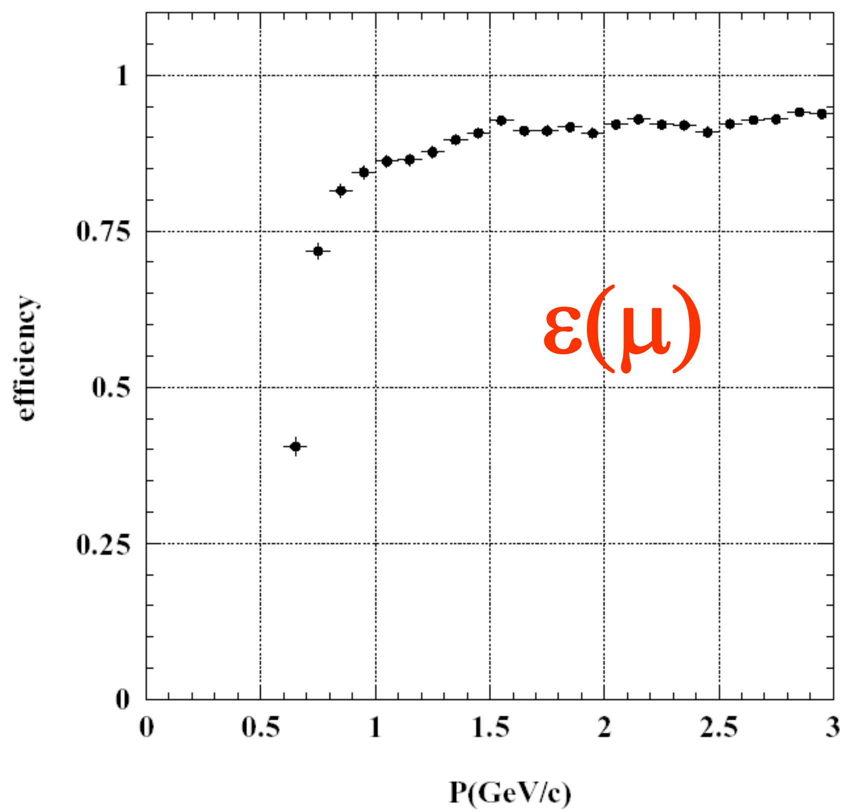


Fig. 109. Muon detection efficiency vs. momentum in KLM.

fake probability

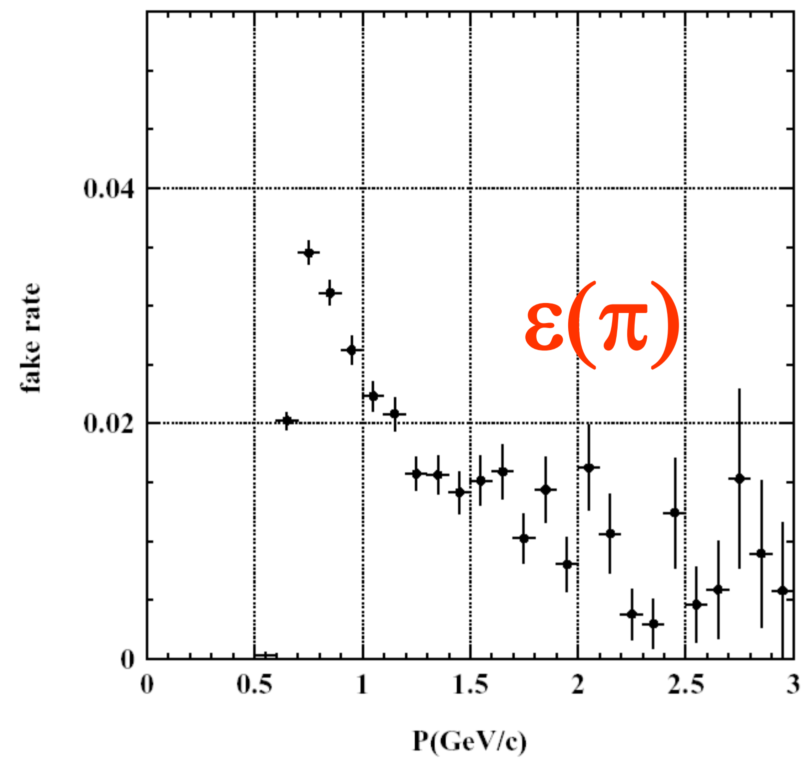
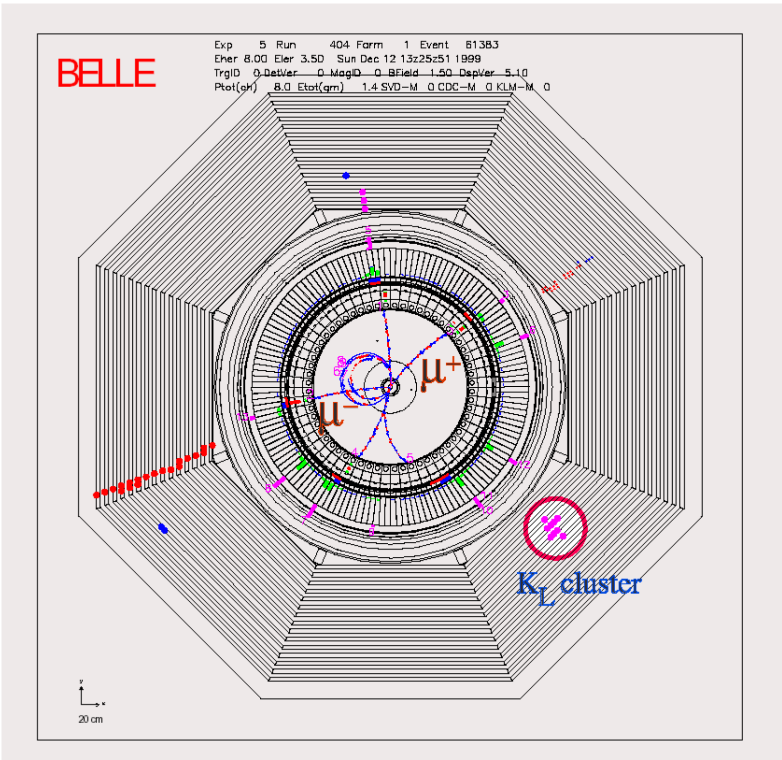


Fig. 110. Fake rate vs. momentum in KLM.

K_L detection performance



K_L detection: resolution in K_L direction

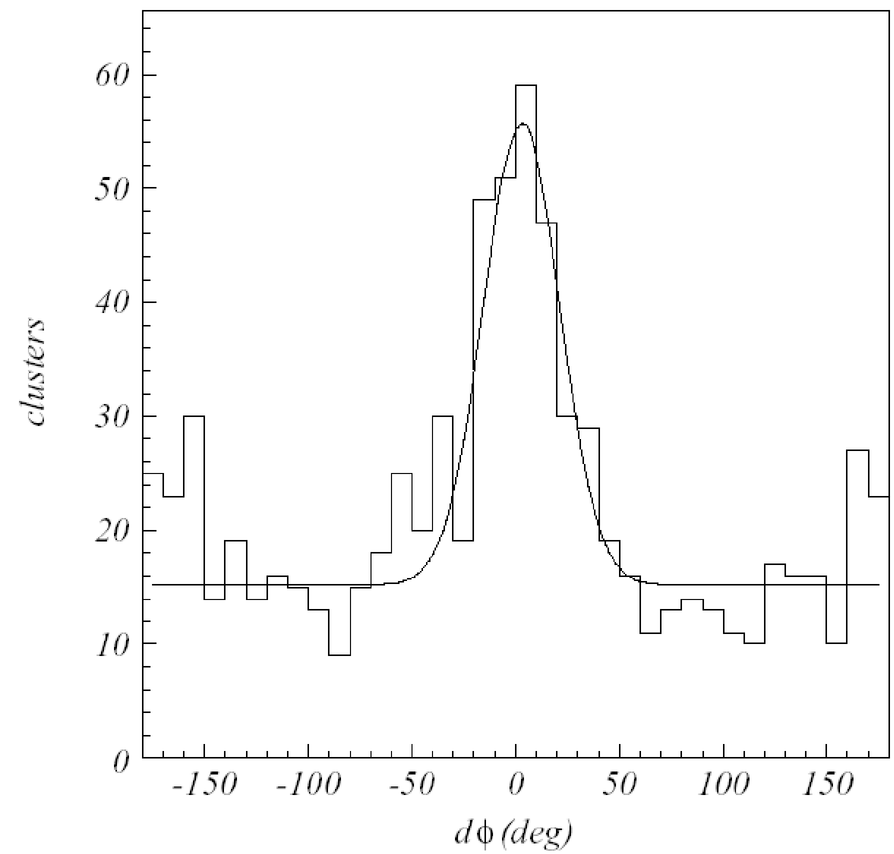
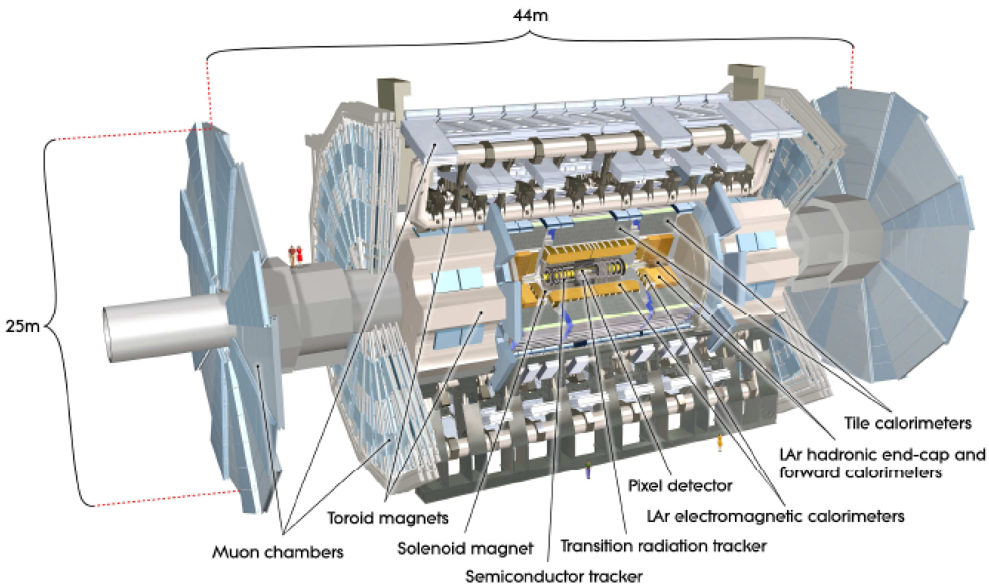
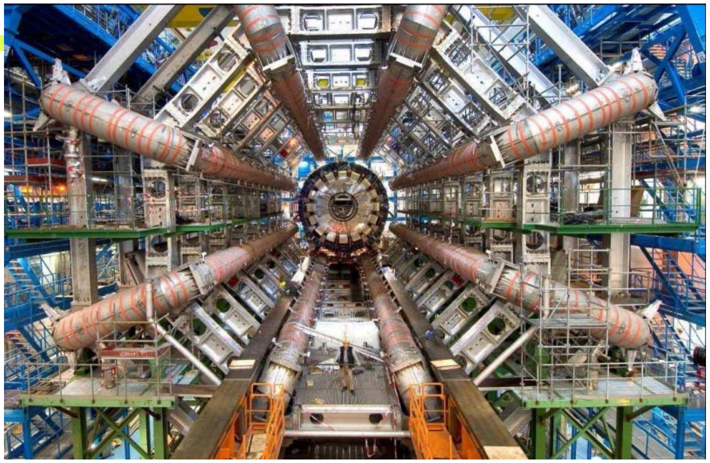


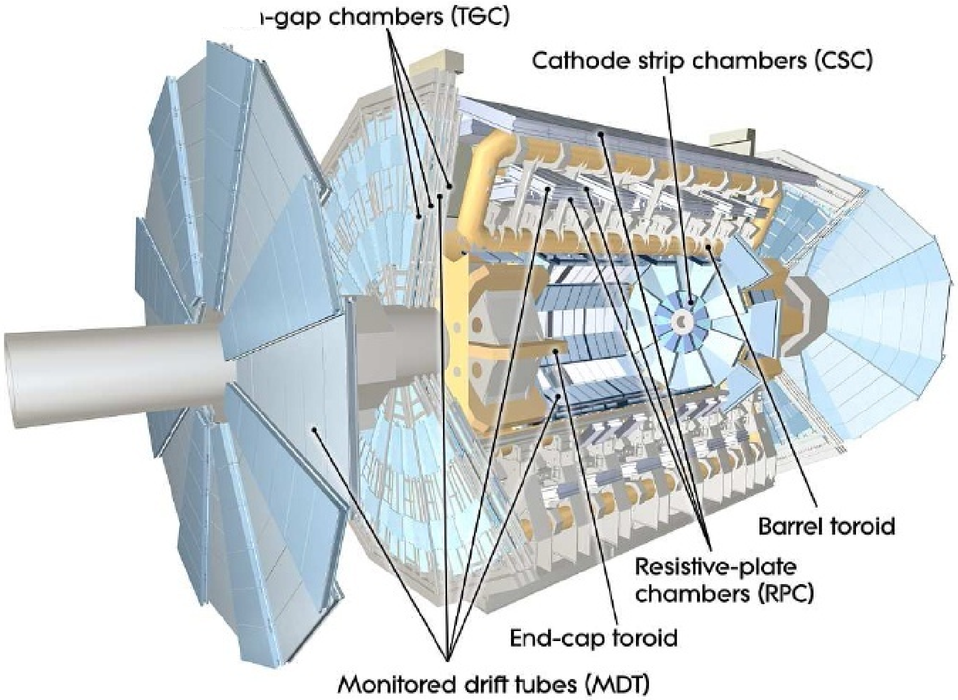
Fig. 107. Difference between the neutral cluster and the direction of missing momentum in KLM.

Identification of muons in ATLAS

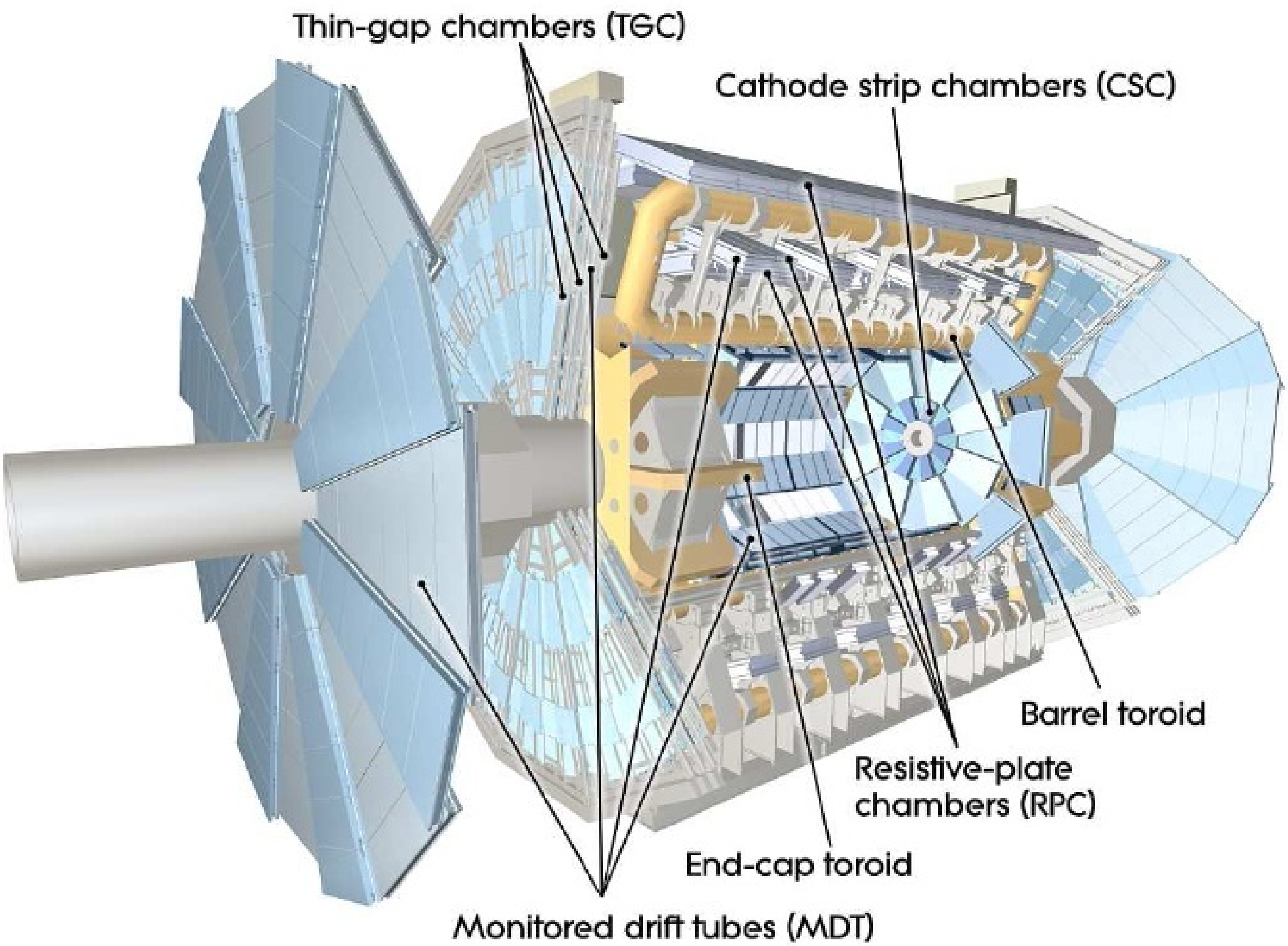




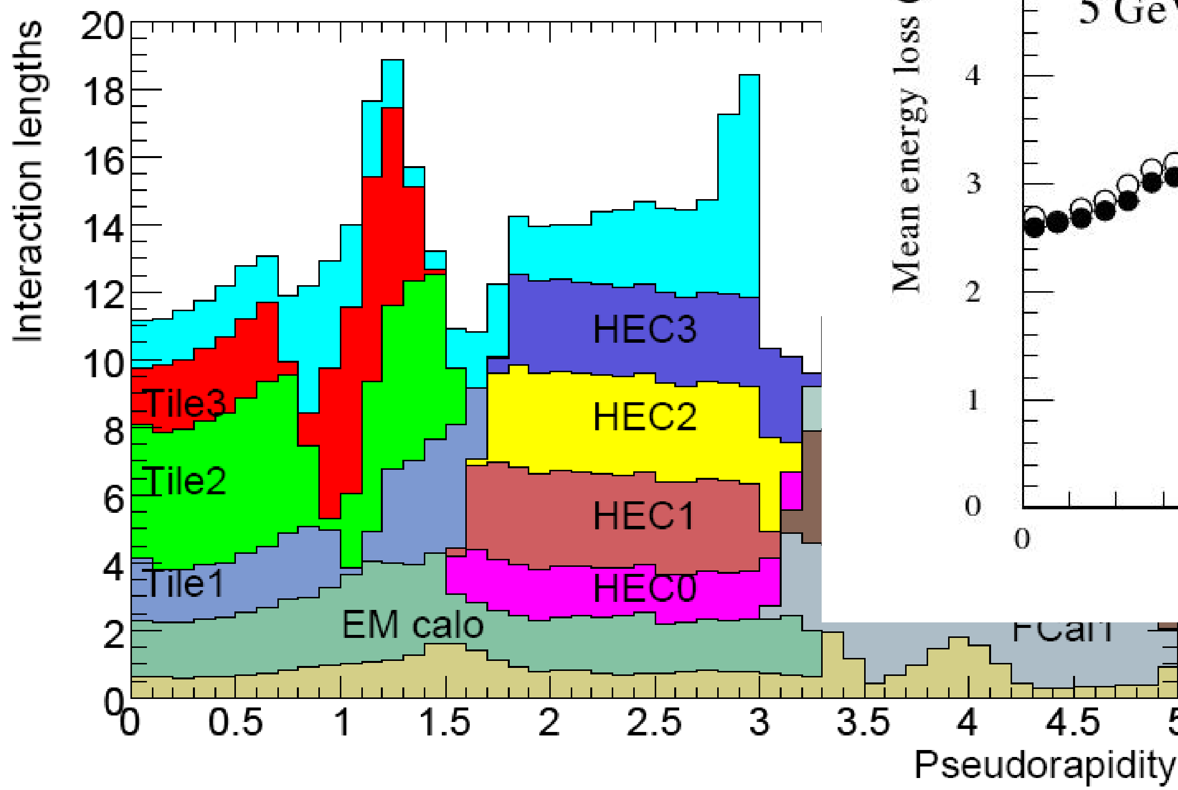
- Identify muons
- Measure their momentum

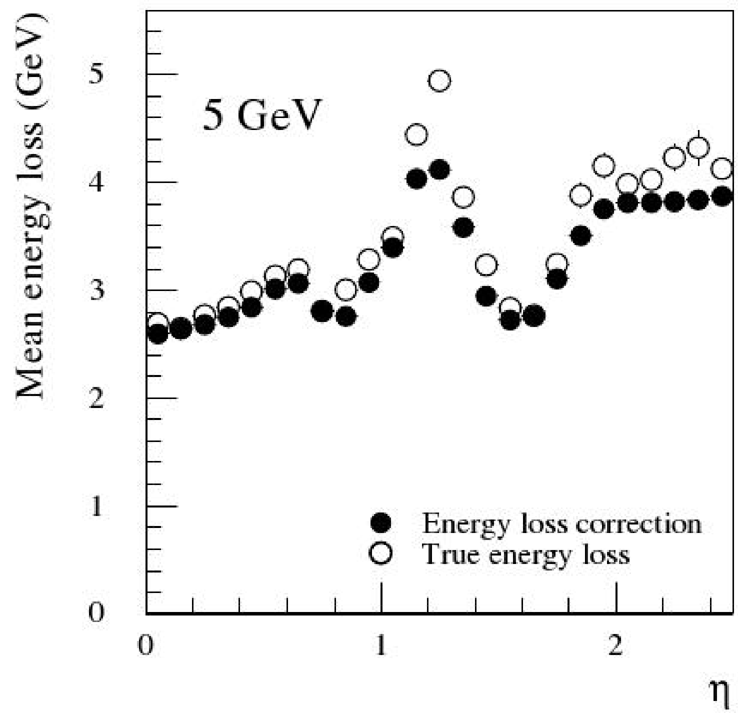


Detection of muons in ATLAS



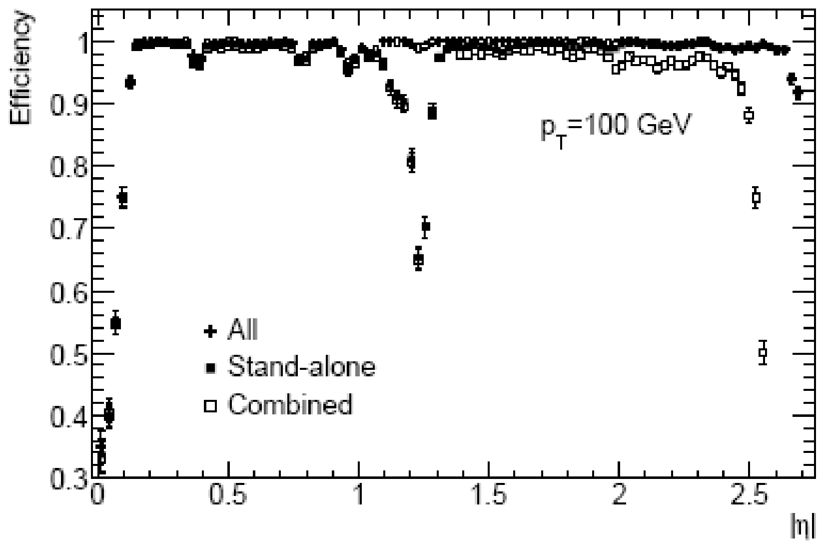
Muon identification in ATLAS

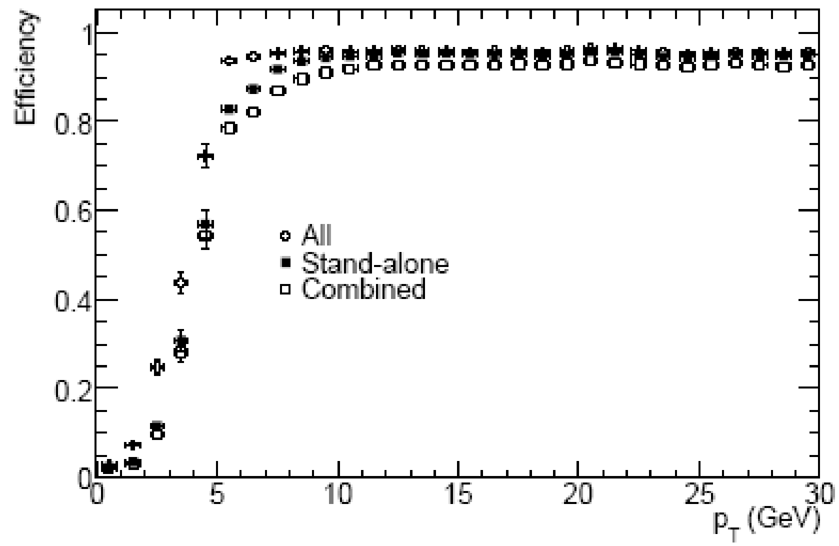




Material in front of the muon system

Muon identification efficiency





Efficiency for 100 GeV muons

Efficiency vs p_T

Summary

Particle identification is an essential part of several experiments, and has contributed substantially to our present understanding of elementary particles and their interactions. It will continue to have an important impact in searches for new physics.

A large variety of techniques has been developed for different kinematic regions and different particles.

New concepts and detectors are being studied → this is a very active area of detector R+D.