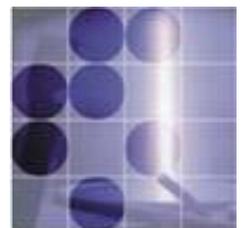
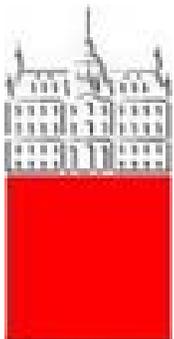


## **Recent advances in particle identification methods**

Peter Križan

*University of Ljubljana and J. Stefan Institute*



# Contents

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Why particle identification?

Ring Imaging Cherenkov counters

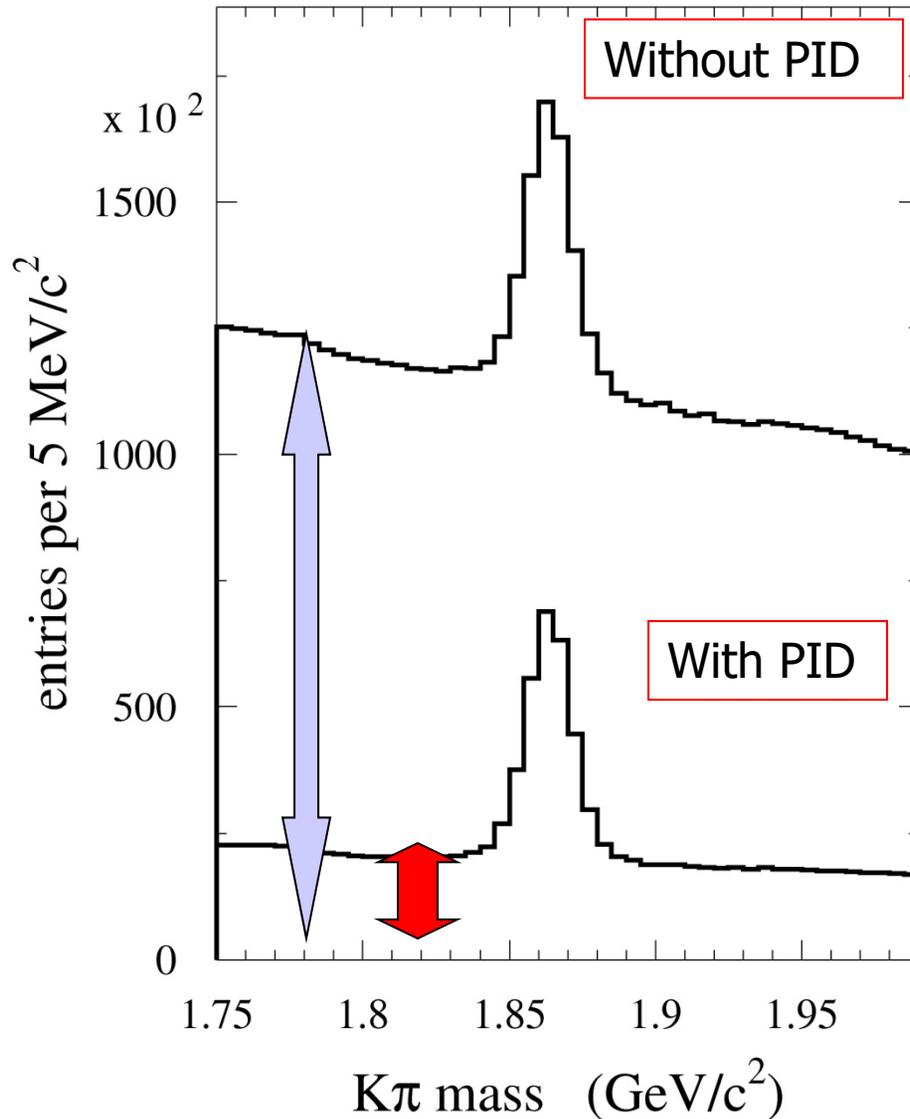
dE/dx and TOF

Transition radiation detectors

Summary

# Why particle ID?

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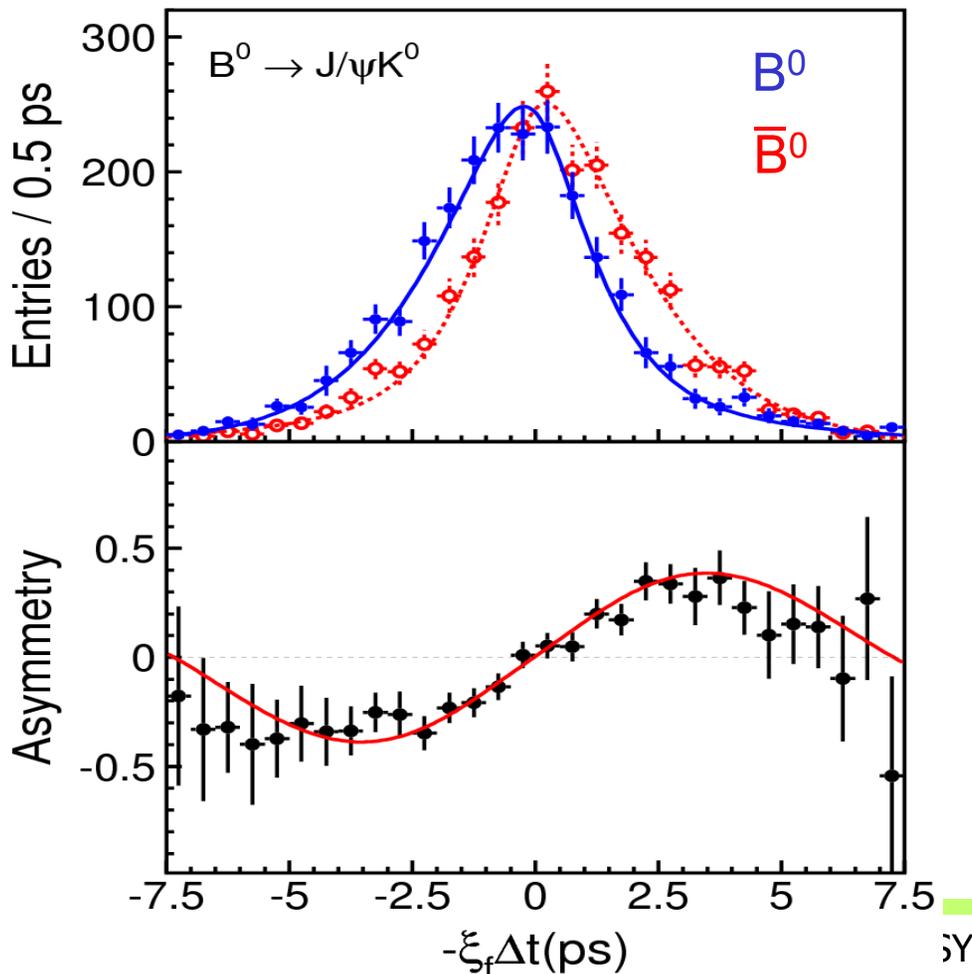


Example 1: BaBar (B factory)

Particle identification reduces the fraction of wrong  $K\pi$  combinations (combinatorial background) by  $\sim 5x$

# Why particle ID?

Particle identification at B factories (Belle and BaBar):  
was essential for the observation of **CP violation in the B meson system**.

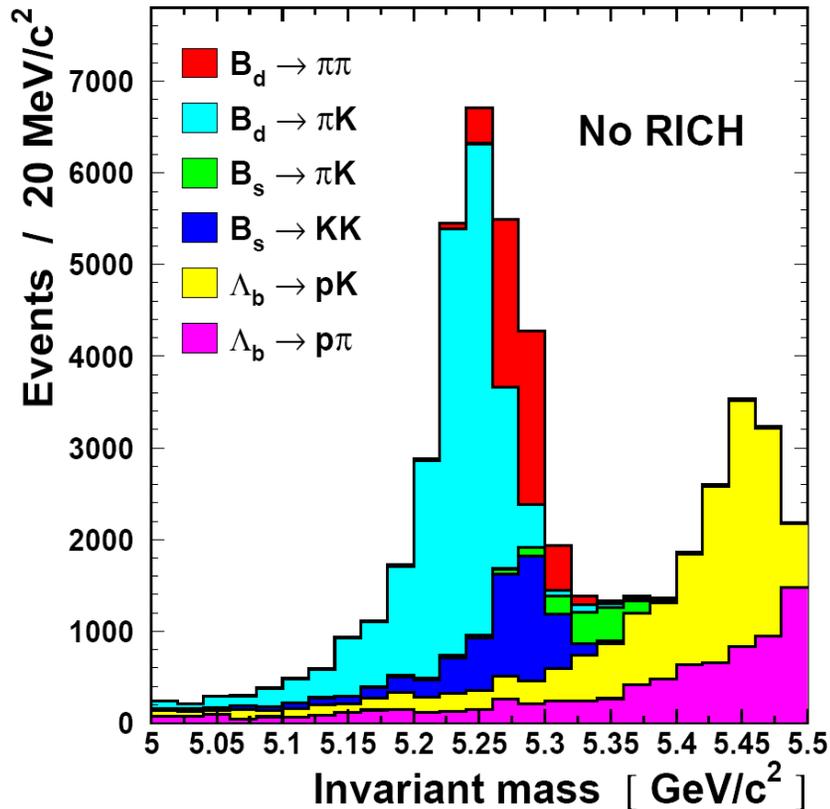


$B^0$  and its **anti-particle**  
**decay differently** to the  
same final state  $J/\psi K^0$

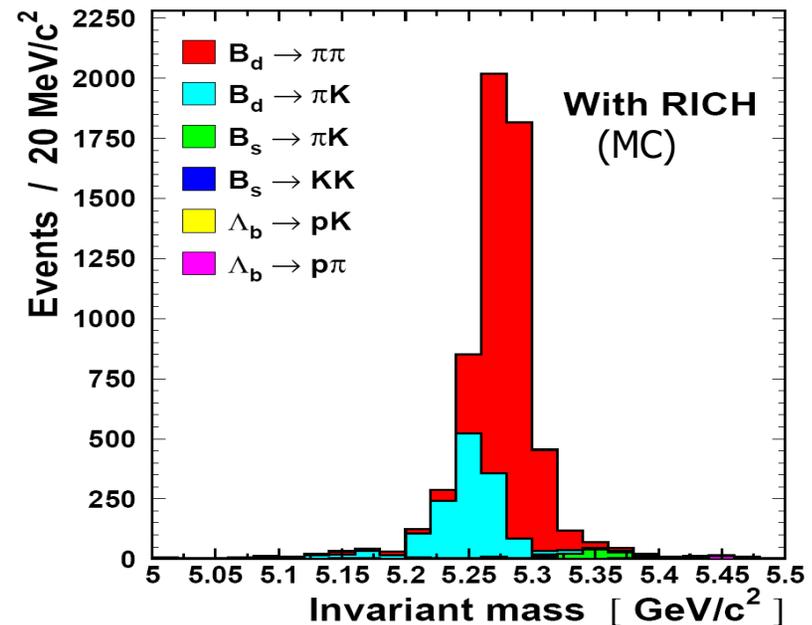
Flavour of the B: from decay  
products of the other B:  
charge of the kaon, electron,  
muon

→ **particle ID is compulsory**

# Why particle ID?

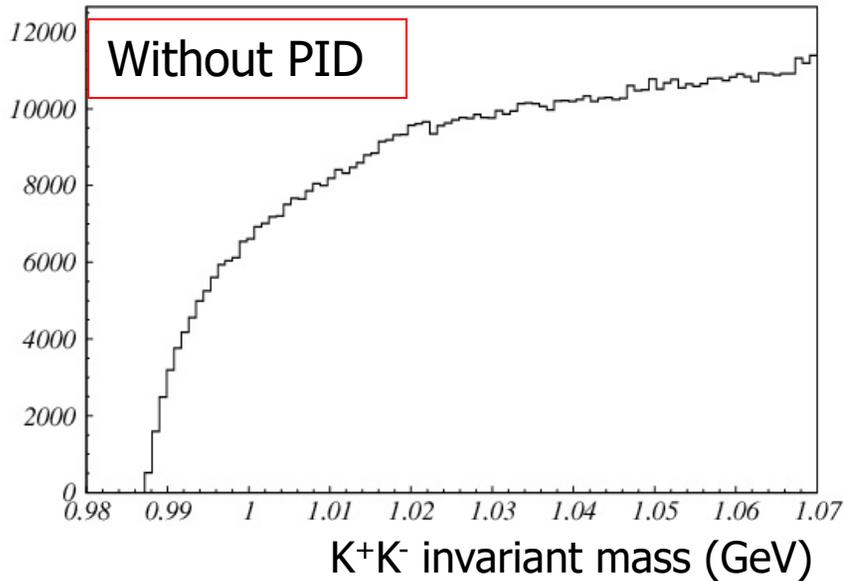


## Example 2: LHCb



Need to distinguish  $B_d \rightarrow \pi\pi$  from other similar topology 2-body decays and to distinguish B from anti-B using K tag.

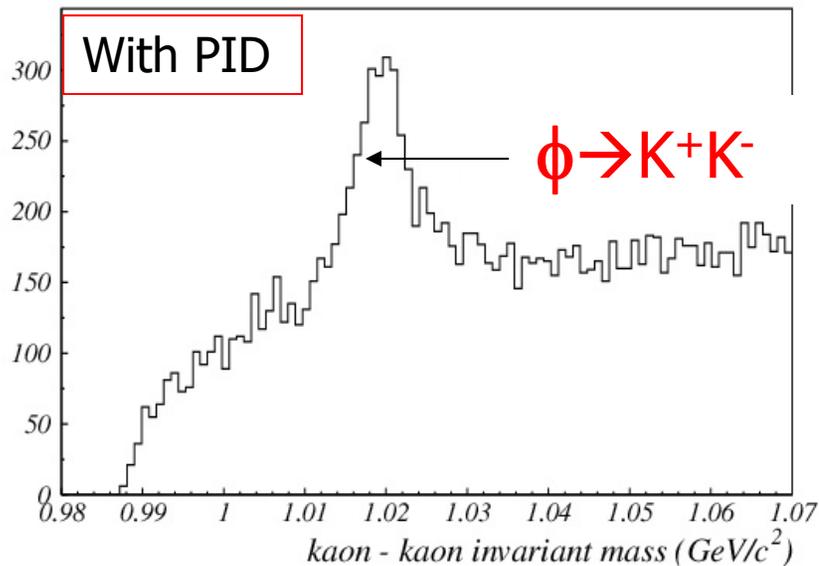
# Why particle ID?



Example 3: HERA-B

K<sup>+</sup>K<sup>-</sup> invariant mass.

The inclusive  $\phi \rightarrow K^+K^-$  decay only becomes visible after particle identification is taken into account.



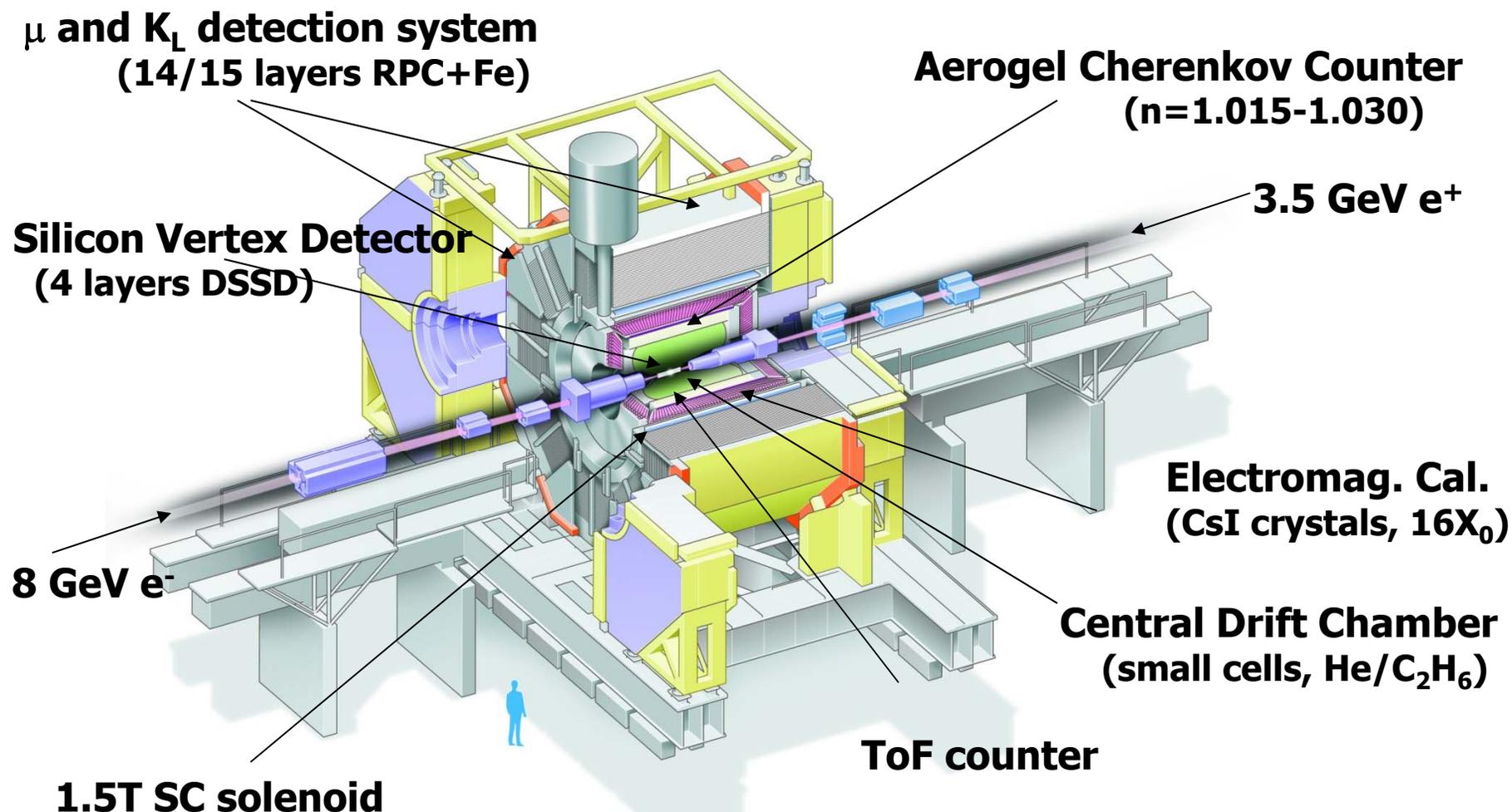
# Why particle ID?

---

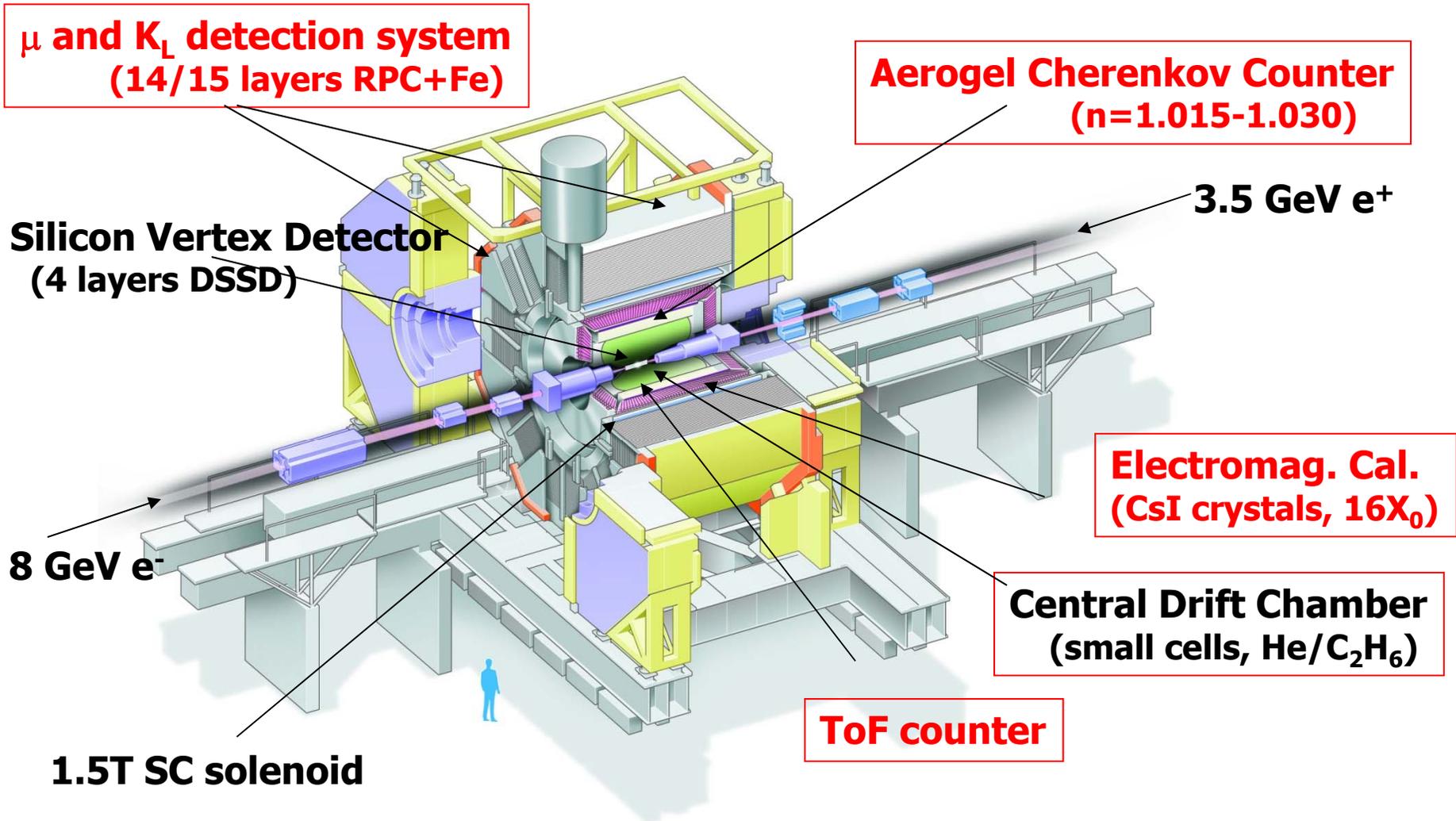
PID is also needed in:

- General purpose LHC experiments: final states with electrons and muons
- Searches for exotic states of matter (quark-gluon plasma)
- Spectroscopy and searches for exotic hadronic states
- Studies of fragmentation functions

# Example: Belle

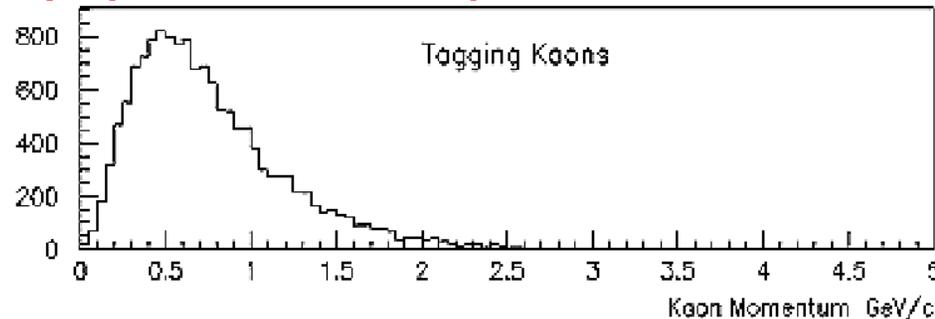


# Particle identification systems in Belle



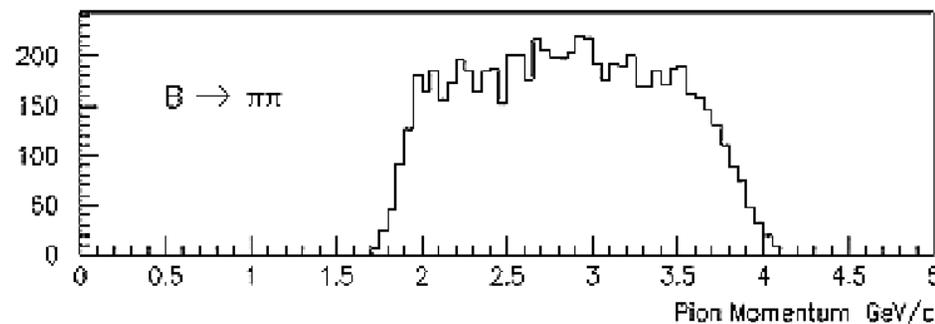
# Particle identification methods depend on the requirements (physics channel, kinematics)

## Example: B factory, pion/kaon separation



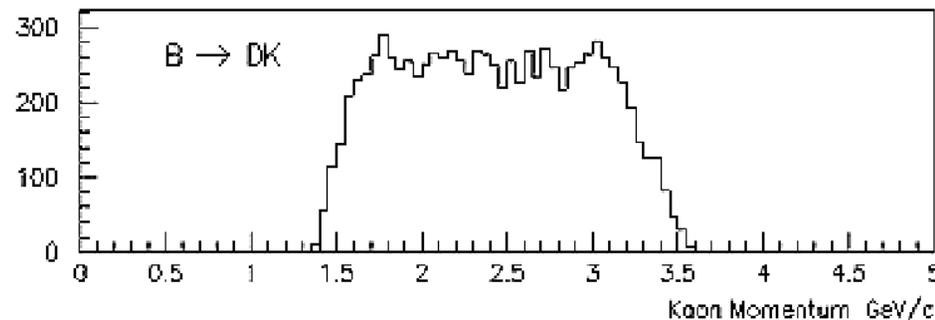
Tagging Kaons

Relatively soft,  
ms dominated  
for tracking



$B \rightarrow \pi\pi$

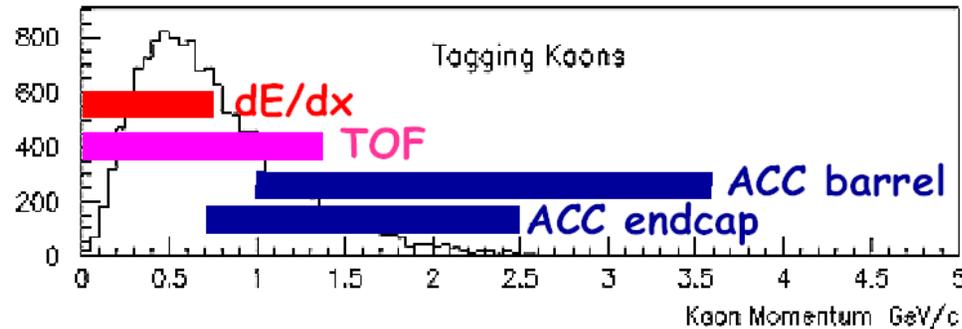
Requires  
dedicated PID



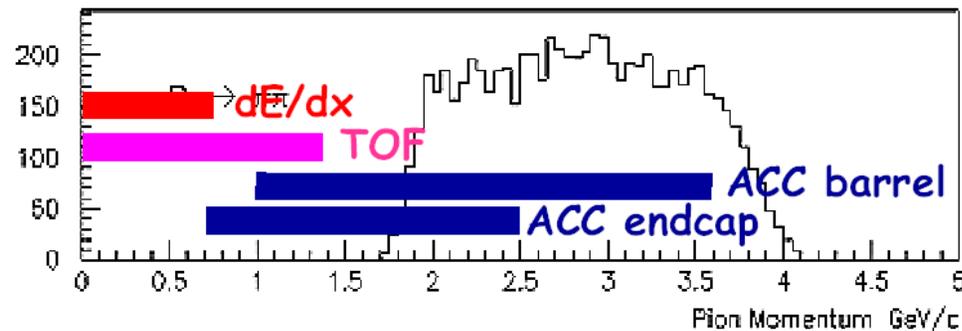
$B \rightarrow DK$

Requires  
dedicated PID

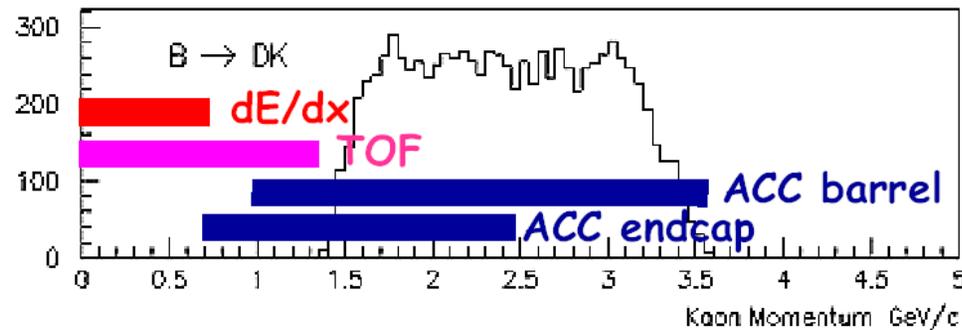
# PID coverage of kaon/pion spectra in Belle



Tagging Kaons

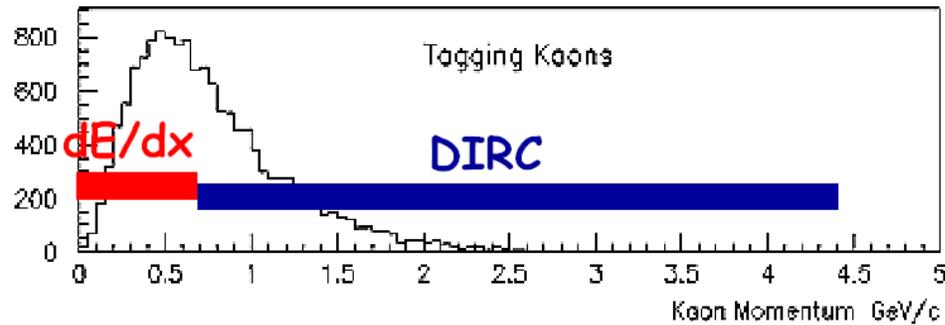
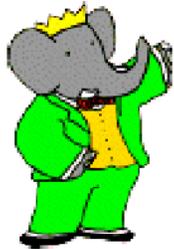


$B \rightarrow \pi\pi$

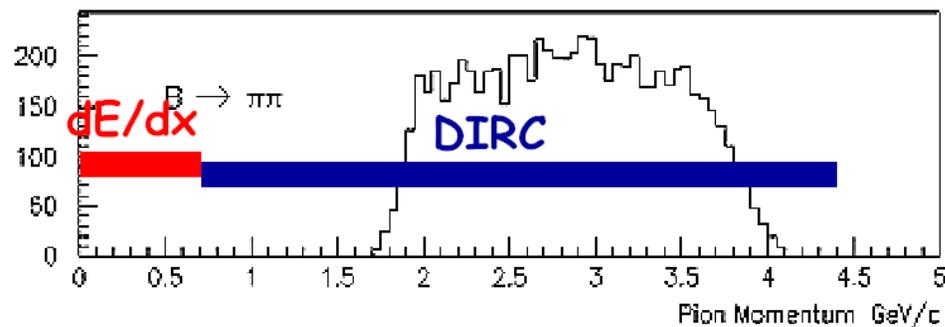


$B \rightarrow DK$

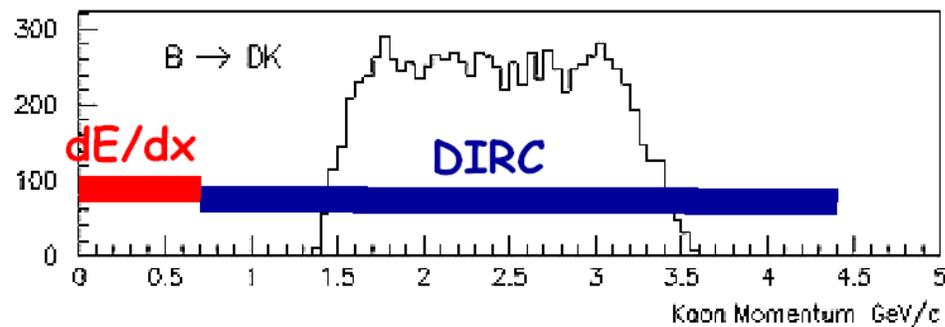
# PID coverage of kaon/pion spectra in BaBar



Tagging Kaons



$B \rightarrow \pi\pi$



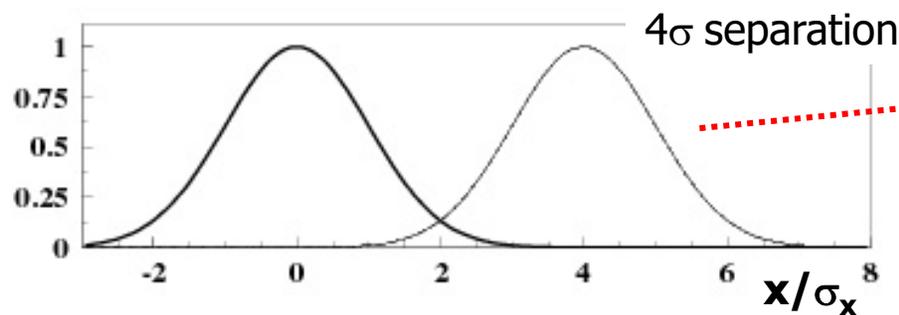
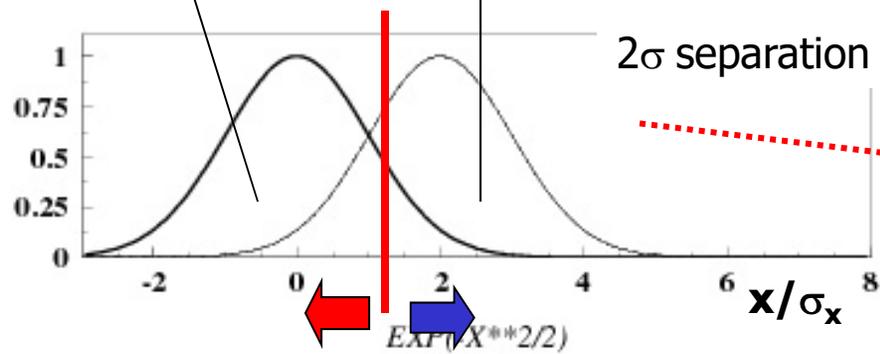
$B \rightarrow DK$

# Efficiency and purity in particle identification

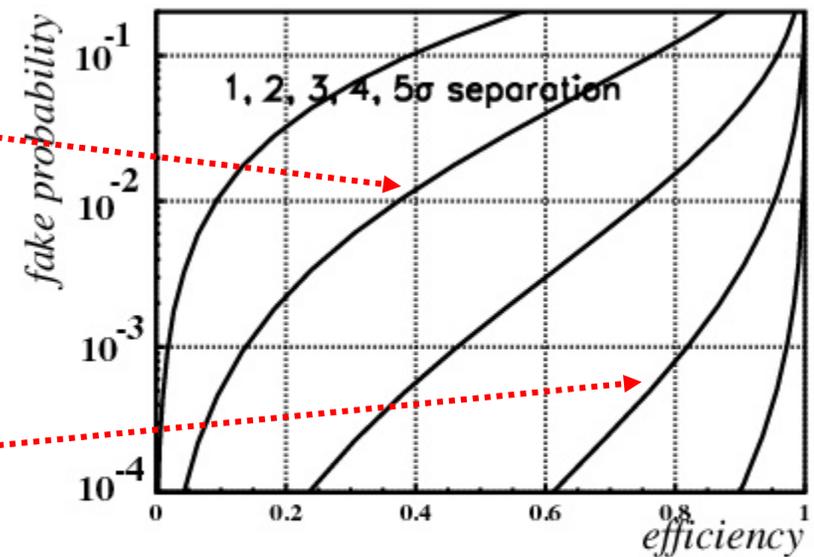
Efficiency and purity are tightly coupled!

Two examples:

particle type 1    type 2



efficiency vs fake probability  
(for Gaussian distributions)



some discriminating variable  $x$ , scaled to the resolution  $\sigma_x$

# Identification of charged particles

---

Particles (e,  $\mu$ ,  $\pi$ , K, p) in the final state are identified by their **mass** or by the **way they interact**.

Determination of **mass**: from the relation between momentum and velocity,  $p = \gamma m v$  (p is known - radius of curvature in the magnetic field)

→ Measure velocity by:

- time of flight
- ionisation losses  $dE/dx$
- Cherenkov photon angle (and/or yield)
- transition radiation

Mainly used for the identification of hadrons.

Identification through **interaction**: electrons and muons

- muon systems
- calorimeters

# Identification of charged particles

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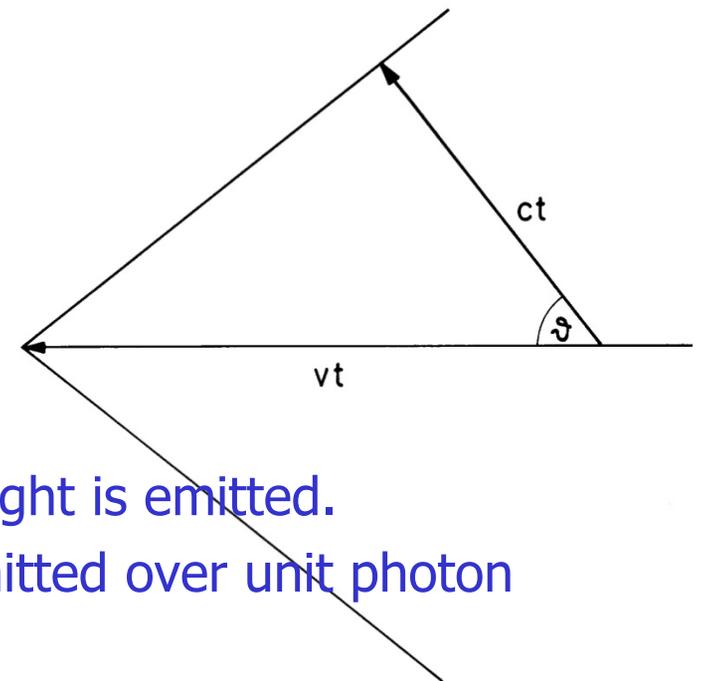
Identification through **interaction**: electrons and muons

- muon systems
- calorimeters

# Cherenkov radiation

A charged track with velocity  $v = \beta c$  exceeding the speed of light  $c/n$  in a medium with refractive index  $n$  emits **polarized light** at a characteristic (Cherenkov) angle,

$$\cos\theta = c/nv = 1/\beta n$$



Two cases:

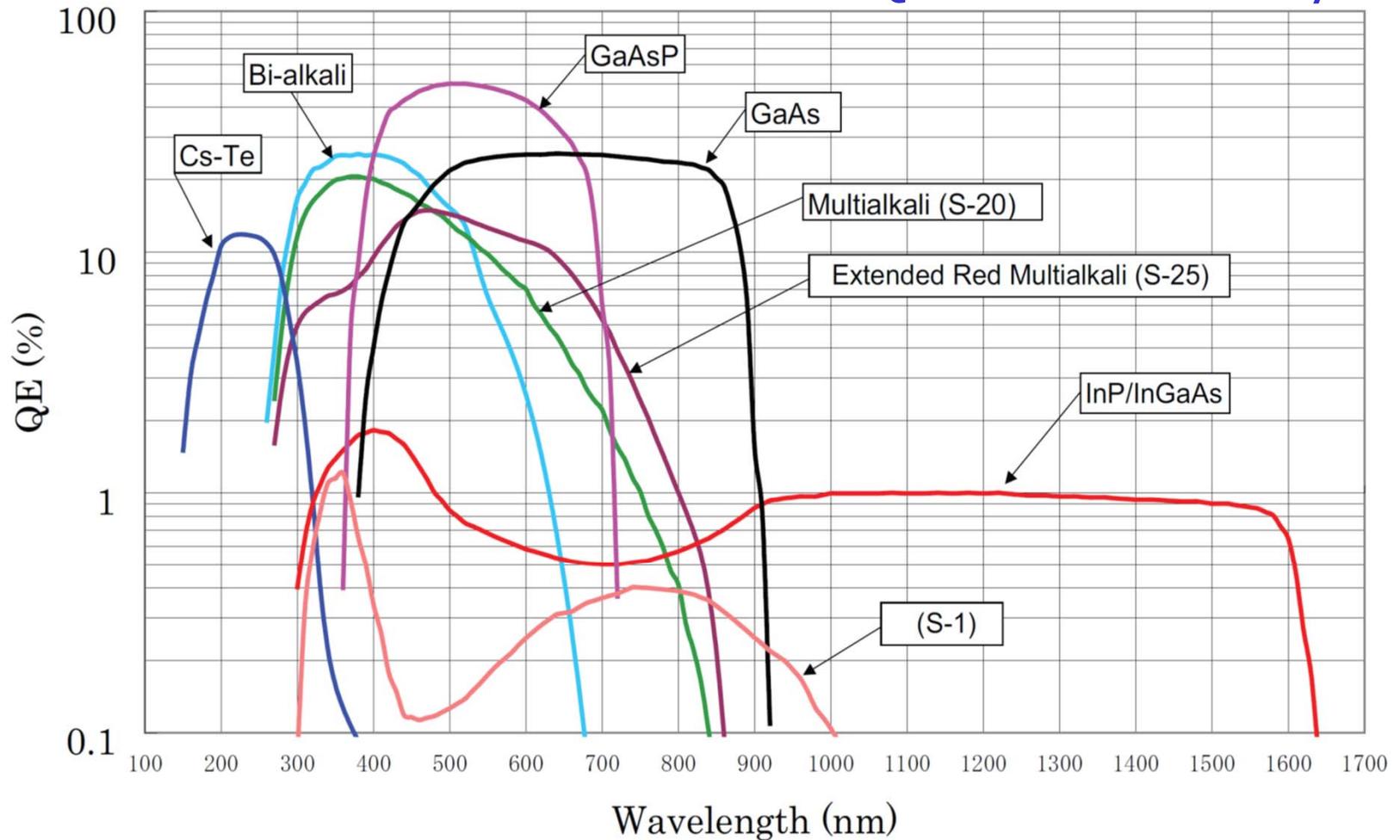
- $\beta < \beta_t = 1/n$ : below threshold **no** Cherenkov light is emitted.
- $\beta > \beta_t$ : the number of Cherenkov photons emitted over unit photon energy  $E = h\nu$  in a radiator of length  $L$ :

$$\frac{dN}{dE} = \frac{\alpha}{\hbar c} L \sin^2 \theta = 370(\text{cm})^{-1} (\text{eV})^{-1} L \sin^2 \theta$$

→ Few detected photons

# Photon detection efficiency

## Quantum efficiency



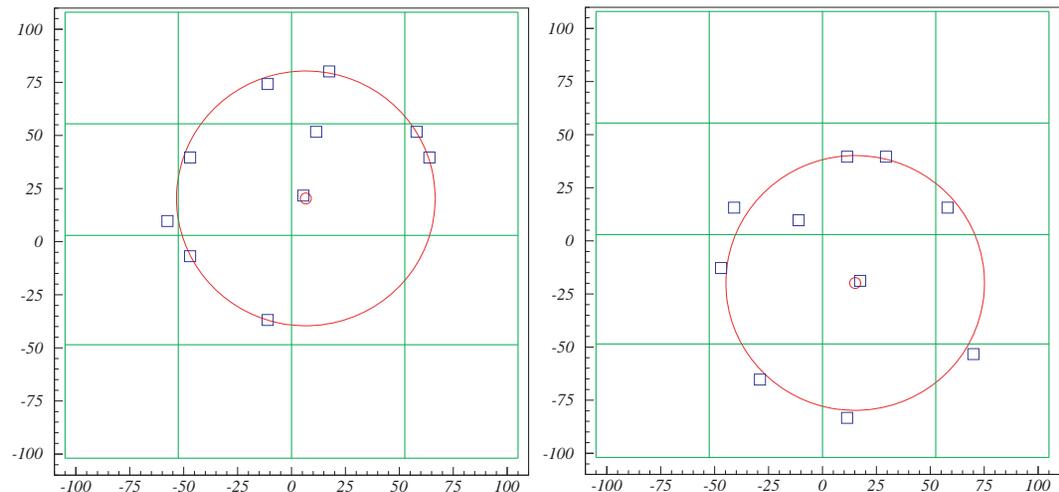
## Number of detected photons

Example: in 1m of air ( $n=1.00027$ ) a track with  $\beta=1$  emits  **$N=41$  photons** in the spectral range of visible light ( $\Delta E \sim 2$  eV).

If Čerenkov photons were detected with an average detection efficiency of  $\varepsilon=0.1$  over this interval,  **$N=4$  photons** would be measured.

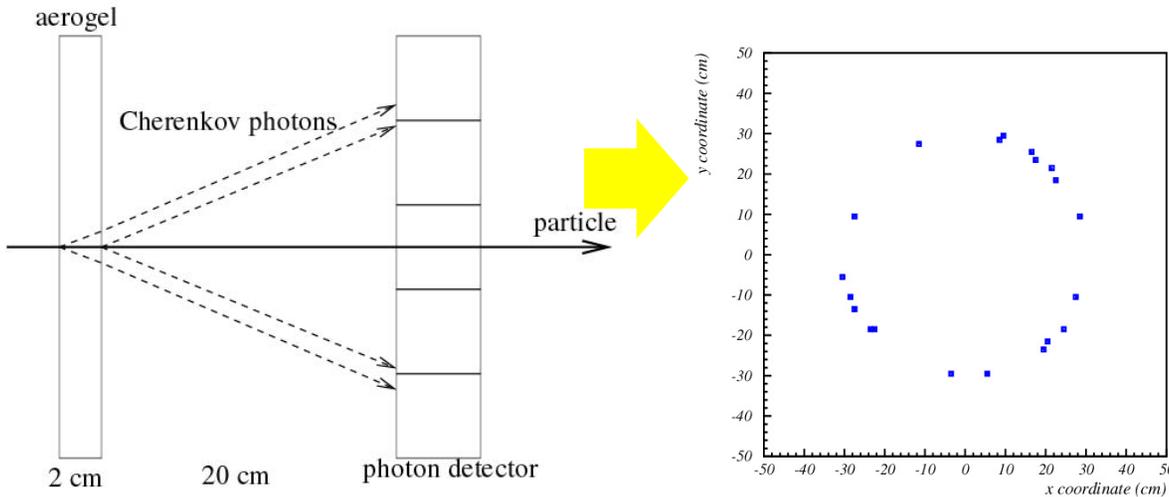
**Few** photons detected

→ Important to have a **low noise** detector



# Measuring the Cherenkov angle

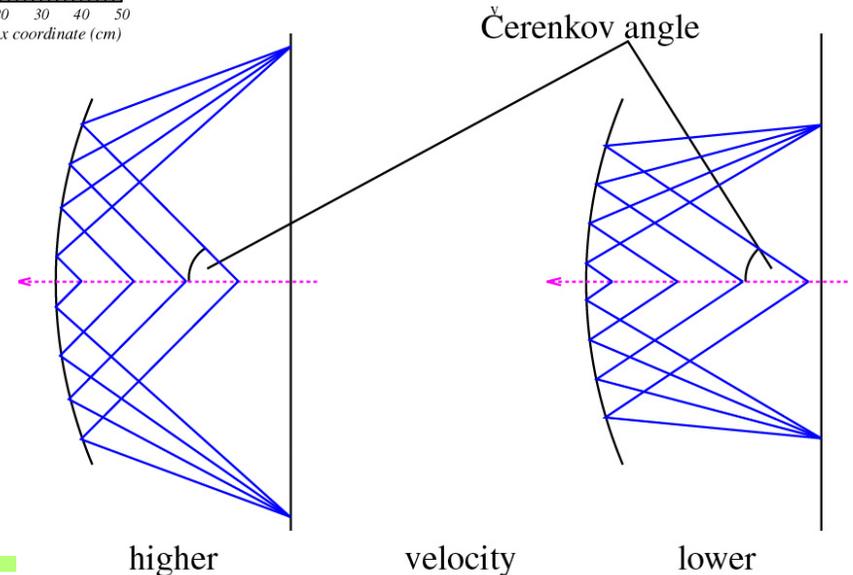
Particles above threshold: measure  $\theta$



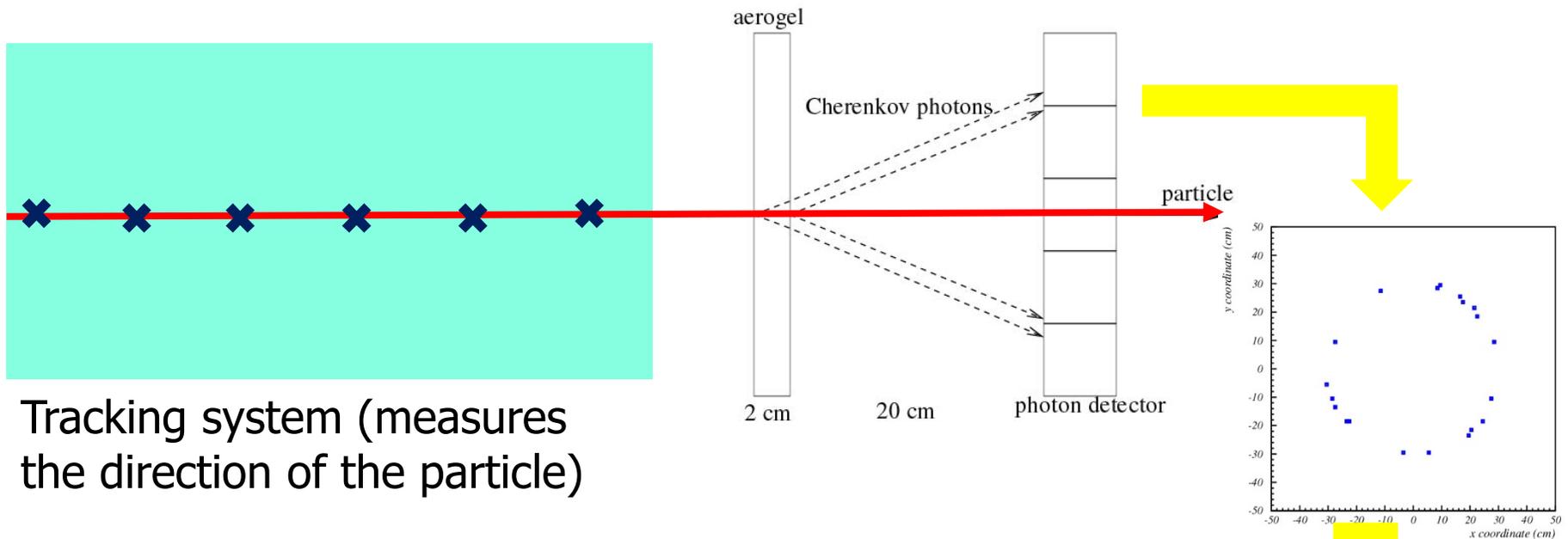
Idea: transform the direction into a coordinate  $\rightarrow$  ring on the detection plane  $\rightarrow$  Ring Imaging Cherenkov (RICH) counter

Proximity focusing RICH

RICH with a focusing mirror



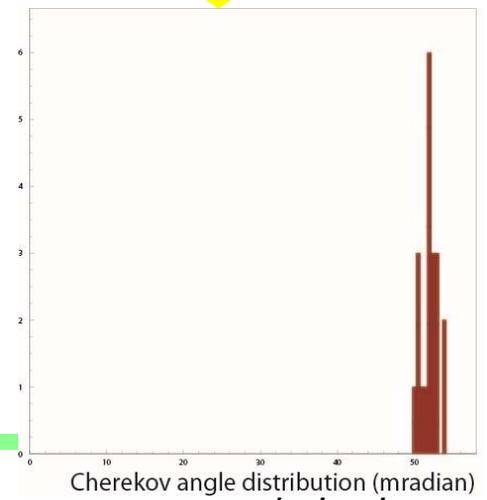
# Measuring the Cherenkov angle



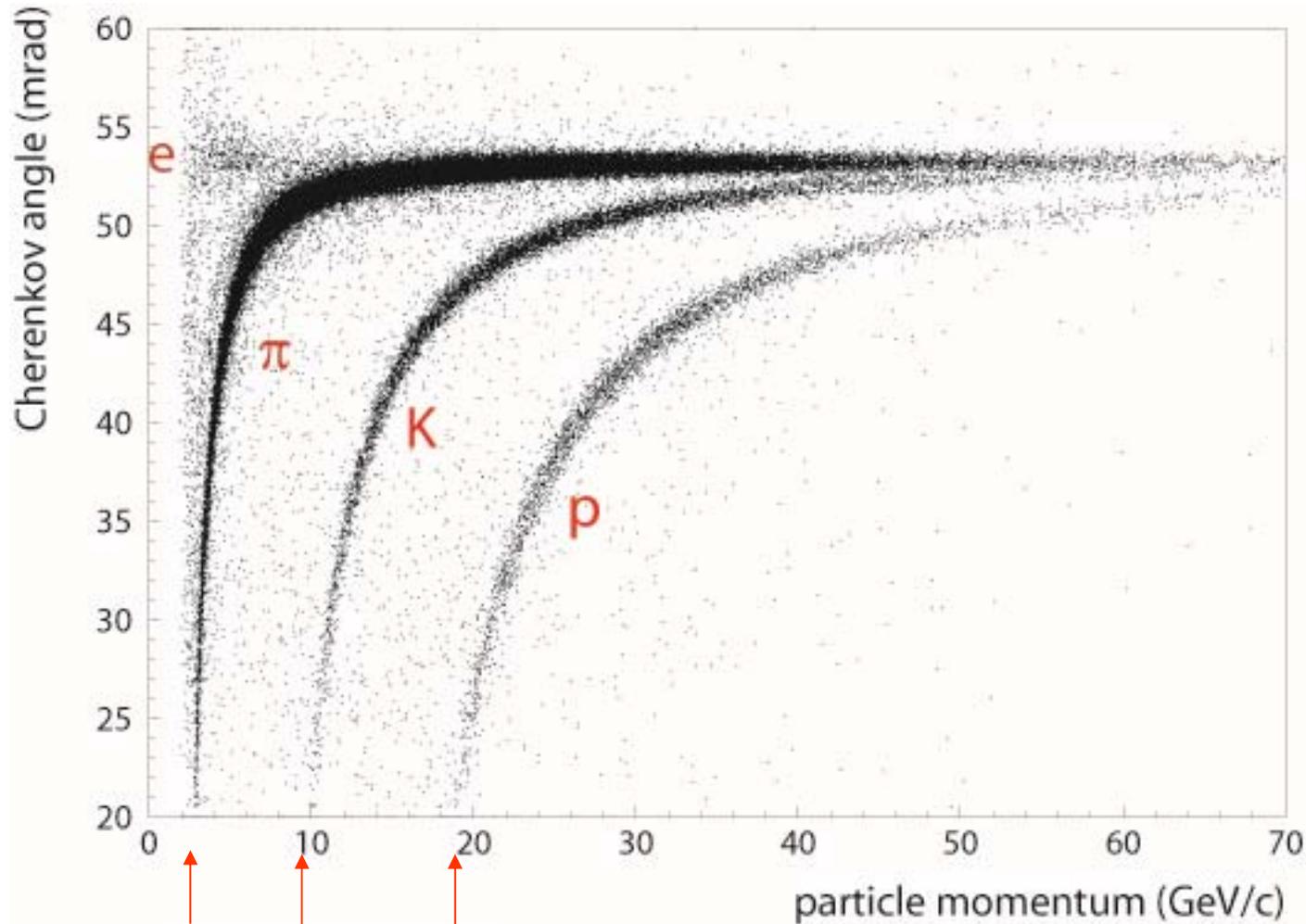
Tracking system (measures the direction of the particle)

Tracking system tells us where the particle hit the radiator, and at which angle.

Use this information to calculate the **Cherenkov angle** for **each individual** detected photon



# Measuring Cherenkov angle



Radiator:  
 $C_4F_{10}$  gas

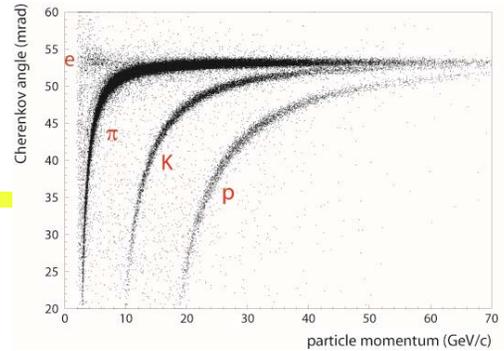
↑ π    ↑ K    ↑ p

Feb. 20 thresholds

EDIT@DESY

Peter Križan, Ljubljana

# Resolution of a RICH counter



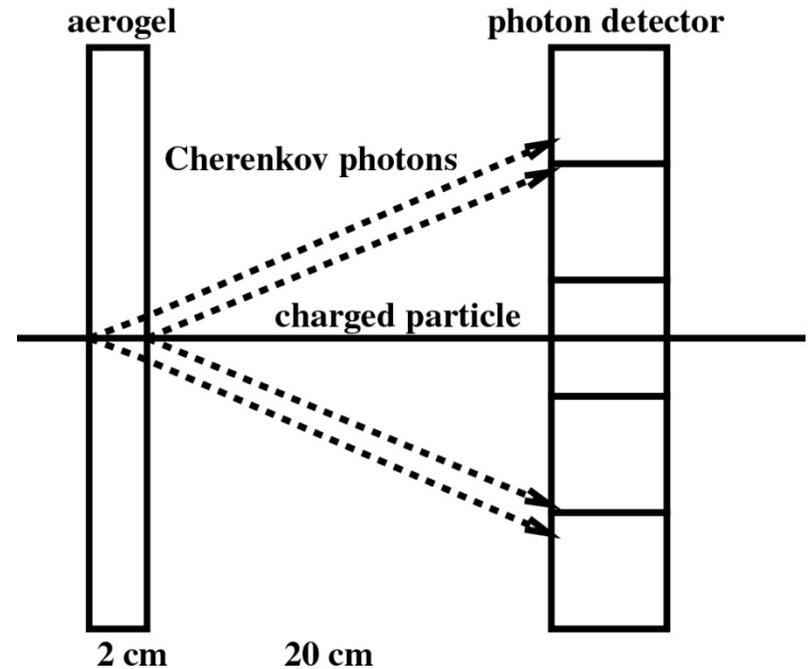
Determined by:

- Photon impact point resolution ( $\sim$ photon detector granularity)
- Emission point uncertainty (not in a focusing RICH)
- Dispersion:  $1/\beta = n(\lambda) \cos\theta$
- Errors of the optical system
- Uncertainty in track parameters

Resolution per track:

$$\sigma_{track} = \frac{\sigma_0}{\sqrt{N_{pe}}}$$

$\sigma_0$  ← single photon resolution  
 $N_{pe}$  ← # of detected photons



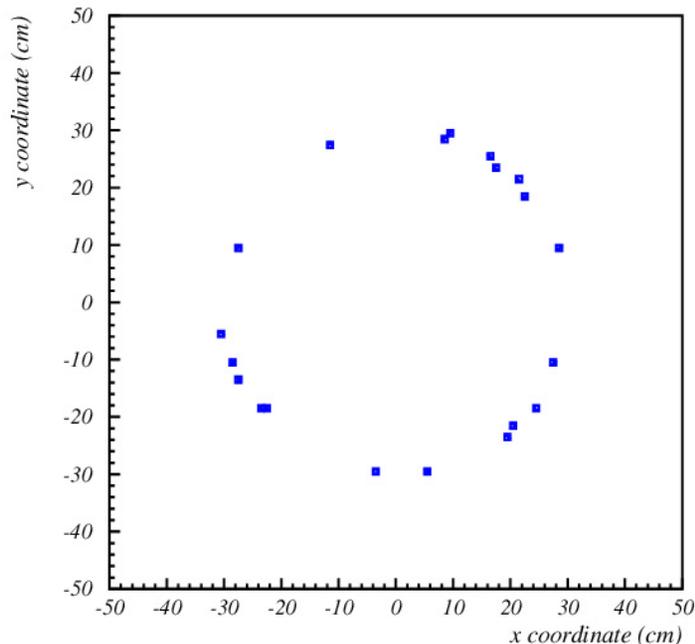
(in the case of low background)

# Photon detection in RICH counters

RICH counter: measure photon impact point on the photon detector surface

→ detection of **single** photons with

- sufficient **spatial resolution**
- **high efficiency** and **good signal-to-noise** ratio (few photons!)
- over a **large area** (square meters)



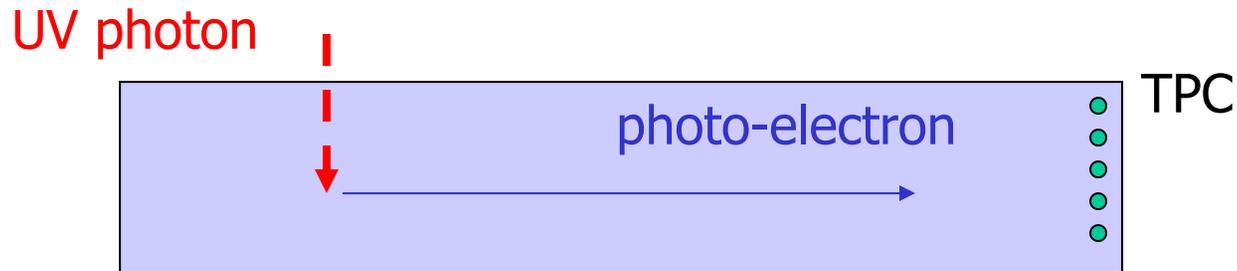
Special requirements:

- **Operation in magnetic field**
- High rate capability
- **Very high spatial resolution**
- **Excellent timing** (time-of-arrival information)

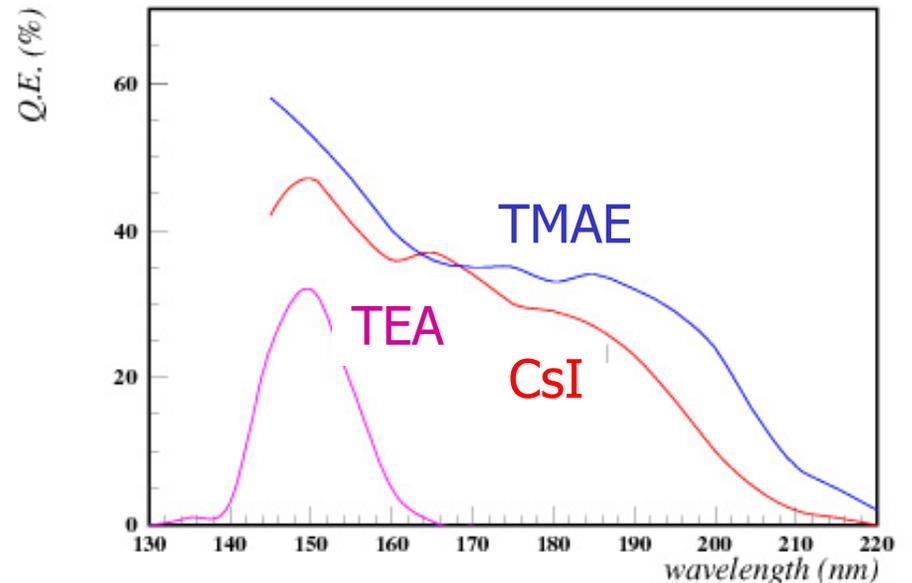
Photon detector is the most crucial element of a RICH counter

# First generation of RICH counters

DELPHI, SLD, OMEGA RICH counters: all employed wire chamber based photon detectors (UV photon  $\rightarrow$  photo-electron  $\rightarrow$  detection of a single electron in a TPC)

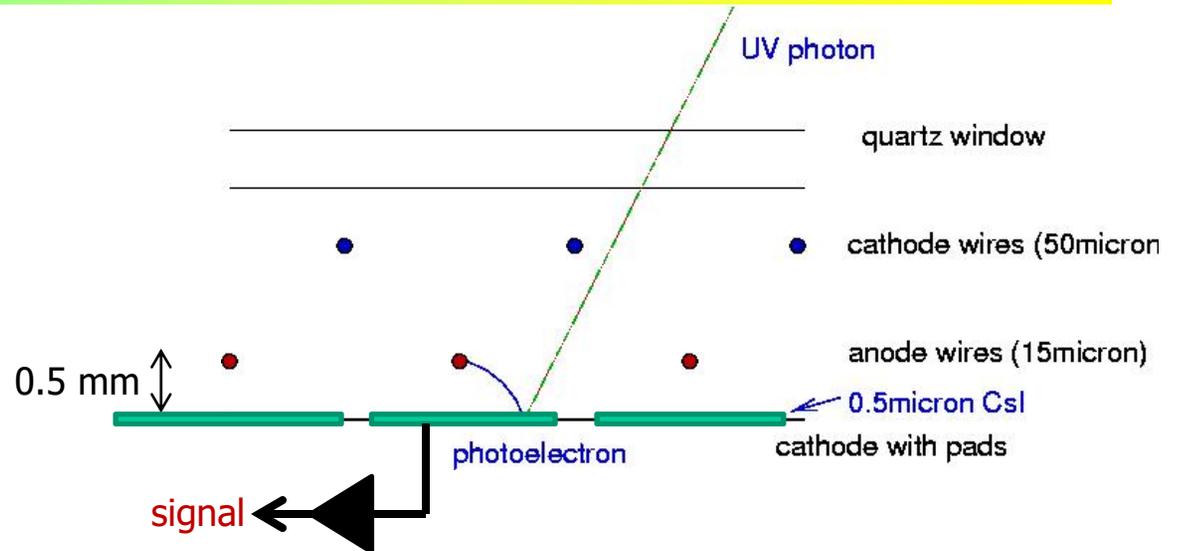


Photosensitive component:  
TMAE added to the gas mixture



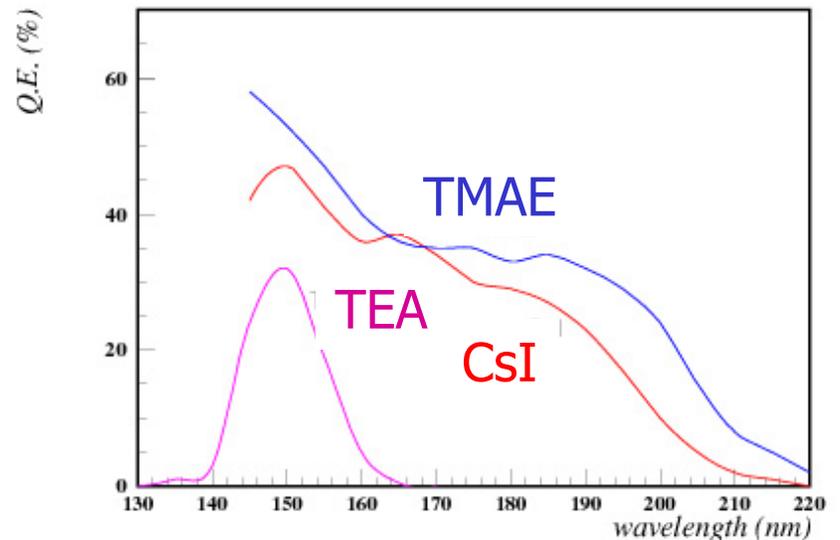
# Fast RICH counters with wire chambers

Multiwire chamber with  
**cathode pad read-out:**  
→ short drift distances,  
fast detector



Photosensitive component:

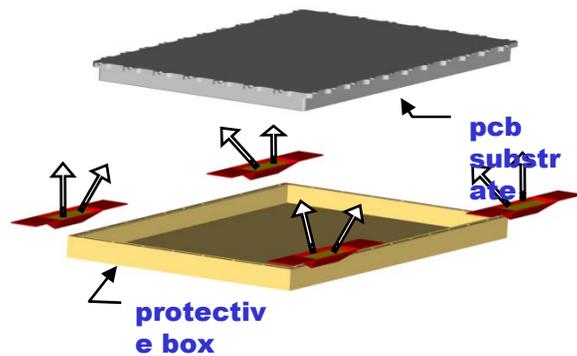
- in the gas mixture (**TEA**):  
CLEOIII RICH
- or a layer on one of the cathodes  
(**CsI** on the printed circuit cathode  
with pads) HADES, COMPASS,  
ALICE



Works in high magnetic field!

# CERN CsI deposition plant

Photocathode produced with a well defined, several step procedure, with CsI vacuum deposition and subsequent heat conditioning

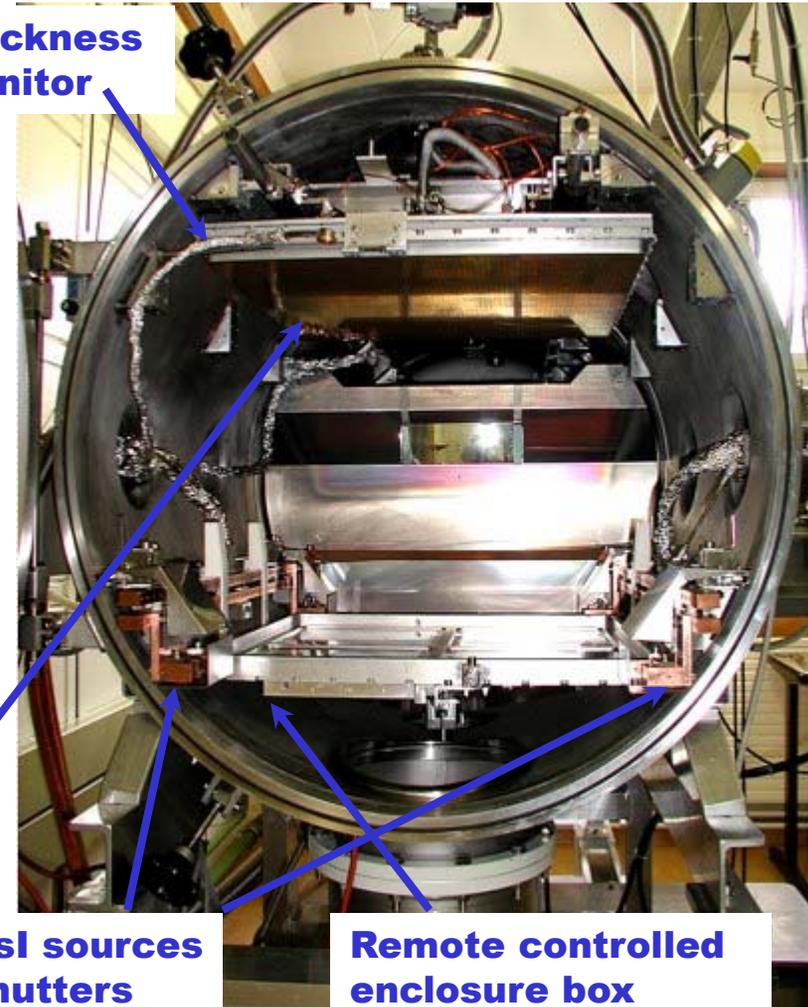


**Thickness monitor**

**PC**

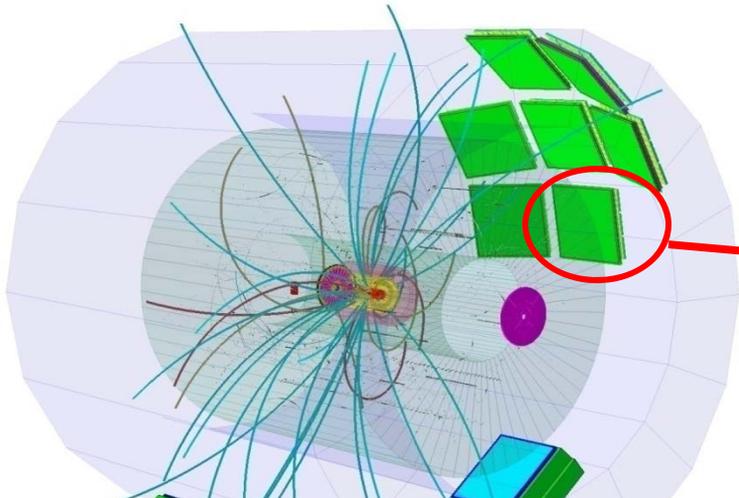
**4 CsI sources + shutters**

**Remote controlled enclosure box**

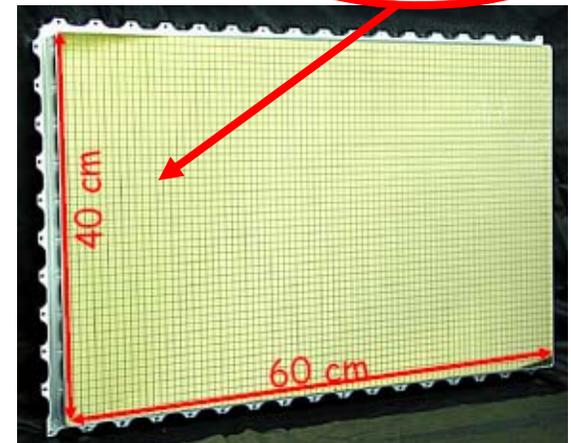
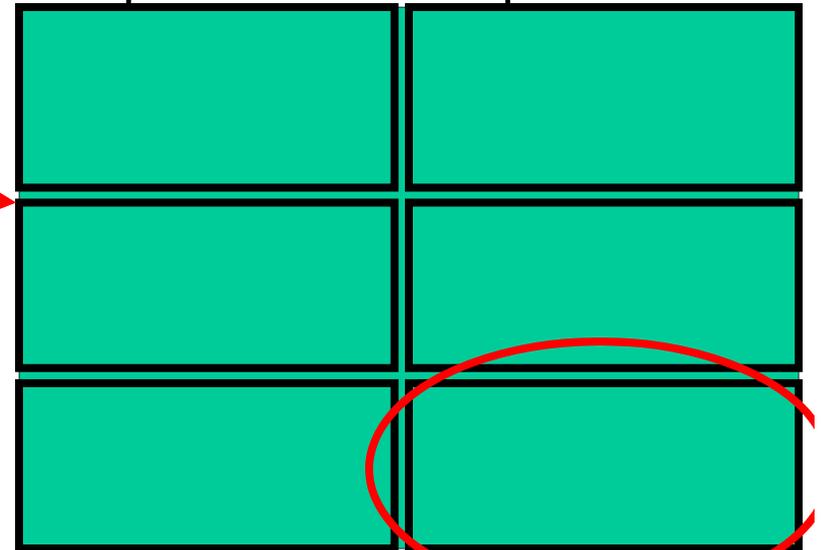


# ALICE RICH = HMPID

The largest scale (11 m<sup>2</sup>) application of CsI photo-cathodes in HEP!



Six photo-cathodes per module



CsI photo-cathode is segmented in **0.8x0.84 cm pads**

# Cherenkov counters with vacuum based photodetectors

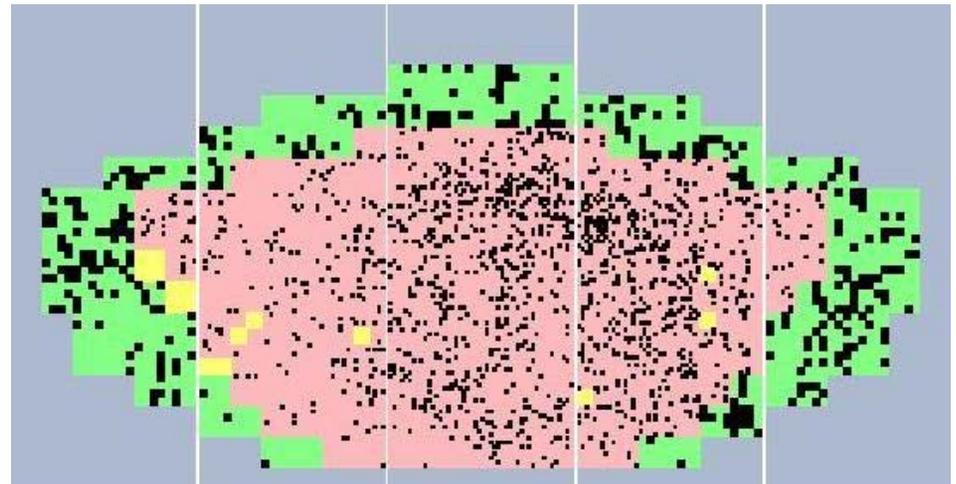
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Many applications: operation at high rates over extended running periods (years) → wire chamber based photon detectors were found to be unsuitable (problems in high rate operation, ageing, only UV photons, difficult handling in  $4\pi$  spectrometers)

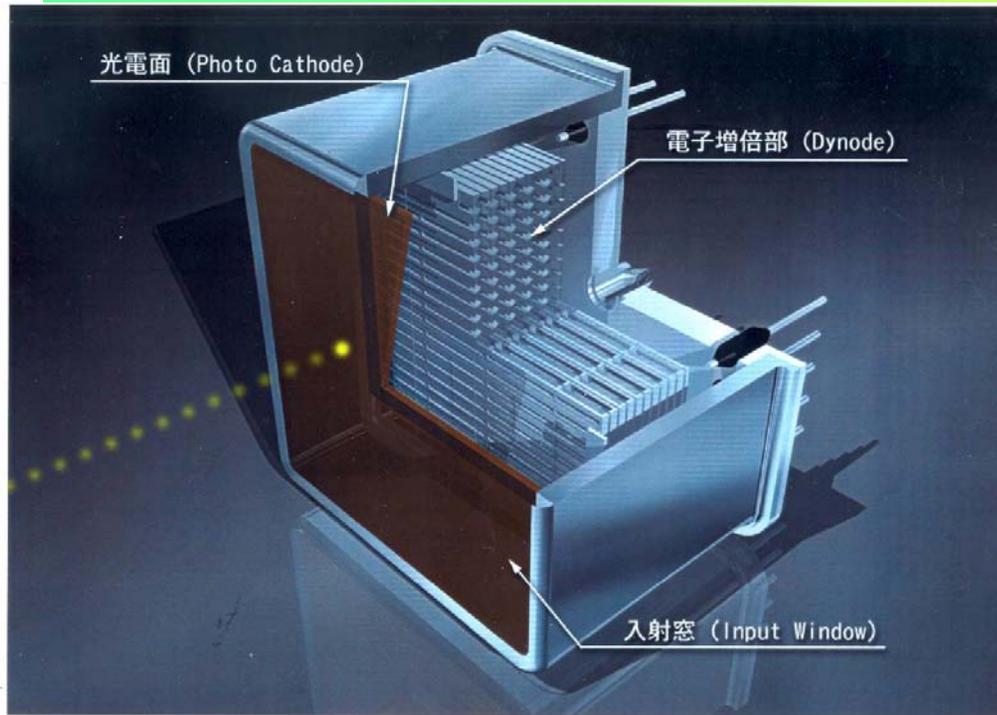
→ Need **vacuum based photon detectors** (e.g. PMTs)

Good spacial resolution (pads with  $\sim 5$  mm size)

→ Solution: **multianode** PMTs (MaPMTs)



# Multianode PMTs



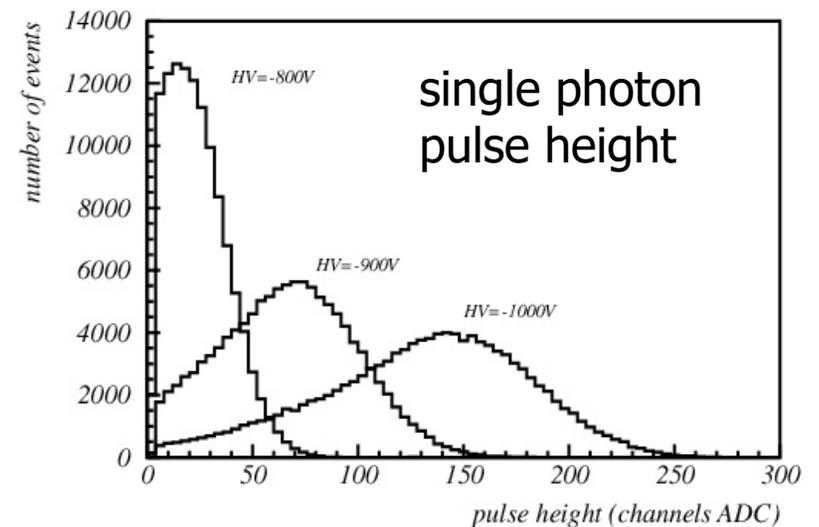
Multianode PMTs (MaPMTs) with metal foil dynodes and 2x2, 4x4 or 8x8 anodes Hamamatsu R5900 (and follow up types 7600, 8500)

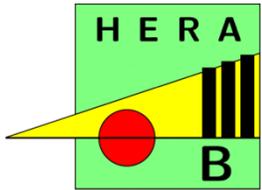
→ Excellent single photon pulse height spectrum

→ Low noise (few Hz/ch)

→ Low cross-talk (<1%)

→ NIM A394 (1997) 27

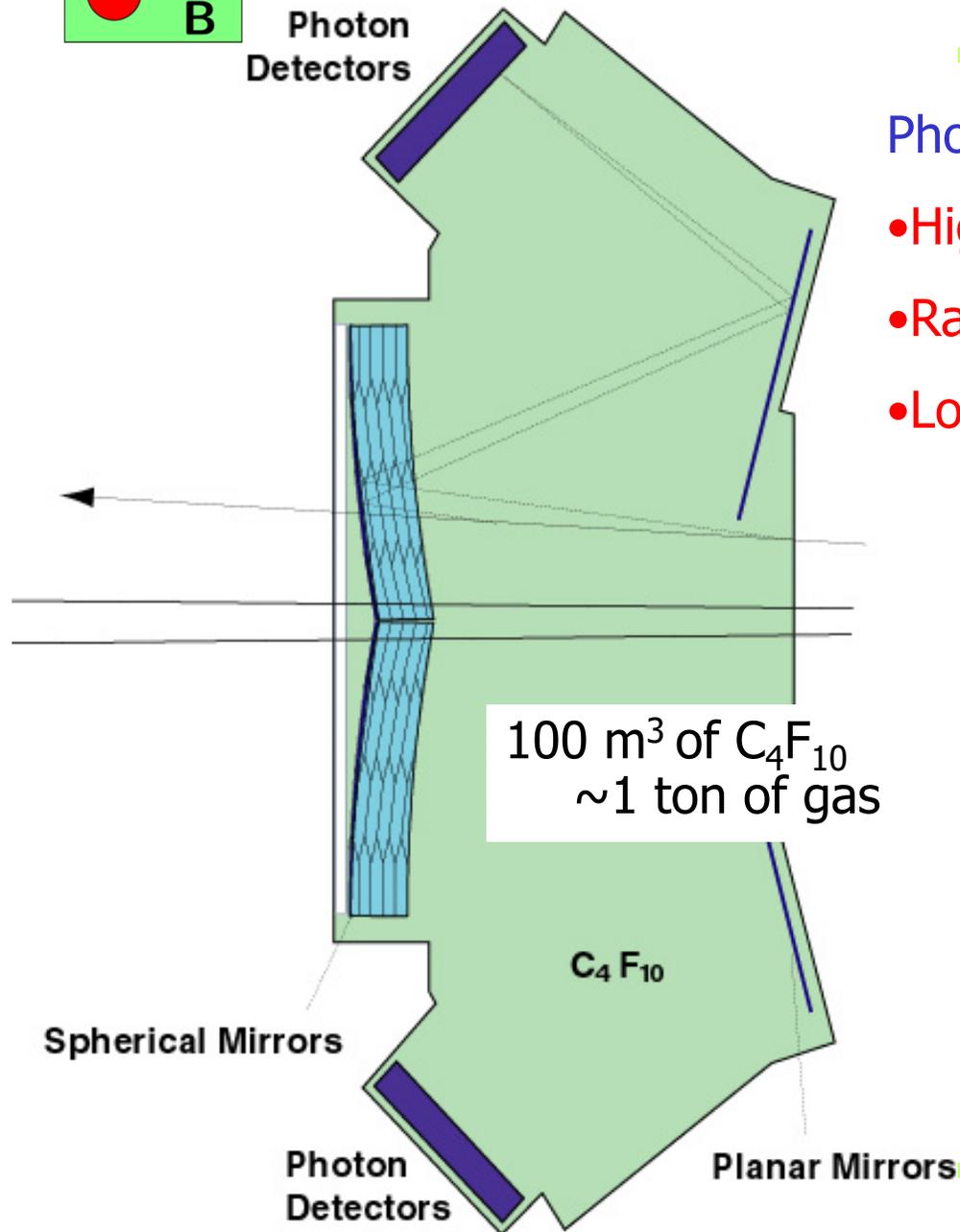


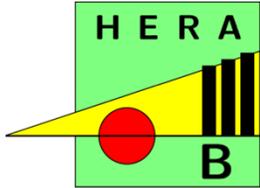


# HERA-B RICH

Photon detector requirements:

- High QE over  $\sim 3\text{m}^2$
- Rates  $\sim 1\text{MHz}$
- Long term stability

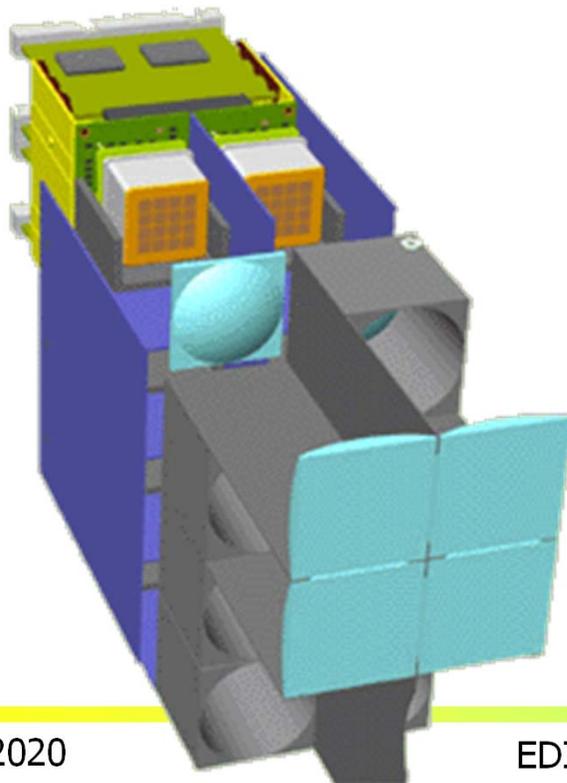




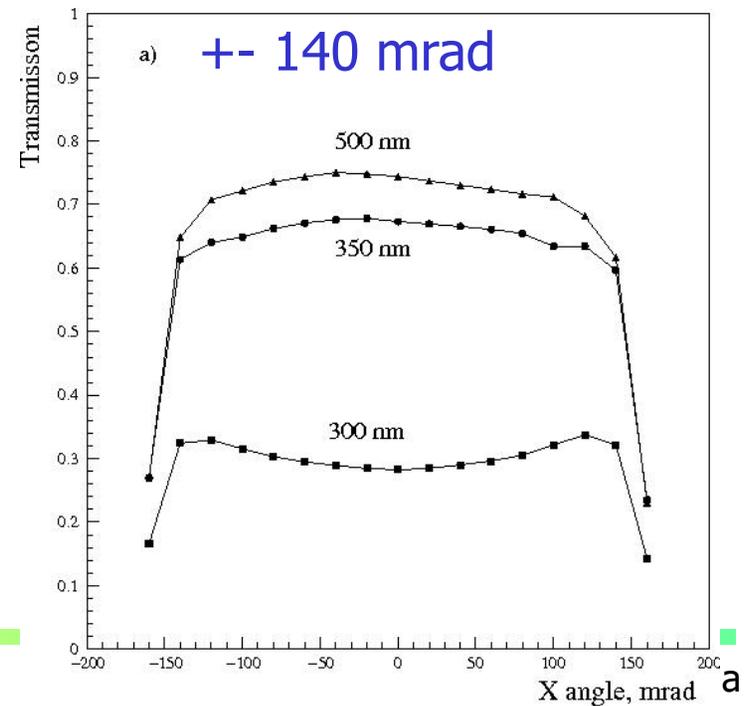
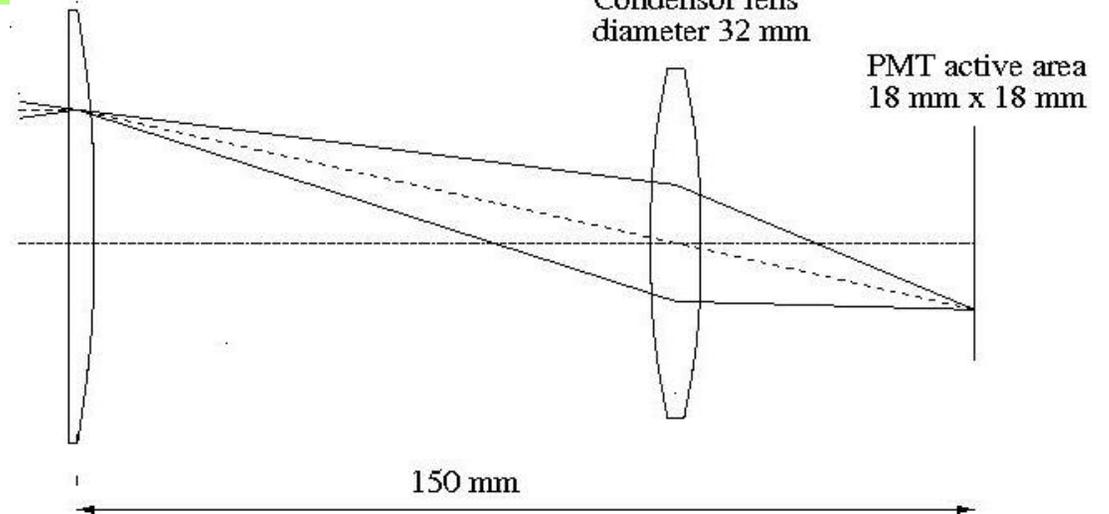
# HERA-B RICH photon detector

Light collection system (imaging!) to:

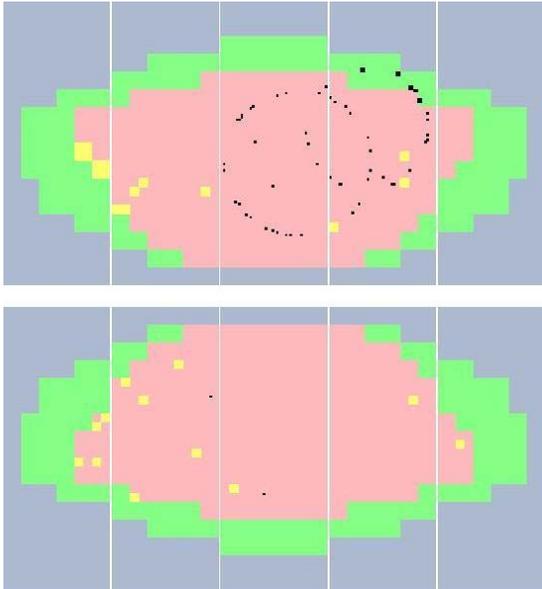
- Eliminate dead areas
- Adapt the pad size



Field lens, 35 mm x 35 mm

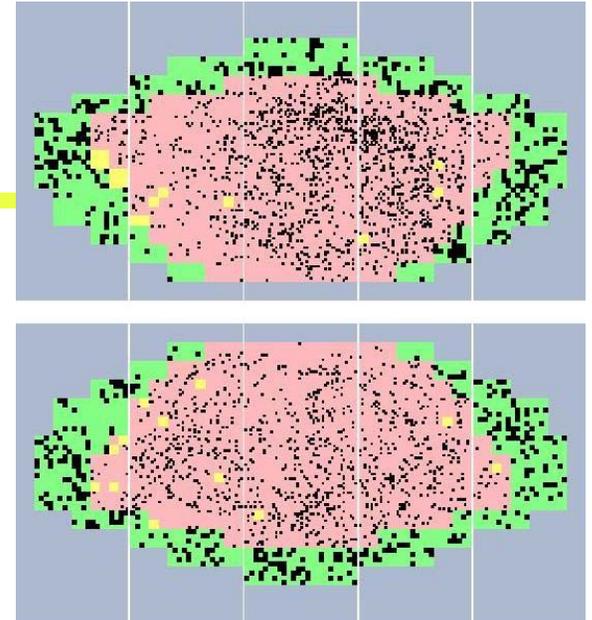


# HERA-B RICH

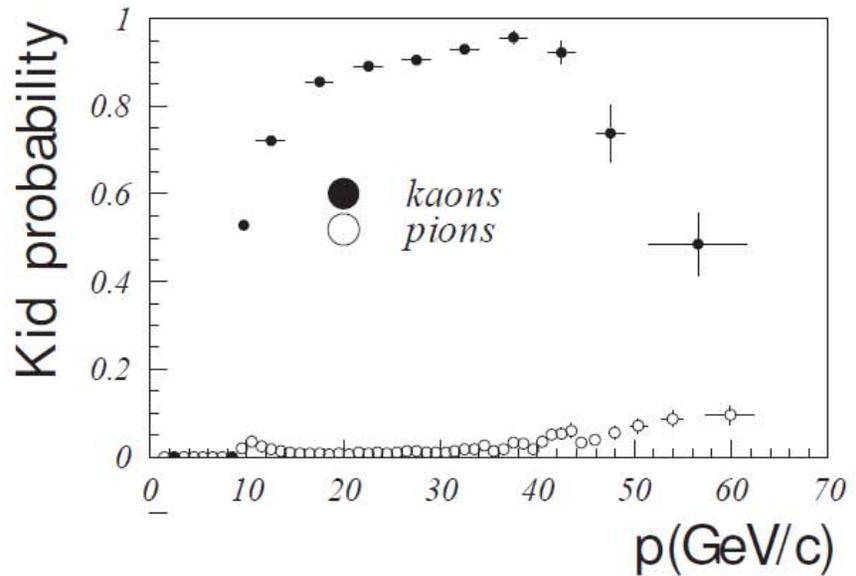
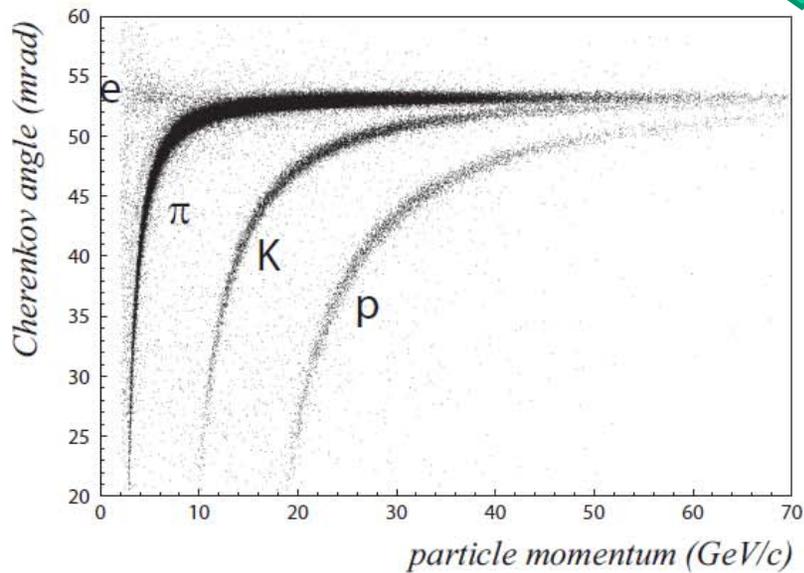


← Little noise, ~30 photons per ring

Typical event →



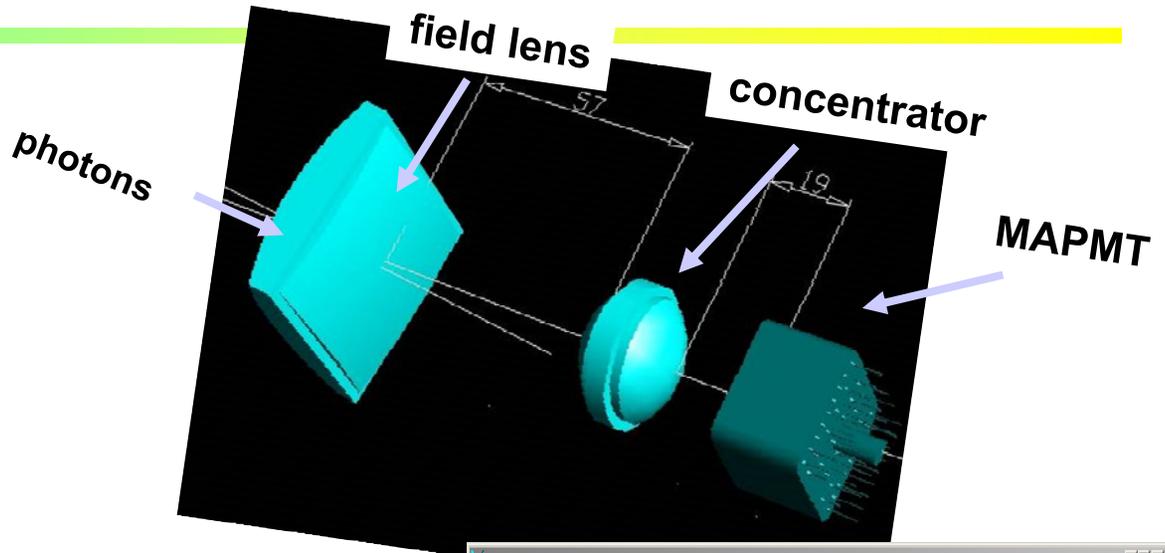
Worked very well!



Kaon efficiency and pion fake probability

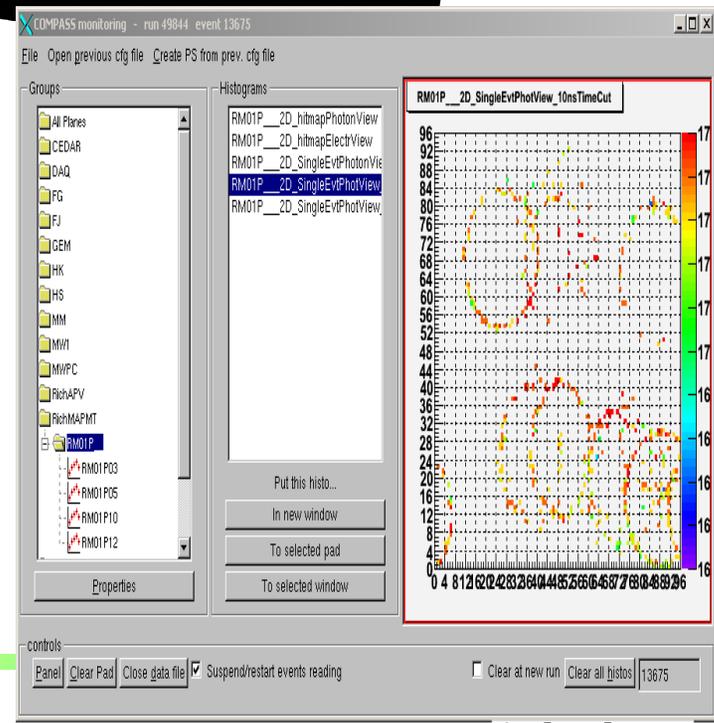
# Photon detector for the COMPASS RICH-1

Upgraded COMPASS RICH-1:  
similar concept as in the  
HERA-B RICH

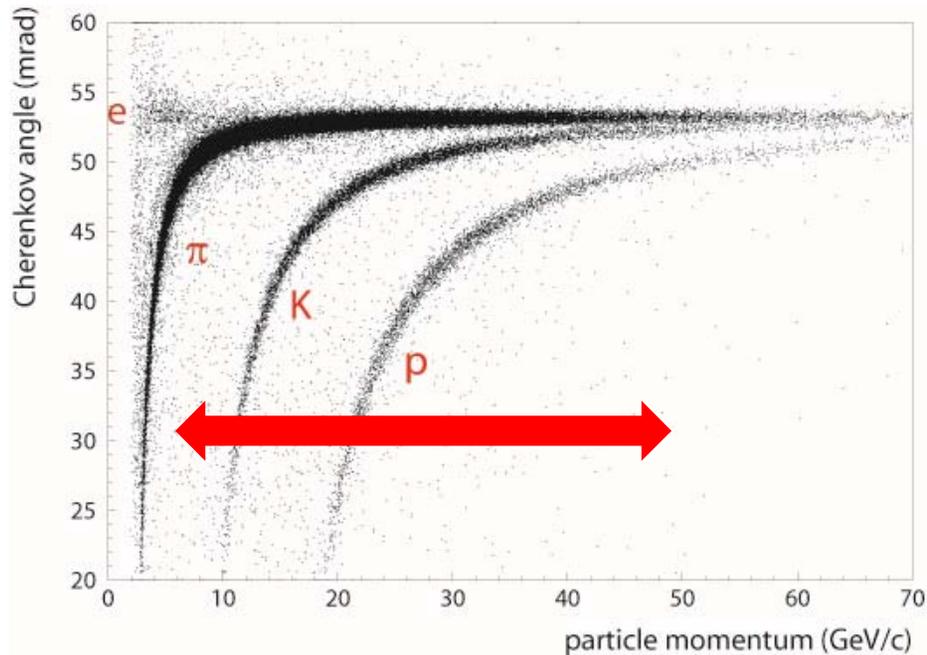


## New features:

- UV extended PMTs & lenses (down to 200 nm) → more photons
- surface ratio = (telescope entrance surface) / (photocathode surface) = 7
- fast electronics with <120 ps time resolution



# Kinematic range of a RICH counter



Example: kinematic range for kaon/pion separation

Kinematic range for separation of two particle types:

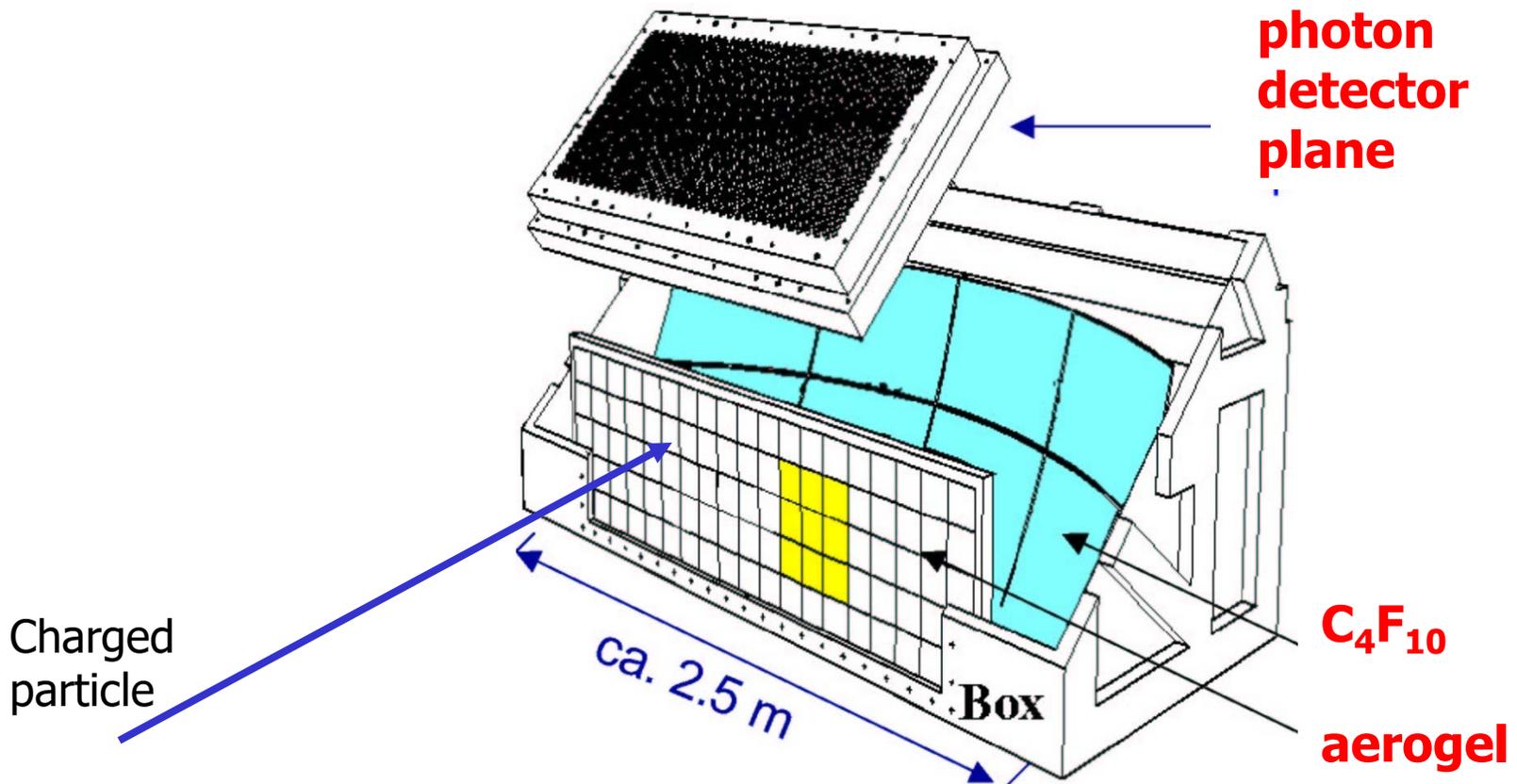
- Lower limit  $p_{\min}$ : sufficiently above lighter particle threshold
- Upper limit  $p_{\max}$ : given by Cherenkov angle resolution – overlap of the two bands

Rule of thumb:  $p_{\max} / p_{\min} < 10$

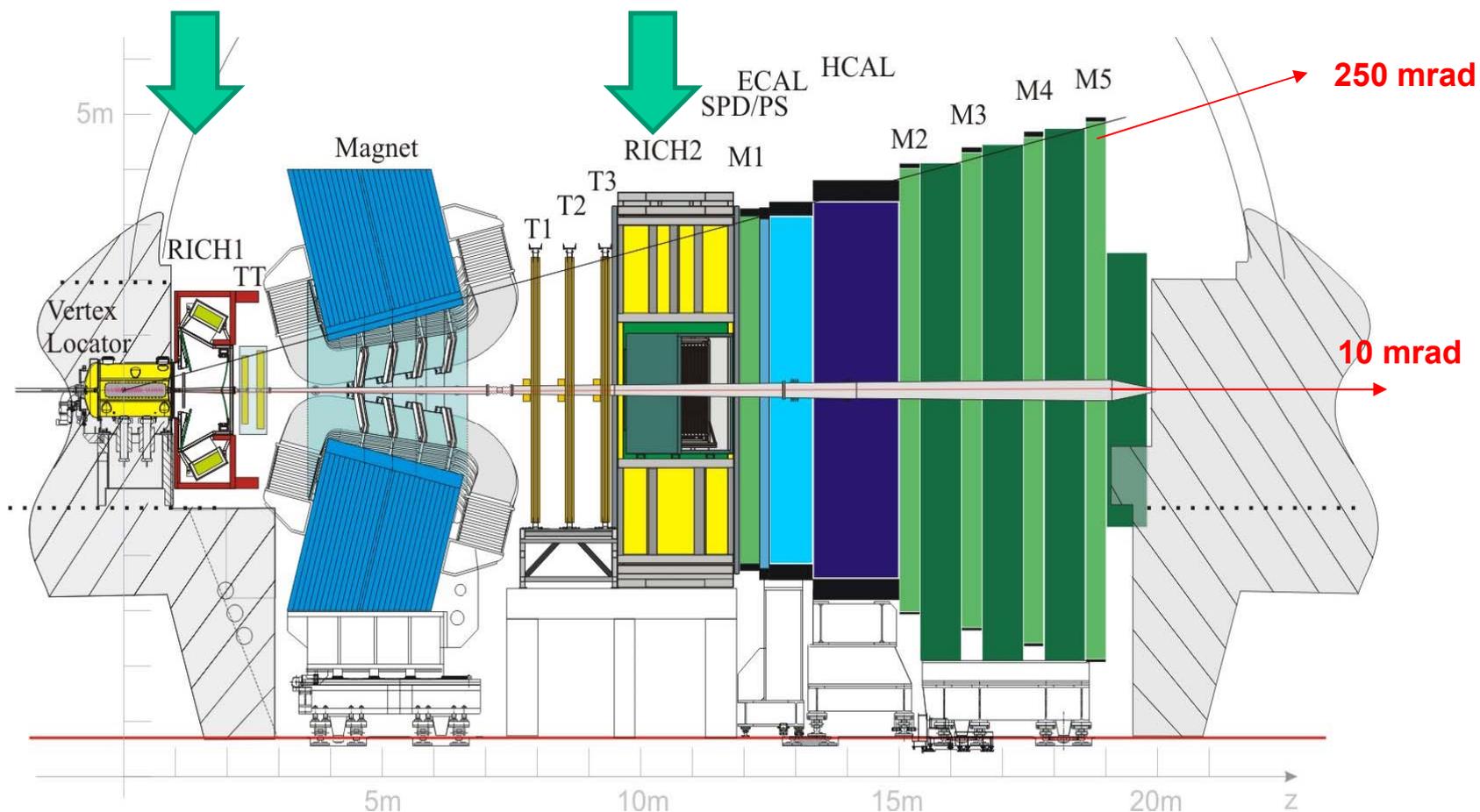
# RICHes with several radiators

Extending the kinematic range → need more than one radiator

- DELPHI at LEP, SLD at SLC (liquid + gas)
- HERMES at HERA (aerogel+gas)



# The LHCb RICH counters



**Vertex reconstruction:**  
**VELO**

**Trigger:**  
**Muon Chambers**  
**Calorimeters**  
**Tracker**

**PID:**  
**RICHes**  
**Calorimeters**  
**Muon Chambers**

**Kinematics:**  
**Magnet**  
**Tracker**  
**Calorimeters**

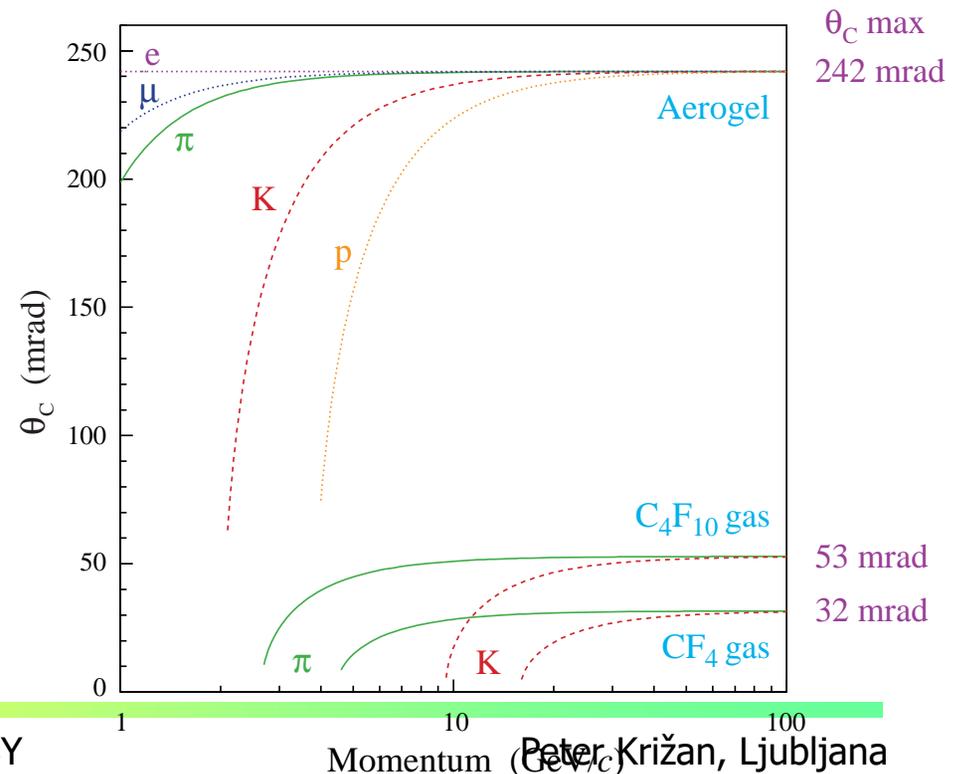
# LHCb RICHes

Need:

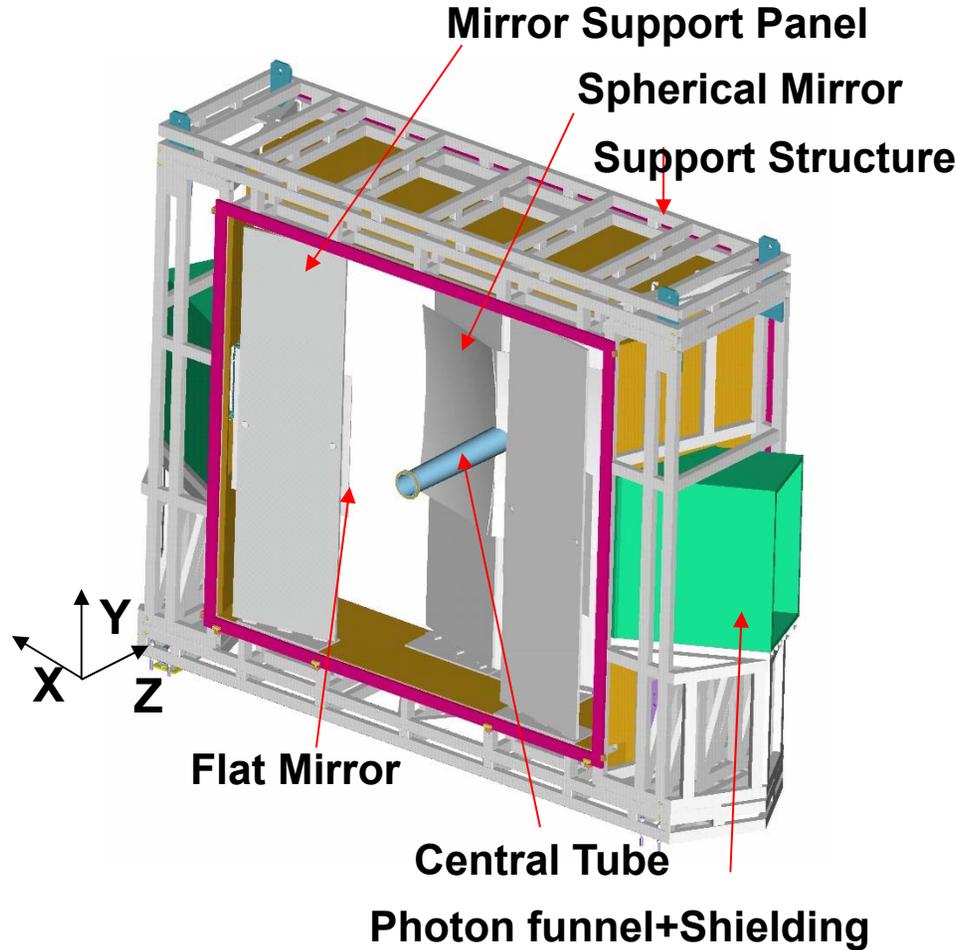
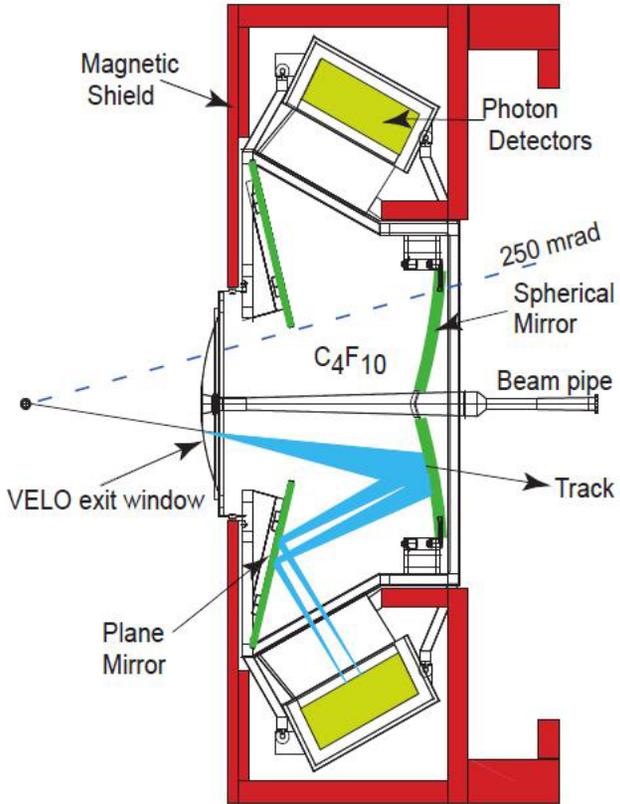
- Particle identification for momentum range  $\sim 2-100 \text{ GeV}/c$
- Granularity  $2.5 \times 2.5 \text{ mm}^2$
- Large area ( $2.8 \text{ m}^2$ ) with high active area fraction
- Fast compared to the 25ns bunch crossing time
- Have to operate in a small B field

→ 3 radiators

- Aerogel
- $\text{C}_4\text{F}_{10}$  gas
- $\text{CF}_4$  gas



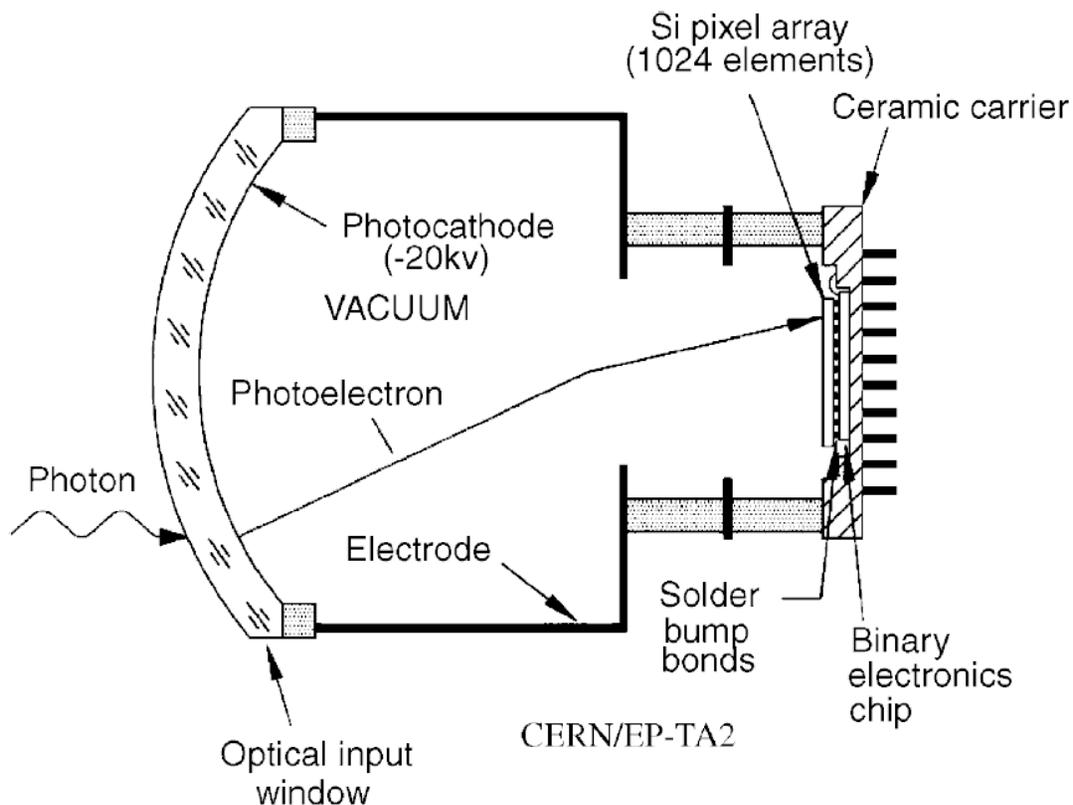
# LHCb RICHes



# LHCb RICHes

Photon detector: hybrid PMT (R+D with DEP) with 5x demagnification (electrostatic focusing).

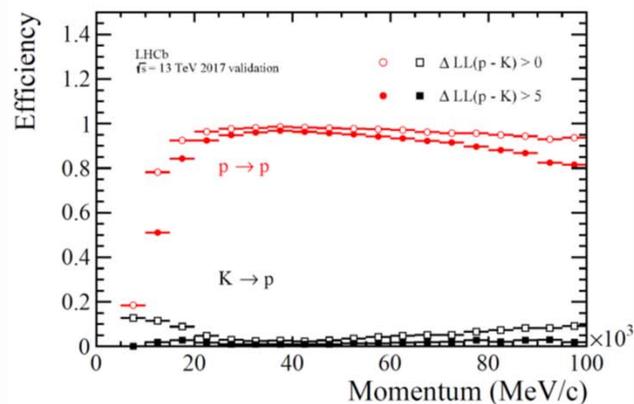
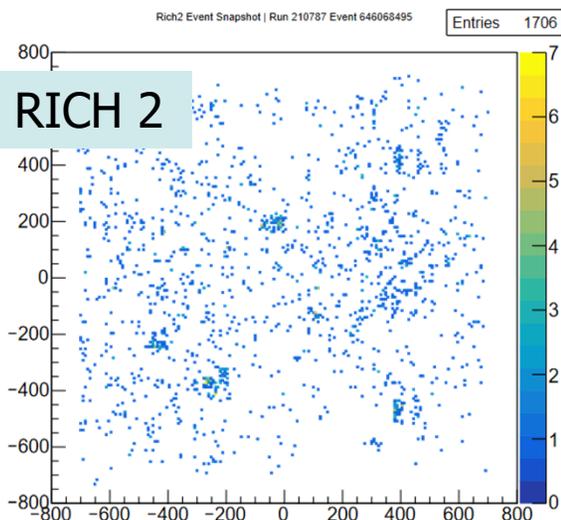
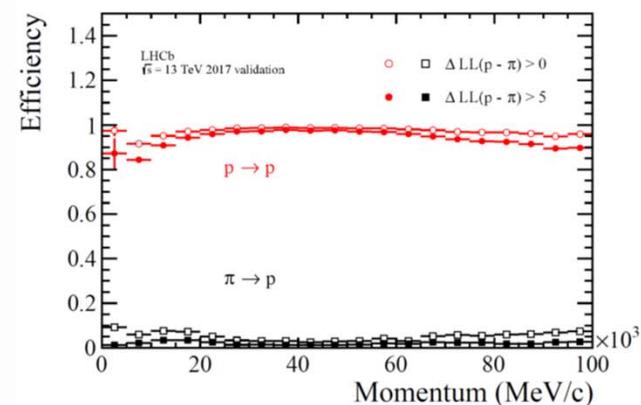
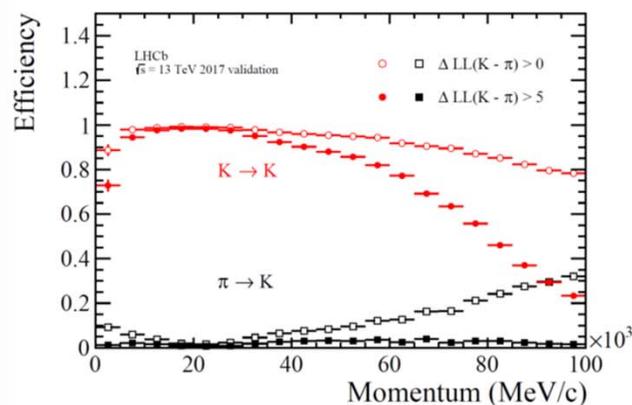
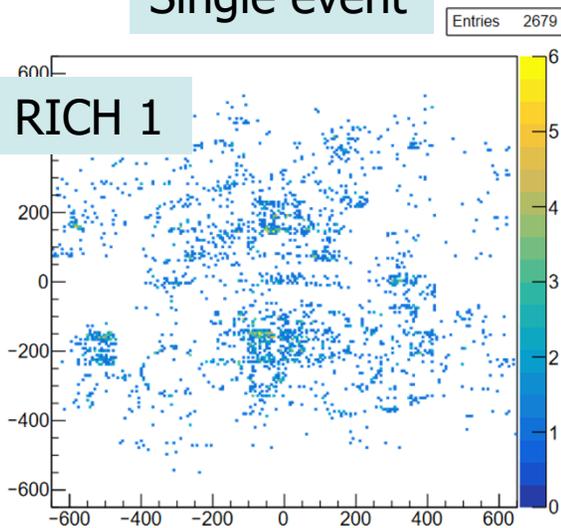
Hybrid PMT: accelerate photoelectrons in electric field ( $\sim 20\text{kV}$ ), detect them in a pixelated silicon detector.



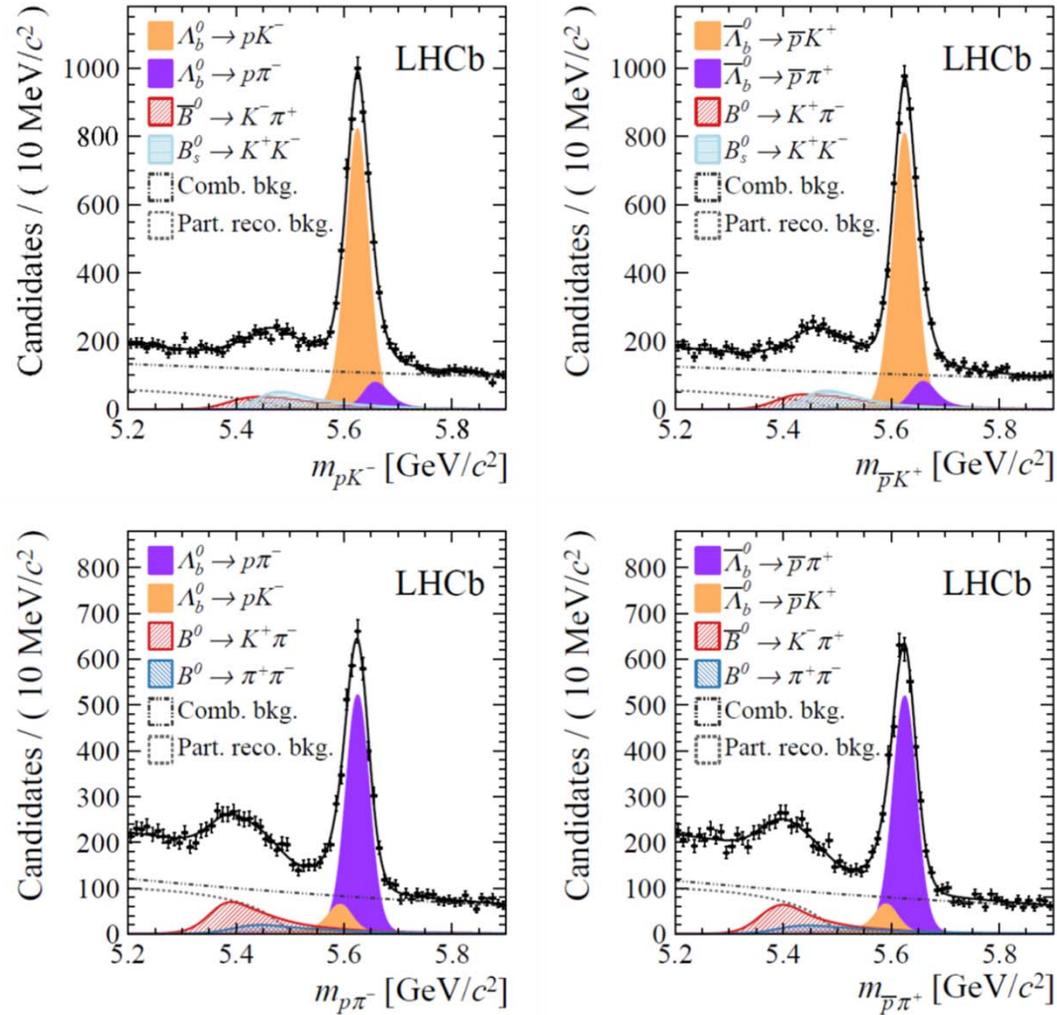
NIM A553 (2005) 333

# Performance of LHCb RICHes

Single event

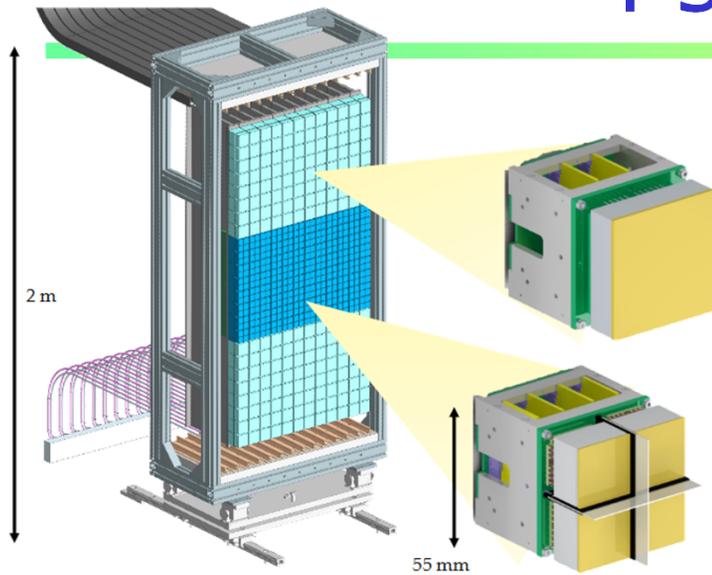


# LHCb RICHes: performance

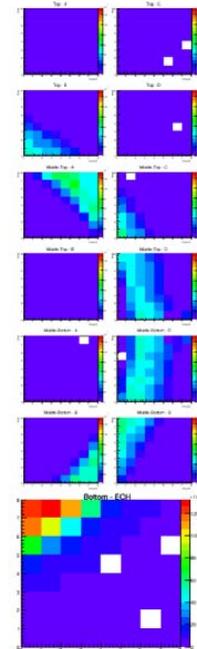
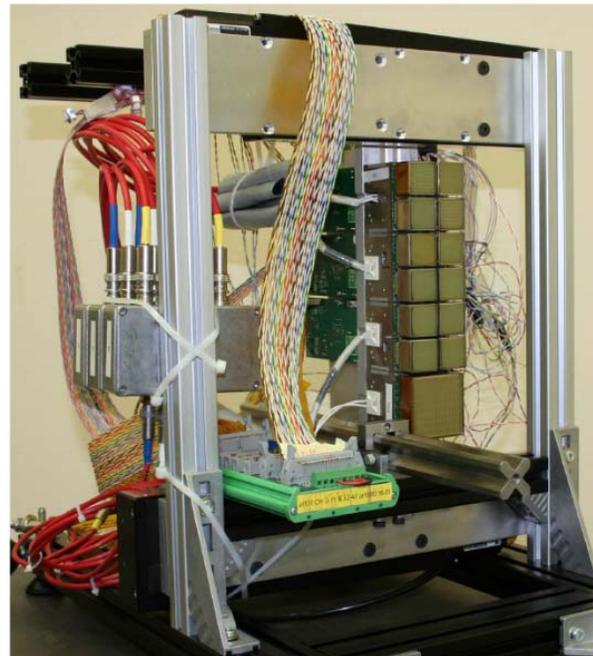
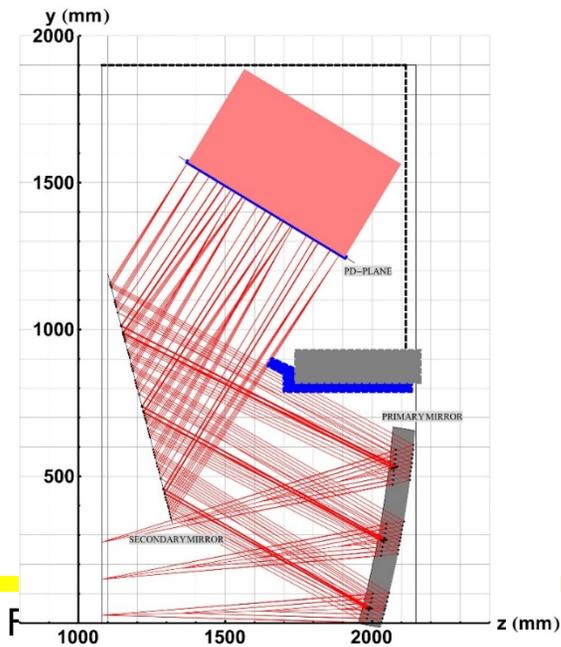


“Search for CP violation in  $\Lambda_b^0 \rightarrow pK^-$  and  $\Lambda_b^0 \rightarrow p\pi^-$  decays”  
 [LHCb-PAPER-2018-025]

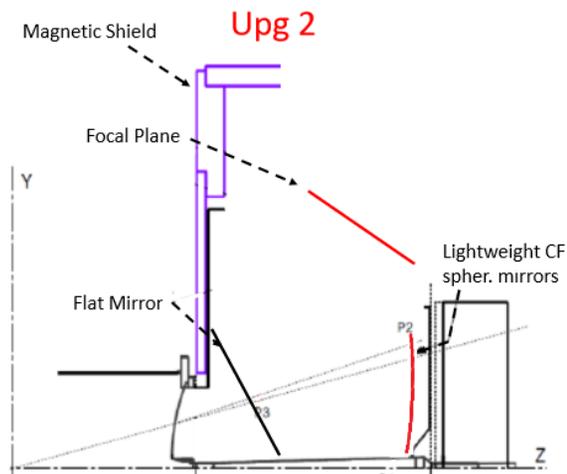
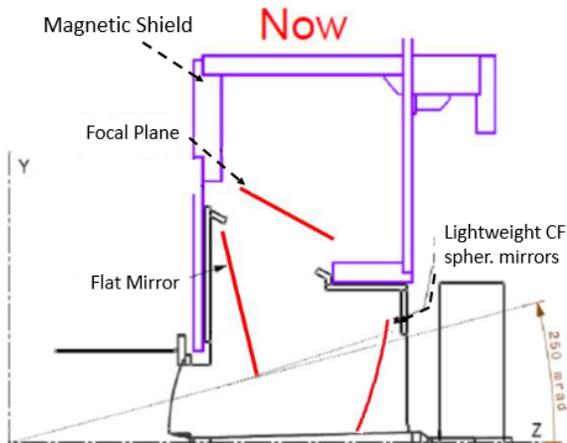
# LHCb Upgrade (under way)



- ❑ New photon detectors: MaPMTs Hamamatsu R13743 (H12700) and R13742 (R11265)
- ❑ New electronics working at 40 MHz readout rate
- ❑ New optics layout for RICH 1



# Future LHCb Upgrade

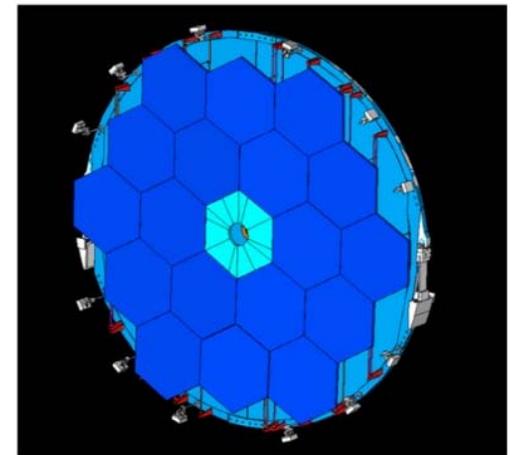
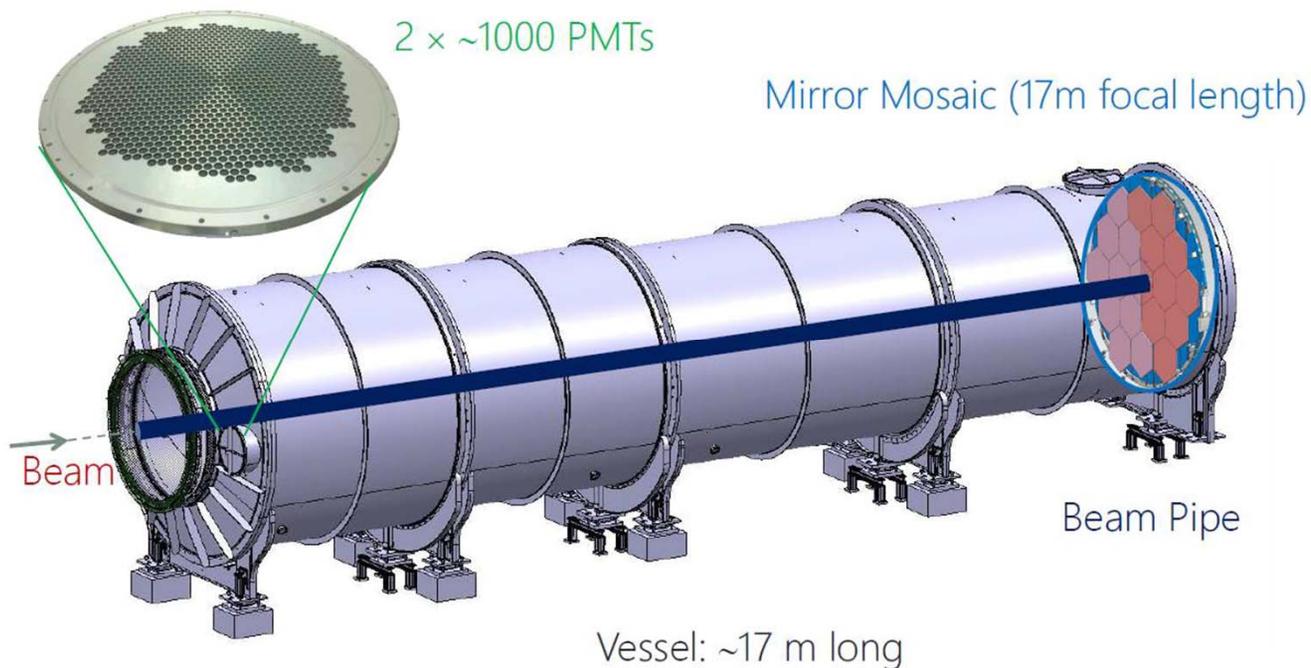


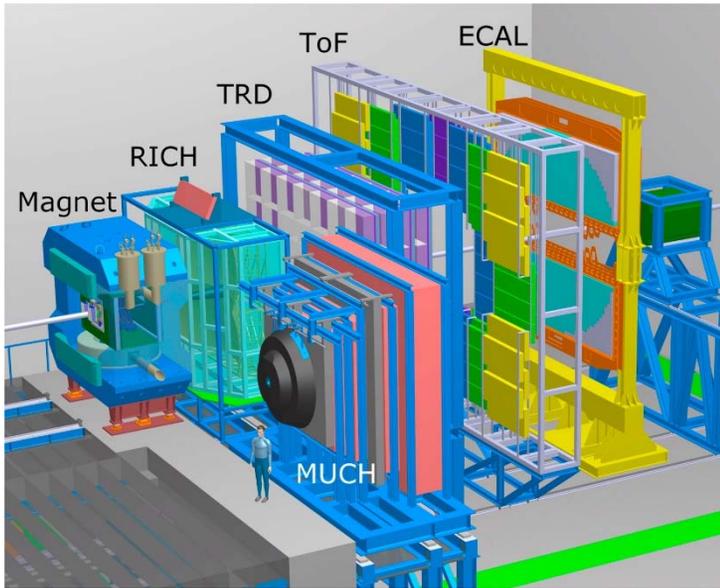
- ❑ Provide PID at p-p luminosity of  $10^{34}$  in the forward region
- ❑ Incremental improvements in:
  - Improve Cherenkov angle resolution
    - More photons in the green → lower chromatic error
  - Reduced event complexity with timing
  - Enhanced number of photons

Radiator	C <sub>4</sub> F <sub>10</sub>			CF <sub>4</sub>	
	RICH 1 Current (HPD)	RICH 1 UPG1	RICH 1 UPG2	RICH 2 UPG1	RICH 2 UPG2
Average Photoelectron Yield	30	40	<b>60–30</b>	22	<b>30</b>
Single Photon Errors (mrad)					
Chromatic	0.84	0.58	<b>0.24–0.12</b>	0.31	<b>0.1</b>
Pixel	0.9	0.44	<b>0.15</b>	0.20	<b>0.07</b>
Emission Point	0.8	0.37	<b>0.1</b>	0.27	<b>0.05</b>
Overall	1.47	0.82	<b>0.3–0.2</b>	0.46	<b>0.13</b>

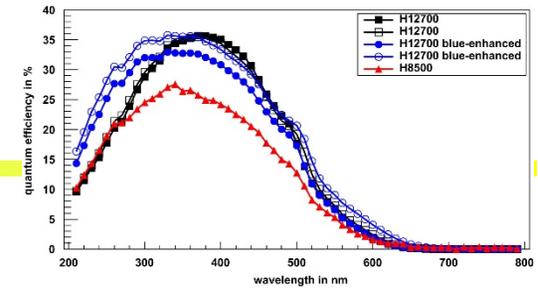
# NA62 RICH

- ❑ Momentum range 15-35 GeV/c
- ❑ 17m long, 200m<sup>3</sup> cylindrical vacuum proof tank with Neon radiator
- ❑ Photon detectors: 2000 PMTs (16mm, 8mm active, with Winstone cone light guides)
- ❑ Mirror alignment  $\sim 30 \mu\text{rad}$
- ❑ Single photon resolution:  $\sim 140 \mu\text{rad}$
- ❑ Operational since 2014

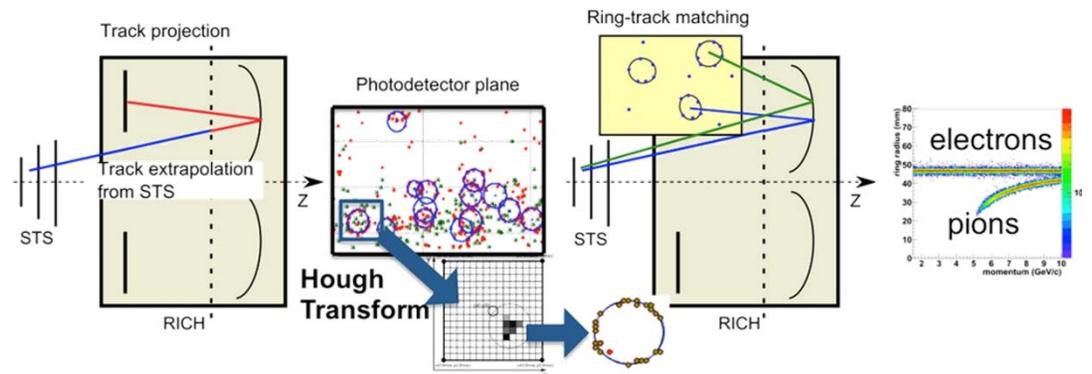
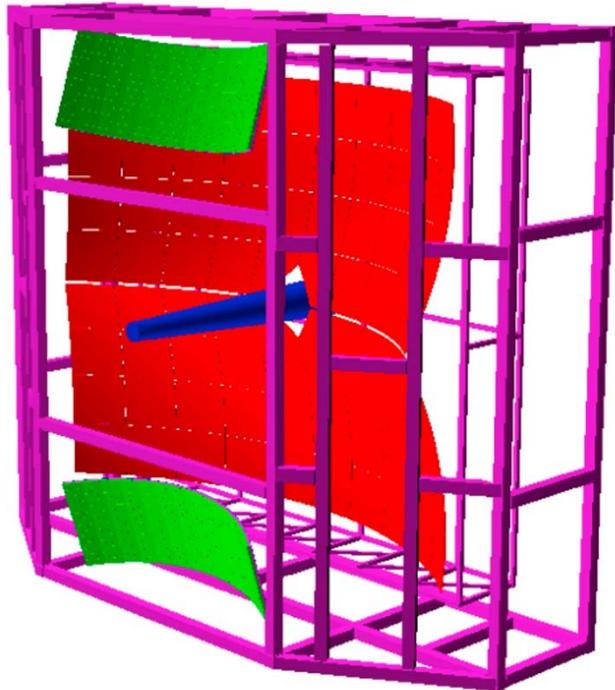




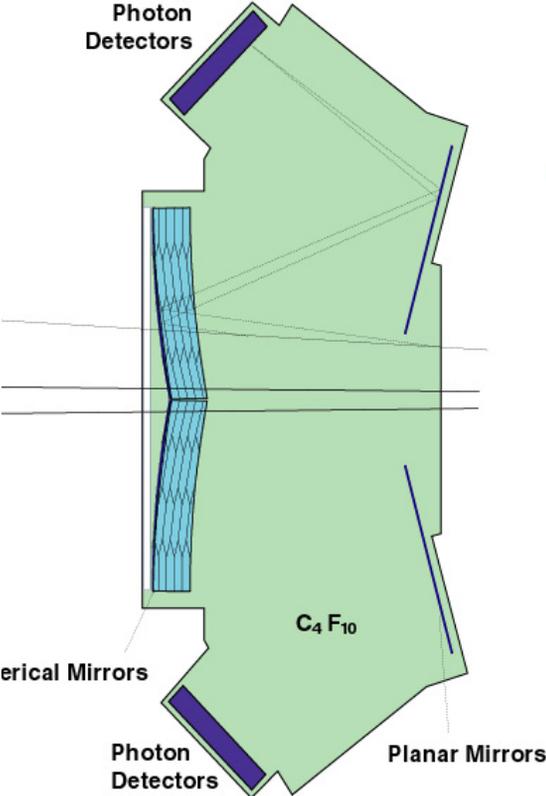
# CBM



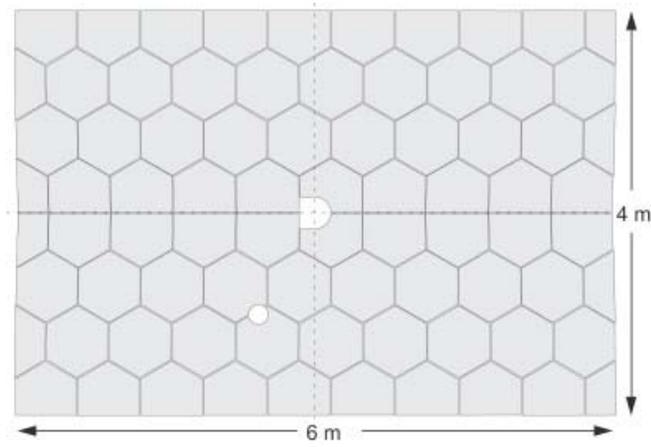
- ❑ RICH with CO<sub>2</sub> radiator
- ❑ MaPMTs: Hamamatsu H12700
- ❑ Cylindrical photon detection surface
- ❑ Extensive testing of MaPMTs for radiation damage
- ❑ Up to 1000 tracks per event
- ❑ Momentum up to 8 GeV/c
- ❑ Pion suppression factor ~5000 (with TRD)



# Mirror alignment

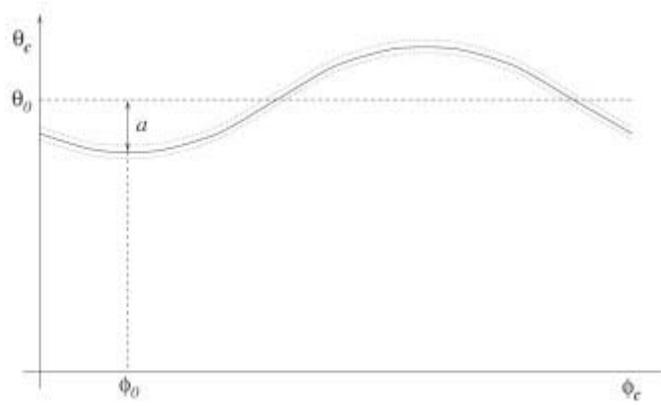
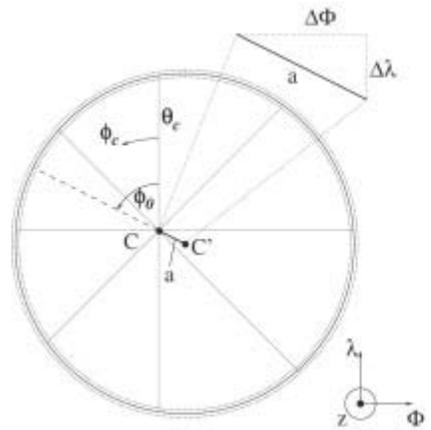


Gas radiator RICHes: large mirrors  $\rightarrow$  tens of segments  $\rightarrow$  need relative alignment



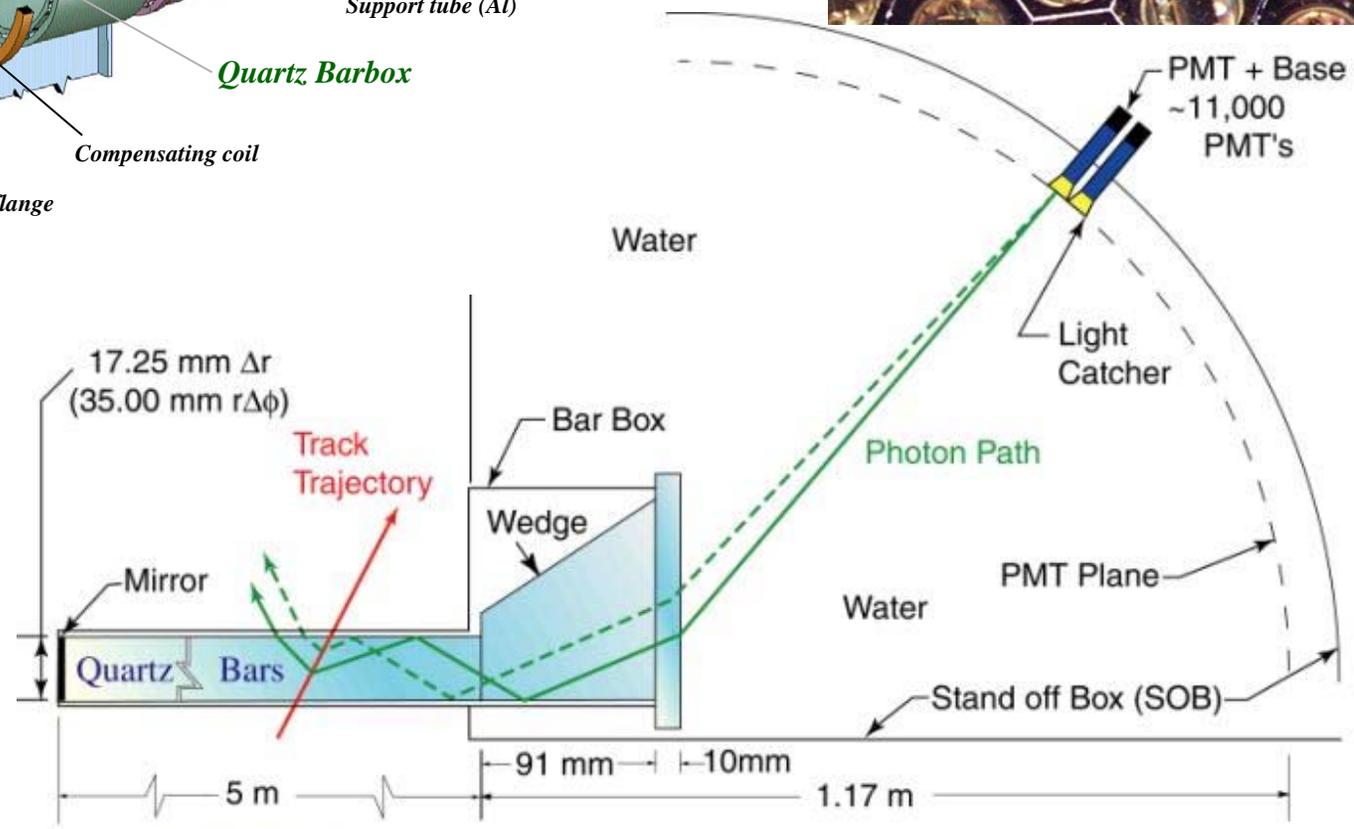
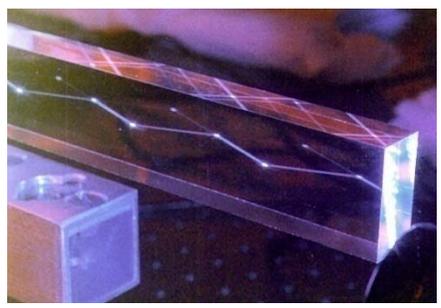
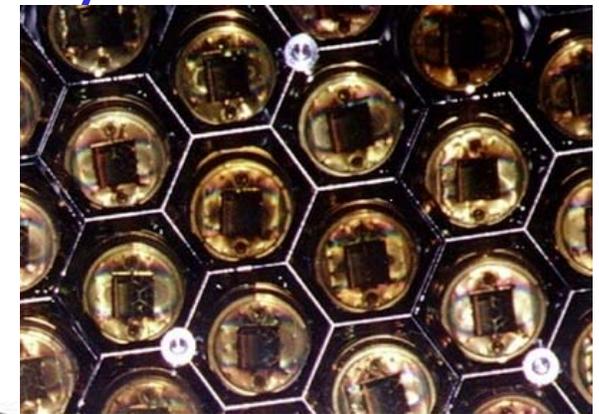
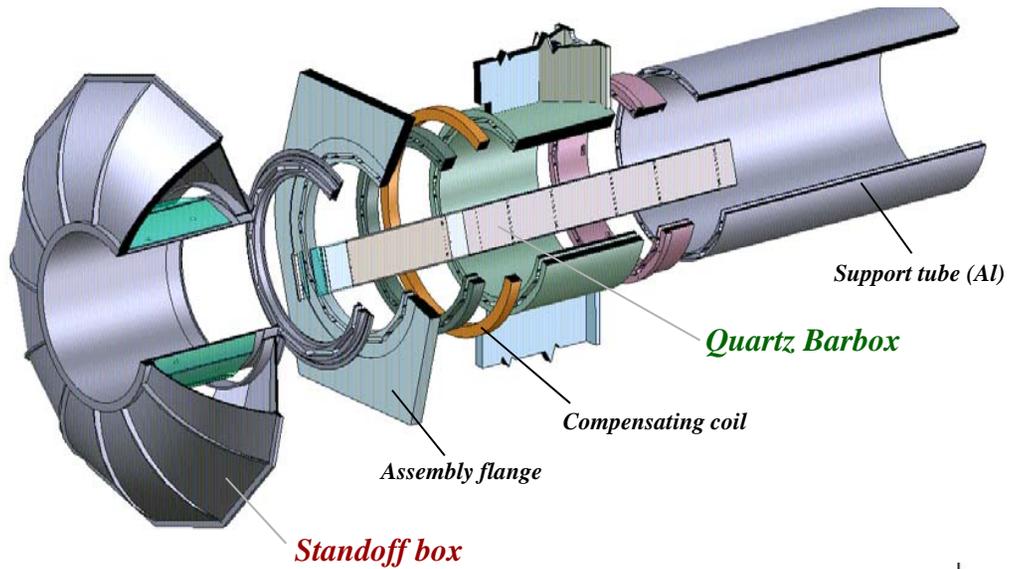
- Spherical mirror: 80 hexagonal segments
- Planar mirrors: 2x18 rectangular segments

Aligning pairs of spherical and planar segments.



Misalignment: Cherenkov angle depends on the azimuthal angle around the track

# DIRC (@BaBar) - detector of internally reflected Cherenkov light

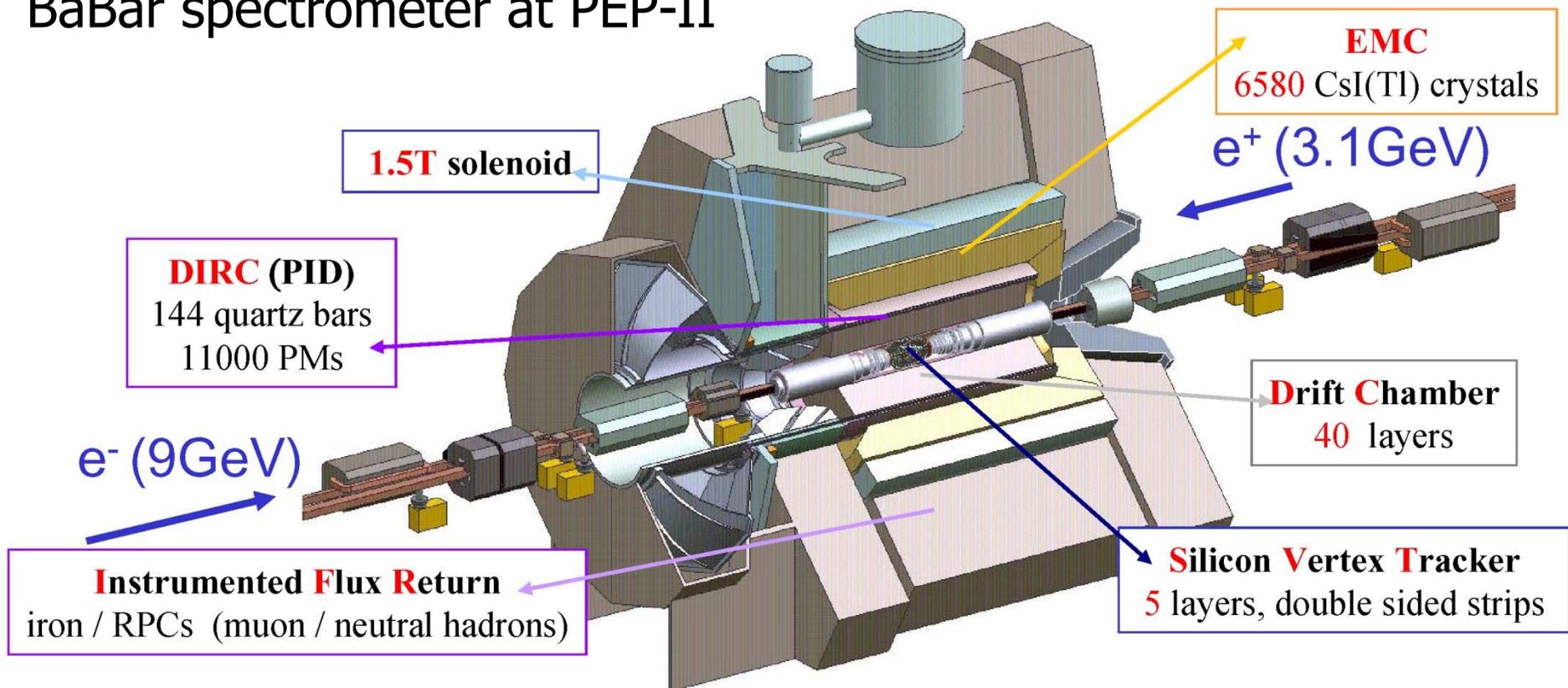


4 x 1.225 m Bars  
 glued end-to-end

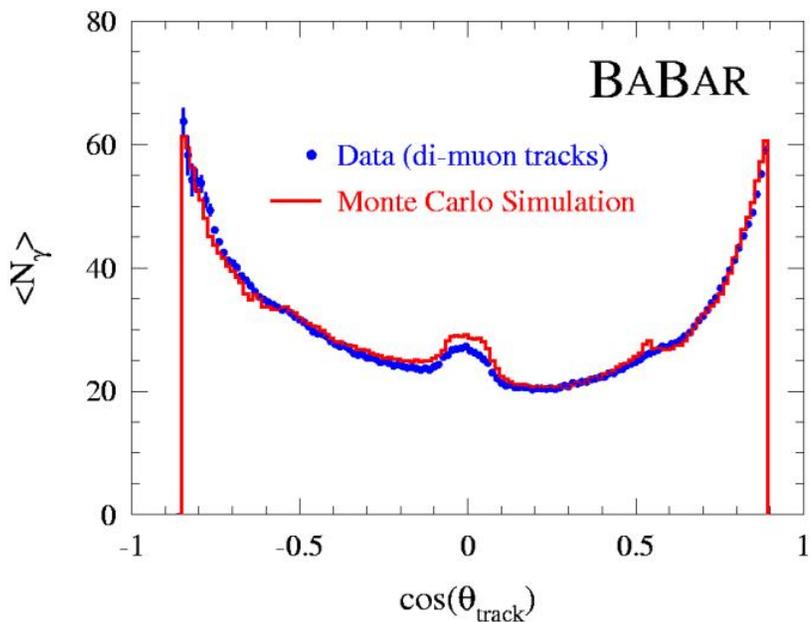


# DIRC - detector of internally reflected Cherenkov light

BaBar spectrometer at PEP-II

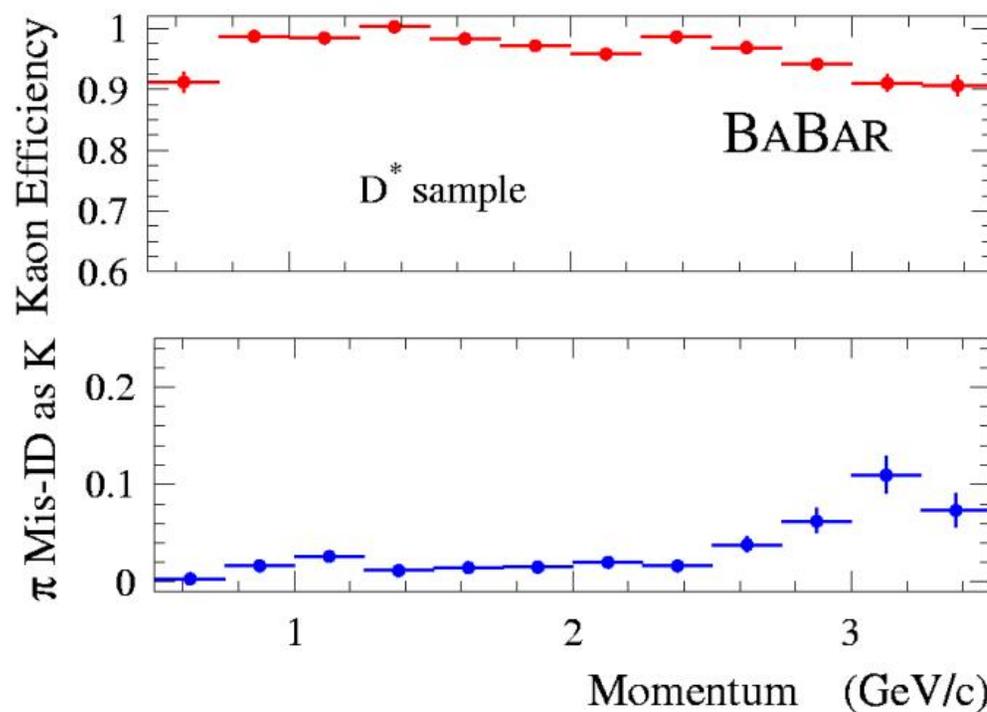


# DIRC performance

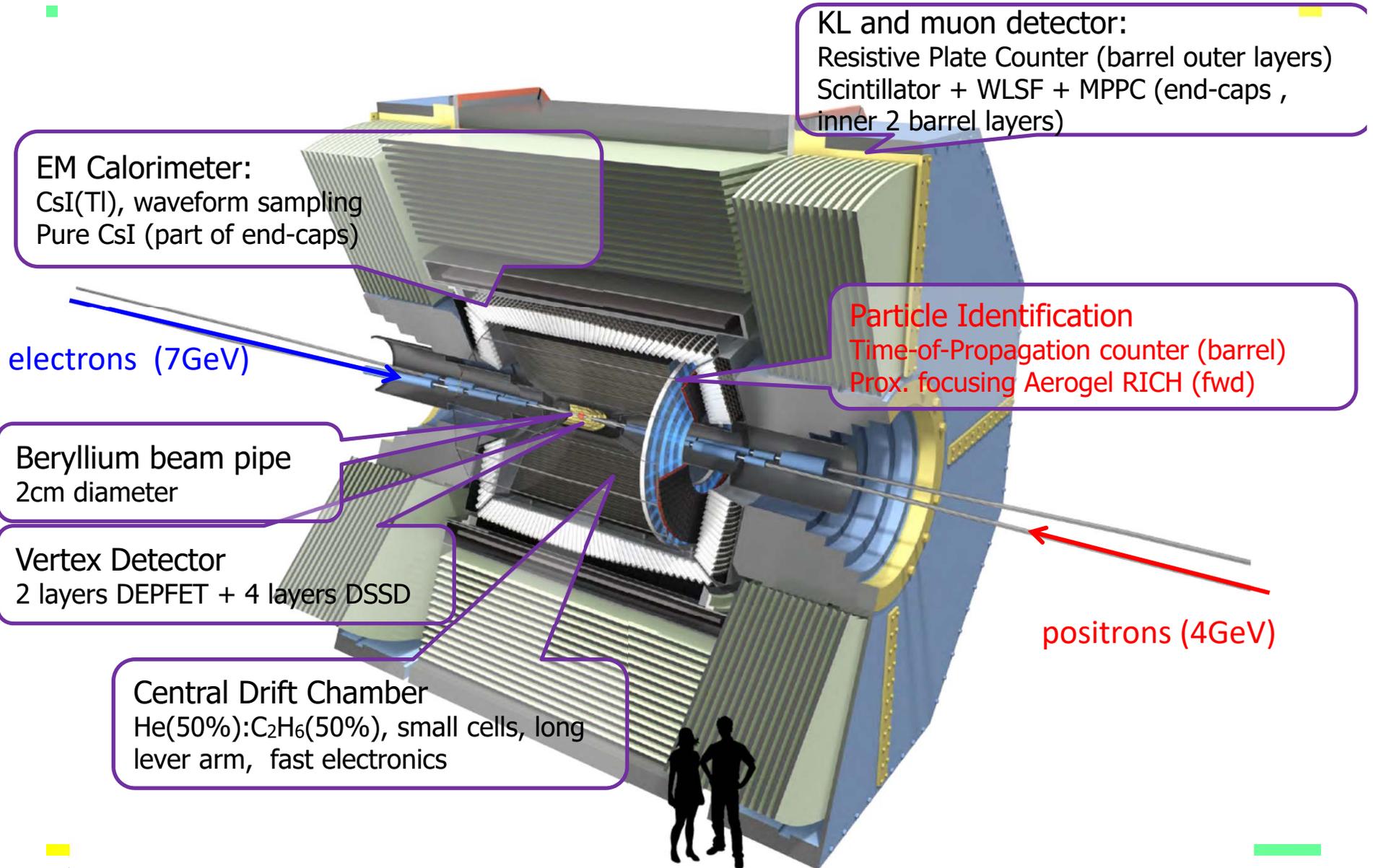


← Lots of photons!

Excellent  $\pi/K$  separation



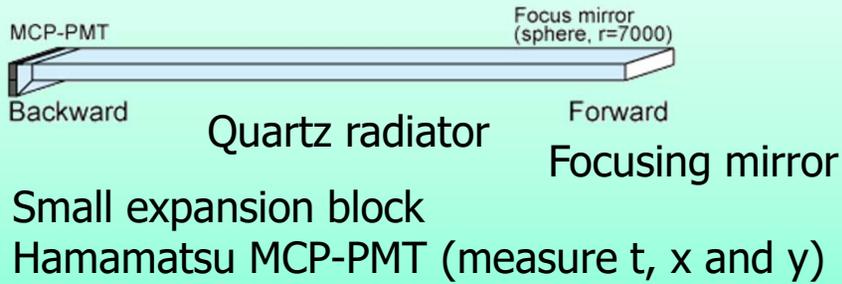
# Hadron PID in the Belle II Detector



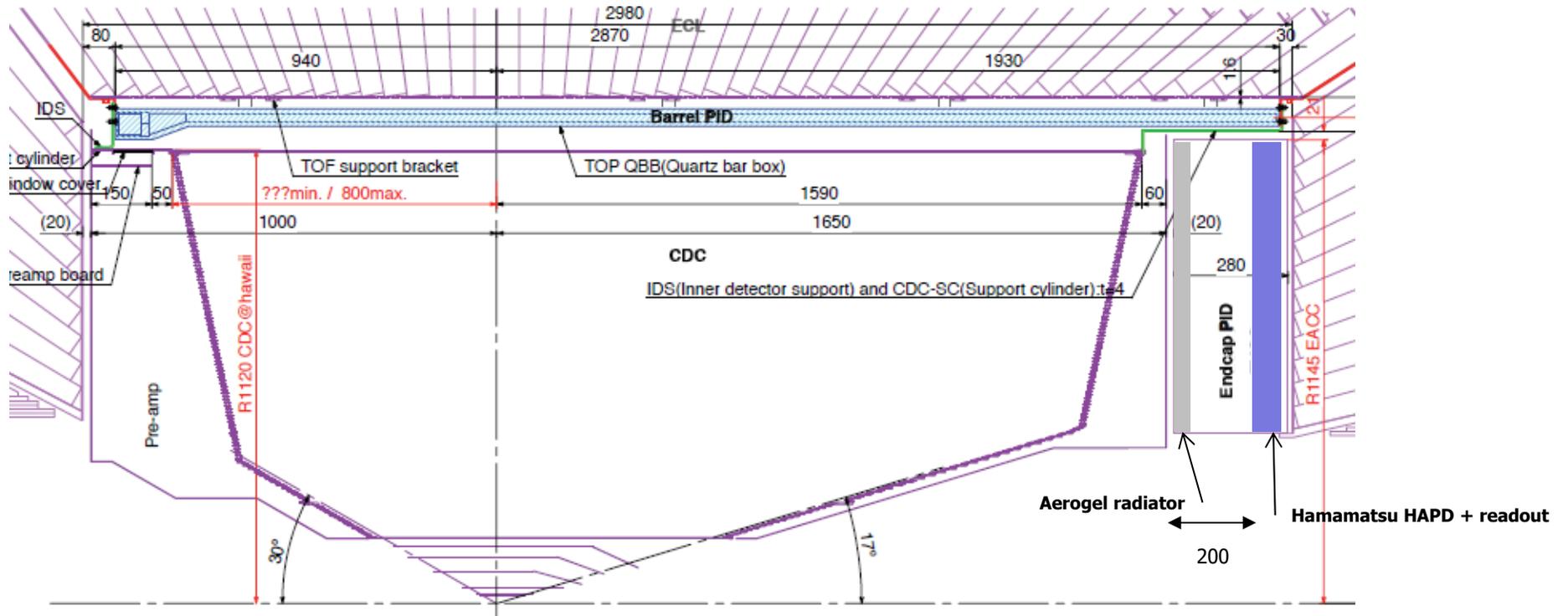
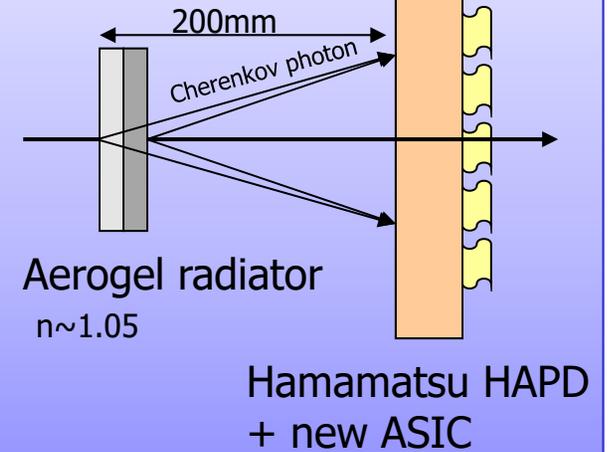


# Belle II Cherenkov detectors

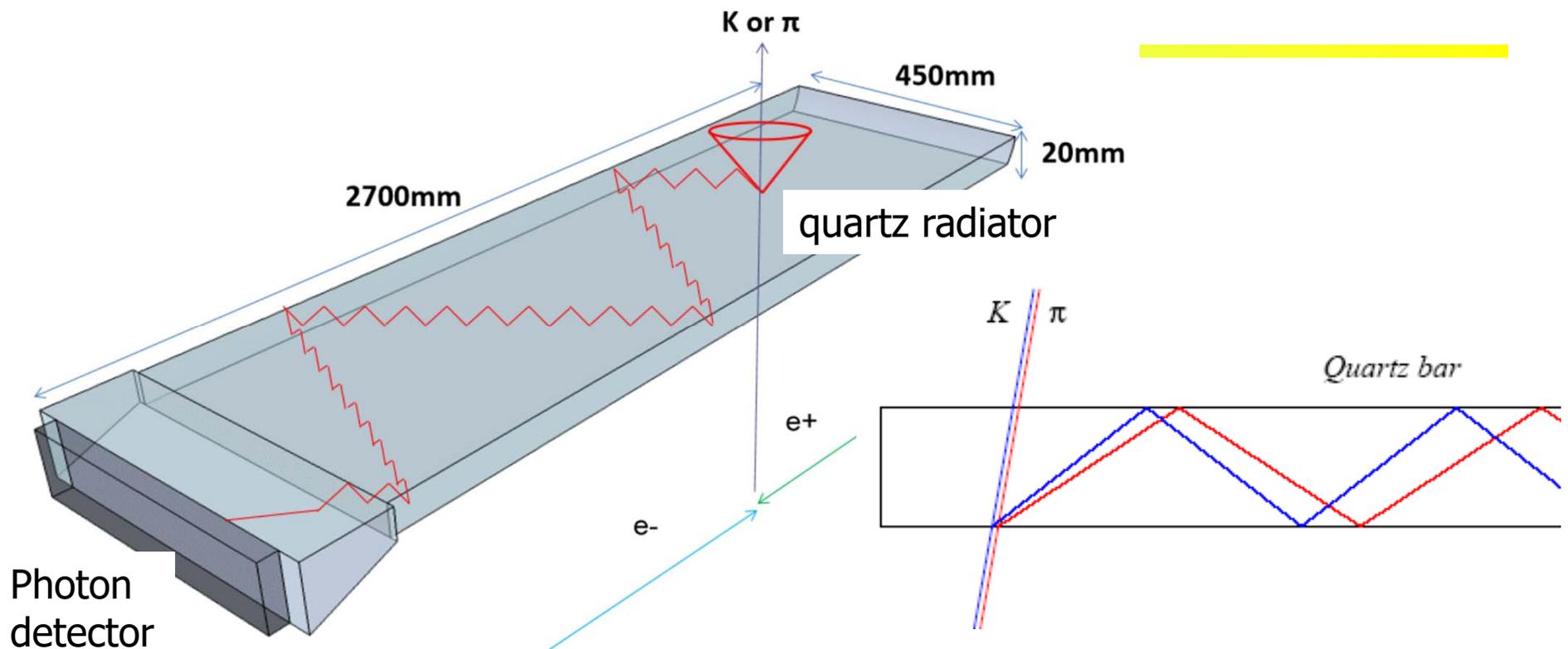
## Barrel PID: Time of Propagation Counter (TOP)



## Endcap PID: Aerogel RICH (ARICH)

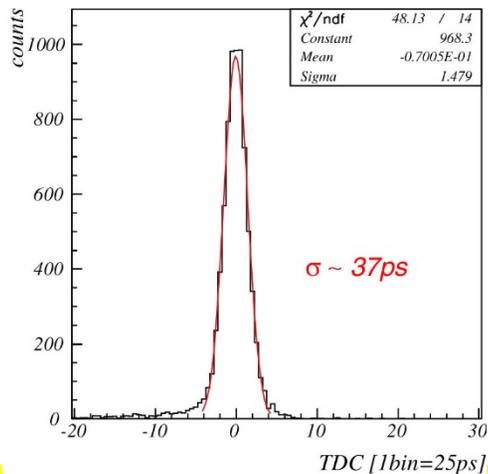
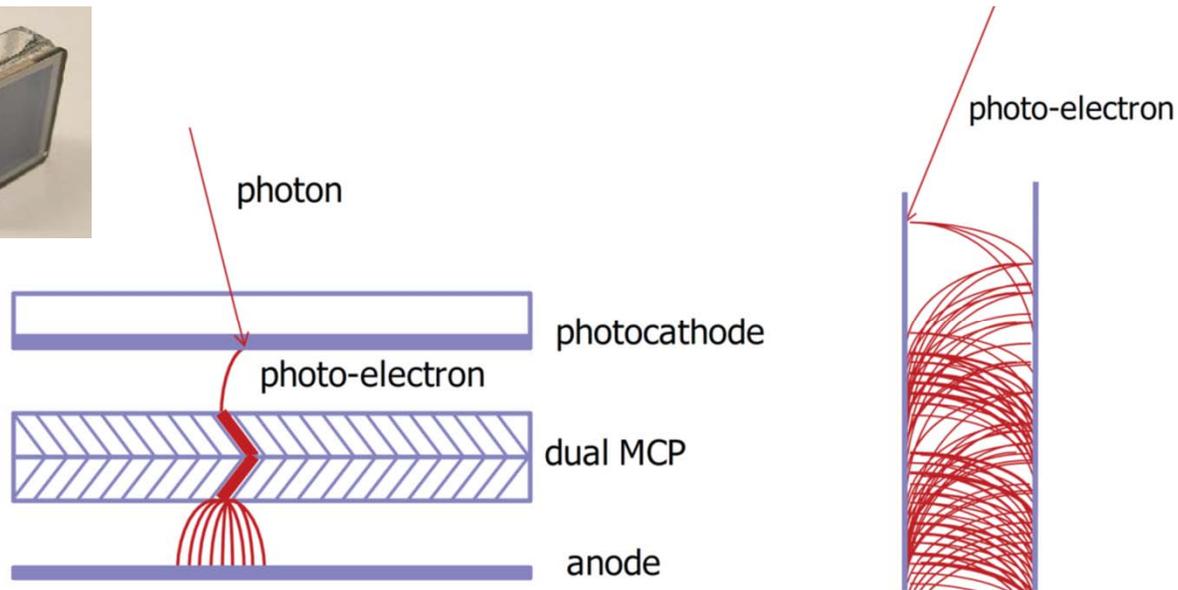


# Belle II Barrel PID: Time of propagation (TOP) counter



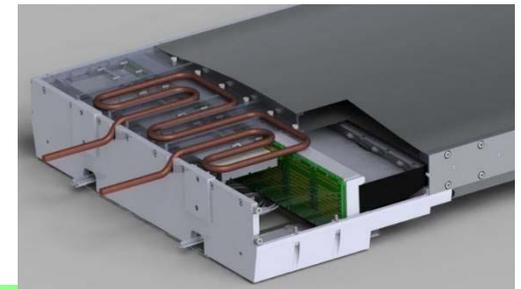
- Similar to the DIRC, Cherenkov ring imaging with precise time measurement.
- Reconstruct Cherenkov angle from two hit coordinates and the time of propagation of the photon
  - Quartz radiator (2cm thick)
  - Photon detector (MCP-PMT)
    - Excellent time resolution  $\sim 40$  ps
    - Single photon sensitivity at 1.5 T

# MCP PMTs for a very fast timing

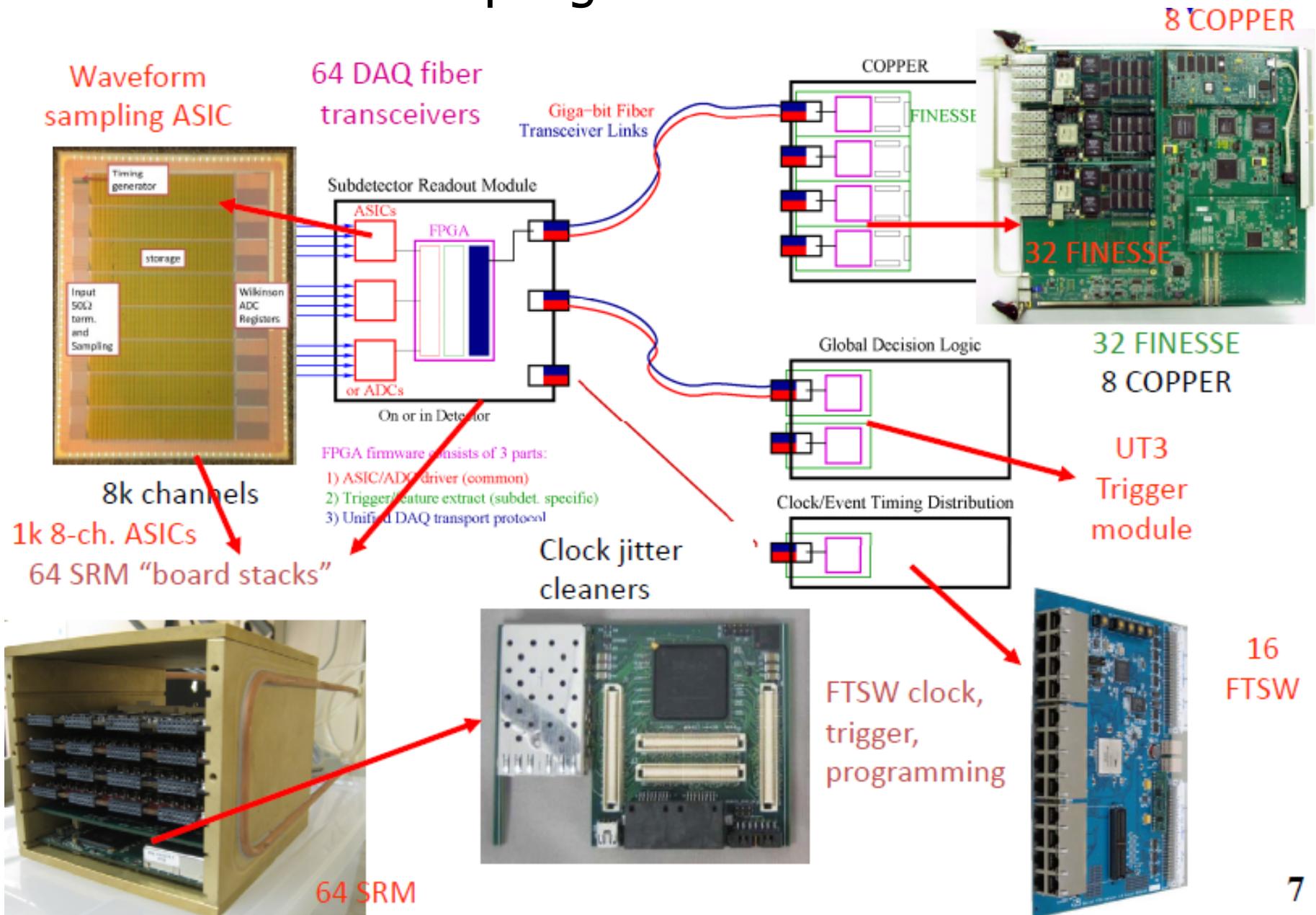


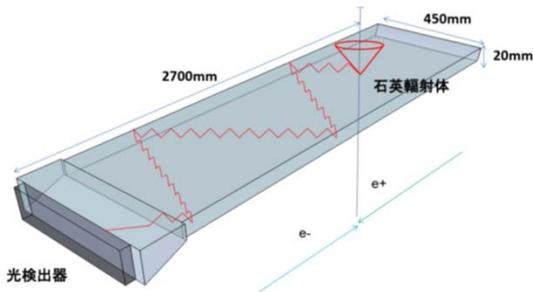
Micro-channel plate PMTs:

Single photon resolution: typically 20ps – 40ps



# TOP Waveform sampling readout

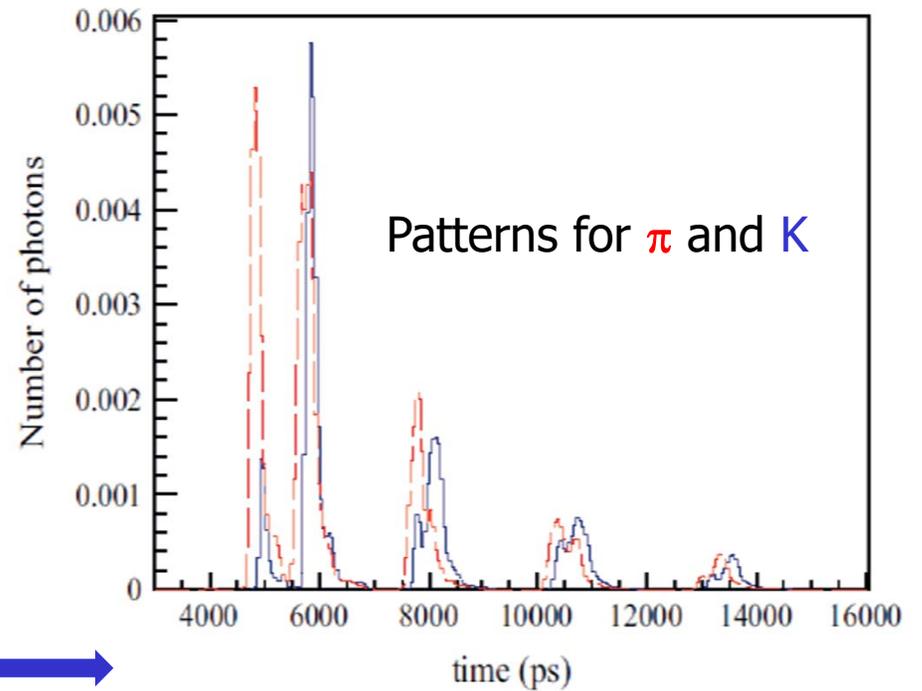
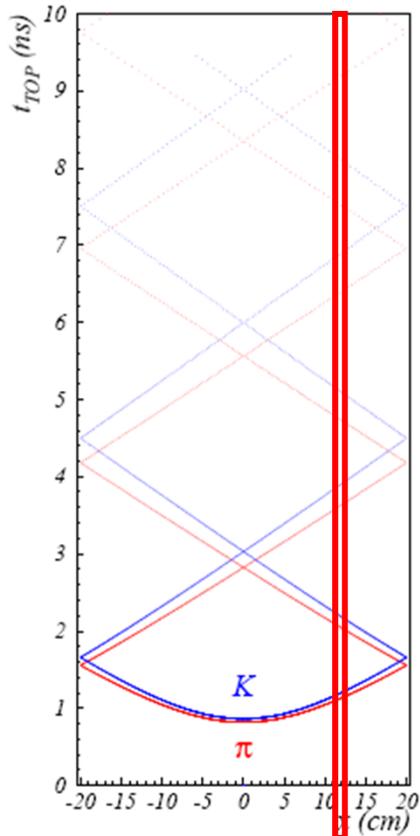




# TOP image reconstruction

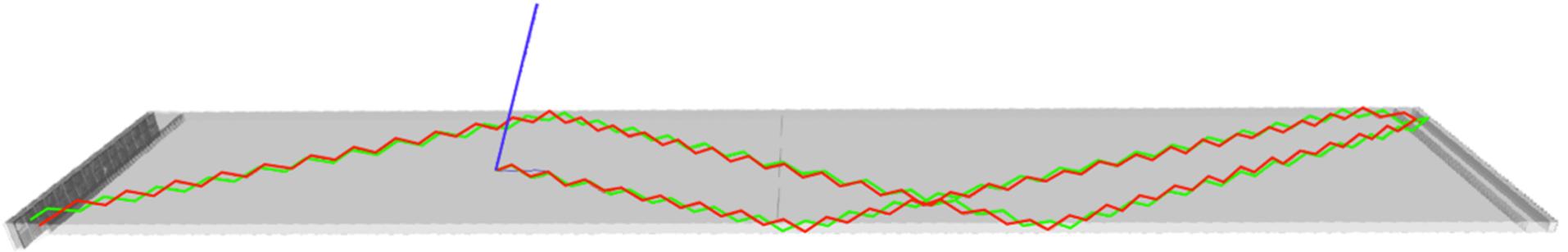
Pattern in the coordinate-time space ('ring') of a pion and kaon hitting a quartz bar

Time distribution of signals recorded by one of the PMT channels (slice in x): different for  $\pi$  and K (~shifted in time)

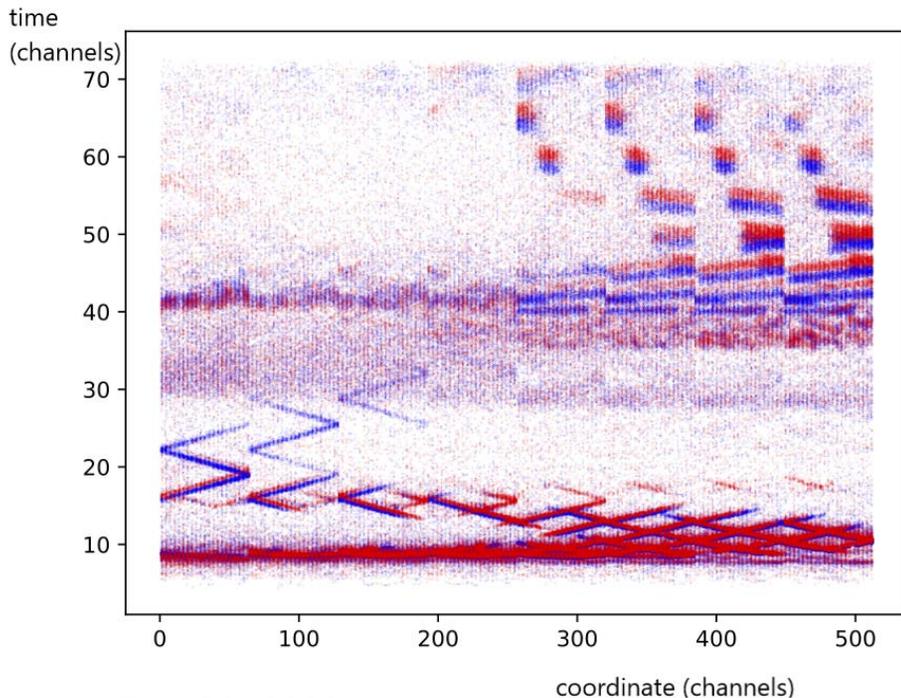


The name of the game: analytic expressions for the 2D likelihood functions  
 → M. Starič et al., NIMA A595 (2008) 252-255

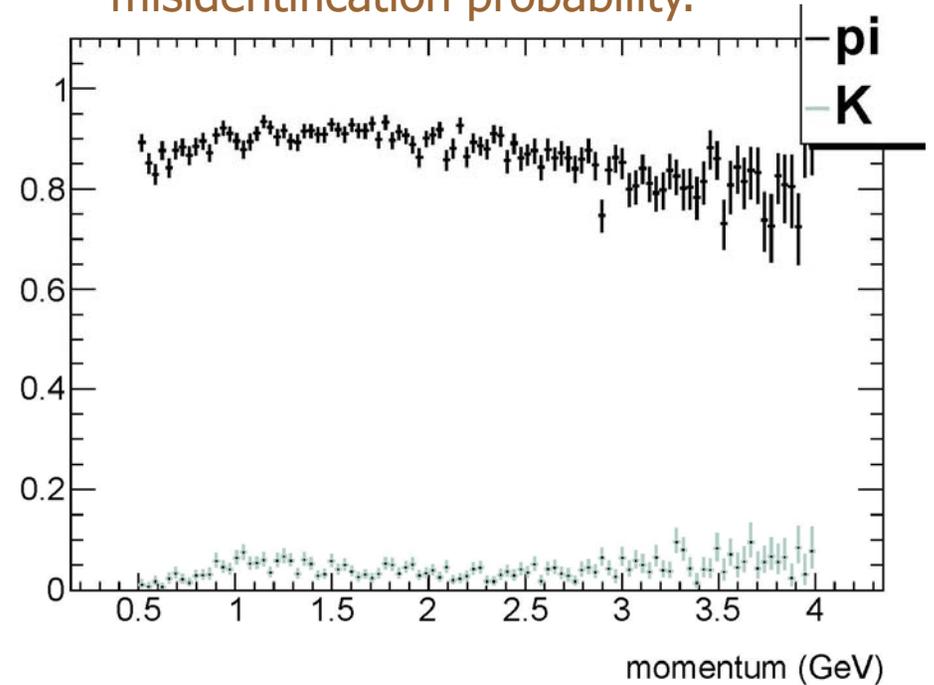
# Separation of kaons and pions



**Pions vs kaons in TOP:**  
different patterns in the time vs  
PMT impact point coordinate

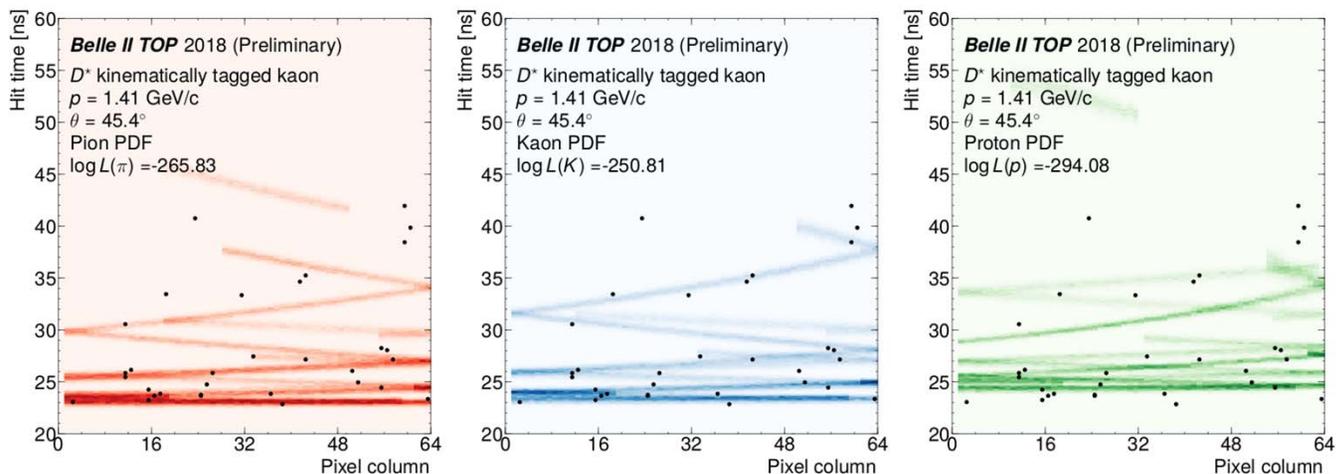


**Pions vs kaons:**  
Expected PID efficiency and  
misidentification probability.

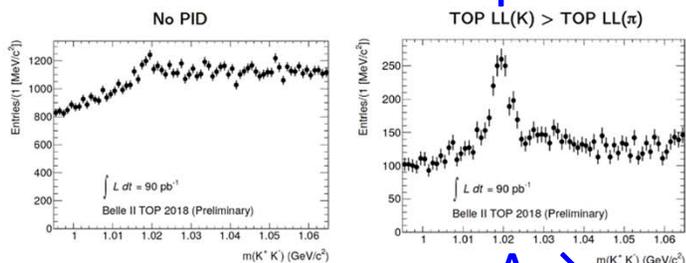


# TOP first events

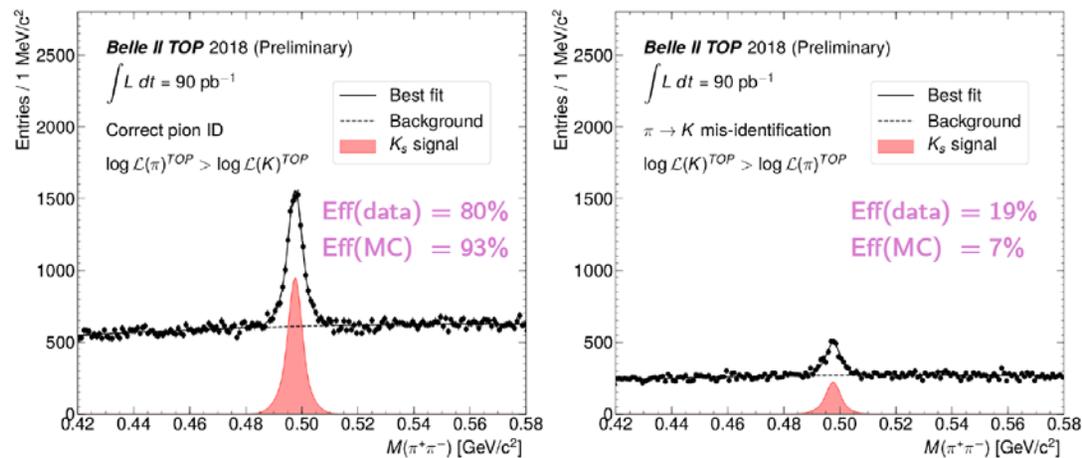
The early data demonstrated that the TOP principle is working



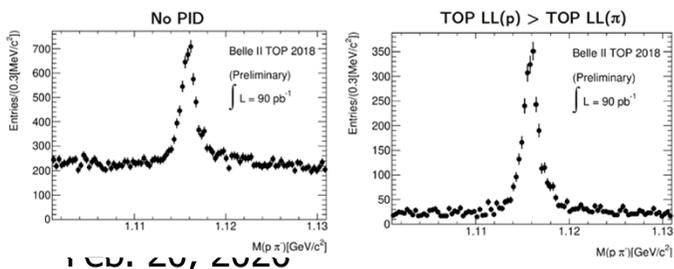
$\phi \rightarrow K^+K^-$  with both the tracks in the TOP acceptance  $\phi \rightarrow KK$



$K_s \rightarrow \pi\pi$

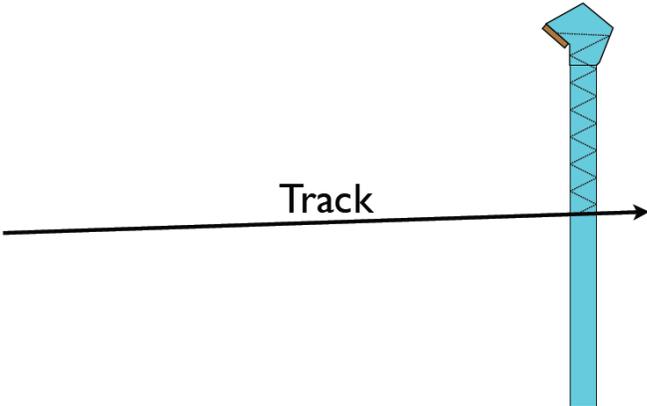


$\Lambda \rightarrow p\pi$  with the proton candidate in the TOP acceptance  $\Lambda \rightarrow p\pi$



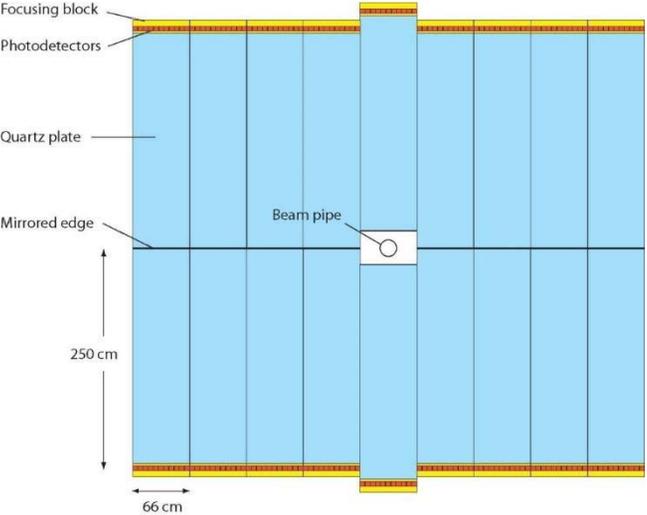
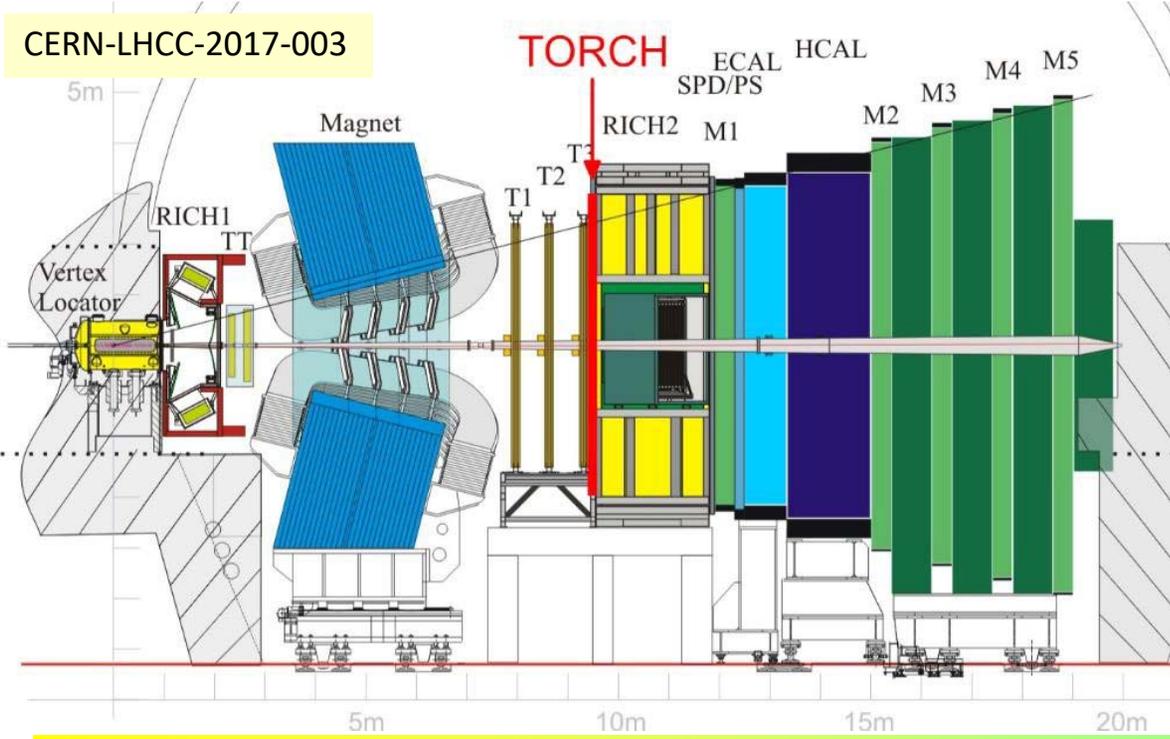
# LHCb PID upgrade: TORCH

A special type of Time-of-Propagation counter for the LHCb upgrade



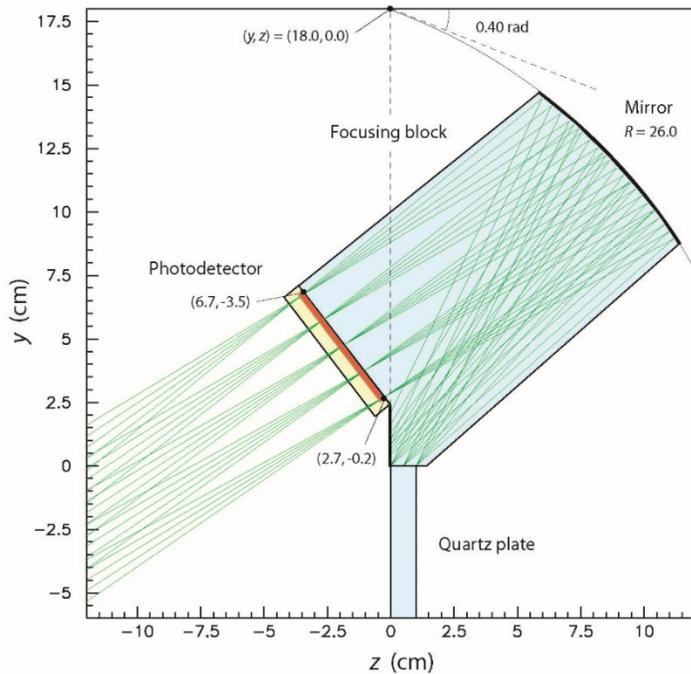
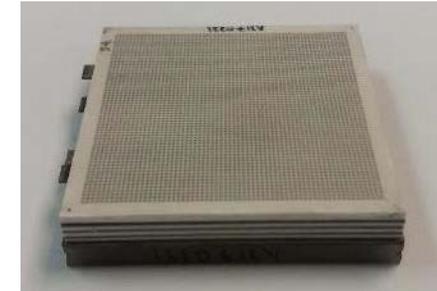
- TORCH area 5 x 6 m<sup>2</sup>
- 18 module system
- 11 MCPs per module

CERN-LHCC-2017-003

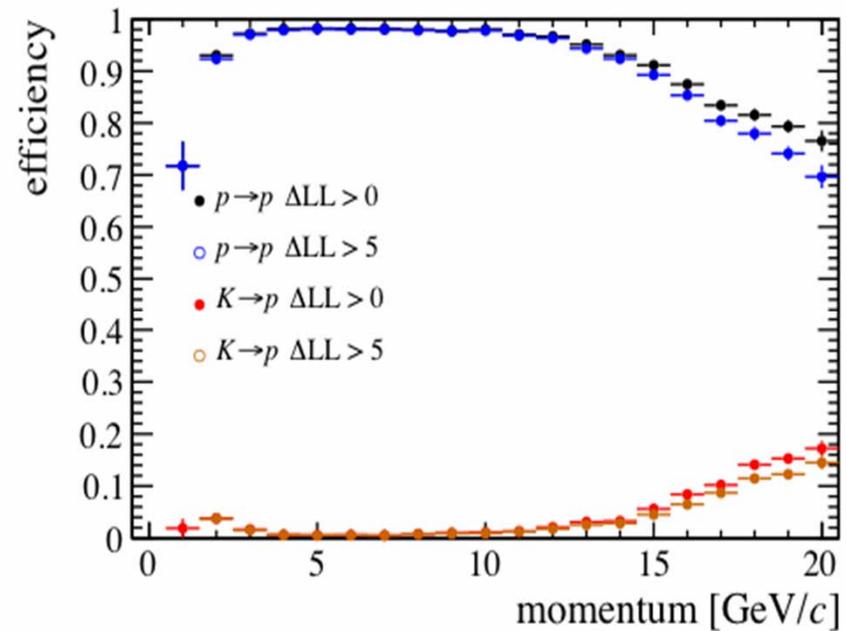
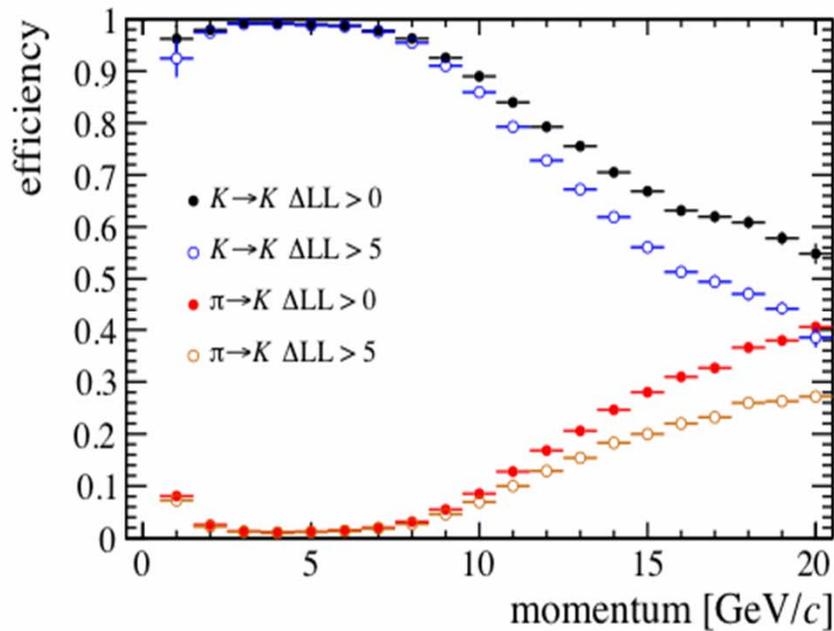


# LHCb PID upgrade: TORCH

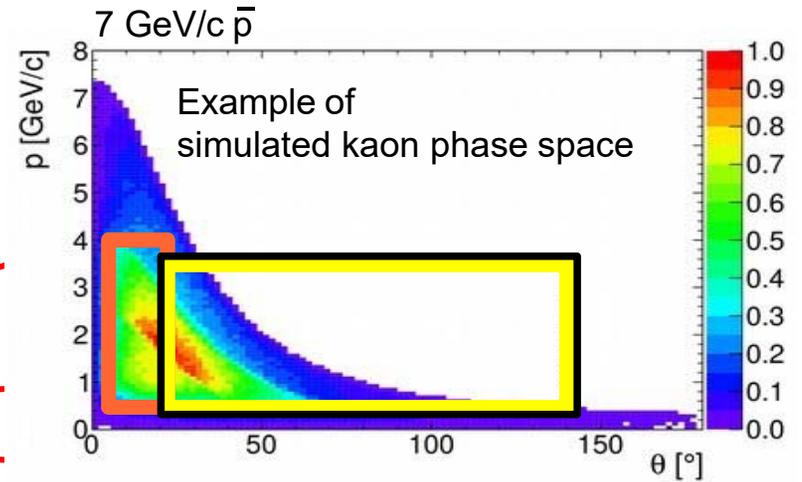
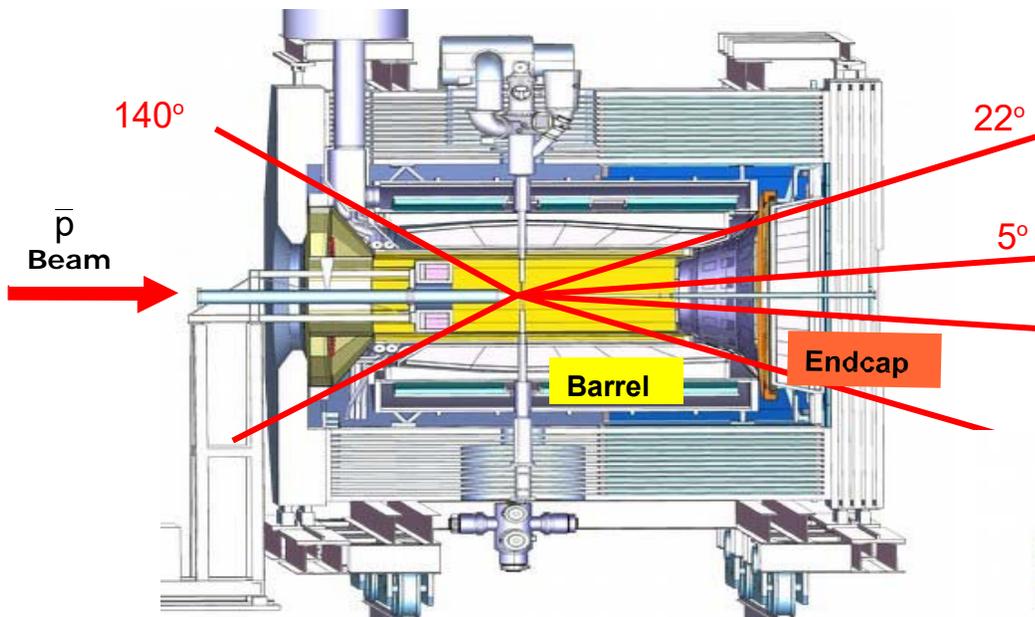
Focusing block with light sensors (MCP PMTs from Photek)



Expected performance



# PANDA DIRC counters

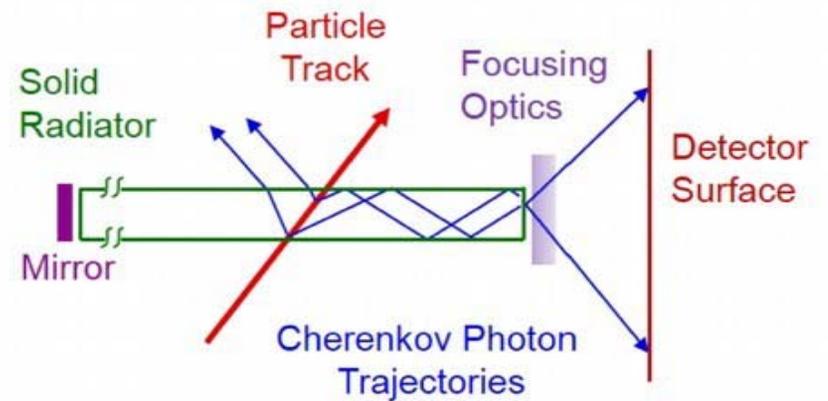


## Barrel DIRC

Goal: 3 s.d.  $\pi/K$  separation up to 3.5 GeV/c

## Endcap Disc DIRC

Goal: 3 s.d.  $\pi/K$  separation up to 4 GeV/c



Magnitude of photon angles in radiator preserved

# PANDA Barrel DIRC

**Design:** based on BABAR DIRC and SuperB FDIRC with key improvements

- Barrel radius ~48 cm; expansion volume depth: 30 cm.

- 48 narrow radiator bars, synthetic fused silica  
17 mm (T) x 53 mm (W) x 2400 mm (L).

- Compact photon detector:**

30 cm fused silica expansion volume 8192 channels of  
MCP-PMTs

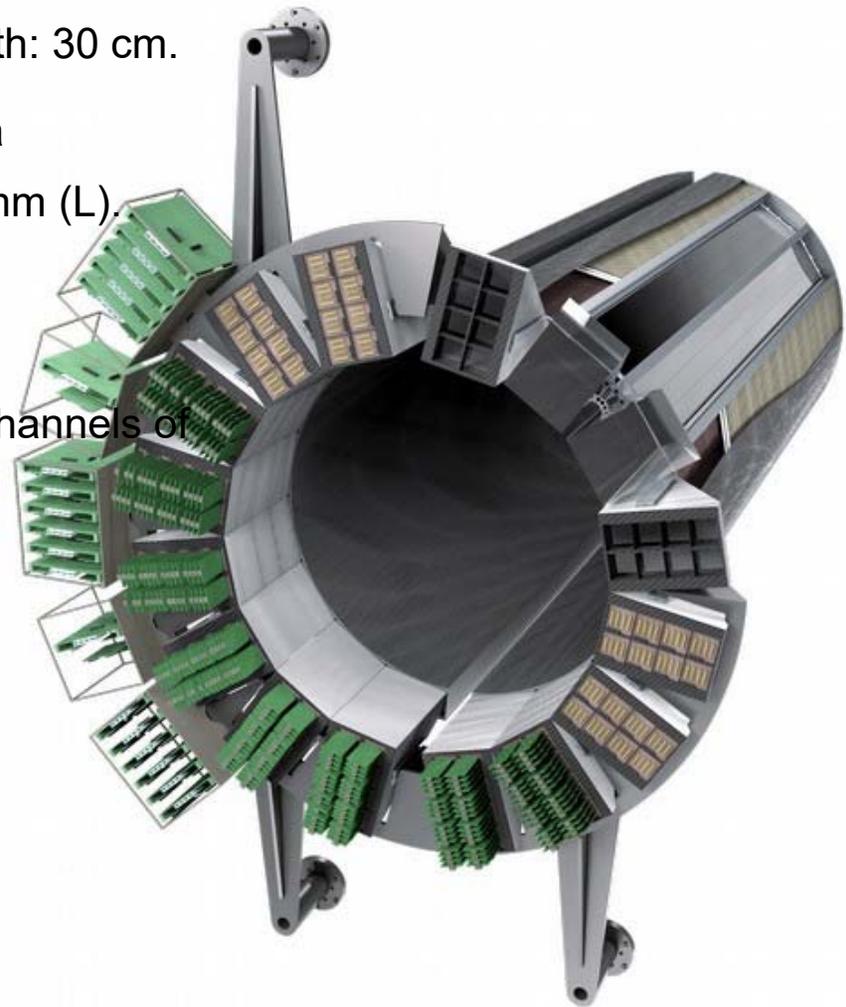
in ~1T B field

- Focusing optics:** spherical lens system

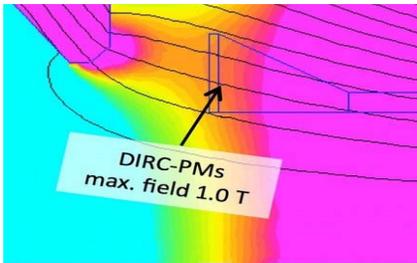
- Fast photon detection:**

fast TDC plus TOT electronics,

→ 100-200 ps timing



# Photon detector



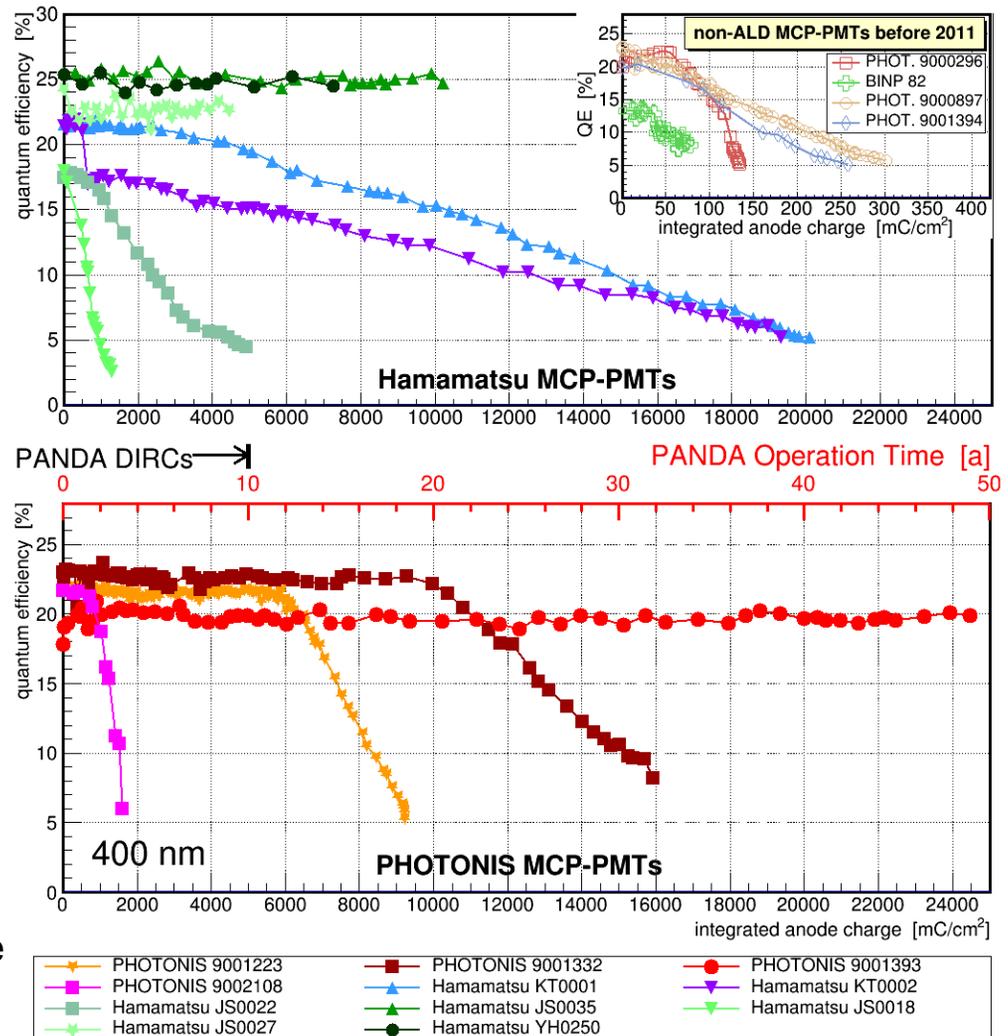
Requirements:

- few mm spatial resolution
- ~100 ps timing resolution

Bar-box:

8 MCP-PMT, 512 pixels (total 8 k readout channels) with pixel size 6 x 6 mm<sup>2</sup> work in 1T magnetic field survive 10 years of PANDA (aging)

Most sensors with ALD coated MCPs have lifetime > 5 C/cm<sup>2</sup>



# Panda Disc DIRC

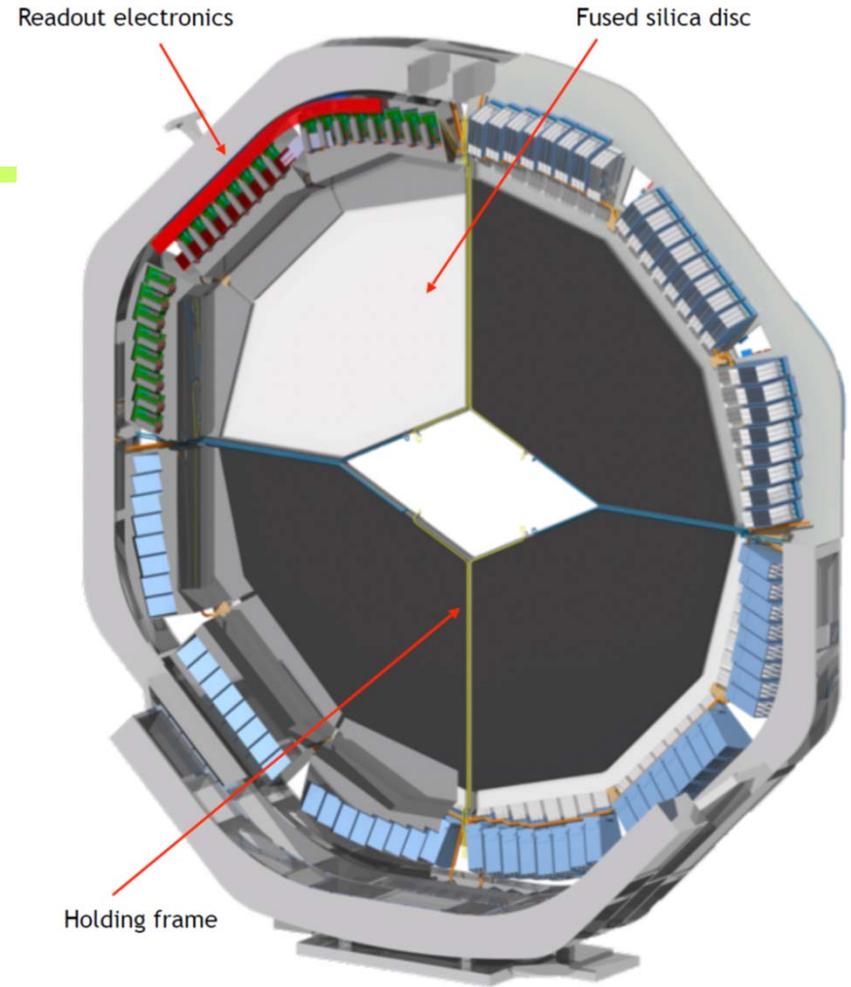
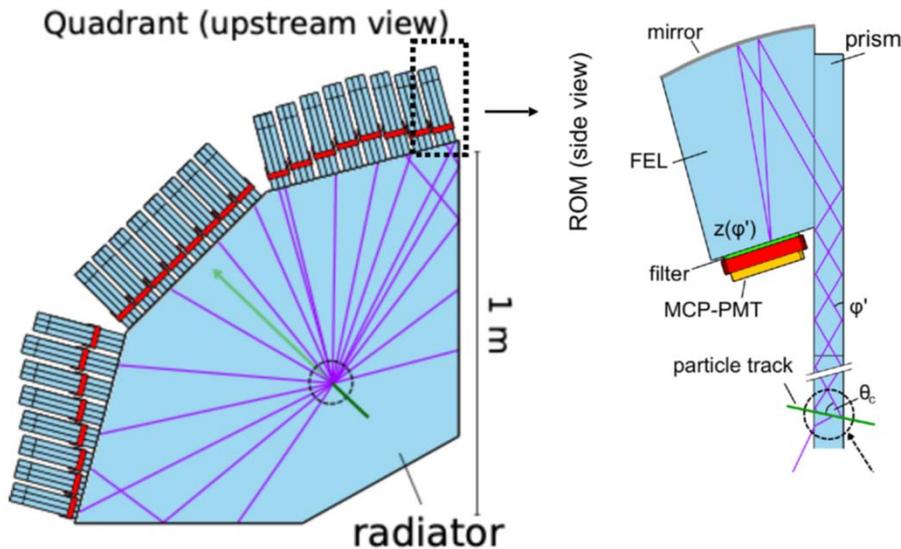
Radiator: fused silica 20 mm thick,  
 $R = 1\text{m}$

$\pi/K$  separation up to 4 GeV/c

Focusing optics

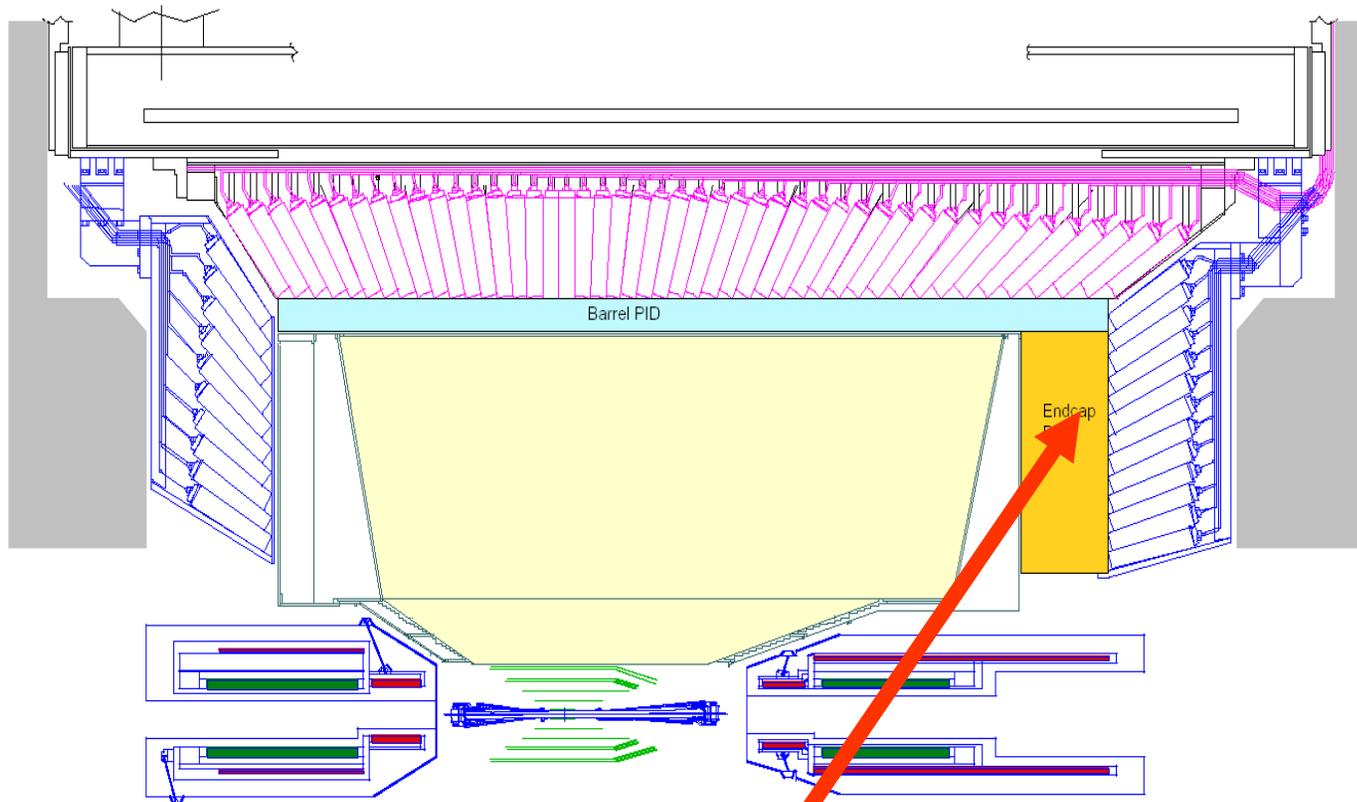
Photon detector in  $\sim 1\text{T}$  field:

- 96 MCP PMTs with a highly segmented anode, TofPET2 readout ASIC





# Belle II PID system



Two novel particle ID devices, both RICHes:

Barrel: Time-of-propagation counter (TOP) counter

Endcap: proximity focusing RICH



# Endcap: Proximity focusing RICH

K/ $\pi$  separation at 4 GeV/c:  
 $\theta_c(\pi) \sim 308$  mrad ( $n = 1.05$ )  
 $\theta_c(\pi) - \theta_c(K) \sim 23$  mrad

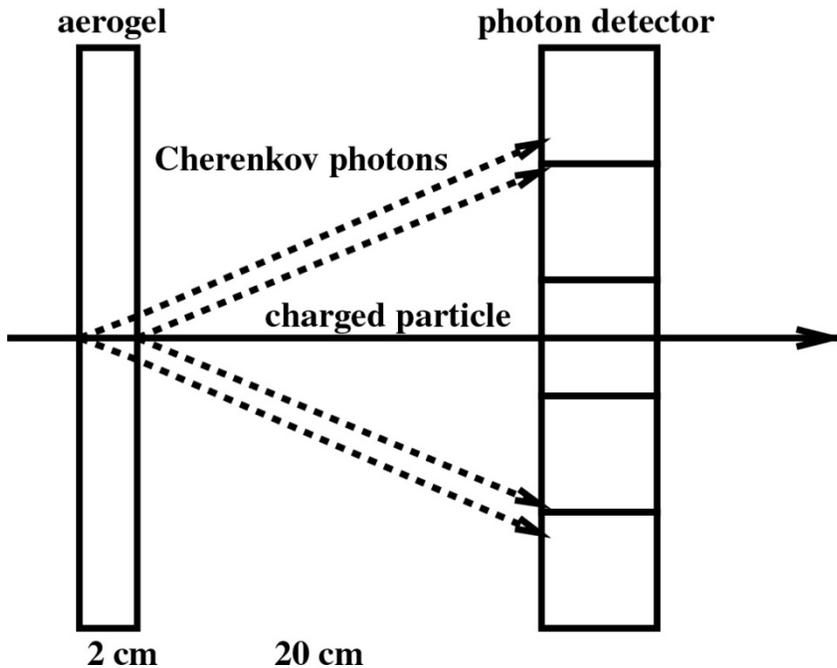
For single photons:  $\delta\theta_c(\text{meas.}) = \sigma_0 \sim 14$  mrad,  
typical value for a 20mm thick radiator and  
6mm PMT pad size

Per track:

$$\sigma_{\text{track}} = \frac{\sigma_0}{\sqrt{N_{pe}}}$$

Separation:  $[\theta_c(\pi) - \theta_c(K)] / \sigma_{\text{track}}$

$\rightarrow 5\sigma$  separation with  $N_{pe} \sim 10$



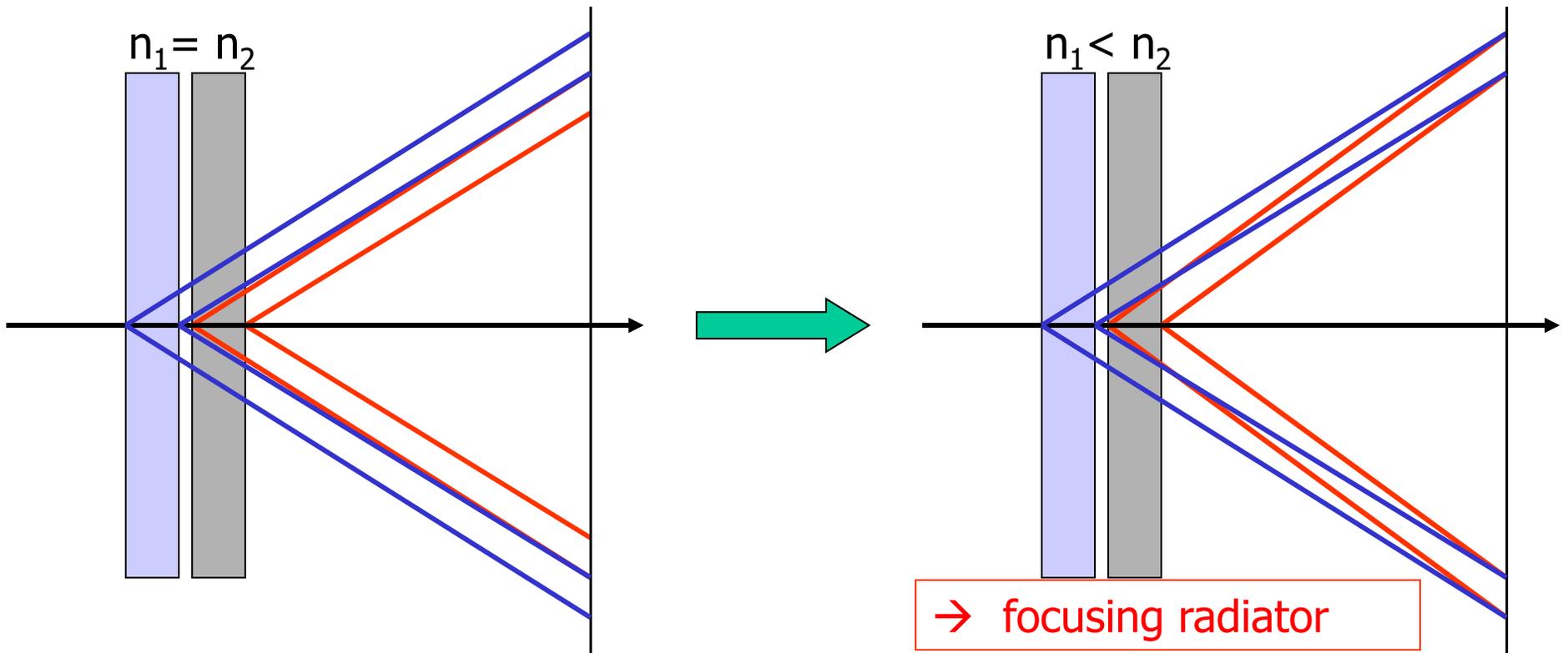


# Radiator with multiple refractive indices

Small number of photons from aerogel  $\rightarrow$  need a thick layer of aerogel.  
How to improve the resolution by keeping the same number of photons?

$\rightarrow$  stack two tiles with different refractive indices:  
“focusing” configuration

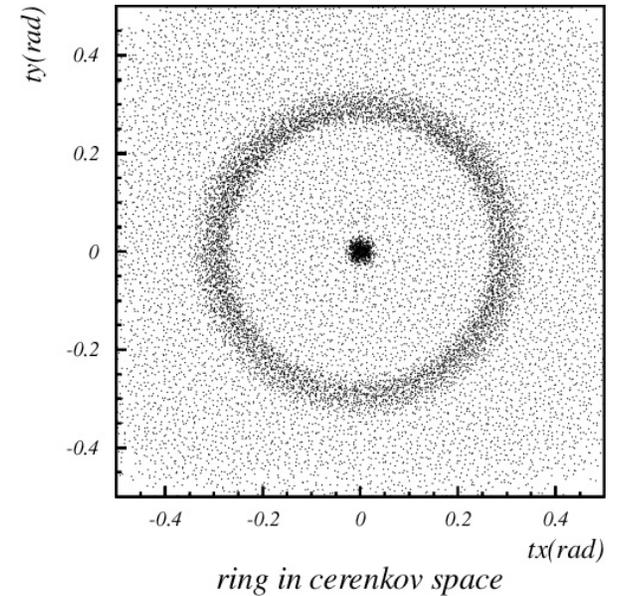
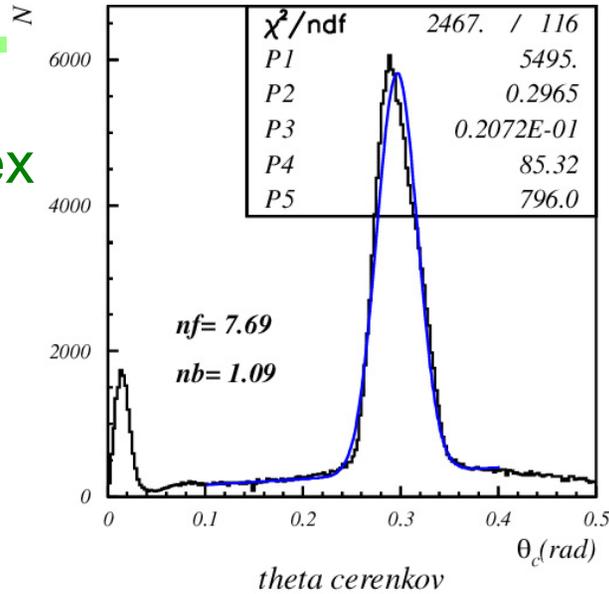
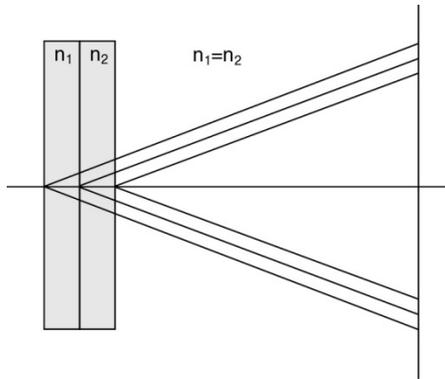
normal



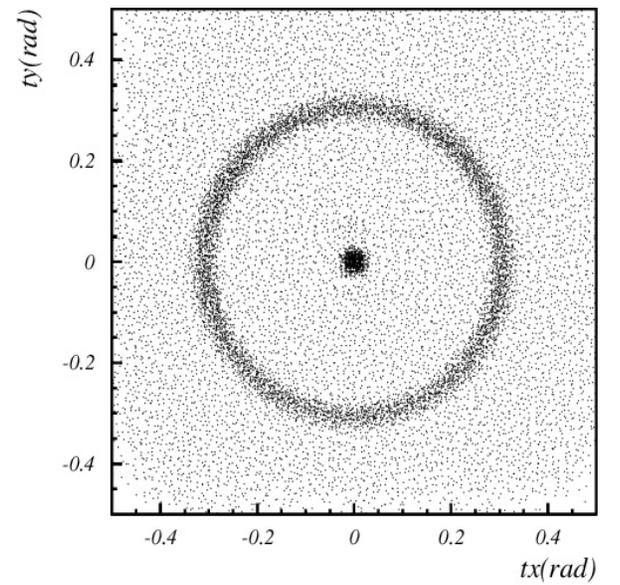
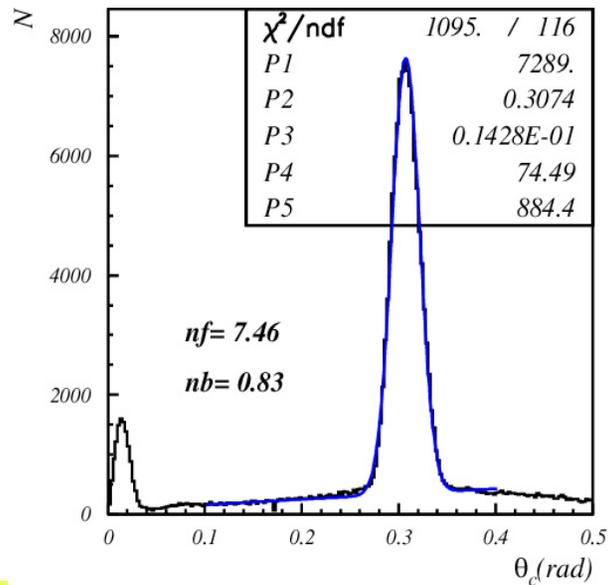
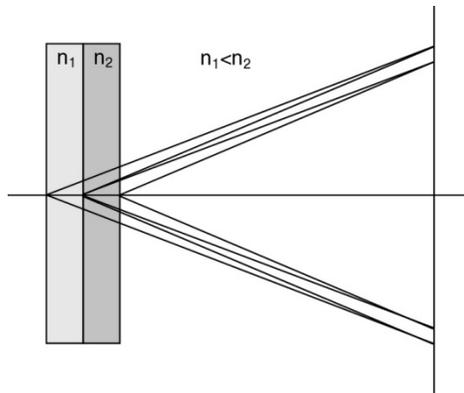


# Focusing configuration – data

4cm aerogel single index



2+2cm aerogel

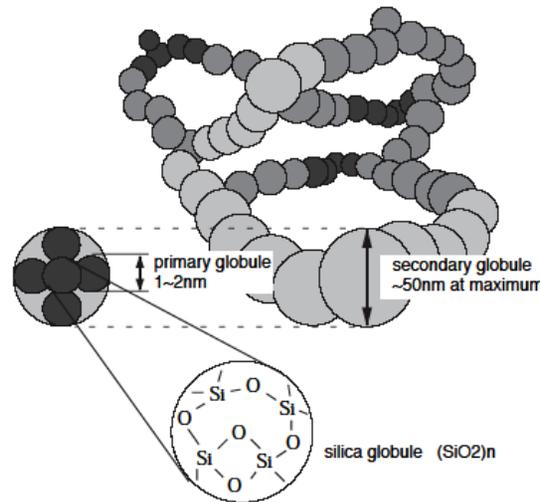
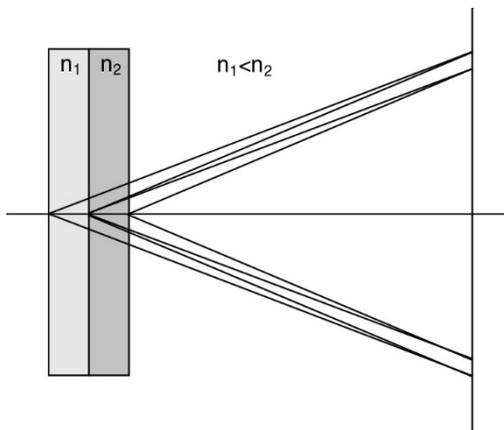


→ NIM A548 (2005) 383, NIMA 565 (2006) 457

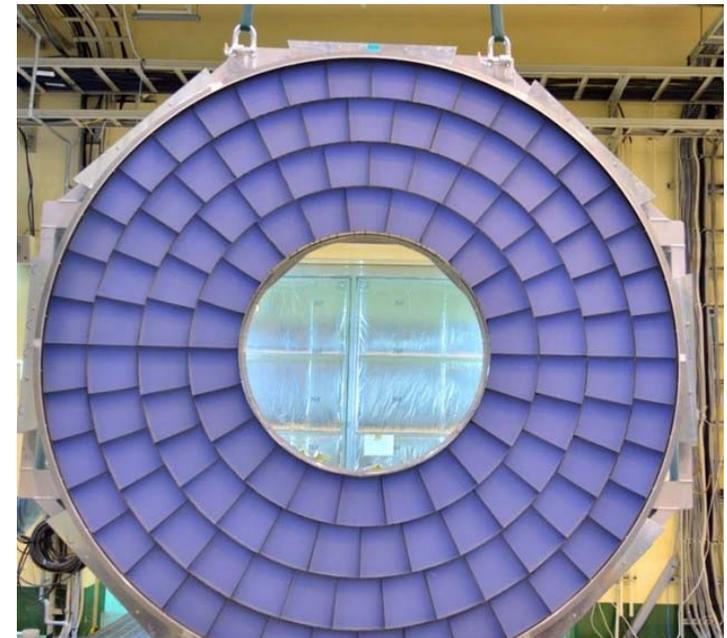
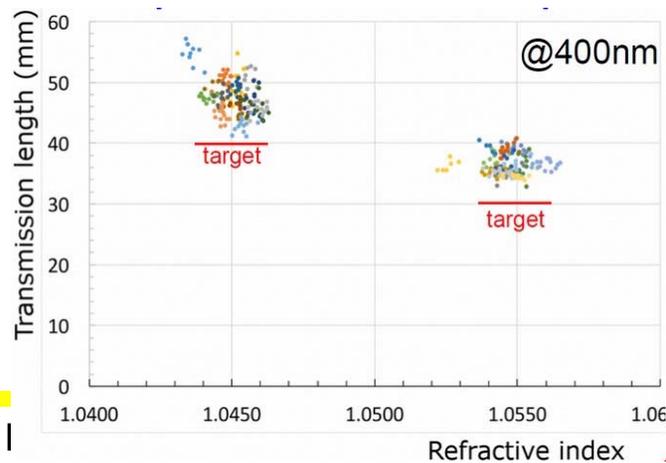
4x4 array of flat panel MAPMTs

# Radiator with multiple refractive indices 2

Such a configuration is only possible with aerogel (a form of  $\text{Si}_x\text{O}_y$ ) – material with a **tunable** refractive index between **1.01** and **1.07**.



Requires aerogel with high transparency



Detector plane covered with 2 x 124 tiles water-jet cut tiles (~ 17x17cm)

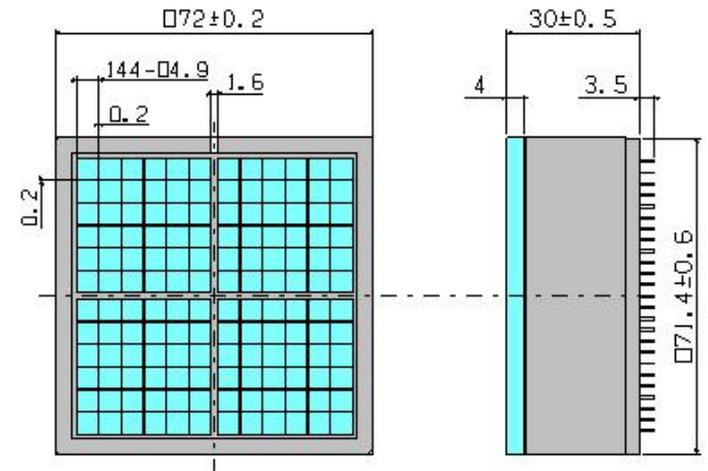
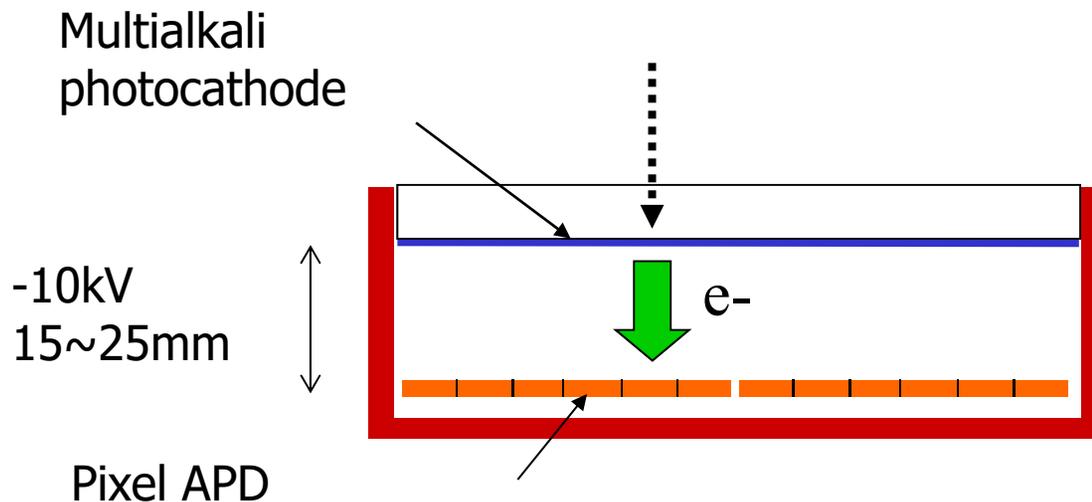


# Photon detectors for the aerogel RICH requirements and candidates

Need: Operation in a high magnetic field (1.5 T)  
Pad size  $\sim 5\text{-}6\text{mm}$

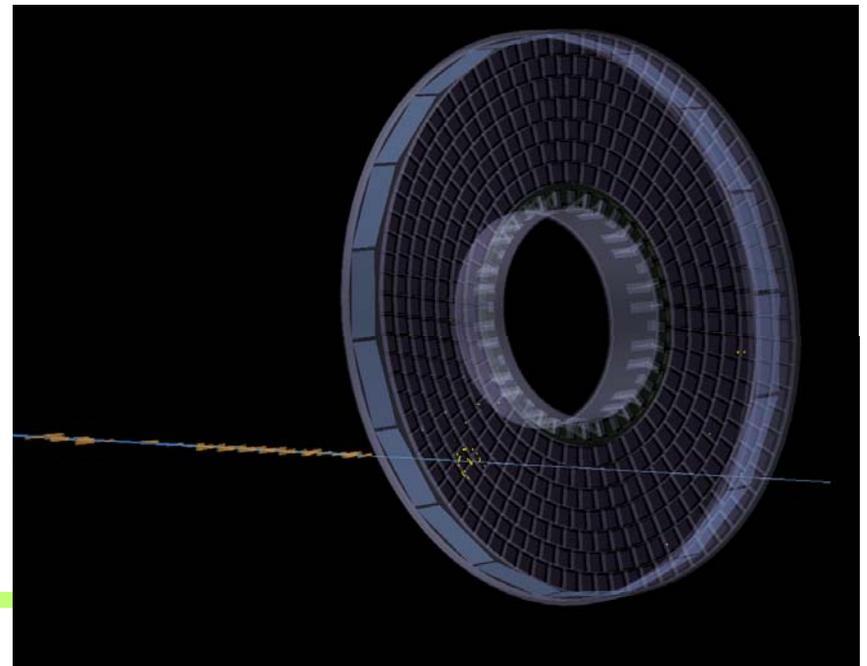
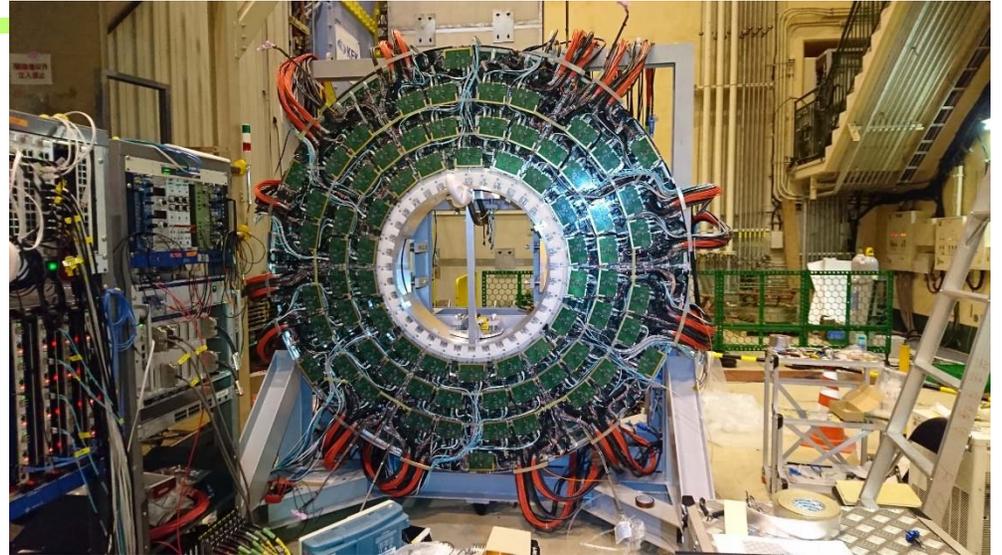
Final choice: large active area HAPD (hybrid avalanche photon detector) of the proximity focusing type

Other candidates: MCP PMT (Photonis 85011), SiPMs



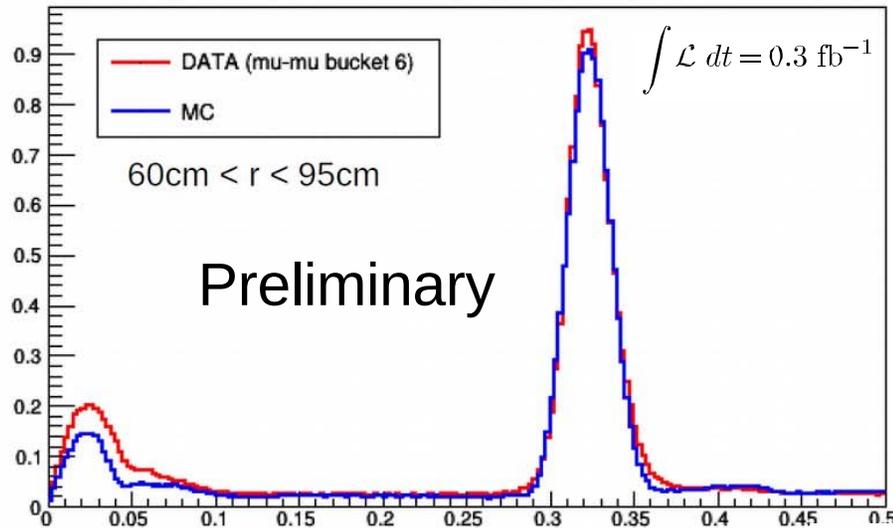
HAPD R&D project in collaboration with HPK.

# The big eye of ARICH – and one of the first rings



# Performance in the early Belle II data

## Cherenkov angle distribution in $e^+e^- \rightarrow \mu^+\mu^-$



**DATA**

$$N_{sig} = 11.38/\text{track}$$

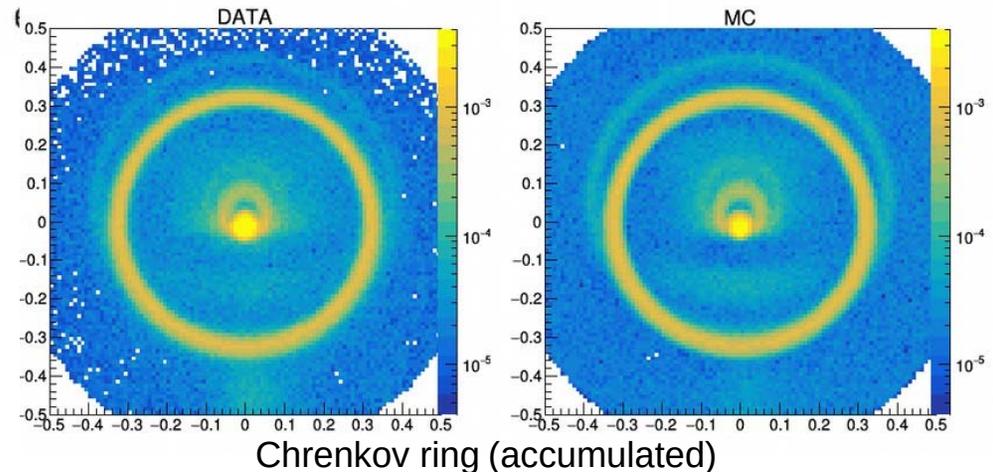
$$\sigma_c = 12.7 \text{ mrad}$$

**MC**

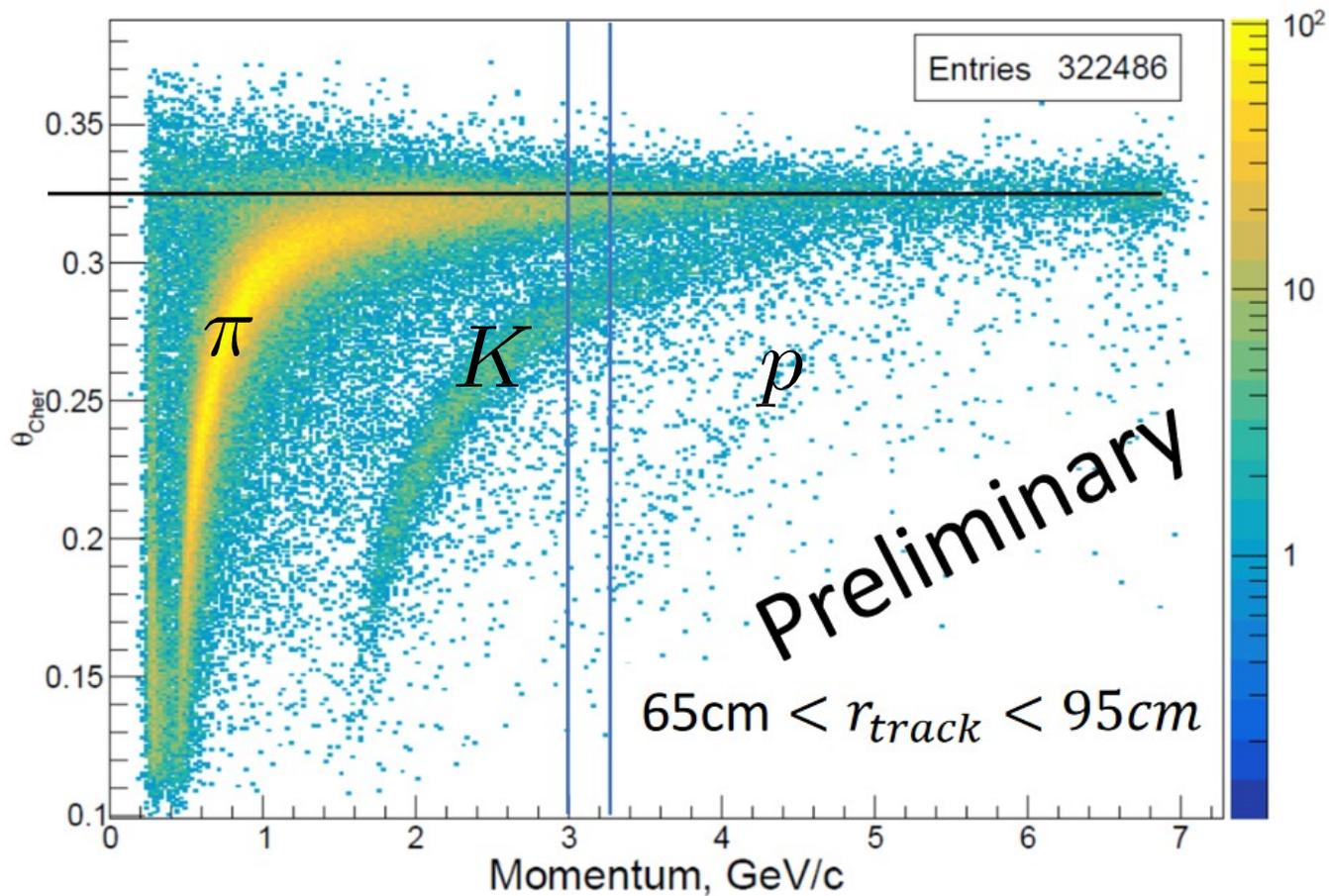
$$N_{sig} = 11.27/\text{track}$$

$$\sigma_c = 12.75 \text{ mrad}$$

Overall a very good  
DATA/MC agreement !



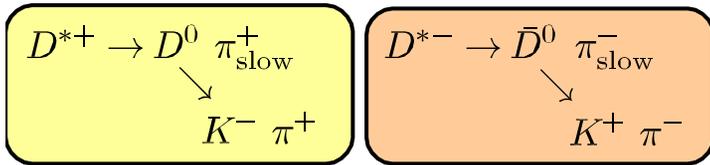
# Cherenkov angle vs momentum in hadronic events



Average Cherenkov angle for tracks from hadronic events

# Estimation of $\pi/K$ separation power using $D^{*\pm}$ decays

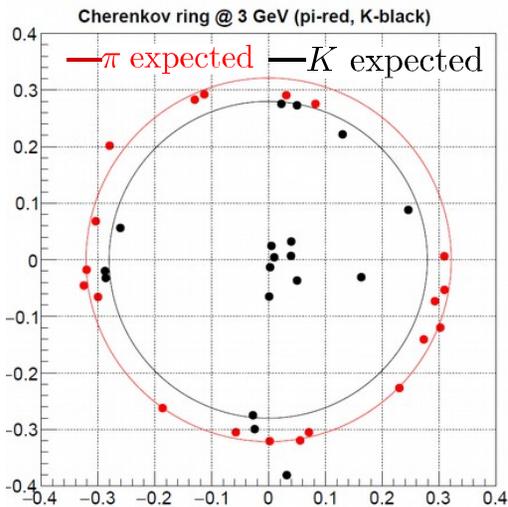
- Identify  $K$ ,  $\pi$  based on track charge in association with the charge of  $\pi_{\text{slow}}$



- Apply selection criteria on

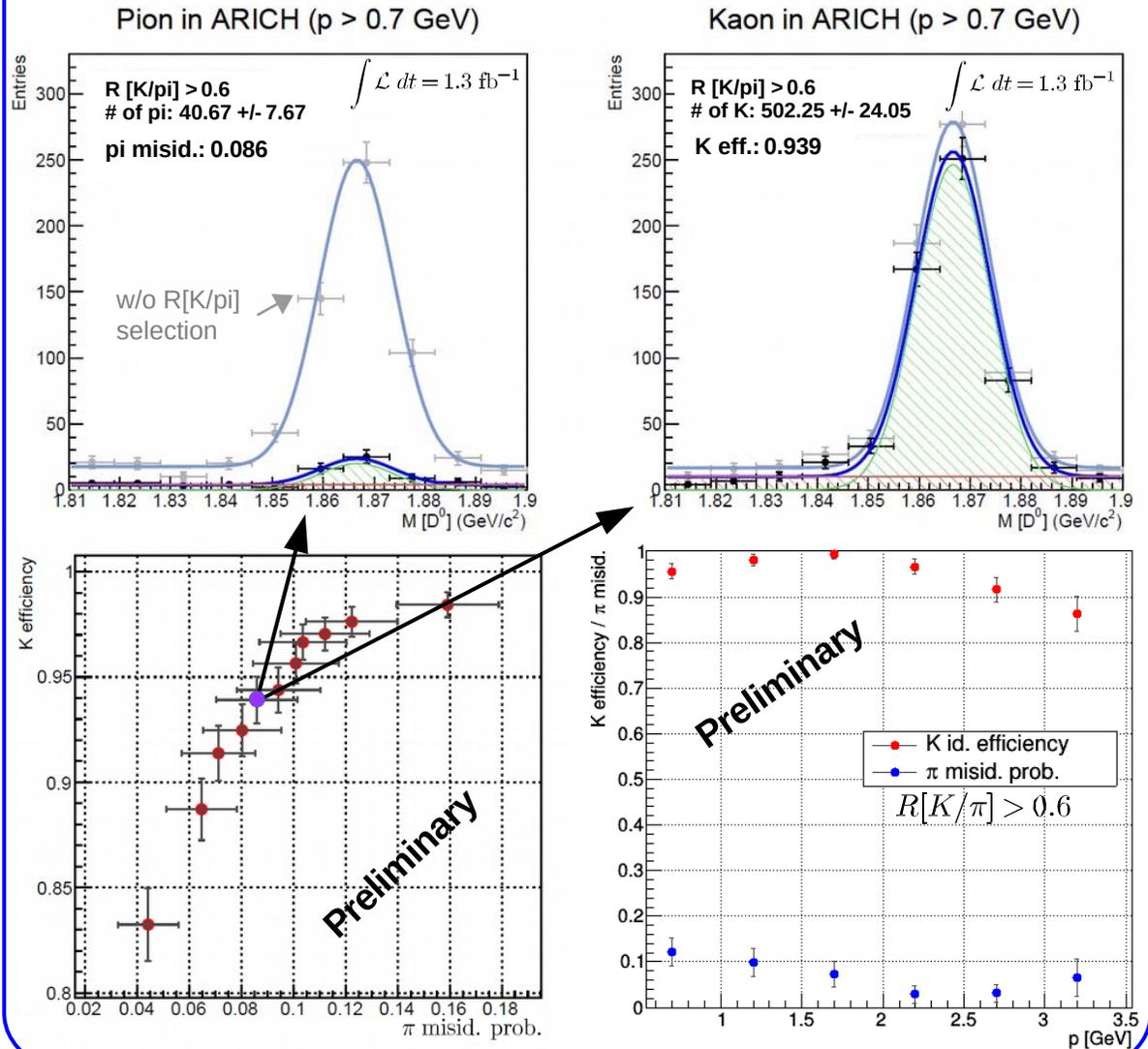
$$R[K/\pi] = \frac{\mathcal{L}_K}{\mathcal{L}_K + \mathcal{L}_\pi}$$

$\mathcal{L}$  - likelihood for given id. hypothesis



- Only coarse/preliminary calibrations included  
→ further improvements expected

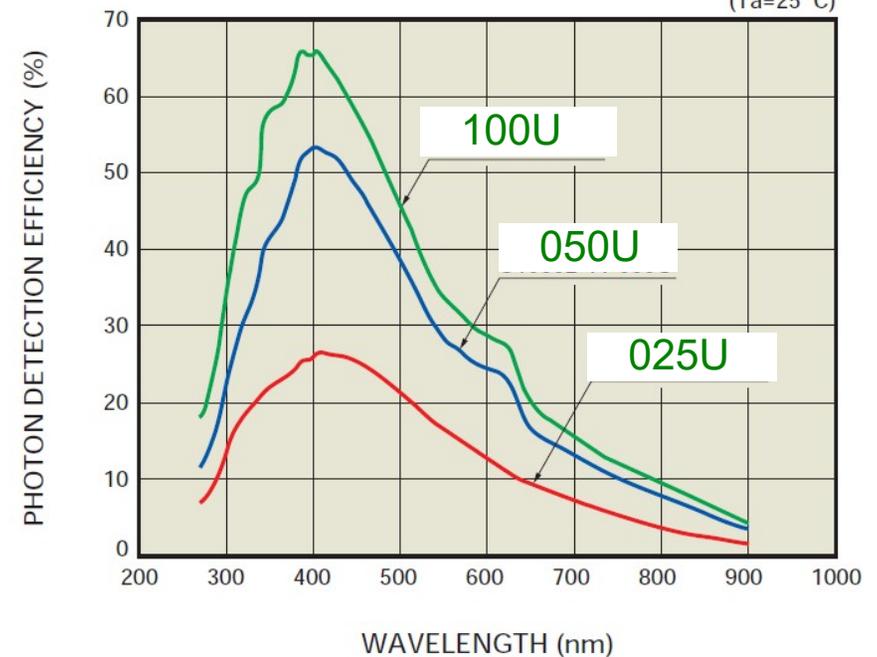
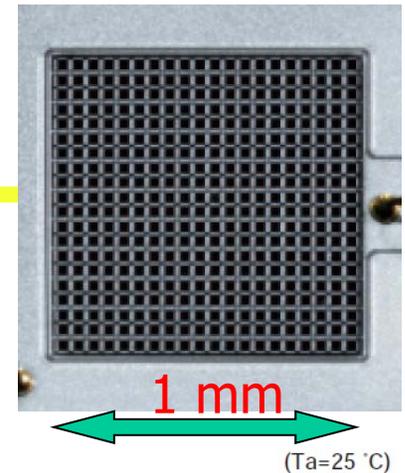
## $D^0$ mass peak



# SiPMs as photon detectors?

SiPM is an array of APDs operating in Geiger mode. Characteristics:

- low operation voltage  $\sim 10\text{-}100\text{ V}$
- gain  $\sim 10^6$
- peak PDE up to 65%(@400nm)  
 $\text{PDE} = \text{QE} \times \epsilon_{\text{geiger}} \times \epsilon_{\text{geo}}$  (up to 5x PMT!)
- $\epsilon_{\text{geo}}$  – dead space between the cells
- time resolution  $\sim 100\text{ ps}$
- works in high magnetic field
- dark counts  $\sim \text{few } 100\text{ kHz/mm}^2$
- radiation damage (p,n)



Not trivial to use in a RICH where we have to detect single photons!

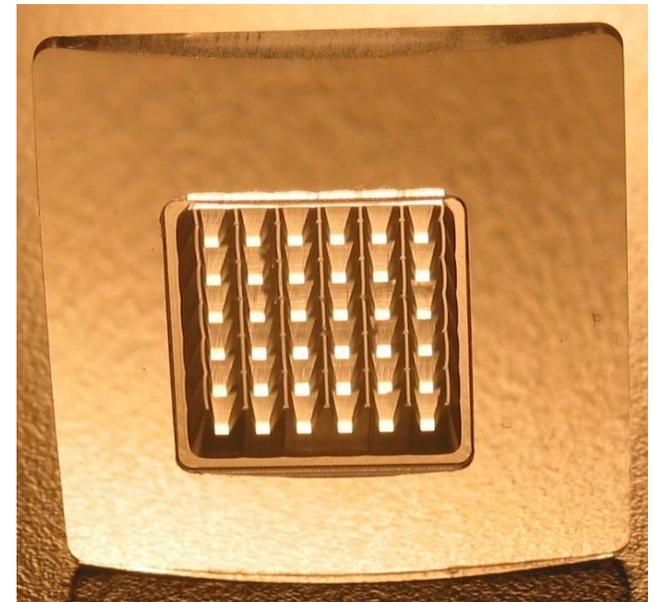
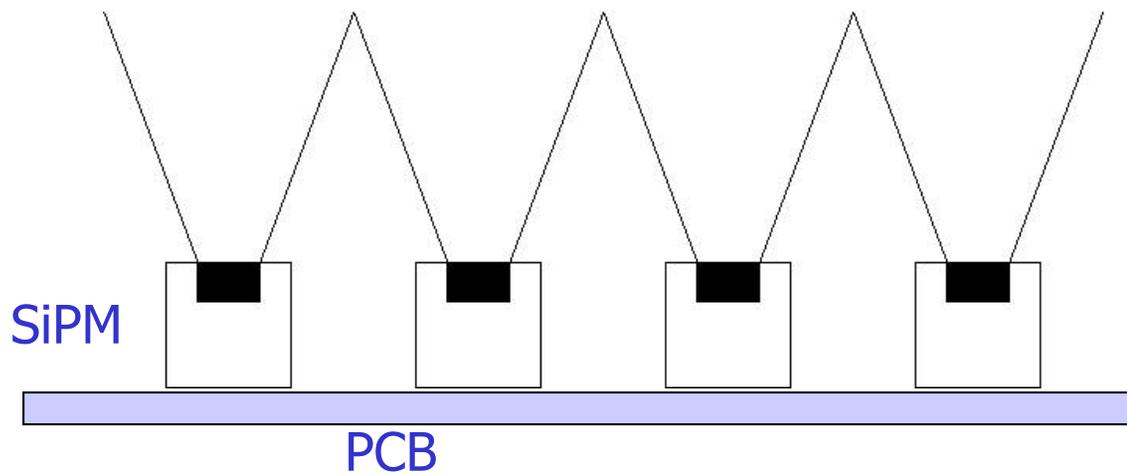
Dark counts have single photon pulse heights (rate 0.1-1 MHz per  $\text{mm}^2$ )

# SiPM as photosensor for a RICH counter

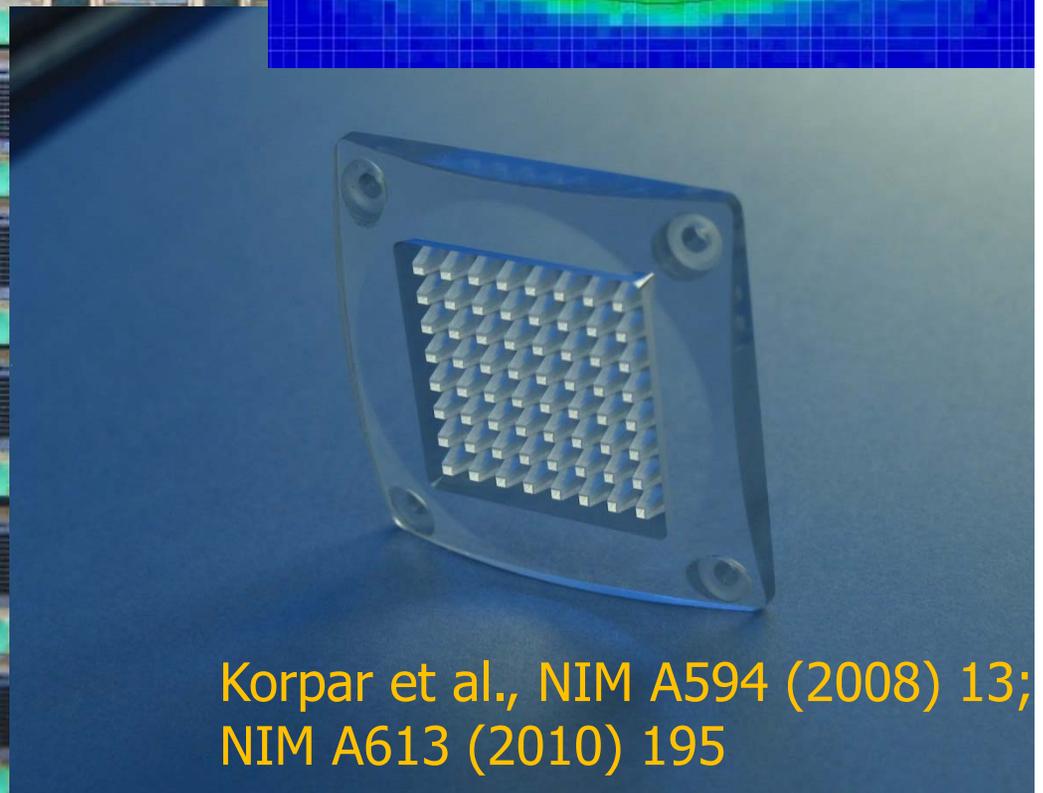
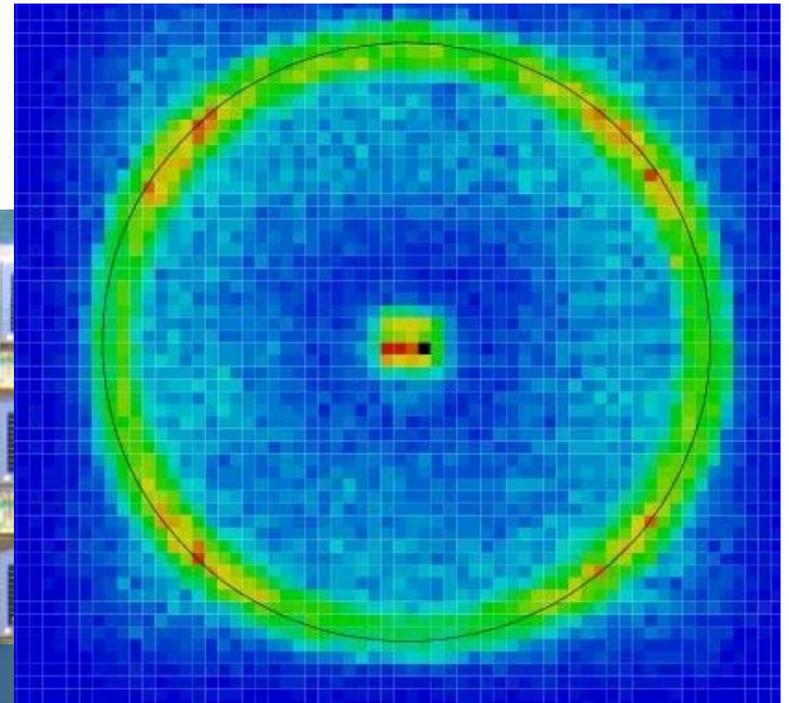
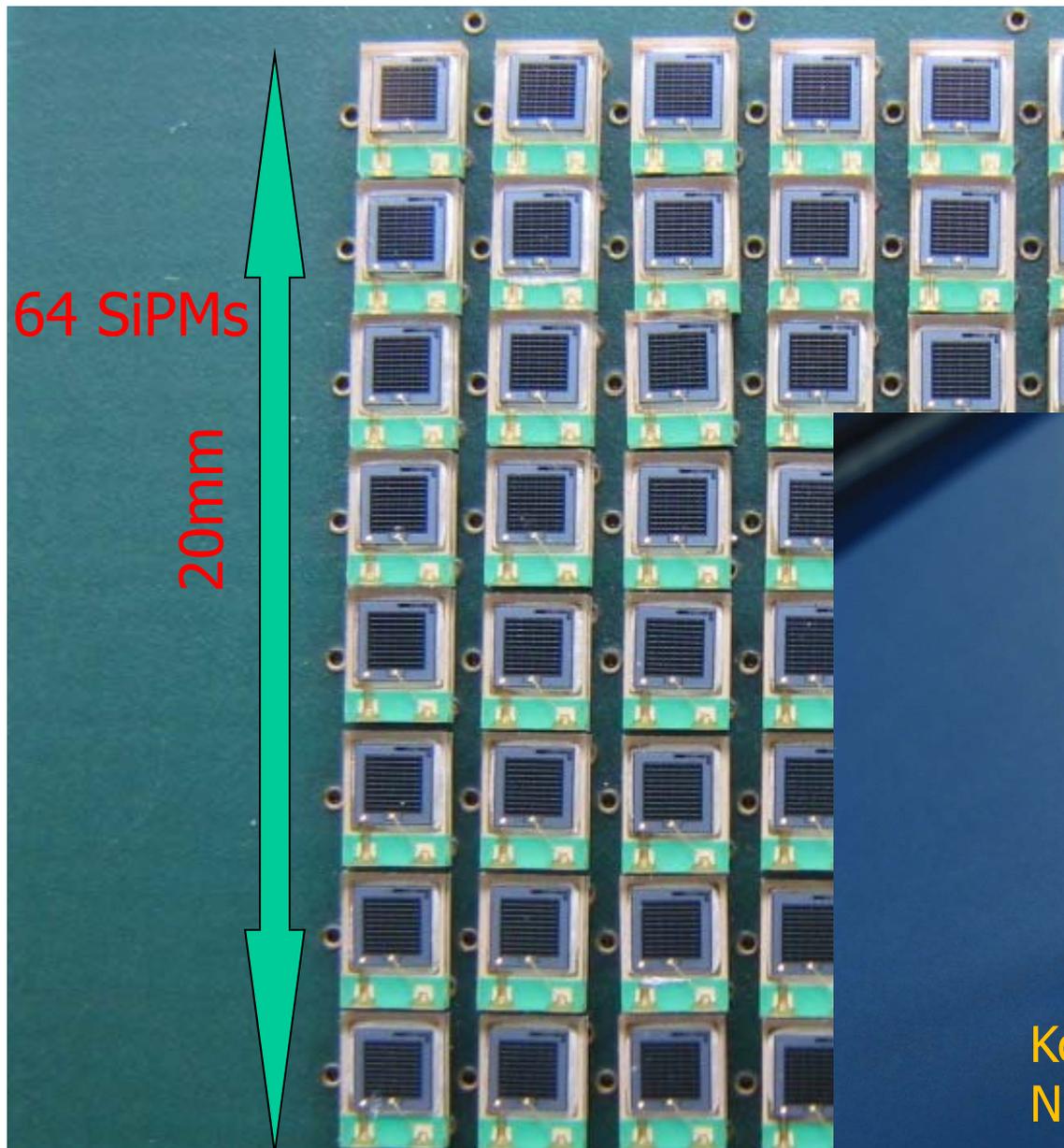
Improve the signal to noise ratio:

- Reduce the noise by a narrow ( $<10\text{ns}$ ) time window (Cherenkov light is prompt!)
- Increase the number of signal hits per single sensor by using light collectors

E.g. light collector with reflective walls or plastic light guide



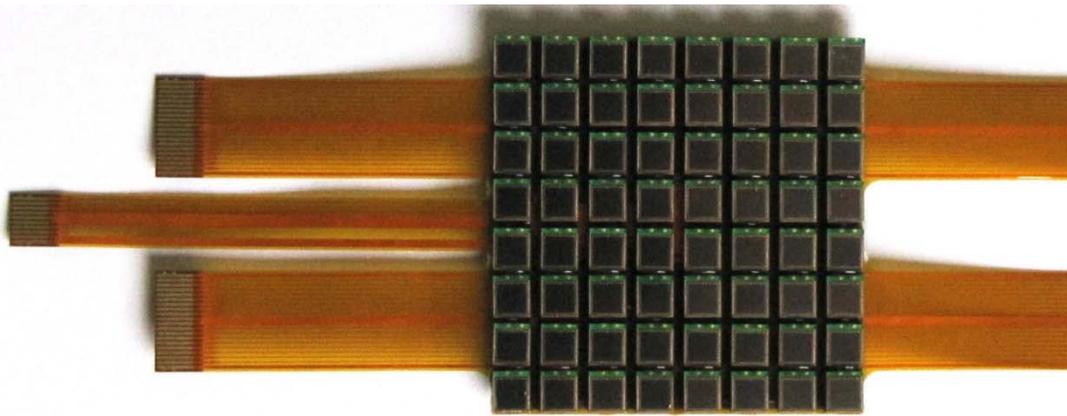
# Photon detector with SiPMs and light guides



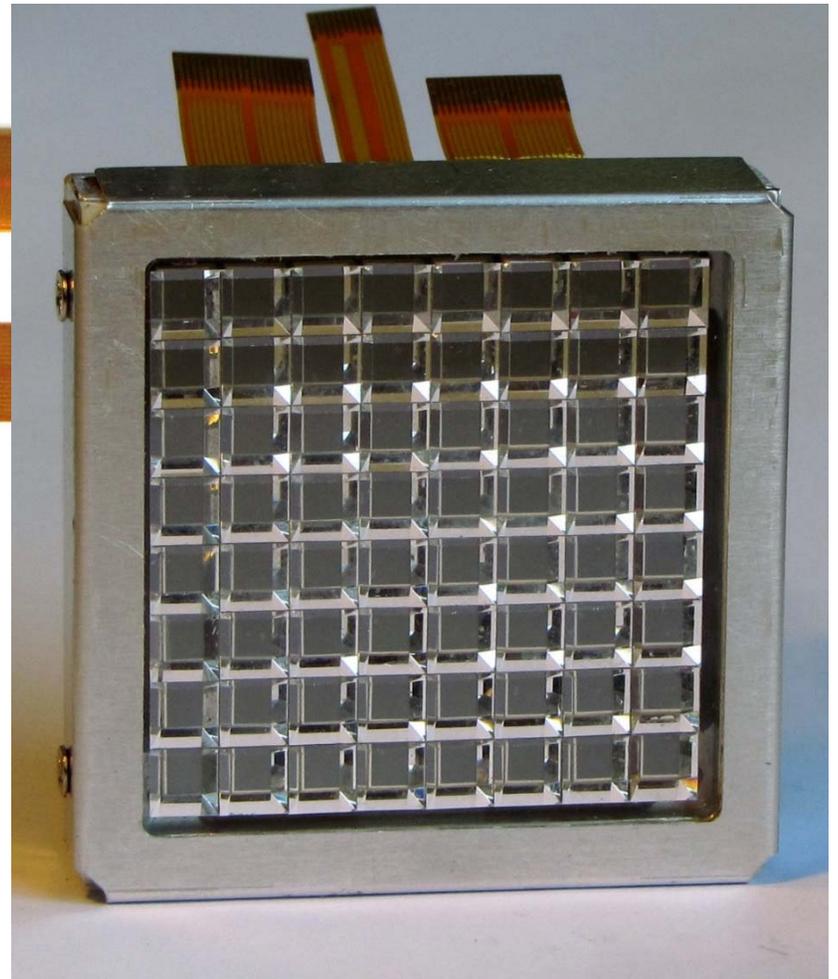
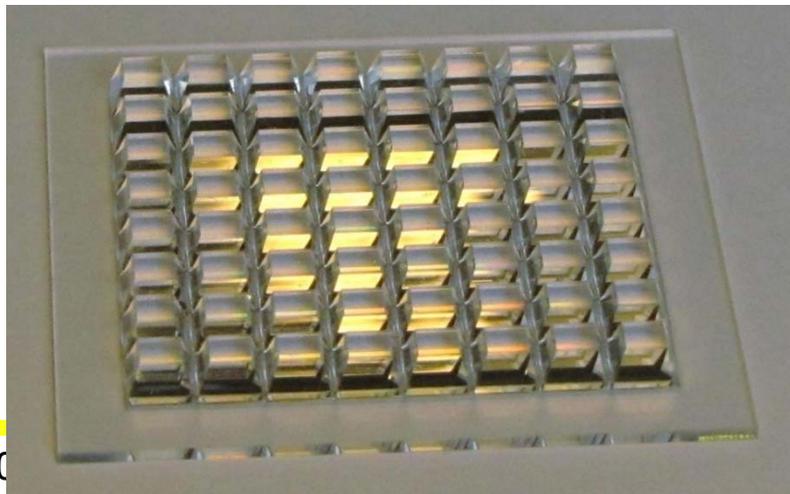
# Next step: use arrays of SiPMs

## Example: Hamamatsu MPPC S11834-3388DF

- 8x8 SiPM array, with 5x5 mm<sup>2</sup> SiPM channels
- Active area 3x3 mm<sup>2</sup>

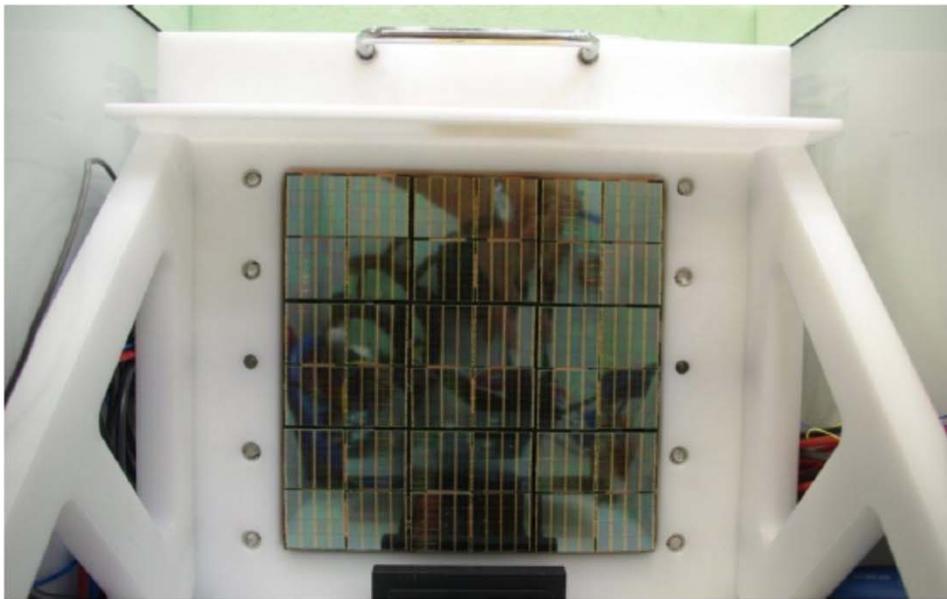
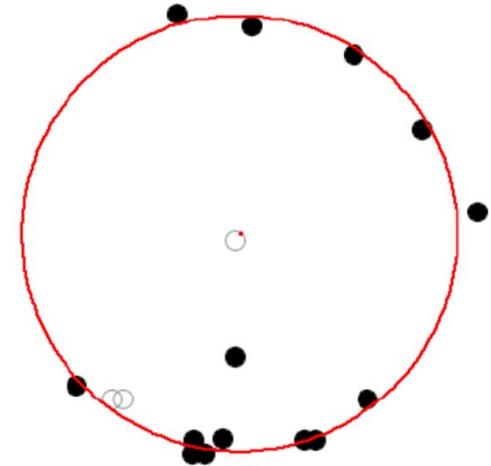


+ array of quartz light collectors



# Digital SiPM

Digital SiPM (Philips): instead of an analog sum of signals from all cells of a single SiPM, use on board electronics for a digital sum + time stamp

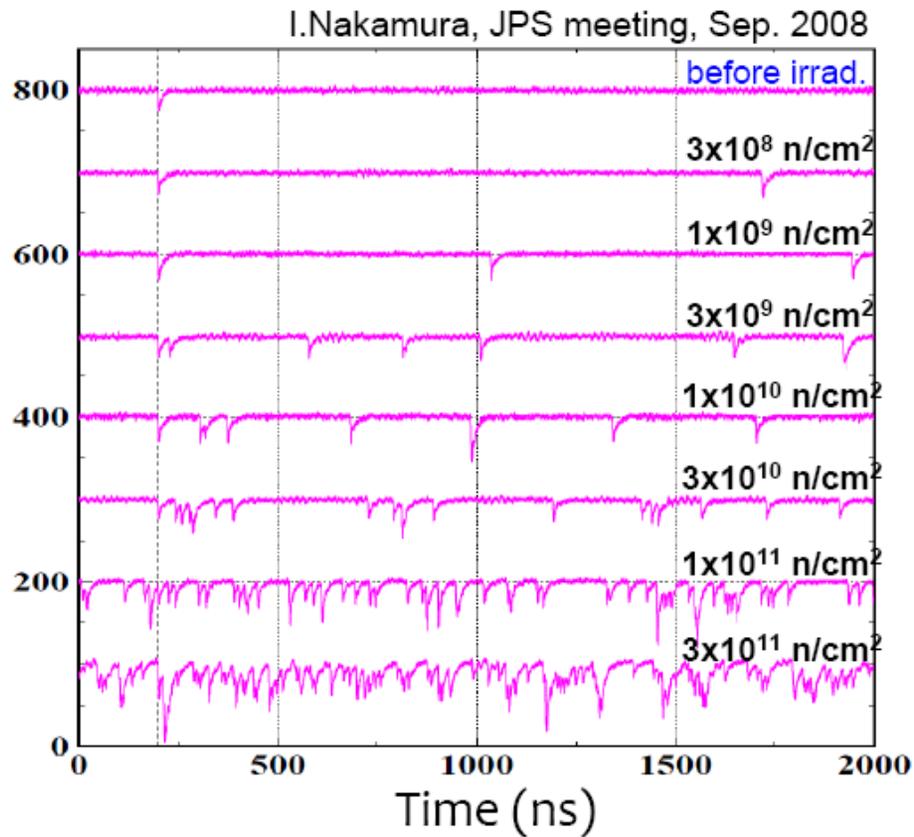


Square matrix **20x20 cm<sup>2</sup>**

- Sensors: DPC3200-22-44
- 3x3 modules = 6x6 tiles = 24x24 dies = 48x48 pixels in total
- 576 time channels
- 2304 amplitude (position) channels
- 4 levels of FPGA readout: tiles, modules, bus boards, test board

→ A.Y. Barnyakov et al., NIM A732 (2013) 352

# SiPMs: Radiation damage



Expected fluence at 50/ab at Belle II:

2-20  $10^{11} \text{ n cm}^{-2}$

→ Worst than the lowest line

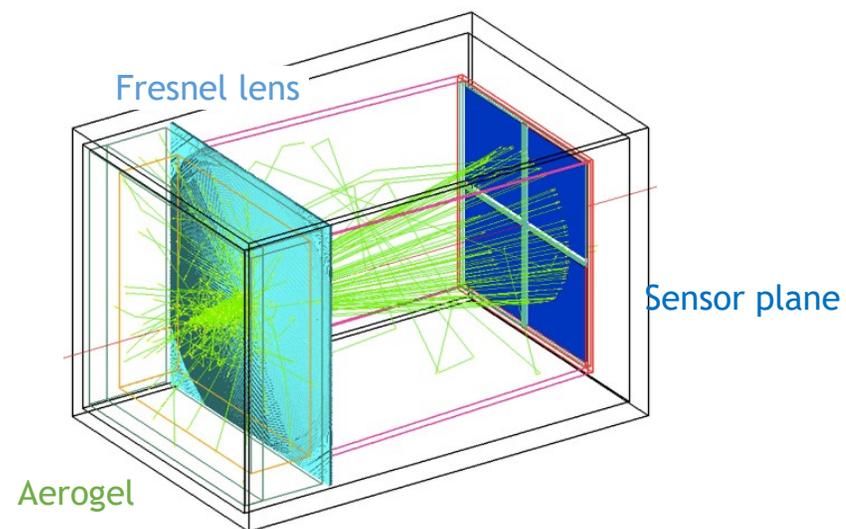
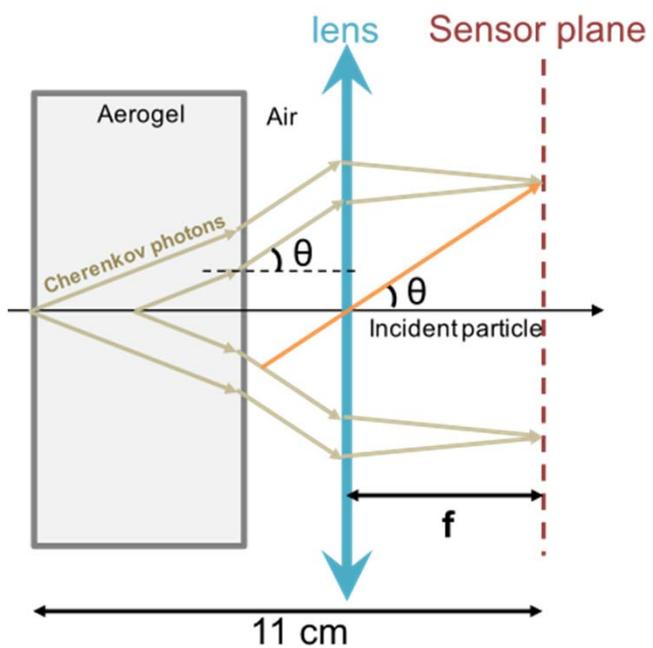
- Need cooling of sensors and wave-form sampling readout electronics
- Annealing?

... and more radiation resistant SiPMs...





# EIC mRICH – Working Principle



Geant4 Simulation

# Identification of charged particles

---

Particles (e,  $\mu$ ,  $\pi$ , K, p) in the final state are identified by their **mass** or by the **way they interact**.

Determination of **mass**: from the relation between momentum and velocity,  $p = \gamma m v$  (p is known - radius of curvature in magnetic field)

→ Measure velocity by:

- time of flight
- ionisation losses  $dE/dx$
- Cherenkov photon angle (and/or yield)
- transition radiation

Mainly used for the identification of hadrons.

Identification through **interaction**: electrons and muons

- muon systems
- calorimeters

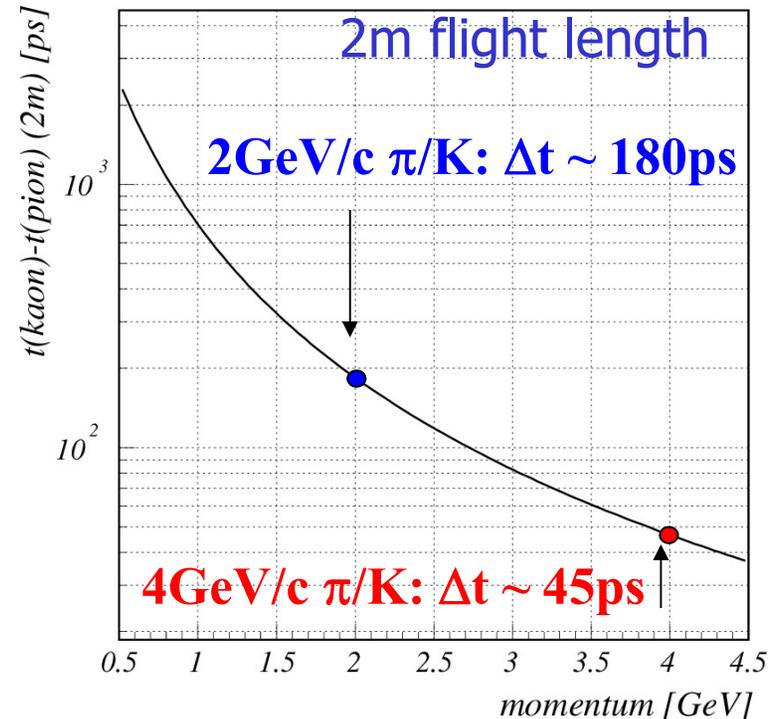
# Time-of-Flight (TOF) counters

Measure velocity by measuring the time between  
-- the interaction and  
-- the passing of the particle through the TOF counter.

Traditionally: plastic scintillator + PMTs

Typical resolution:  $\sim 100$  ps  $\rightarrow \pi/K$  separation up to  $\sim 1$  GeV.

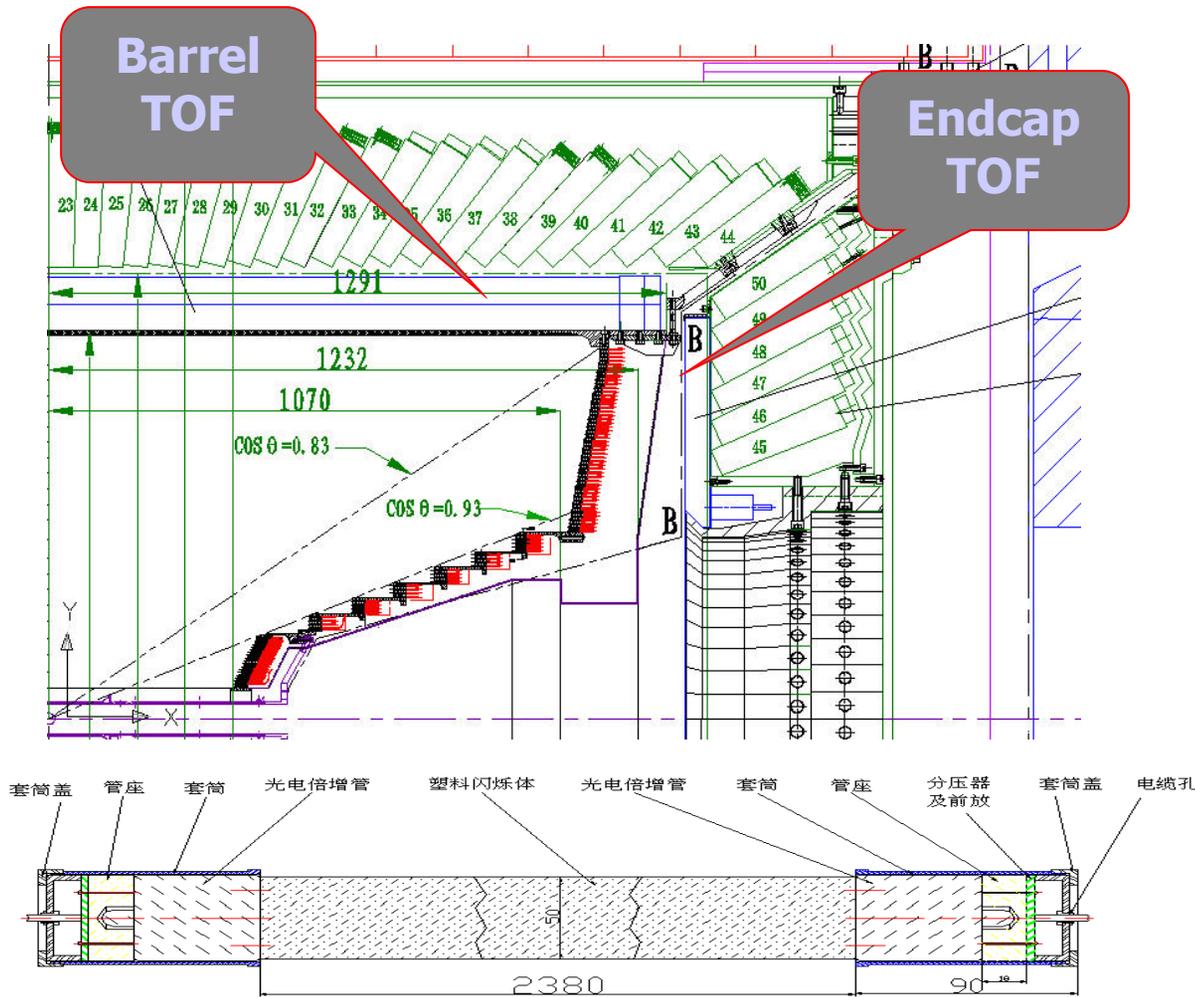
Time difference between  $\pi$  and K:



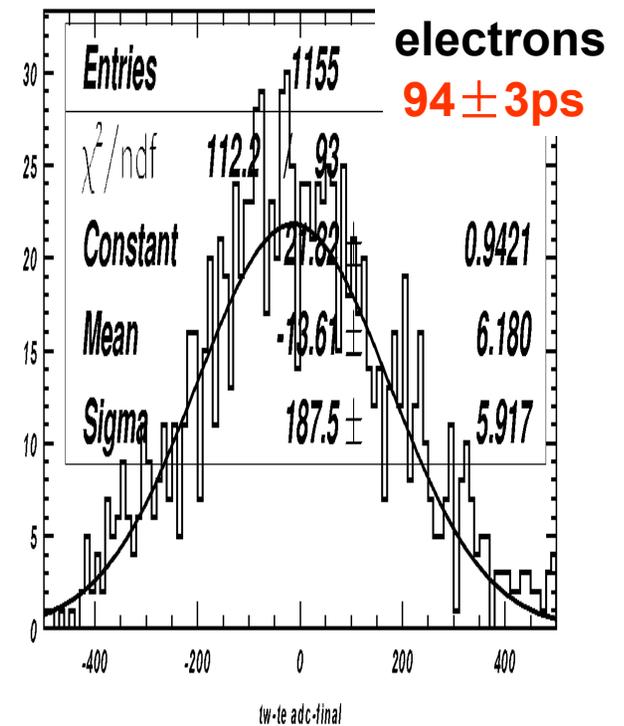
$\rightarrow$  BESSIII



# BESIII: Time-Of-Flight counters



**Beam test results:  
two TOF modules**



**TOF module: high quality plastic scintillator: 2.4 m long, 5cm thick, two PMTs with preamplifiers**

# Time-of-Flight (TOF) counters

Measure velocity by measuring the time between the interaction and the passing of the particle through the TOF counter.

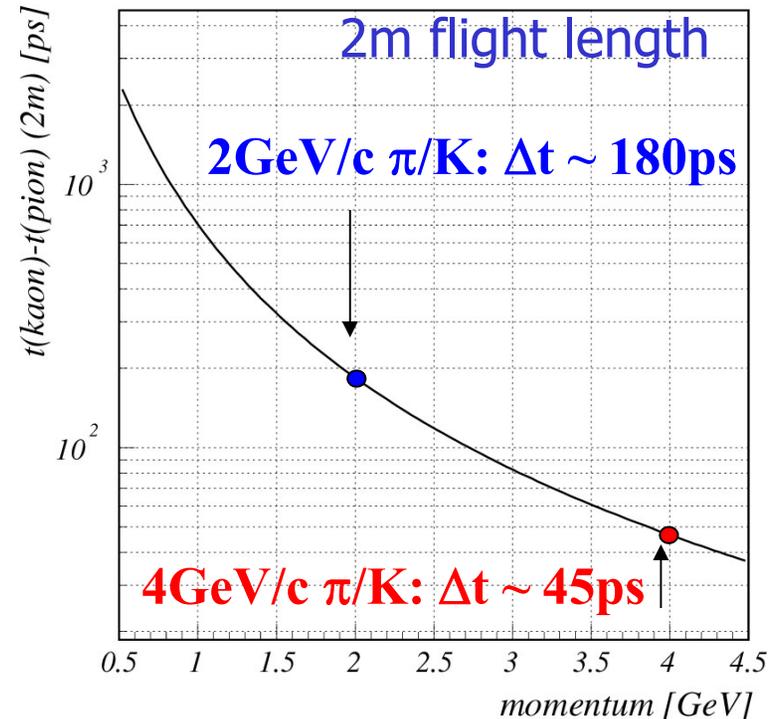
Traditionally: plastic scintillator + PMTs

Typical resolution:  $\sim 100$  ps  $\rightarrow \pi/K$  separation up to  $\sim 1$  GeV.

To go beyond that: need faster detectors:  
 $\rightarrow$  use Cherenkov light (prompt) instead of scintillations  
 $\rightarrow$  use a fast gas detector (Multi gap RPC)

However: make sure you also know the interaction time very precisely...

Time difference between  $\pi$  and K:



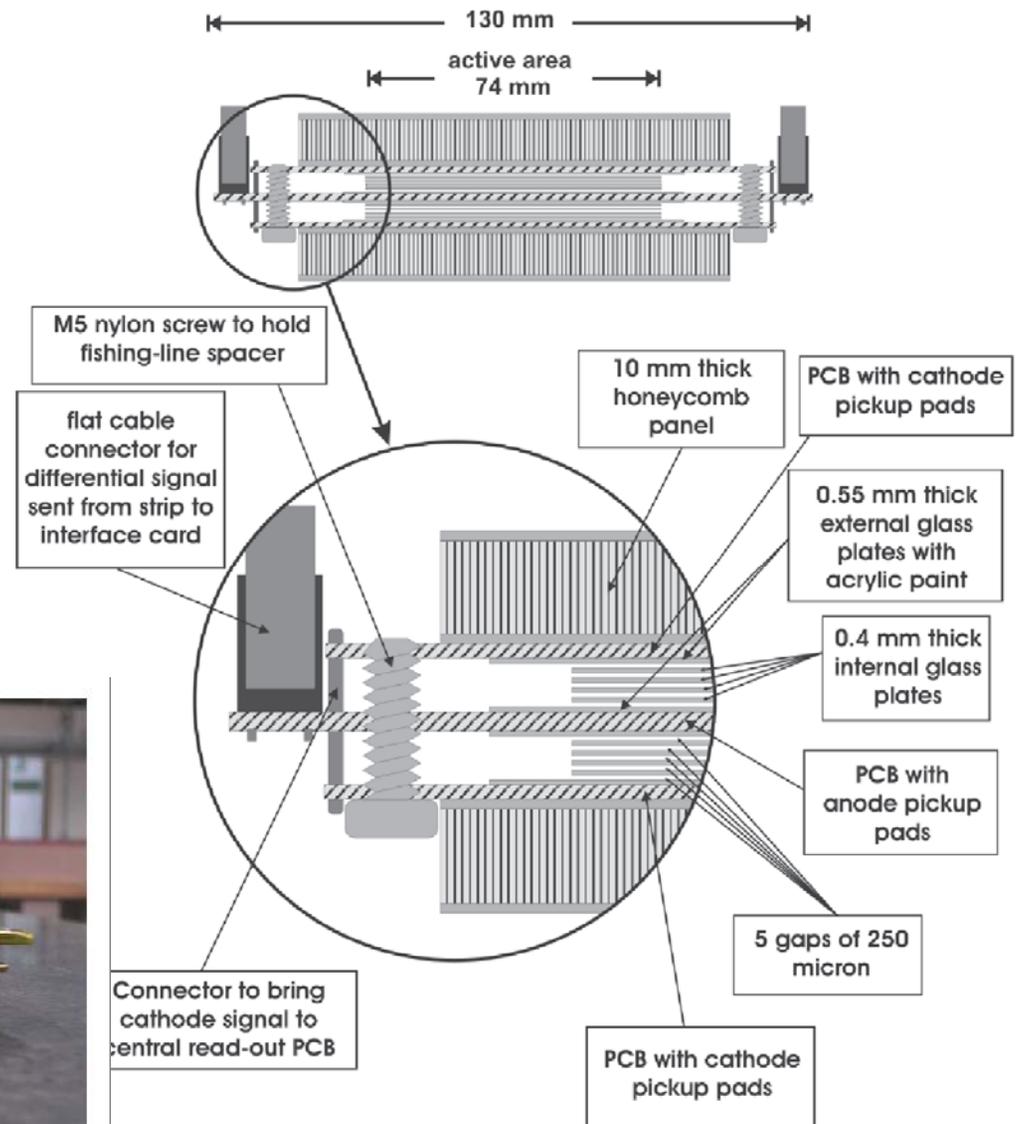
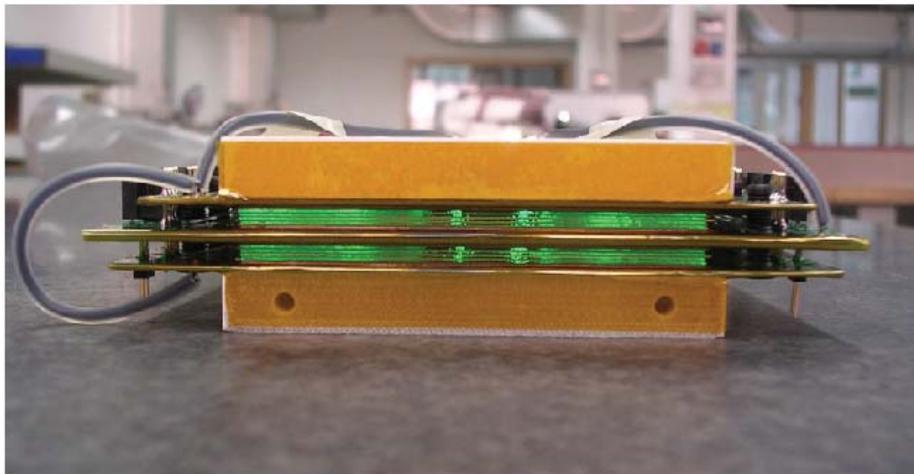
# ALICE TOF

Very fast large area ( $140\text{m}^2$ )  
particle detector:

→ MRPC, multi-gap RPC

$\sigma = 50\text{ps}$  (incl. read-out)

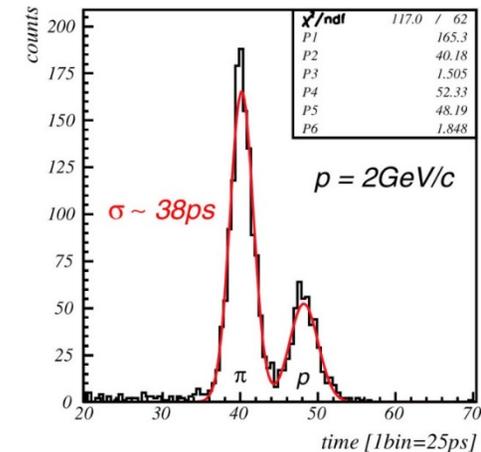
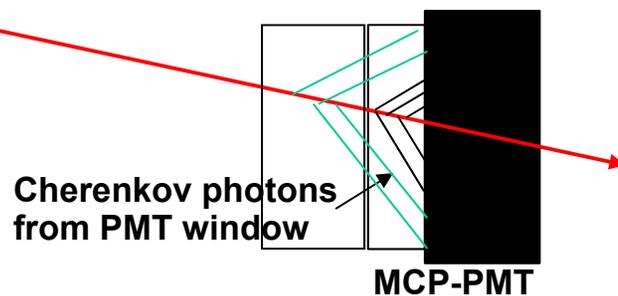
$\pi/K$  separation ( $3\sigma$ ) up to  $2.5$   
 $\text{GeV}/c$  at large track densities



# TOF with Cherenkov light

Idea: detect Cherenkov light with a very fast photon detector (MCP PMT).

Cherenkov light is produced in a quartz plate in front of the MCP PMT and in the PMT window.



Proof of principle: beamt test with pions and protons at 2 GeV/c.

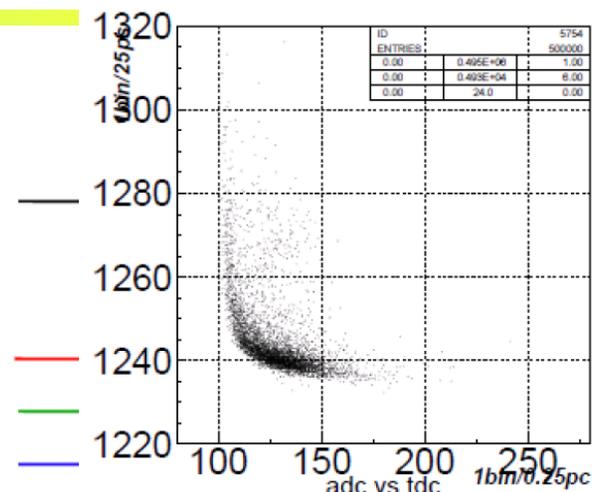
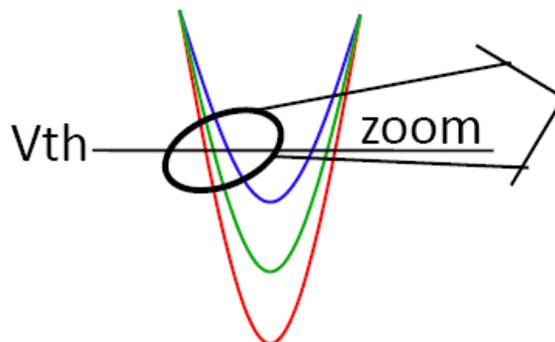
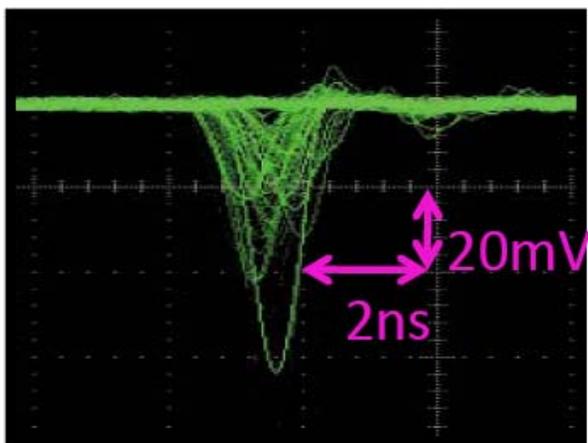
→ resolution  $\sim 5ps$  measured

- K. Inami NIMA 560 (2006) 303
- J. Va'vra NIMA 595 (2008) 270

Open issues:

- read-out
- start time

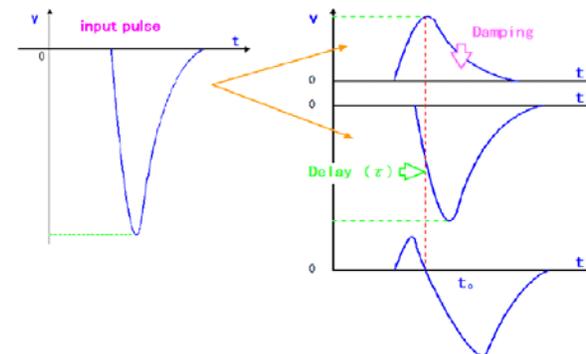
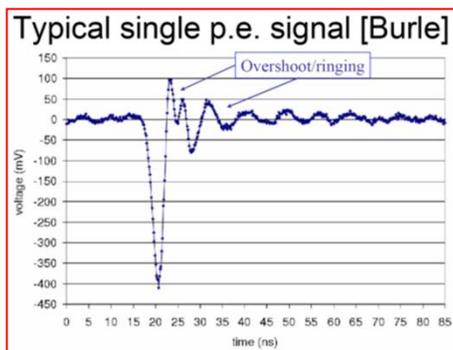
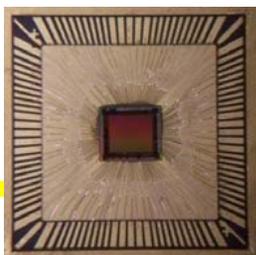
# Read out: for precise timing mitigate time walk



Variation of time determined with a leading edge discriminator: **smaller pulses give a delayed signal.** → Has to be corrected!

- Measure both time (TDC) and amplitude (ADC), correct time of arrival by using a  $\Delta T(\text{ADC})$  correction
- Use constant fraction discriminator (CFD) →
- Wave-form sampling →

e.g. Labrador 3,  
G. Varner, U Hawaii



# Identification of charged particles

---

Particles (e,  $\mu$ ,  $\pi$ , K, p) in the final state are identified by their **mass** or by the **way they interact**.

Determination of **mass**: from the relation between momentum and velocity,  $p = \gamma m v$  (p is known - radius of curvature in magnetic field)

→ Measure velocity by:

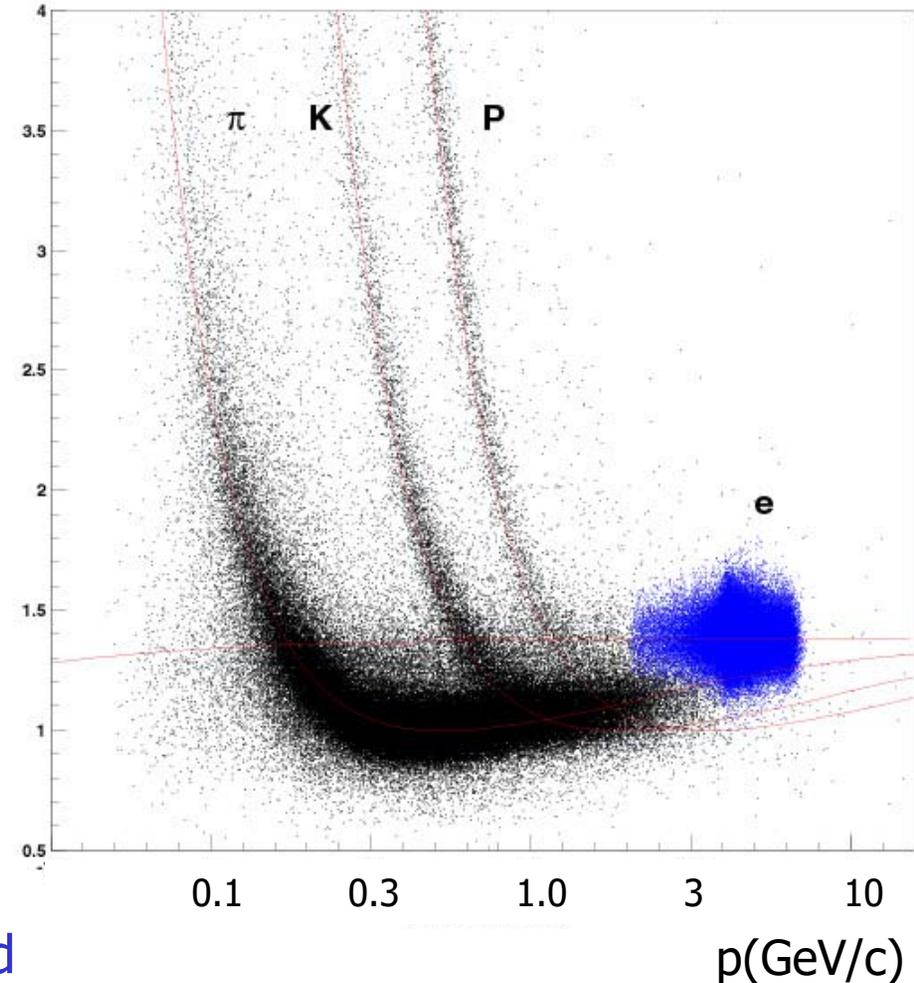
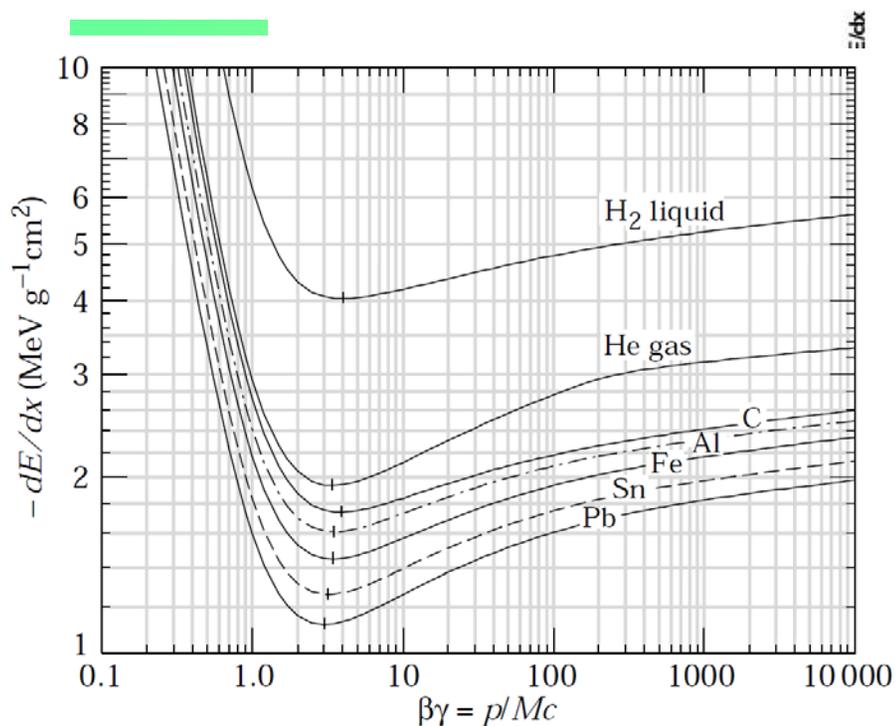
- time of flight
- ionisation losses  $dE/dx$
- Cherenkov photon angle (and/or yield)
- transition radiation

Mainly used for the identification of hadrons.

Identification through **interaction**: electrons and muons

- muon systems
- calorimeters

# Identification with the $dE/dx$ measurement



$dE/dx$  is a function of velocity  $\beta$

For particles with different mass the Bethe-Bloch curve gets displaced if plotted as a function of  $p$

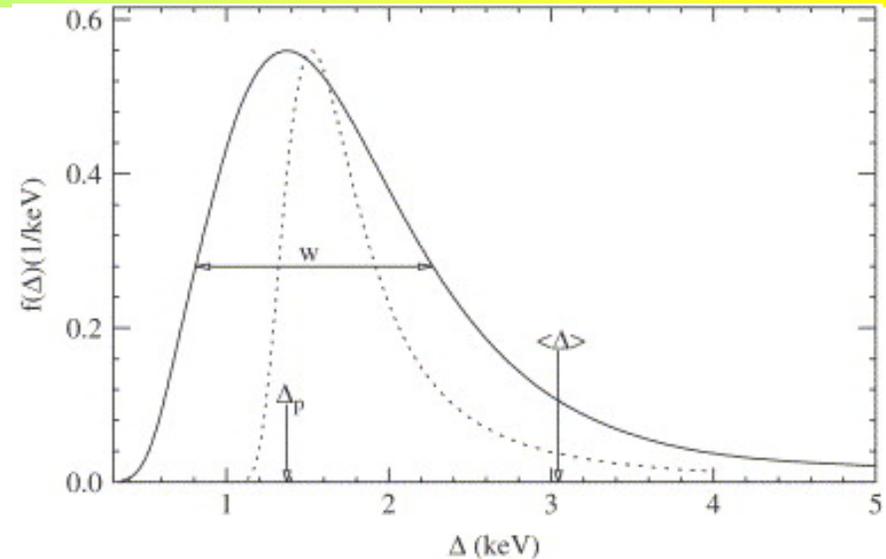
For good separation: resolution should be  $\sim 5\%$

Measure in each drift chamber layer – use truncated mean

# Identification with dE/dx measurement

Problem: long tails (not Gaussian!)

Energy loss distribution for particles with  $\beta\gamma=3.6$  traversing 1.2 cm of Ar gas (solid line).



Parameters describing  $f(\Delta)$  are

$\Delta_p(x; \beta\gamma)$ : the most probable energy loss = the position of the maximum at 1371 eV, and

$W$ : the full-width-at-half-maximum (FWHM) of 1463 eV. The mean energy loss is 3044 eV.

Dotted line: the original Landau function.

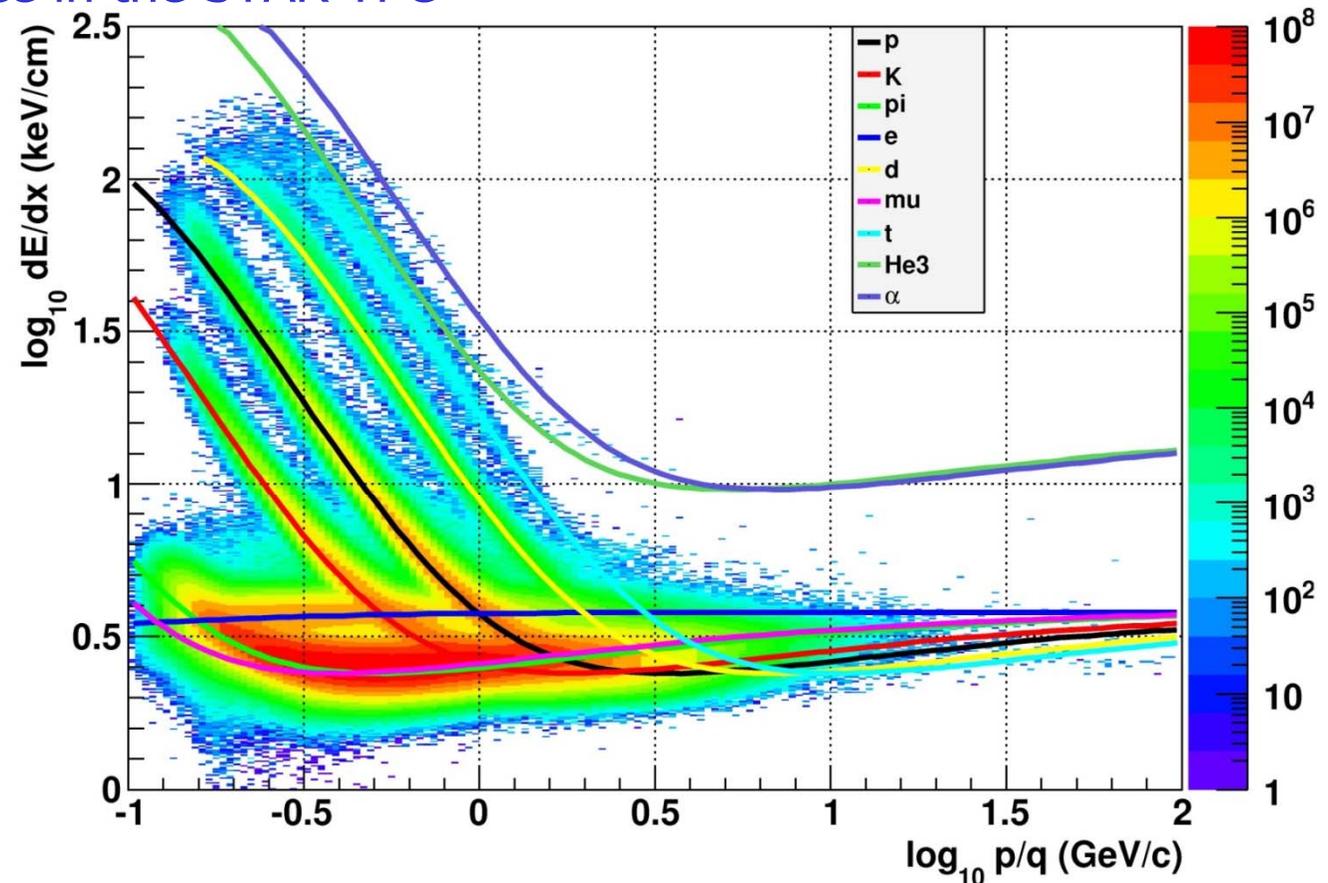
→ Many samples along the track ( $\sim 100$  in ALICE TPC), remove the largest  $\sim 40\%$  values (reduce the influence of the long tail) → truncated mean

→ Hans Bichsel: A method to improve tracking and particle identification in TPCs and silicon detectors, NIM A562 (2006) 154

# Identification with dE/dx measurement

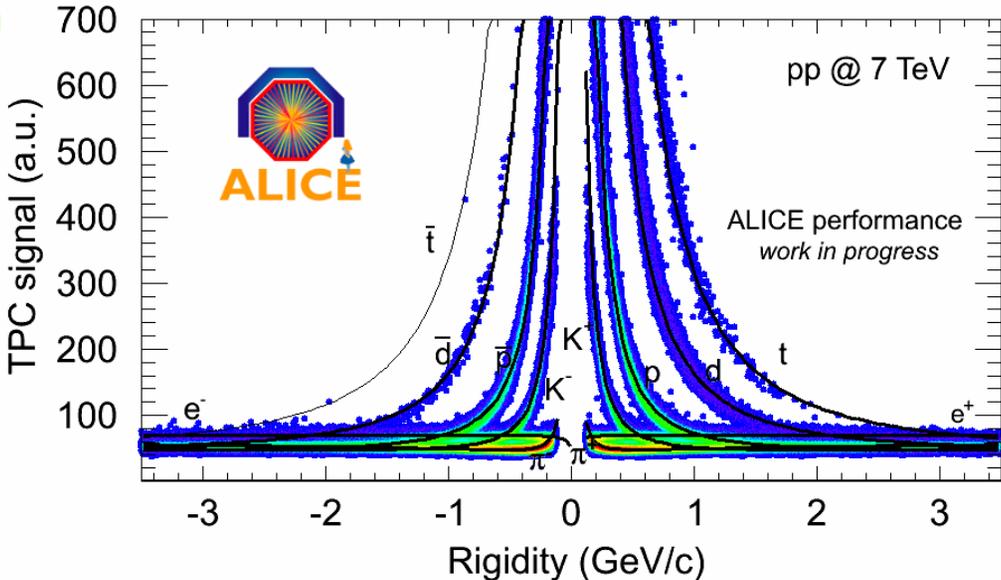
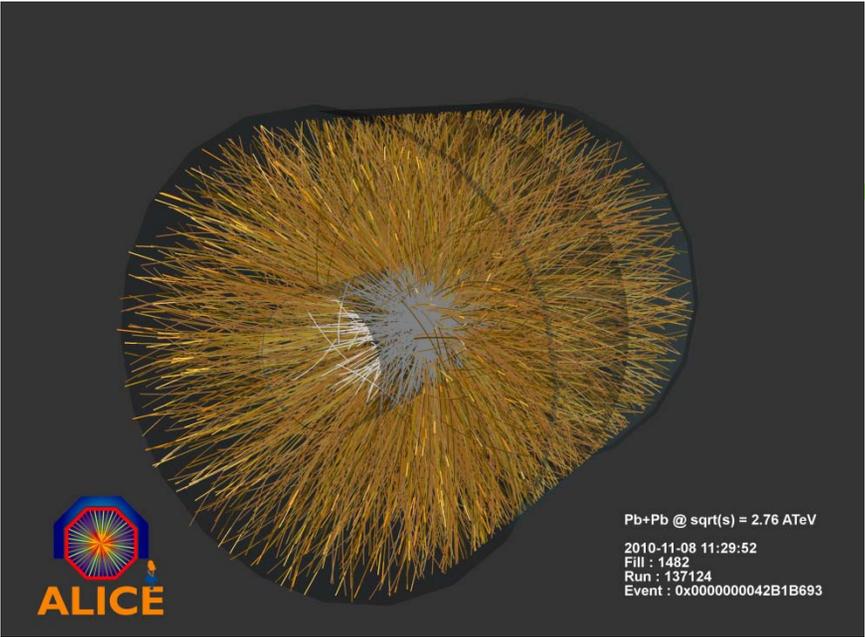
## dE/dx performance in the STAR TPC

gold-gold  
collisions

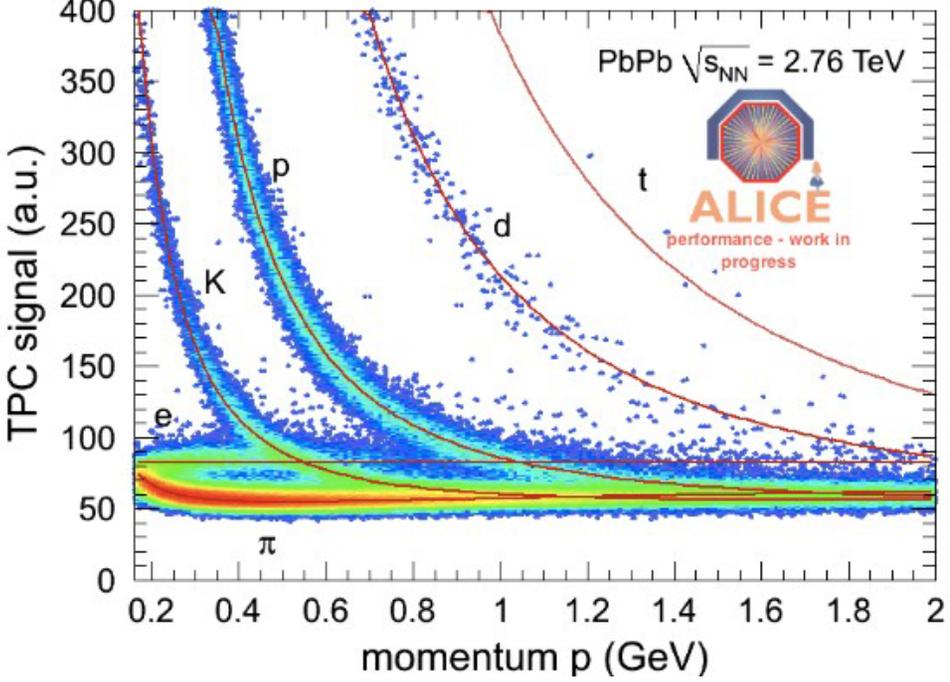


Energy loss in the STAR TPC: truncated mean as a function of momentum. The curves are Bichsel model predictions.

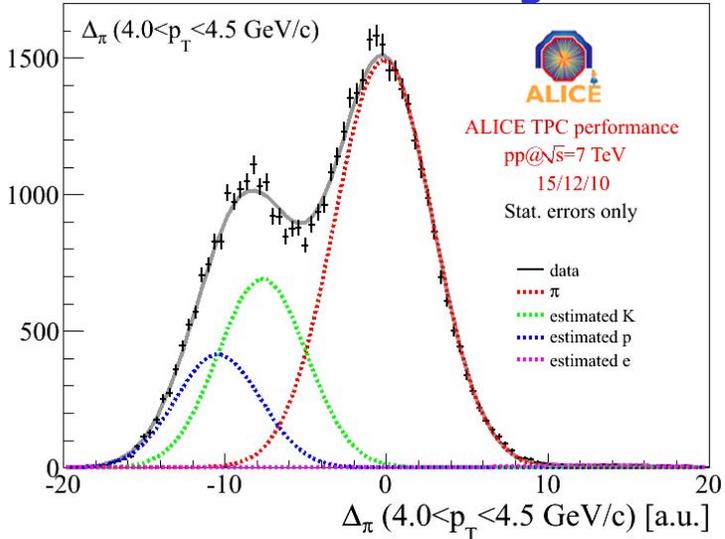
# dE/dx in ALICE



## TPC dE/dx



## relativistic rise region



# Identification of charged particles

---

Particles ( $e$ ,  $\mu$ ,  $\pi$ ,  $K$ ,  $p$ ) in the final state are identified by their **mass** or by the **way they interact**.

Determination of **mass**: from the relation between momentum and velocity,  $p = \gamma m v$  ( $p$  is known - radius of curvature in magnetic field)

→ Measure velocity by:

- time of flight
- ionisation losses  $dE/dx$
- Cherenkov photon angle (and/or yield)
- transition radiation

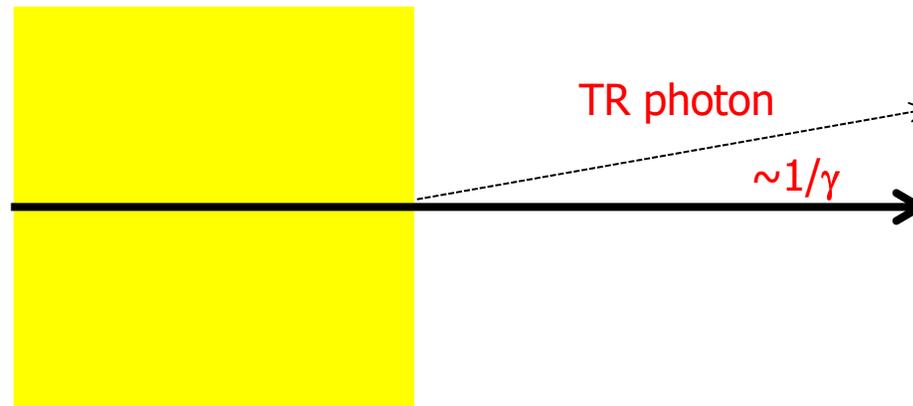
Mainly used for the identification of hadrons.

Identification through **interaction**: electrons and muons

- muon systems
- calorimeters

# Transition radiation

E.M. radiation emitted by a charged particle at the boundary of two media with different refractive indices



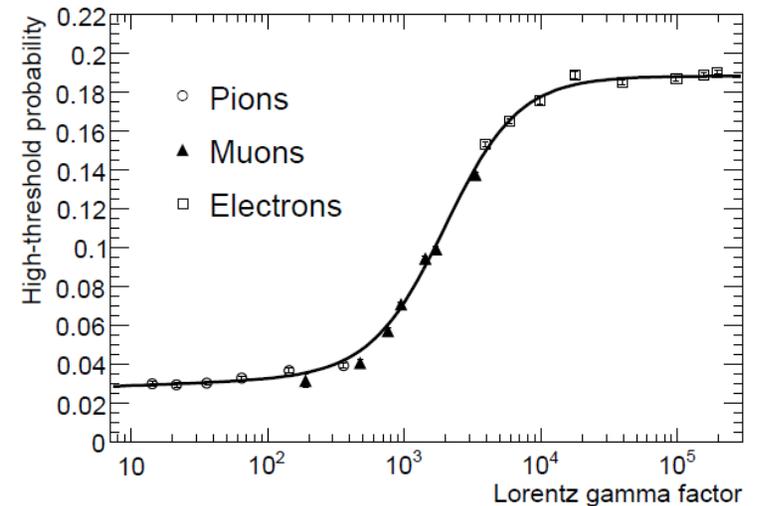
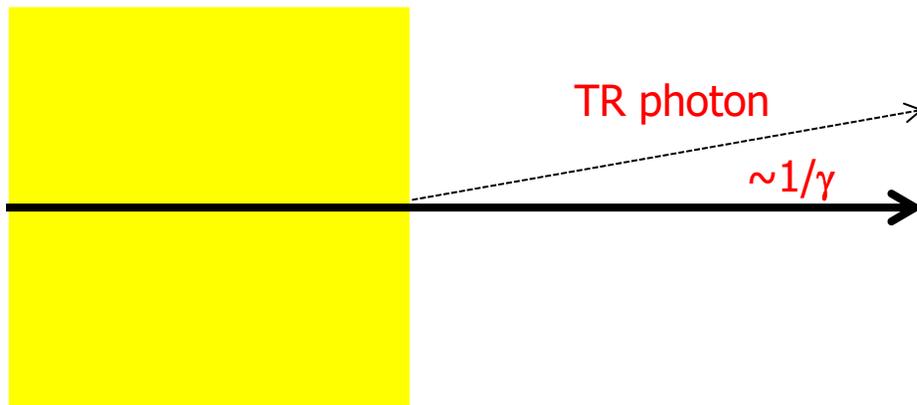
- Accelerated particle emits E.M. radiation
- Transition radiation: particle has a constant velocity, but the phase velocity of the medium changes abruptly at the boundary → radiation

→B. Dolgoshein, NIM A326 (1993) 434-469; J.D. Jackson, Classical Electrodynamics.

→H. Kolanoski, N. Wermes, Teilchendetektoren, Springer.

# Transition radiation

E.M. radiation emitted by a charged particle at the boundary of two media with different refractive indices



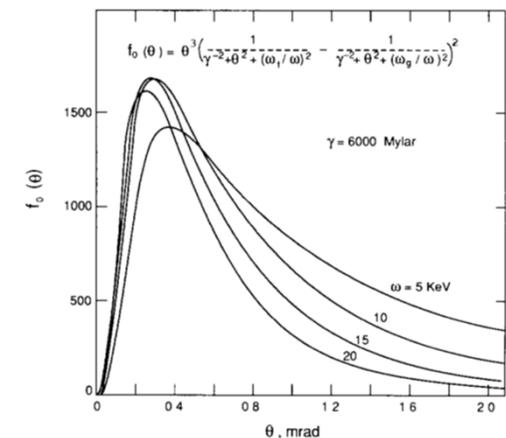
Emission rate depends on  $\gamma$  (Lorentz factor): becomes important at  $\gamma \sim 1000$

- Electrons at 0.5 GeV
- Pions above 140 GeV

Emission probability per boundary  $\sim \alpha = 1/137$

Emission angle  $\sim 1/\gamma$

Typical photon energy:  $\sim 10$  keV  $\rightarrow$  X rays



# Transition radiation - detection

---

Emission probability per boundary  $\sim \alpha = 1/137$

→ Need many boundaries

- Stacks of thin foils or
- Porous materials – foam with many boundaries of individual 'bubbles'

Typical photon energy:  $\sim 10$  keV → X rays

→ Need a wire chamber with a high Z gas (Xe) in the gas mixture  
(=large cross section for photoeffect of X rays)

Emission angle  $\sim 1/\gamma$

→ Hits from TR photons along the charged particle direction

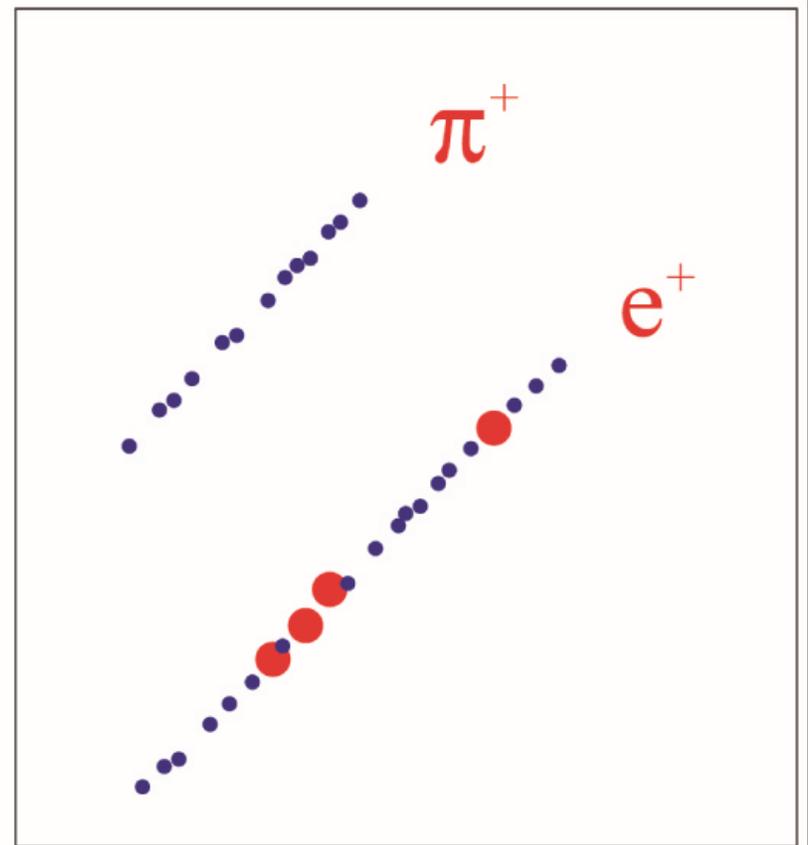
- Separation of X ray hits (high energy deposit on one place) against ionisation losses (spread out along the track)
- Two thresholds: lower for ionisation losses, higher for X ray detection

# Transition radiation - detection

- Hits from TR photons along the charged particle direction
- Separation of X ray hits (high energy deposit on one place) against ionisation losses (spread out along the track)
- Two thresholds: lower to remove noise, higher to separate X ray conversions from ionisation by charged particles

- Small circles: between low and high threshold (ionisation)
- Big circles: high threshold (X ray detection)

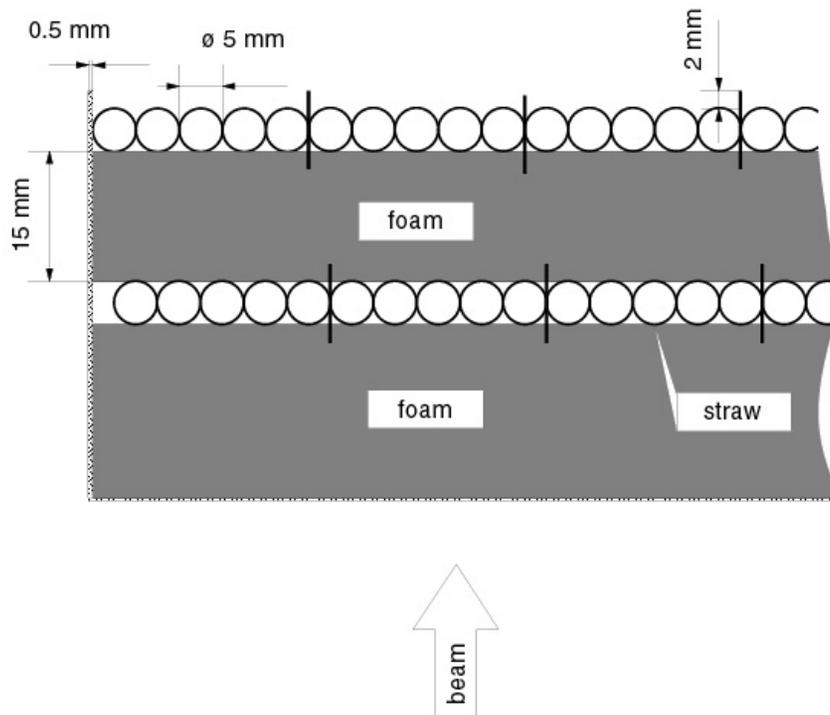
(pion below the TR threshold, e above the TR threshold)



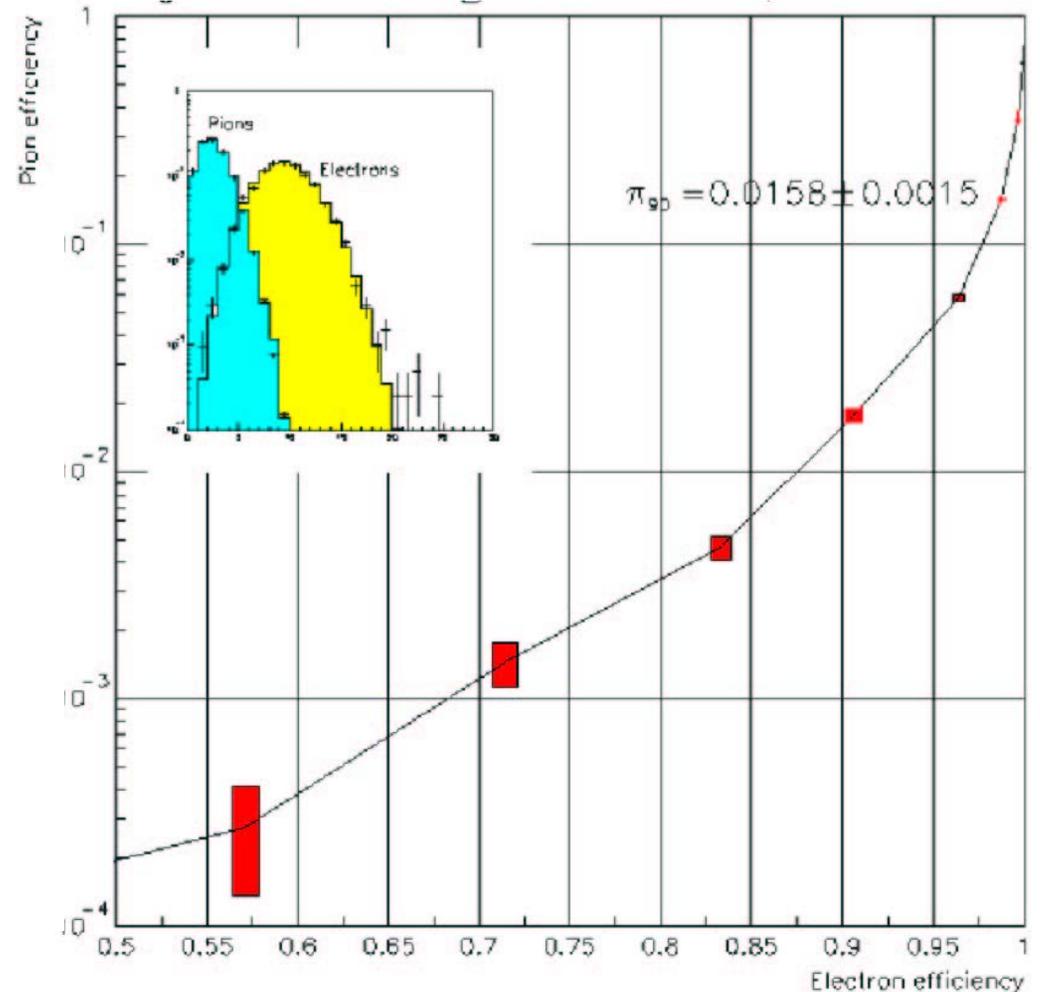
# Transition radiation detectors

Example:

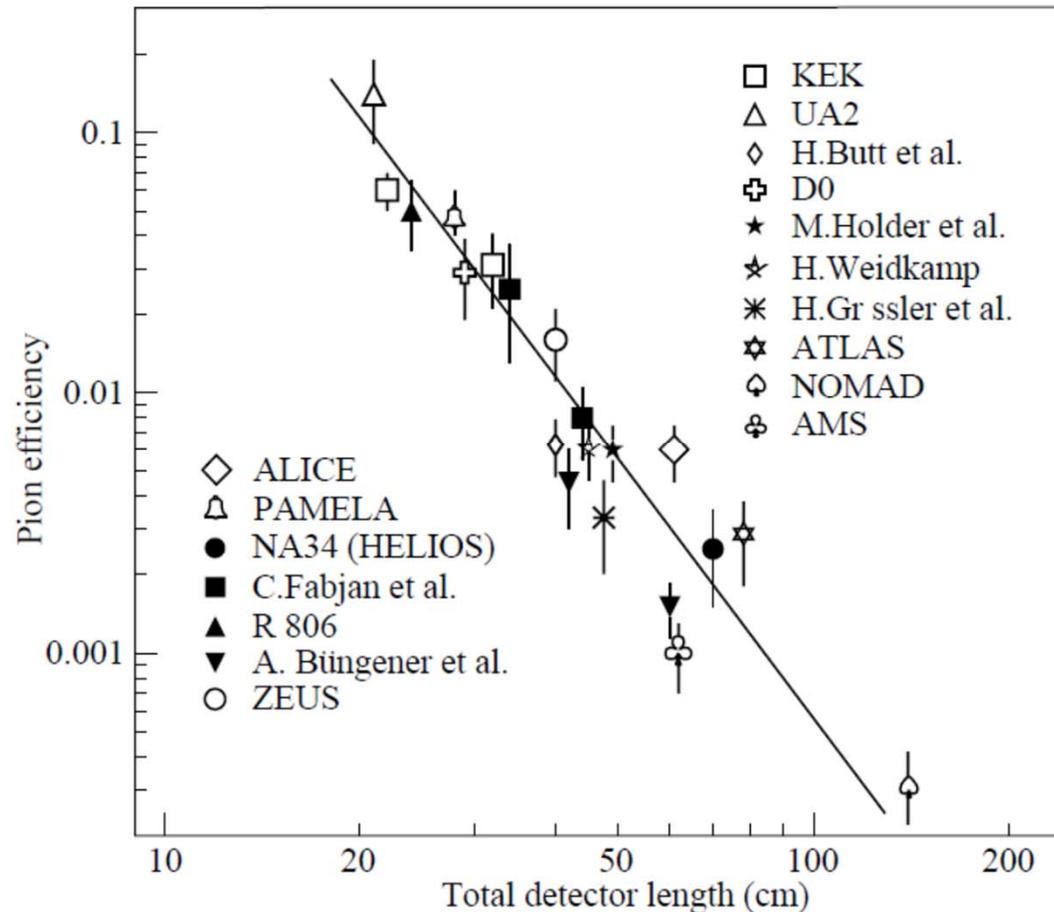
Radiator: organic foam  
between the detector  
tubes (straws made of  
capton foil)



Performance: pion efficiency (fake prob.)  
vs electron efficiency



# Transition radiation detectors - performance



Performance: pion efficiency (fake prob.) vs detector length

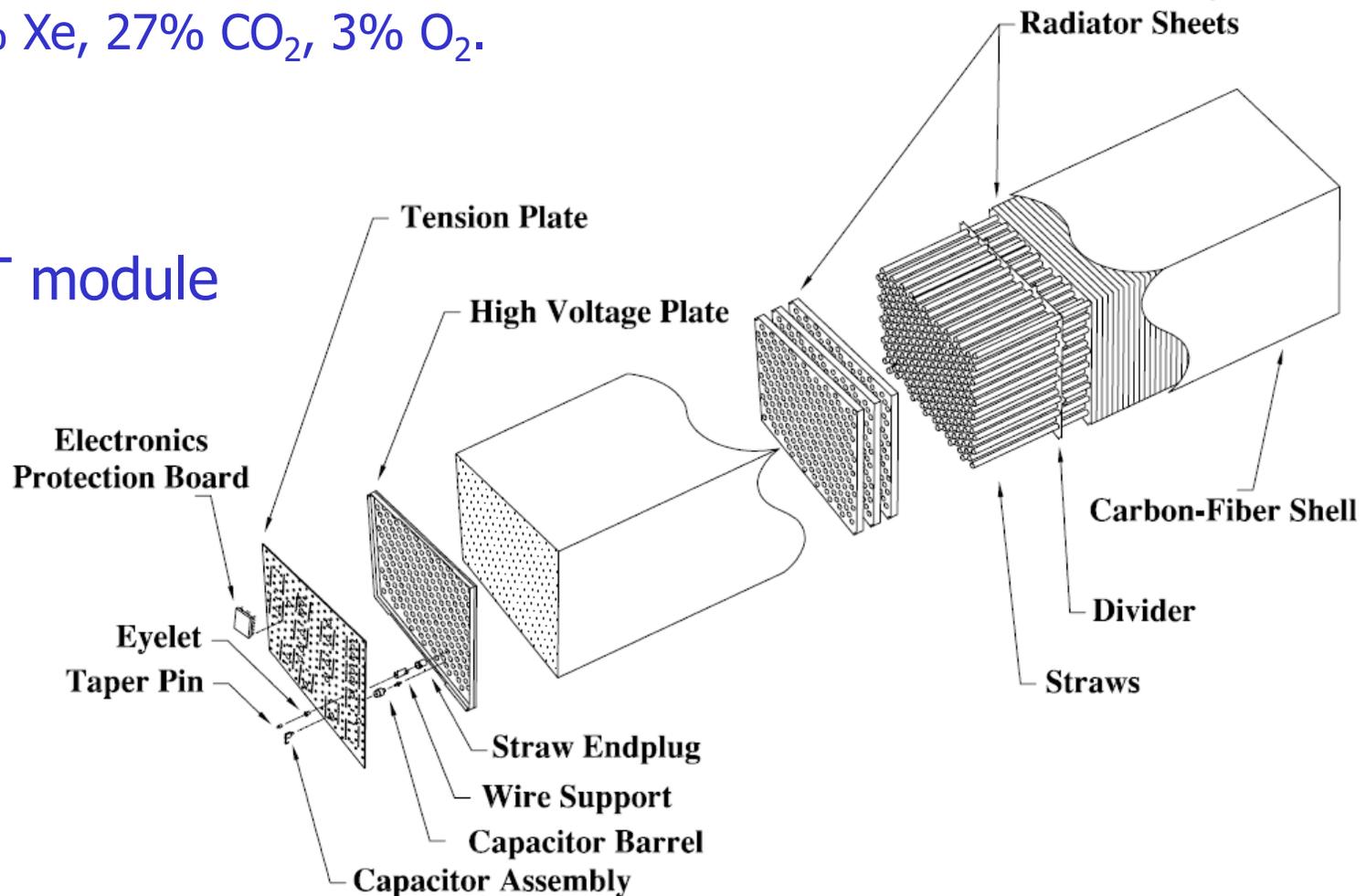


# ATLAS TRT

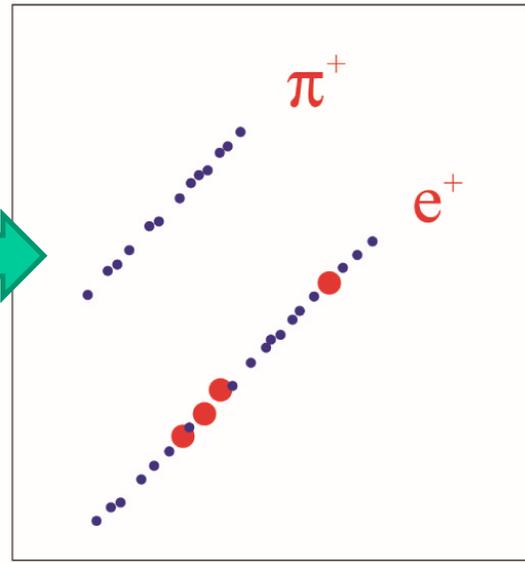
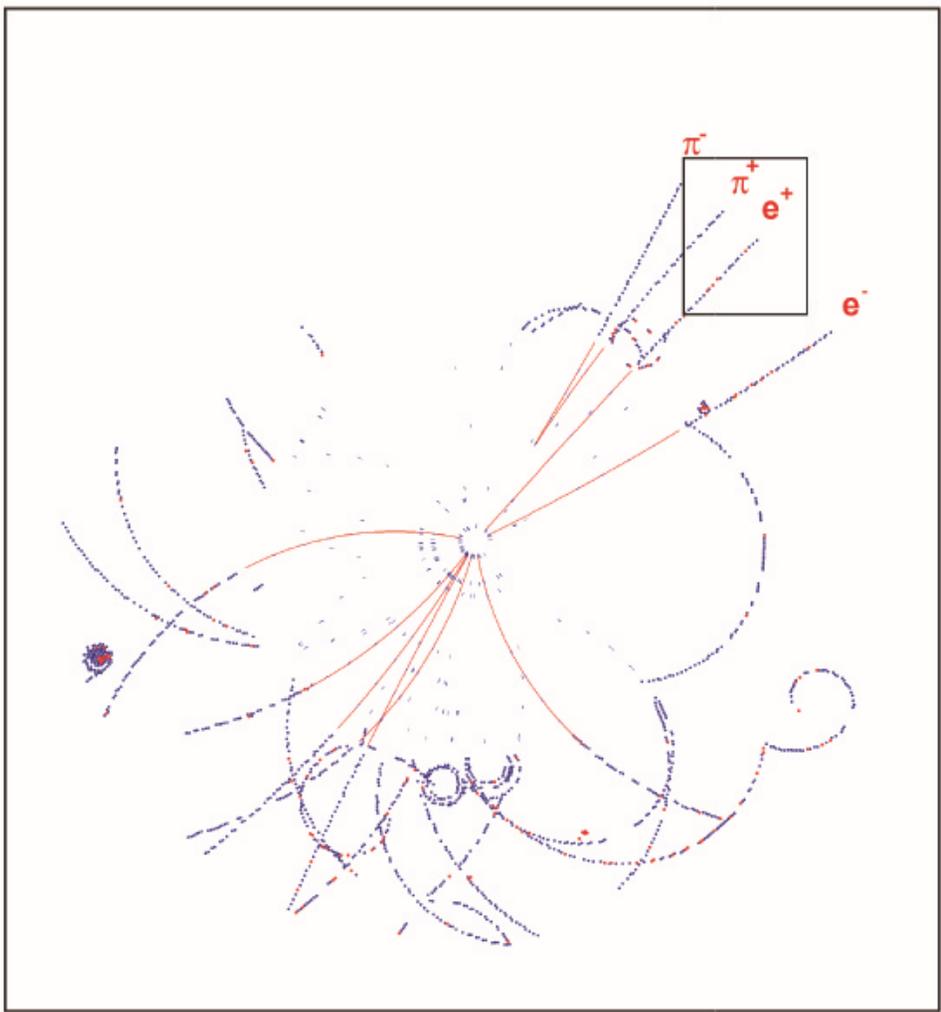
Radiator: 3mm thick layers made of polypropylene-polyethylene fibers with  $\sim 19$  micron diameter, density:  $0.06 \text{ g/cm}^3$

Straw tubes: 4mm diameter with 31 micron diameter anode wires, gas: 70% Xe, 27% CO<sub>2</sub>, 3% O<sub>2</sub>.

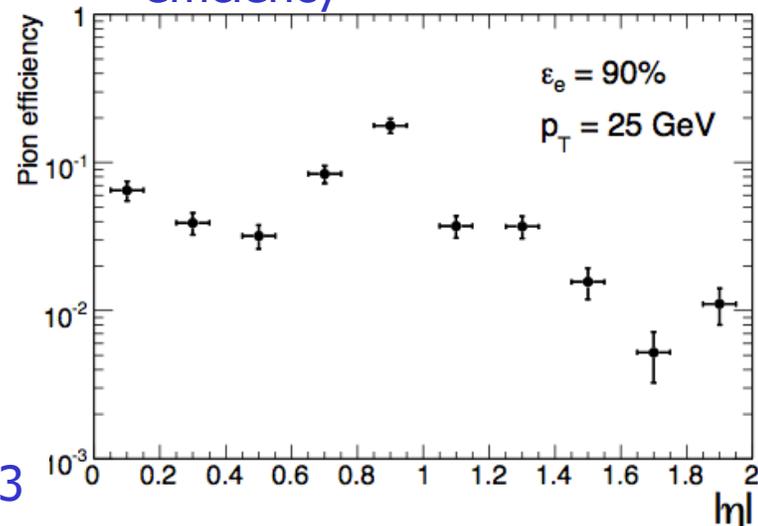
## TRT module



# TRT: pion-electron separation



$\pi$  fake probability at 90% e efficiency



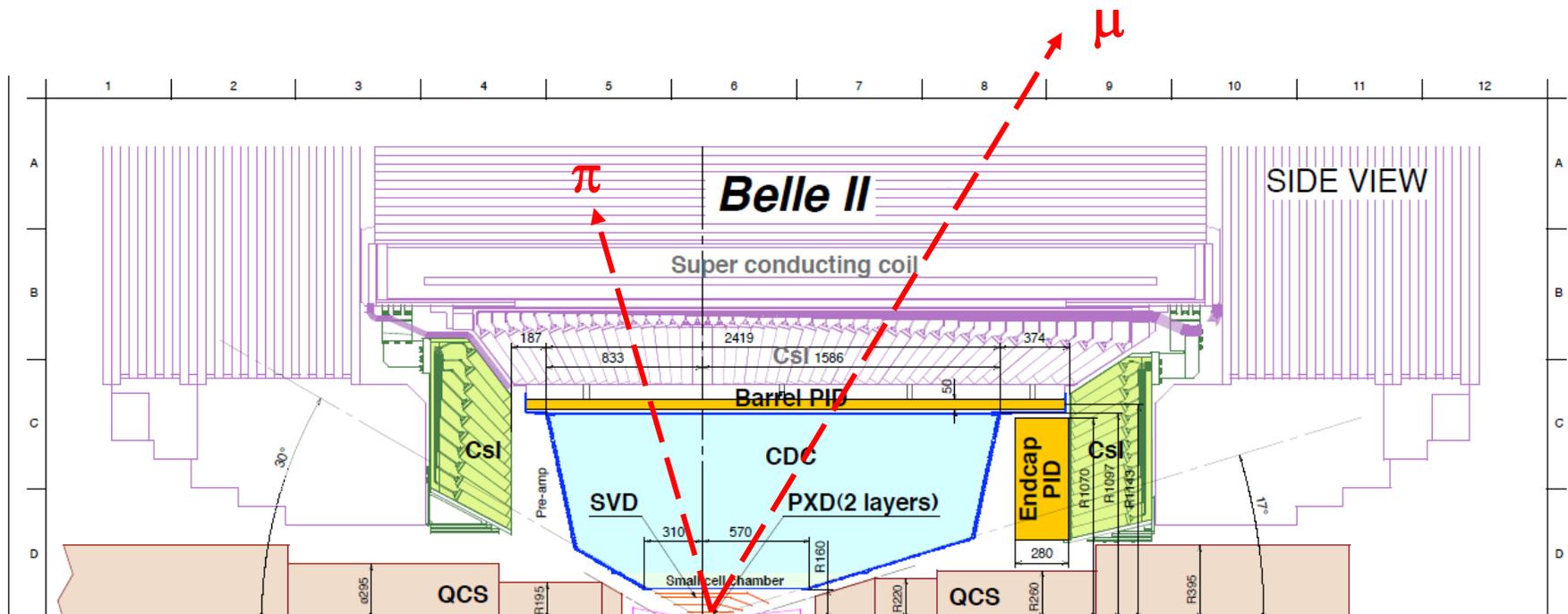
# Identification of muons

## Separate muons from hadrons (pions and kaons):

Exploit the fact that muons interact only electromag., while hadrons interact strongly

→ need a few interaction lengths to stop hadrons (interaction lengths = about 10x radiation length in iron, 20x in CsI).

→ **A particle is identified as a muon if it penetrates the material.**



# Example: muon detection at B factories

## Separate muons from hadrons (pions and kaons):

Need a few interaction lengths to stop hadrons (interaction length = about 10x radiation length in iron, 20x in CsI).

Some numbers: 0.8 interaction length (CsI) + 3.9 interaction lengths (iron)

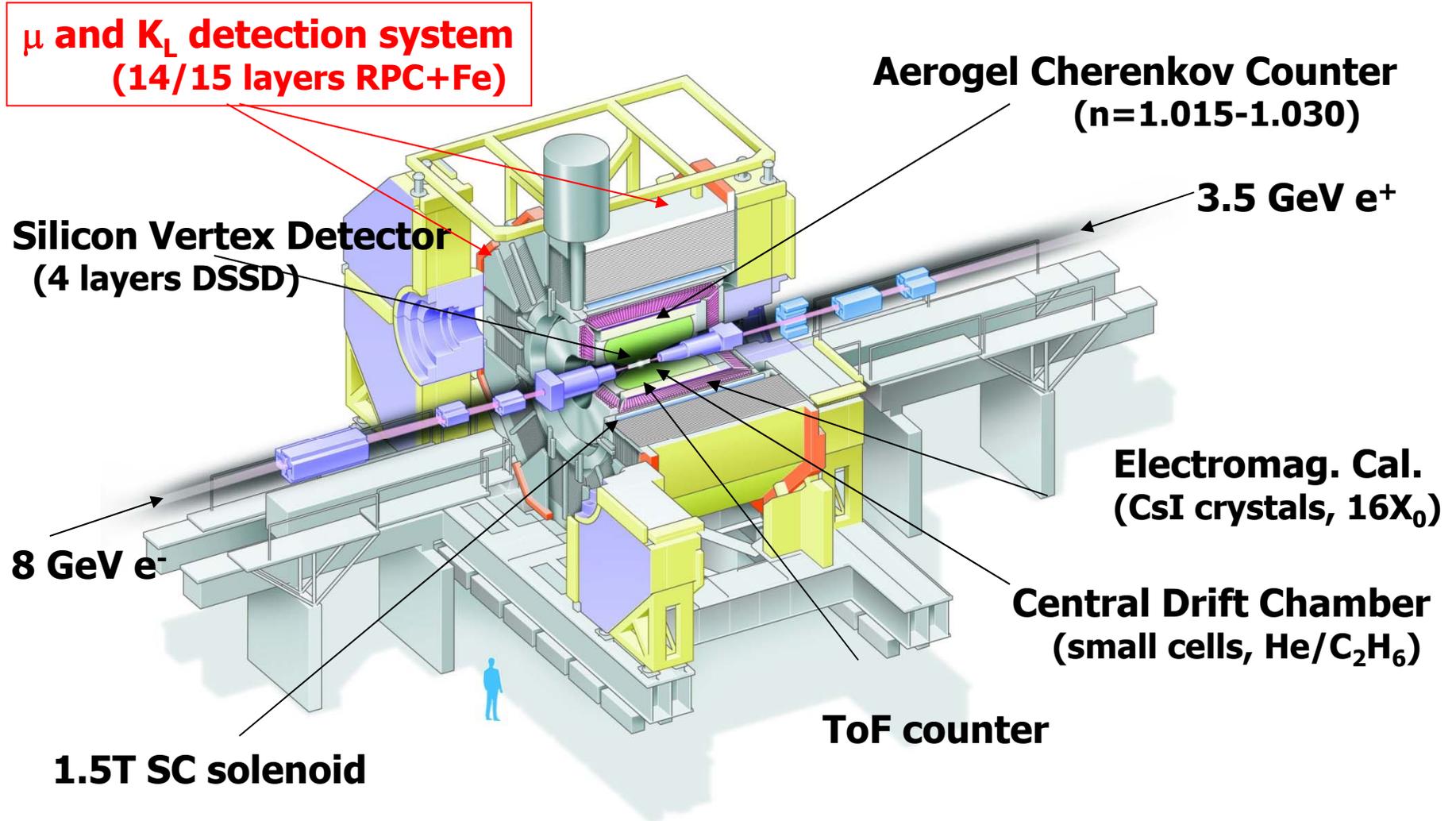
Interaction length: CsI 167 g/cm<sup>2</sup>, iron 132 g/cm<sup>2</sup>,

(dE/dx)<sub>min</sub>: CsI 1.24 MeV/(g/cm<sup>2</sup>), iron 1.45 MeV/(g/cm<sup>2</sup>)

→  $\Delta E_{\min} = (0.15+0.75) \text{ GeV} = 0.90 \text{ GeV}$  → reliable identification of muons possible around ~1 GeV

**Detect K<sub>L</sub> interaction (cluster):** again need a few interaction lengths – the same system can be used for both – bonus!

# Example: Muon and $K_L$ detection at Belle



# Muon and $K_L$ detector

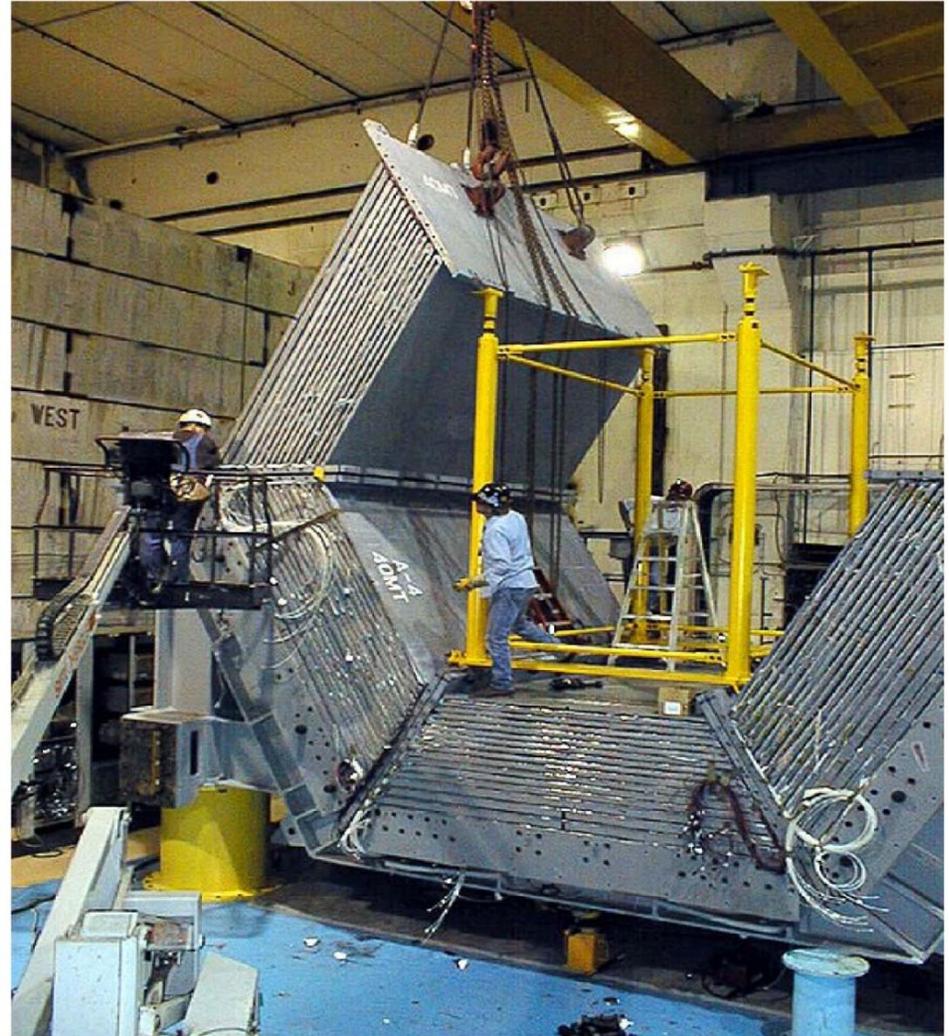
Up to 21 layers of resistive-plate chambers (RPCs) between iron plates of the flux return

Bakelite RPCs at BABAR

Glass RPCs at Belle

(glass was better choice because of ageing effects)

Scintillator strips + RPCs at Belle II



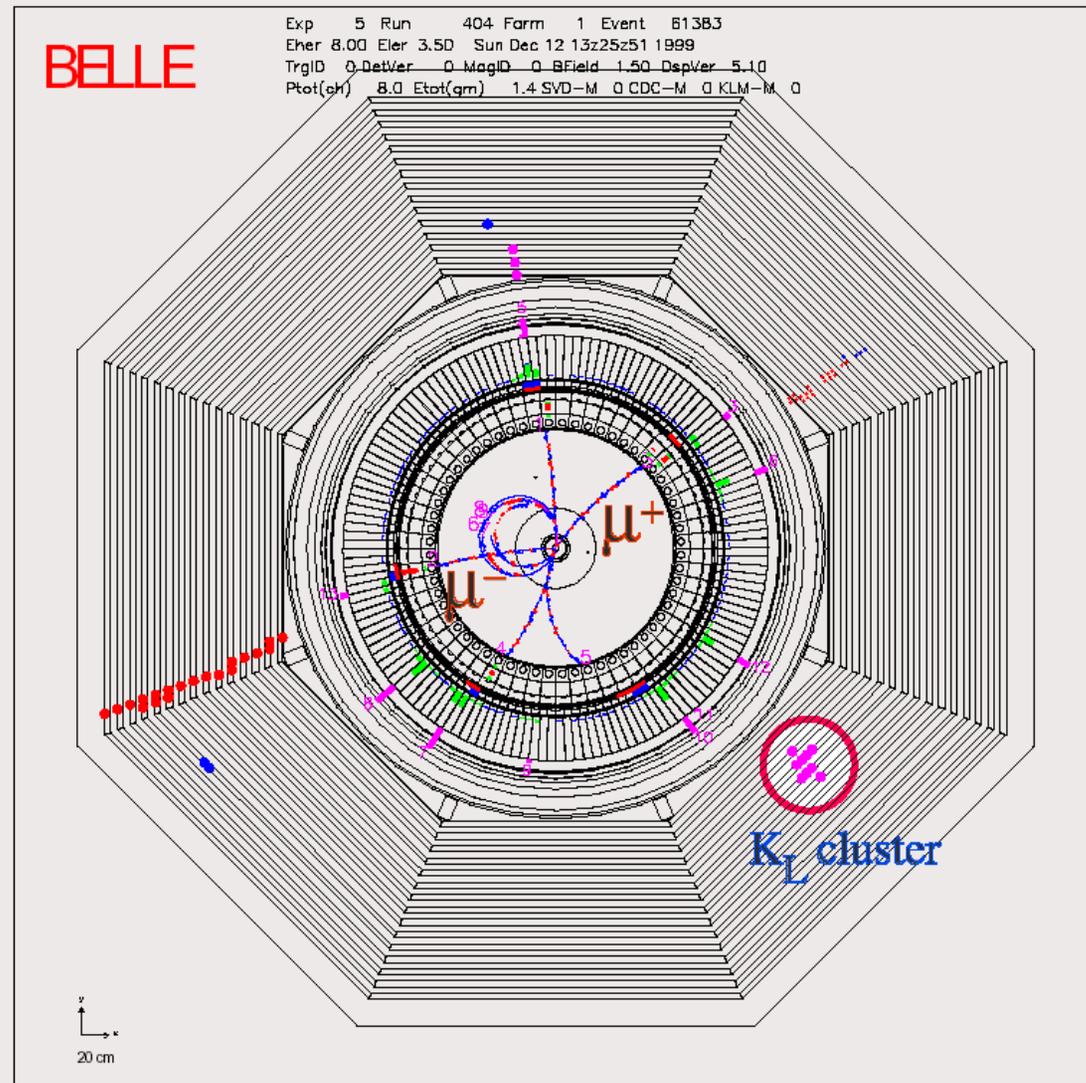
# Muon and $K_L$ detector

Example:

event with

- two muons and a
- $K_L$

and a pion that only partly penetrated



# Muon identification performance

Muon identification: efficient for  $p > 800$  MeV/c

efficiency

fake probability

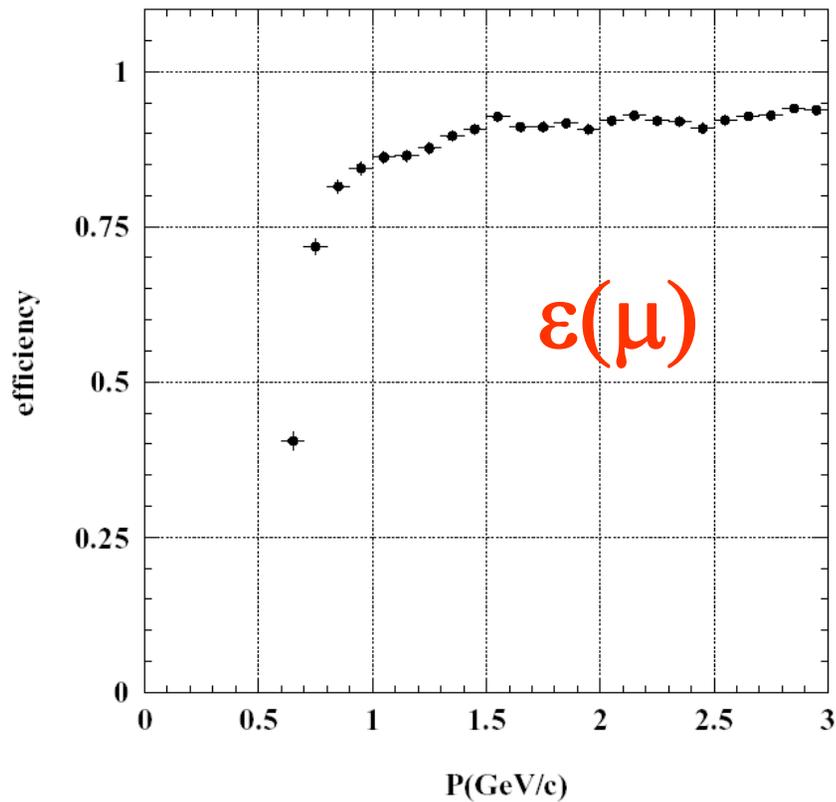


Fig. 109. Muon detection efficiency vs. momentum in KLM.

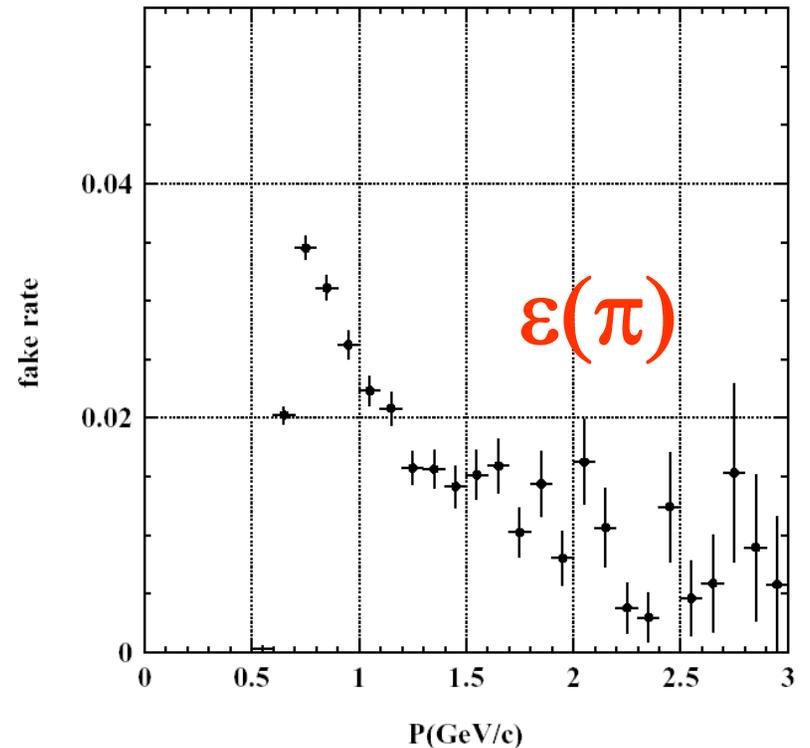
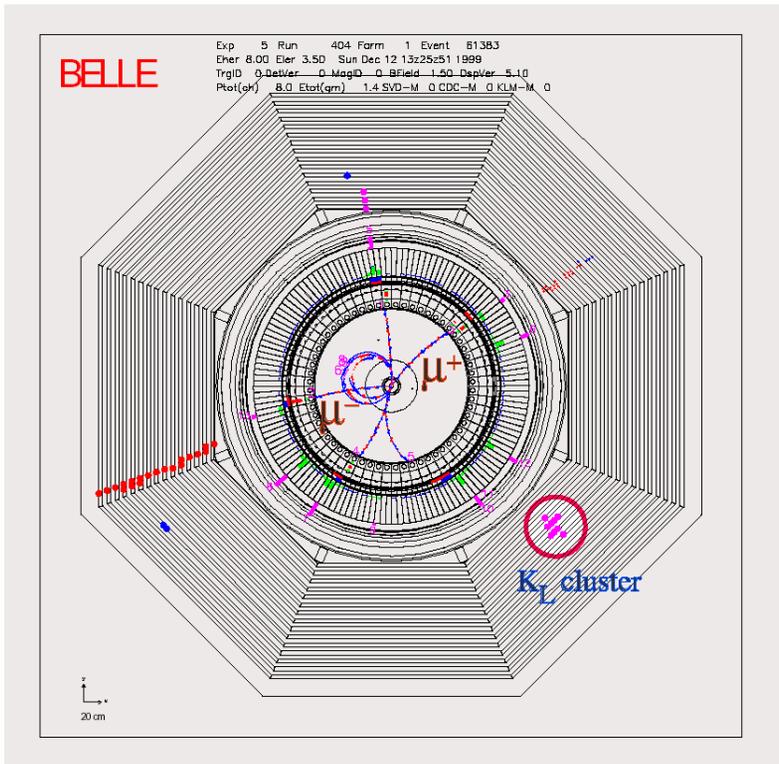


Fig. 110. Fake rate vs. momentum in KLM.

# $K_L$ detection performance



$K_L$  detection: resolution in  $K_L$  direction

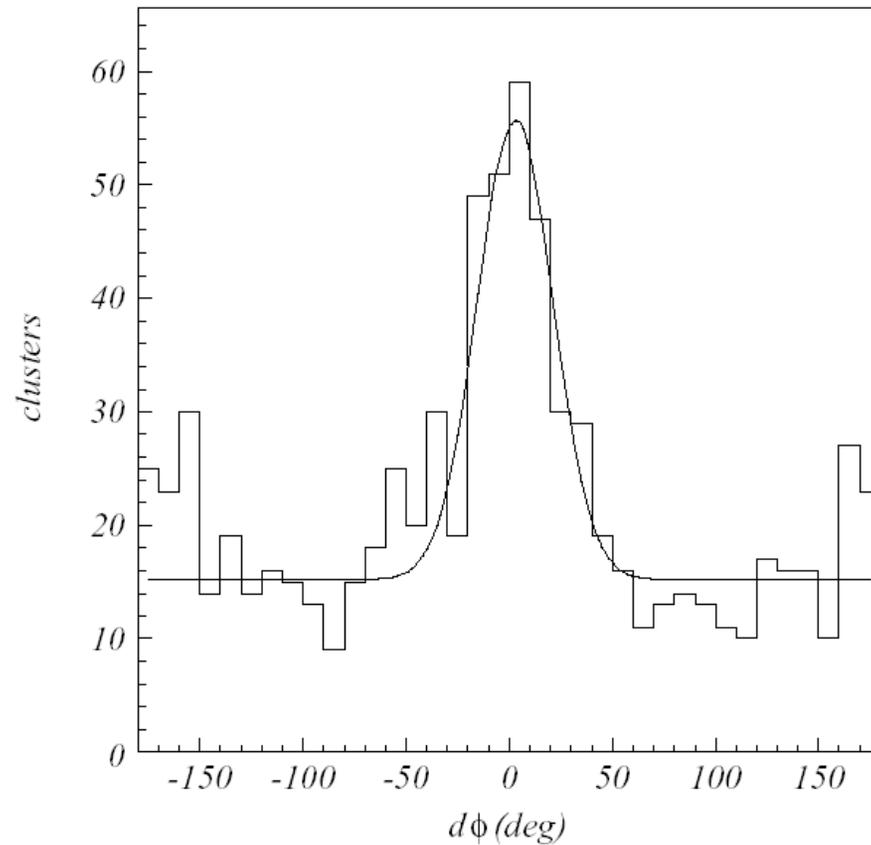
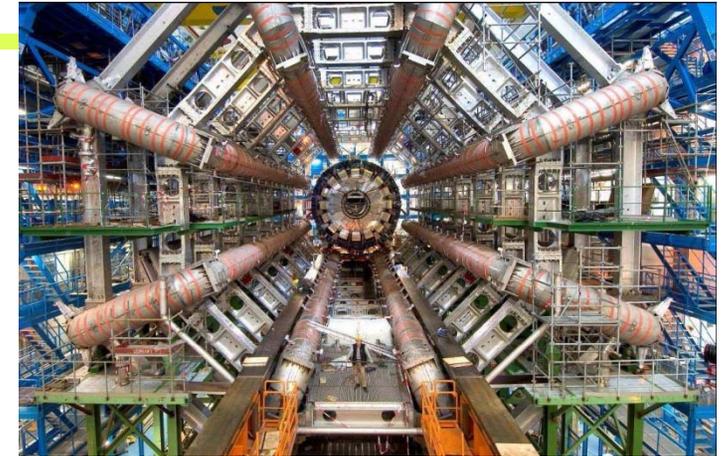
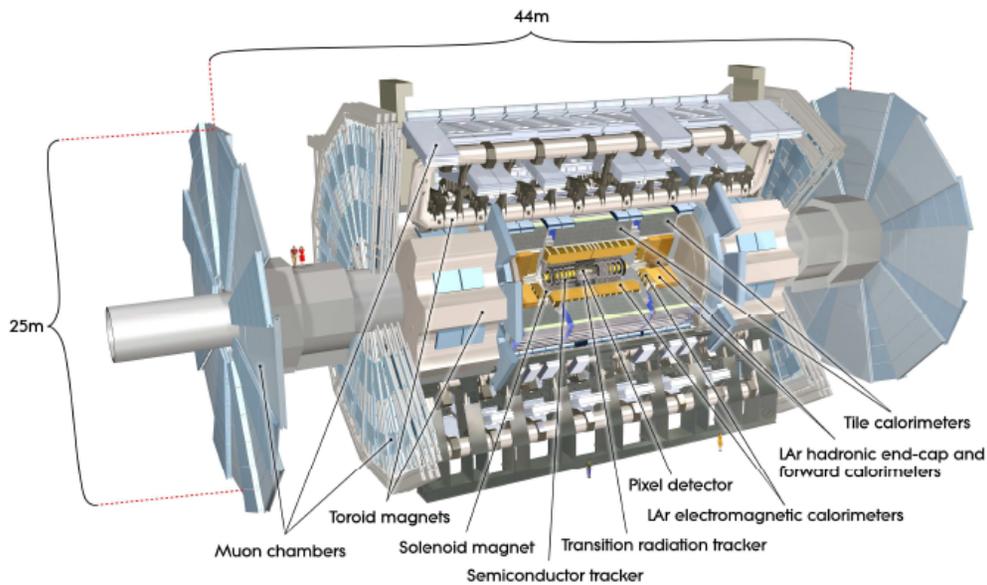
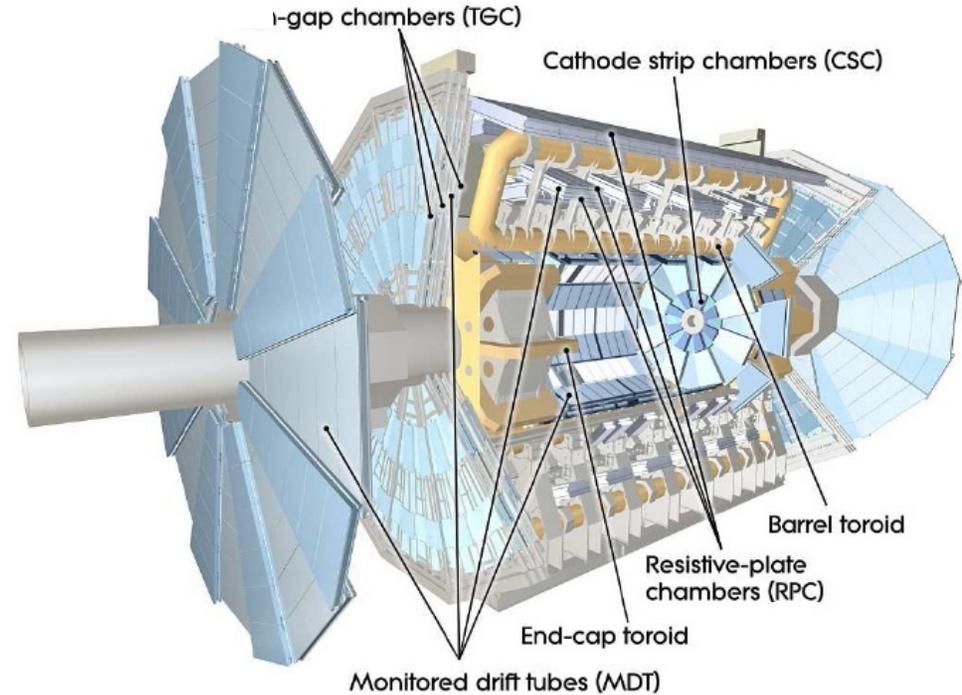


Fig. 107. Difference between the neutral cluster and the direction of missing momentum in KLM.

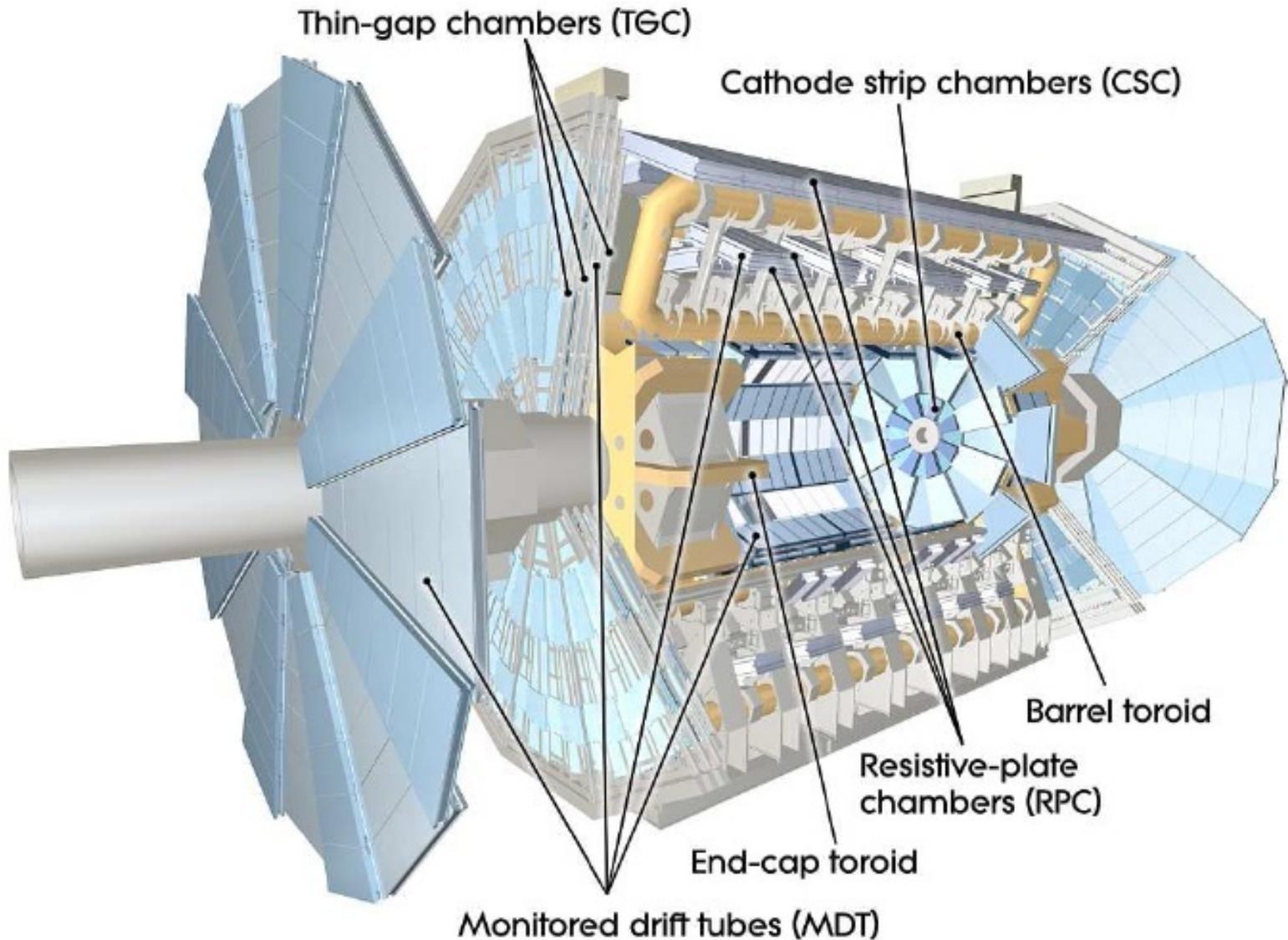
# Identification of muons in ATLAS



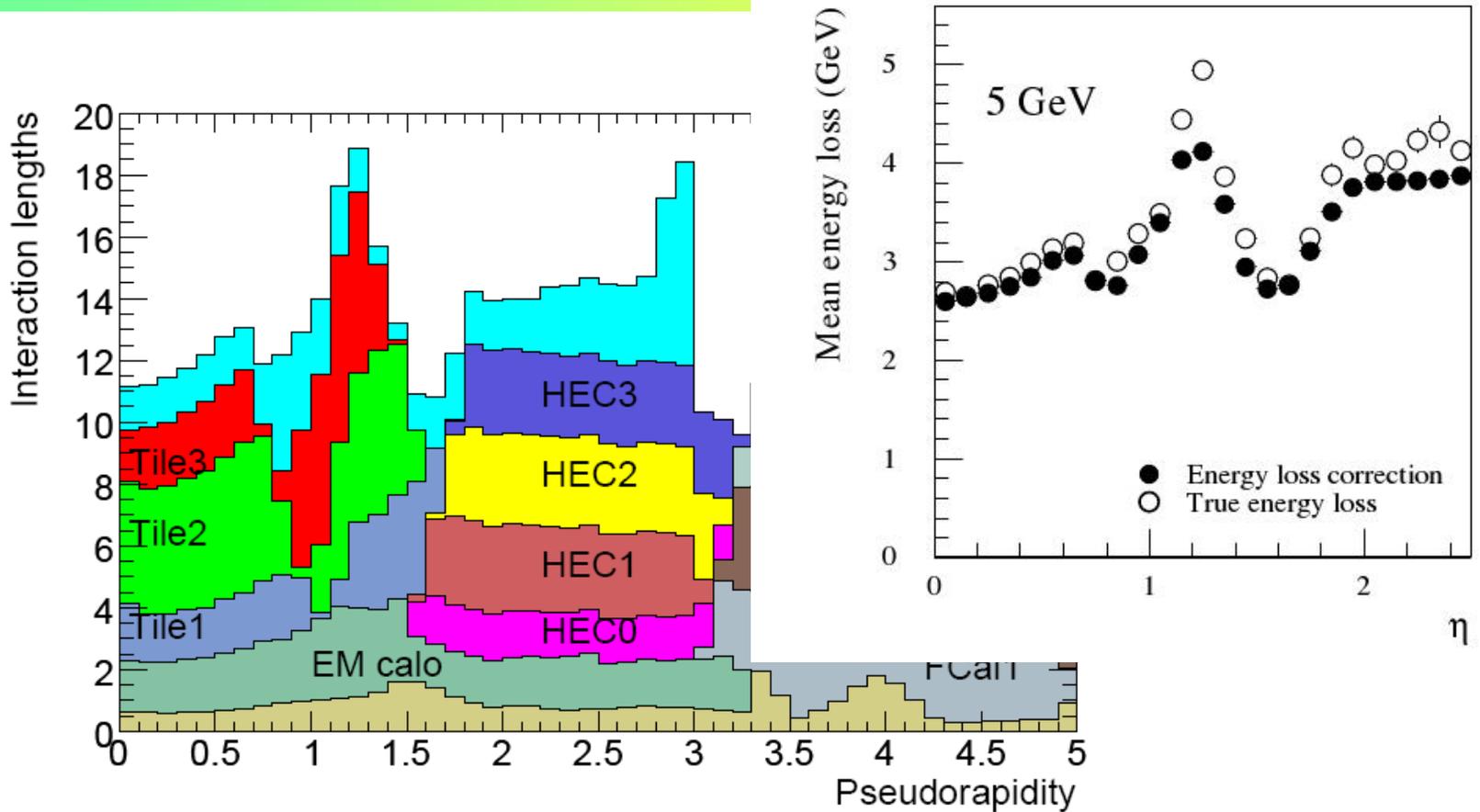
- Identify muons
- Measure their momentum



# Detection of muons in ATLAS

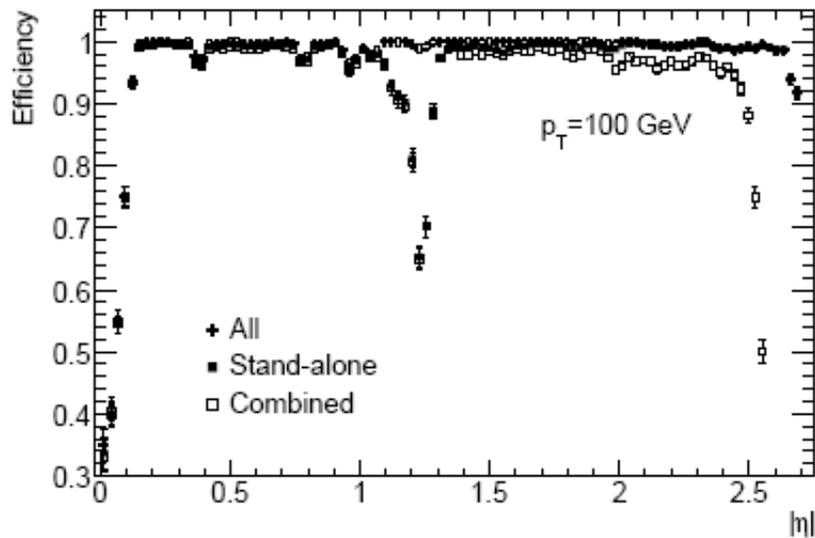


# Muon identification in ATLAS

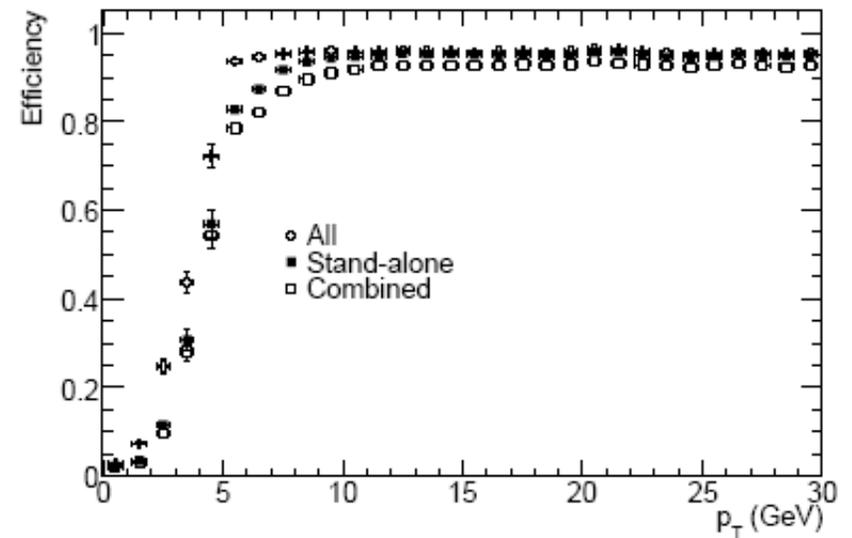


**Material in front of the muon system**

# Muon identification efficiency



Efficiency for 100 GeV muons



Efficiency vs  $p_T$

# Summary

---

Particle identification is an essential part of several experiments, and has contributed substantially to our present understanding of elementary particles and their interactions. It will continue to have an important impact in searches for new physics.

A large variety of techniques has been developed for different kinematic regions and different particles.

New concepts and detectors are being studied → this is a very active area of detector R+D.

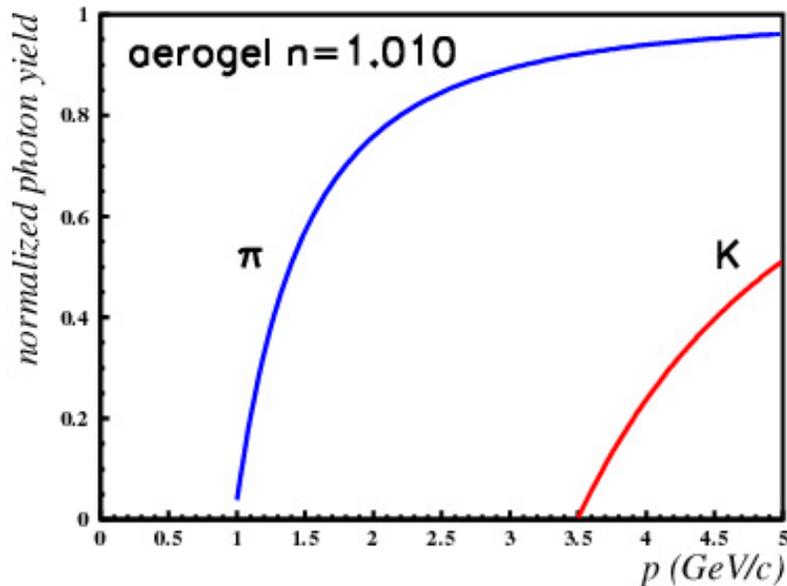
# Back-up slides

---

# Threshold Cherenkov counter

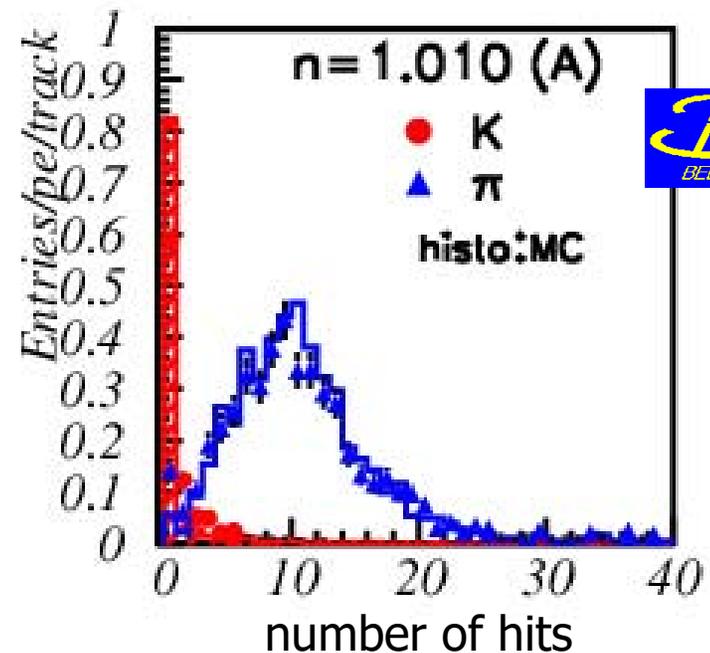
$\cos\theta = c/nv = 1/\beta n$  → Separate **K** (below threshold) from  $\pi$  (above) by properly choosing **n**

Photon yield vs p



→ Good separation between pions (light) and kaons (no light) between  $\sim 1.5$  GeV/c and 3.5 GeV/c

Choice of n: depends on the momentum range.

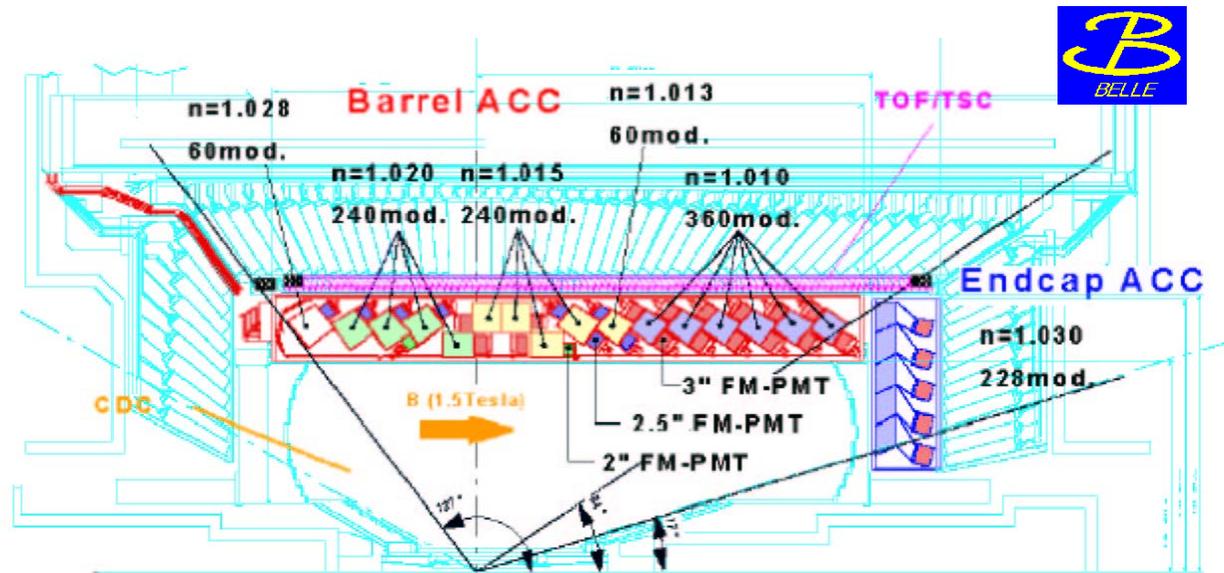
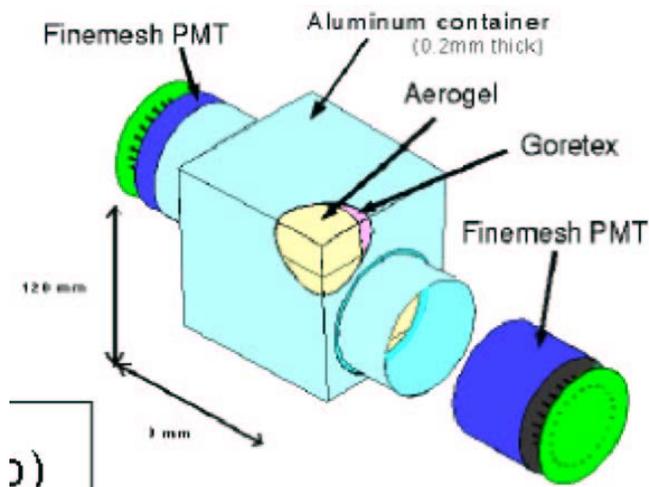


# Belle: Threshold Cherenkov counter, ACC (aerogel Cherenkov counter)

K (below threshold) vs.  $\pi$  (above) by properly choosing n for a given kinematic region

→ more energetic particles fly in the 'forward region' → lower n

Detector unit: a block of aerogel and two fine-mesh PMTs



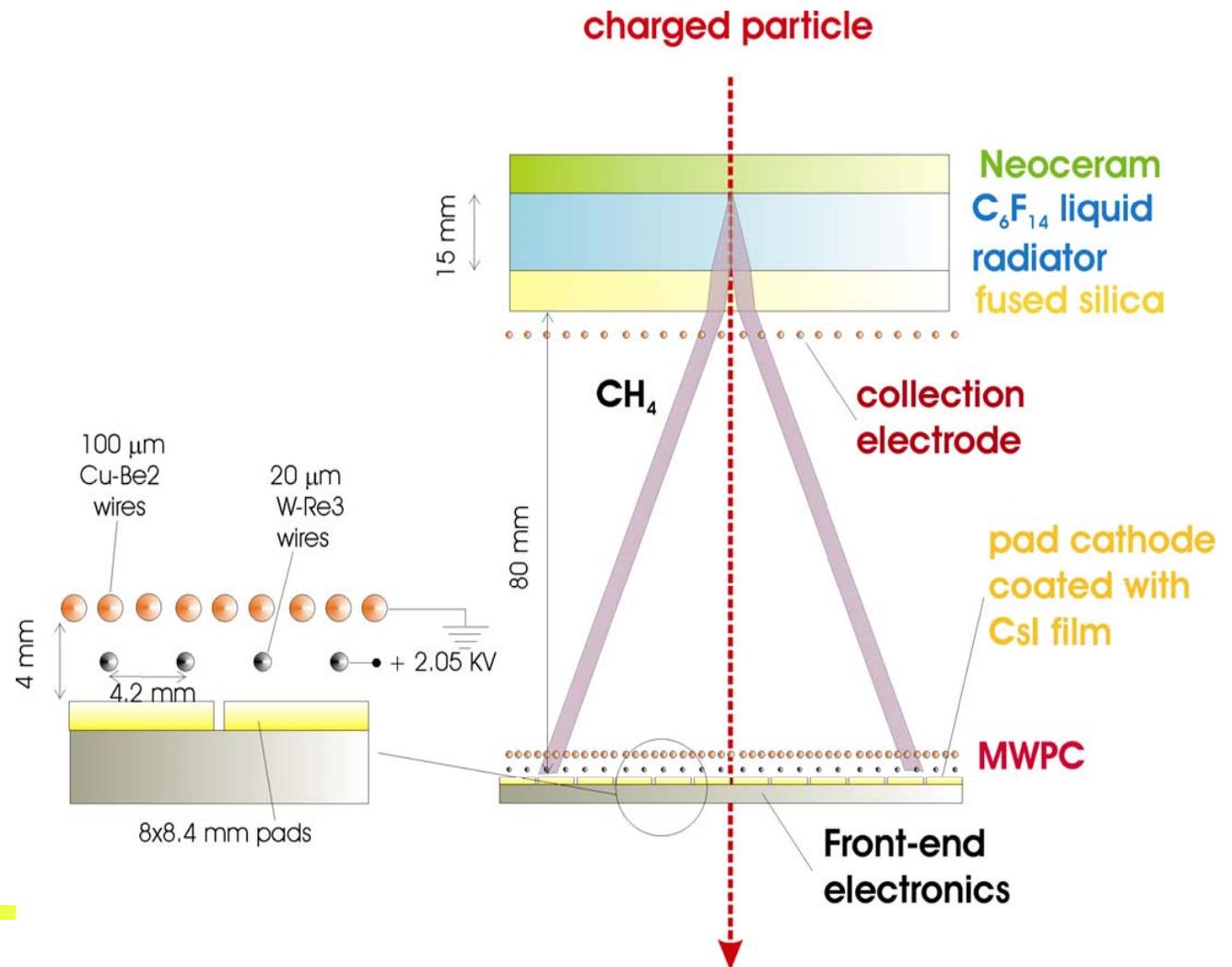
Fine-mesh PMT: works in high B fields (1.5 T)

# CsI based RICH counters: HADES, COMPASS, ALICE

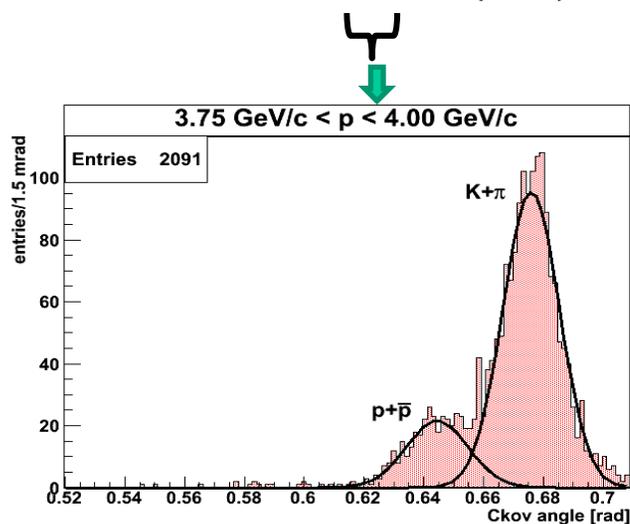
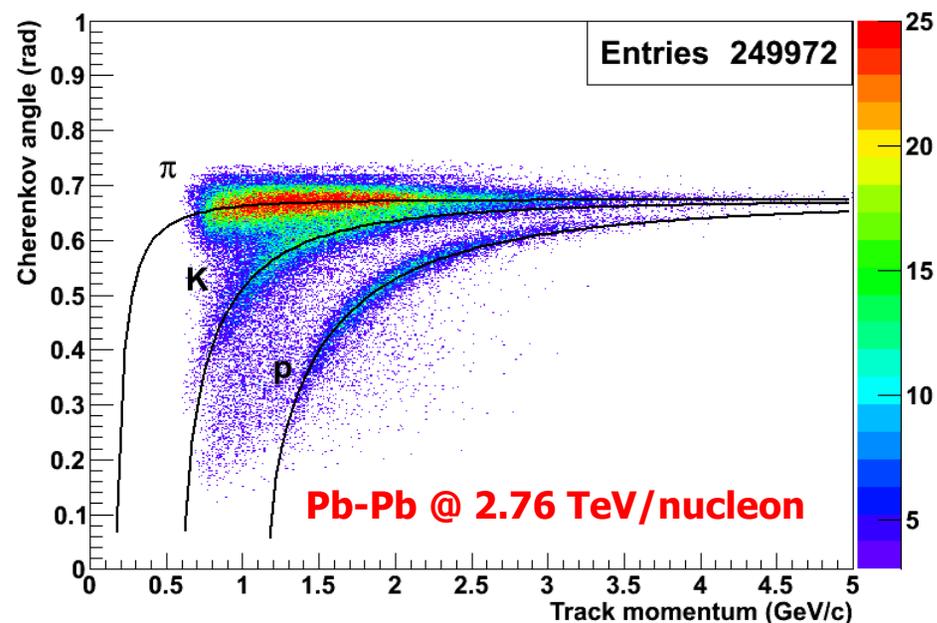
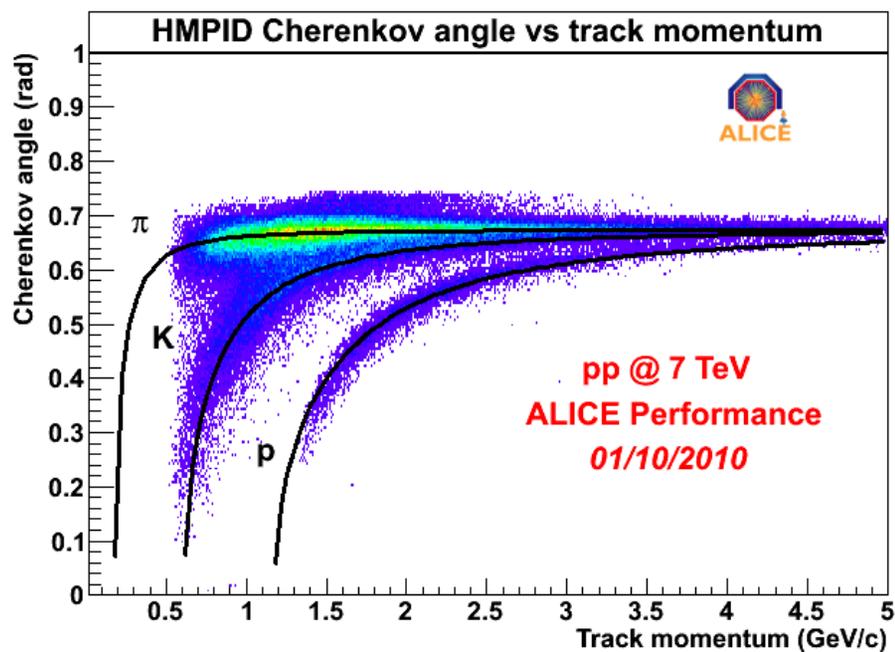
HADES and COMPASS RICH: gas radiator + CsI photocathode – long term experience in operation

ALICE:

- liquid radiator
- proximity focusing



# ALICE HMPID performance



Feb. 20, 2020

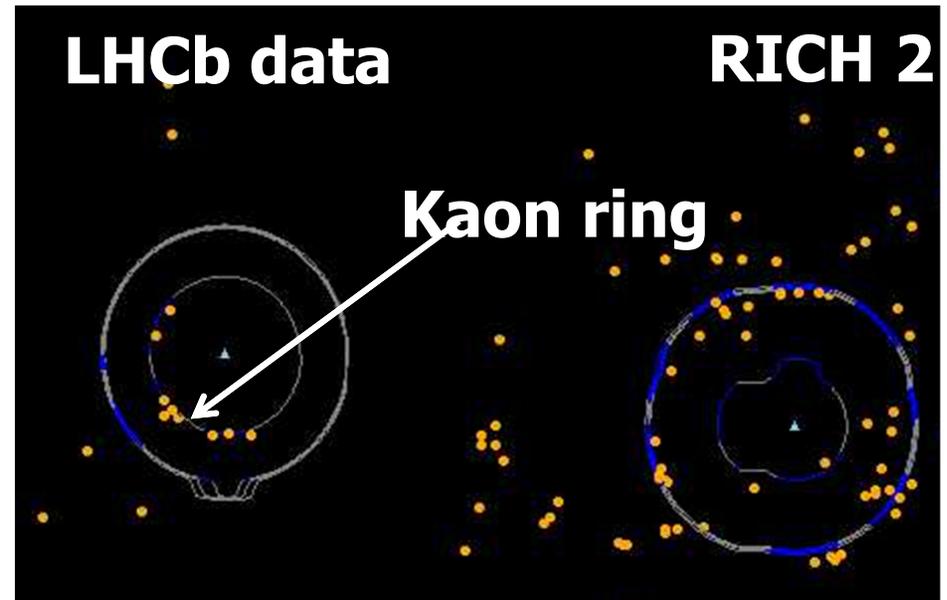
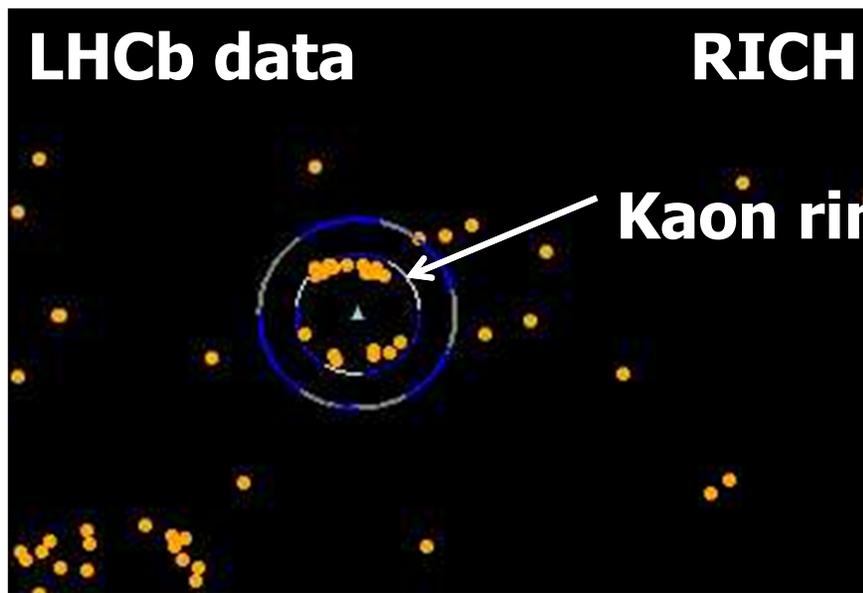
Peter Križan, Ljubljana

# LHCb Event Display

RICH1

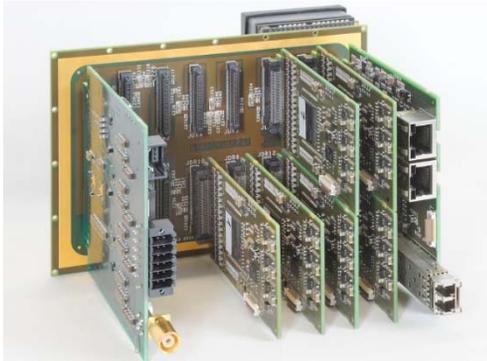
Early data, Nov/Dec 2009  
LHC beams  $\sqrt{s} = 900$  GeV

RICH2

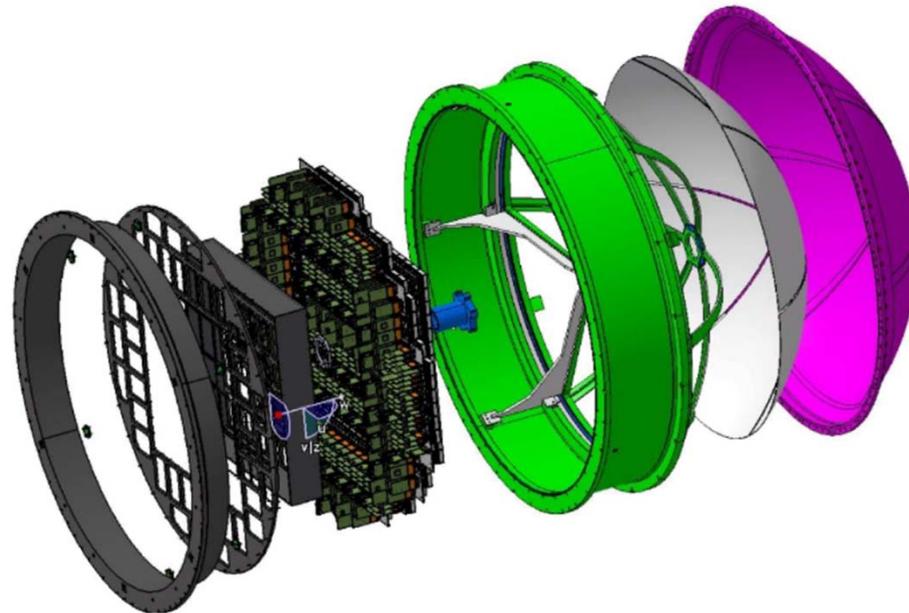
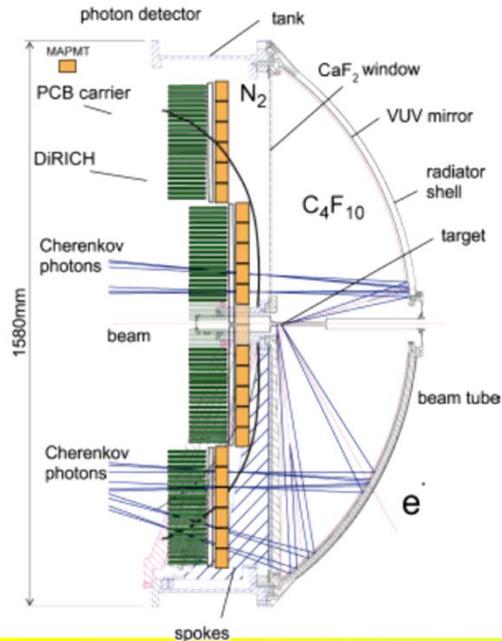


- Orange points → photon hits
- Continuous lines → expected distribution for each particle hypothesis

# HADES Upgrade

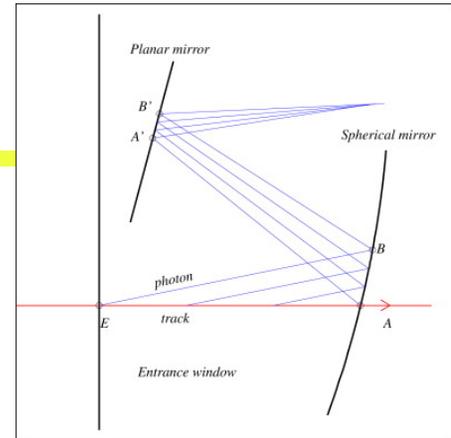


- Replace CsI-MWPCs with MaPMTs
  - Hamamatsu H12700
  - Same as for CBM-RICH
  - Also share electronics

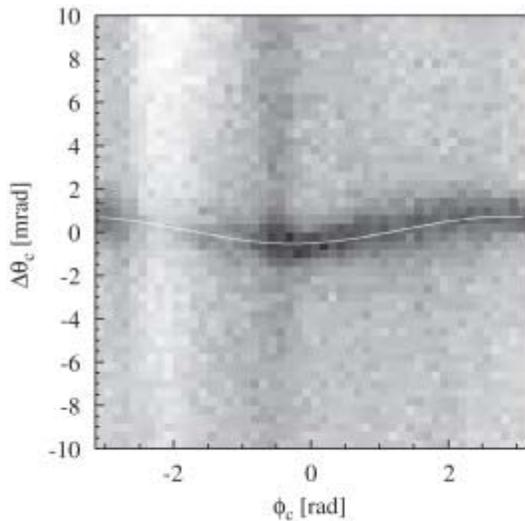
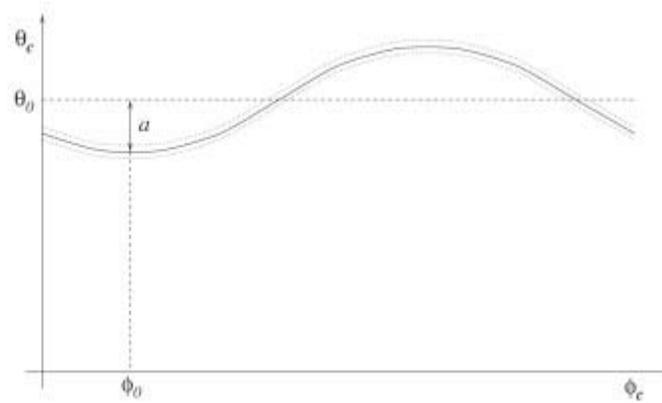
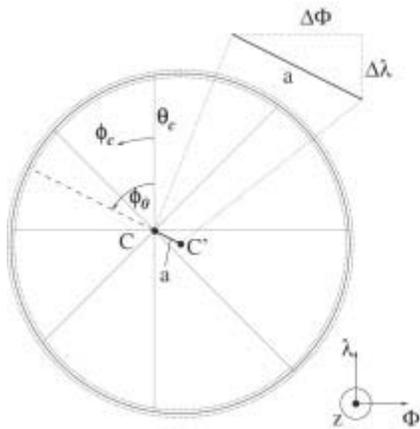


# Mirror alignment

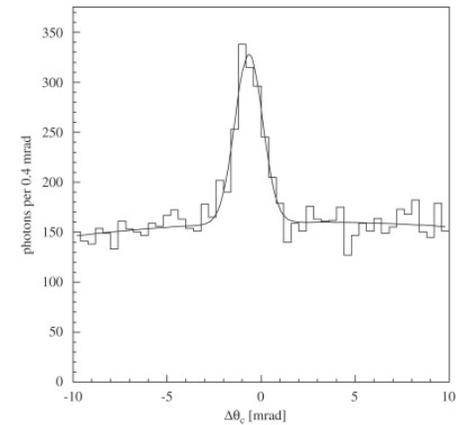
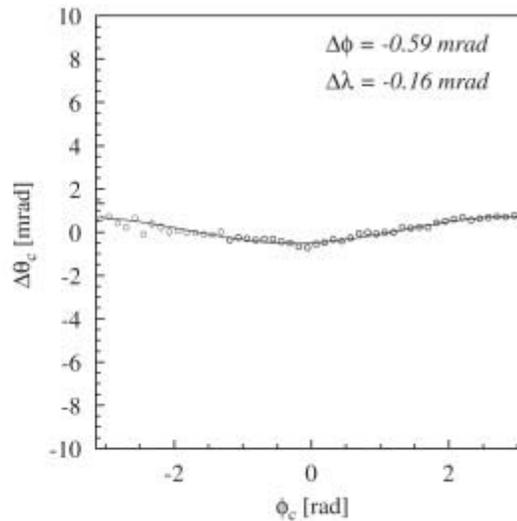
Misalignment: Cherenkov angle depends on the azimuthal angle around the track



Use unambiguous photons.



mirrors 34 14



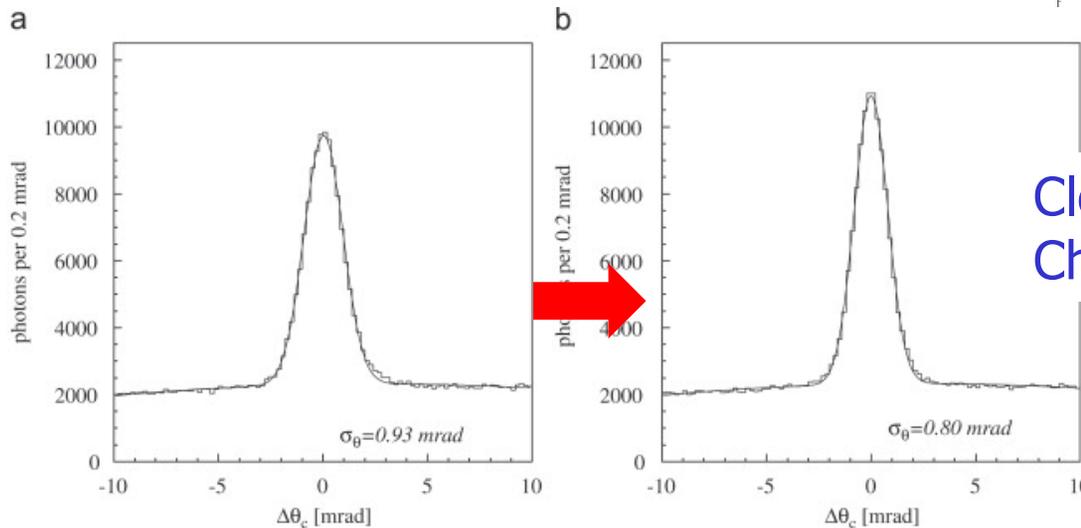
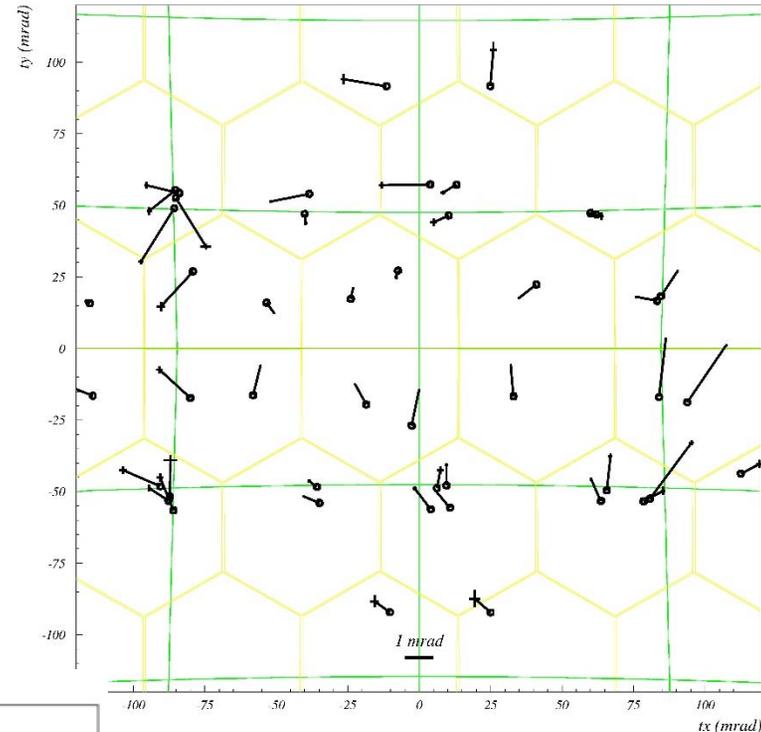
Slice in  $\phi_c$

# Mirror alignment

Initial mirror system alignment:  
with optical methods, theodolite.

Alignment with data: tells you the  
ultimate truth...

Combine all alignment data for all  
(possible) pairs of segments →  
solve a system of linear equations

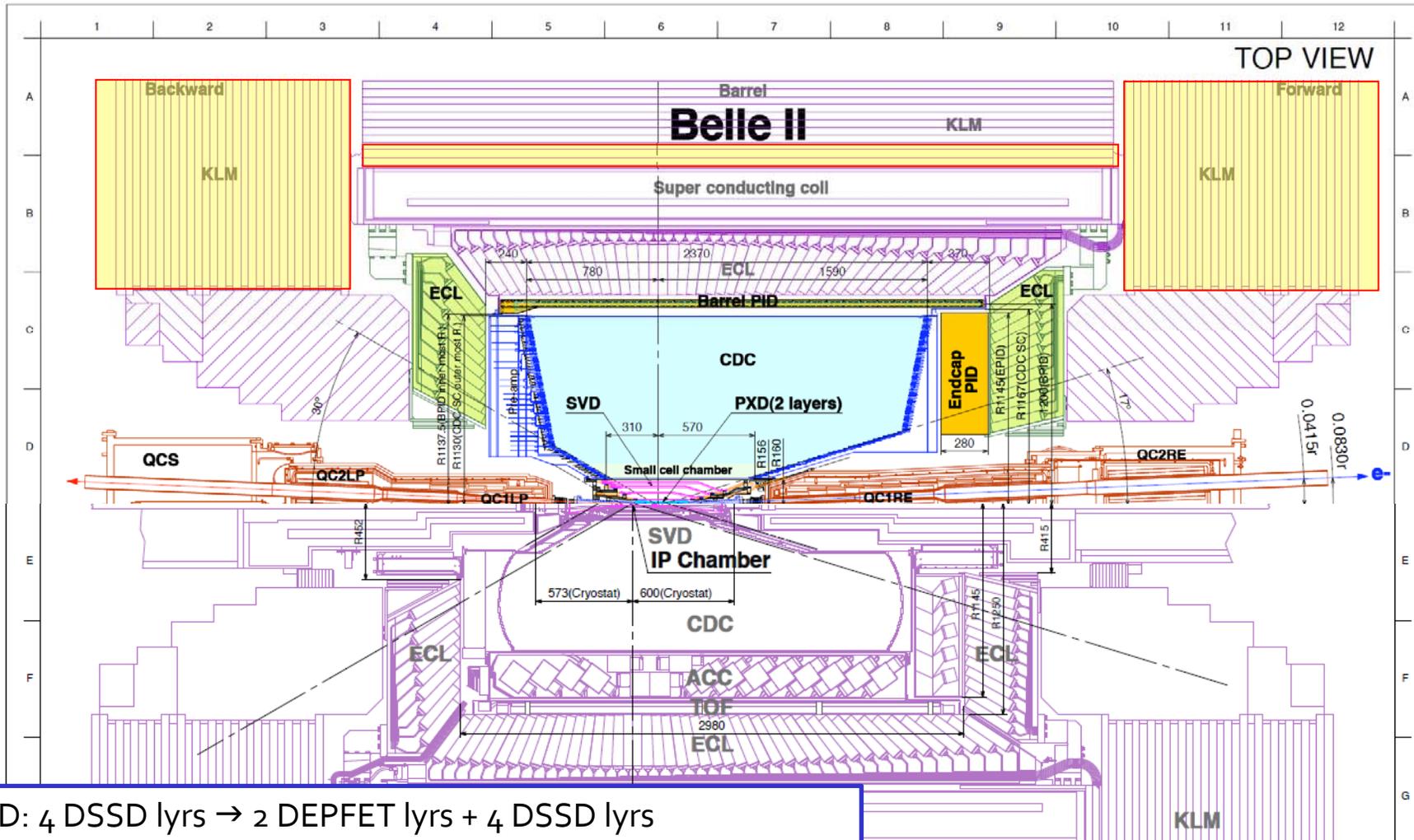


Clear improvement in  
Cherenkov angle resolution

→ NIMA 586 (2008) 174



# Belle II Detector (compared to Belle)

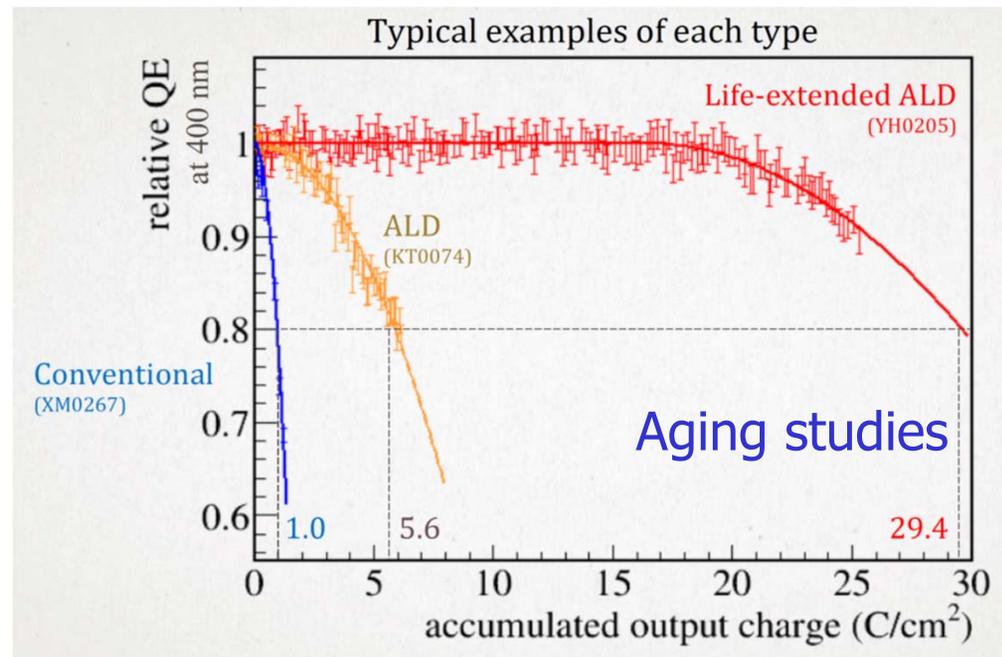
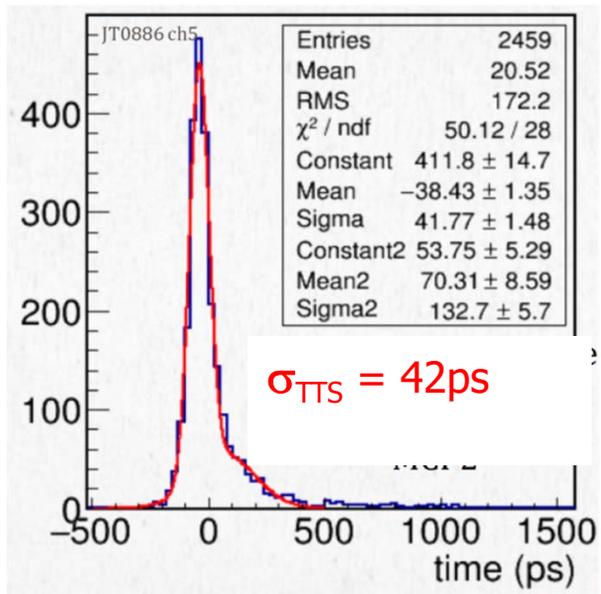


SVD: 4 DSSD lyrs → 2 DEPFET lyrs + 4 DSSD lyrs  
 CDC: small cell, long lever arm  
 ACC+TOF → TOP+A-RICH  
 ECL: waveform sampling (+pure CsI for part of endcaps)  
 KLM: RPC → Scintillator +MPPC (endcaps, barrel inner 2 lyrs)

In colours: new components

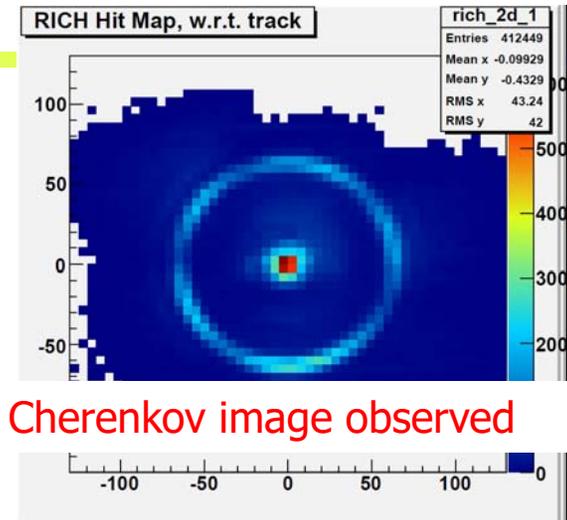
# TOP R+D areas

- Very fast photosensors for operation in 1.5 T field (MCP PMTs)
- R+D to mitigate aging of photocathodes in MCP PMTs (ALD)

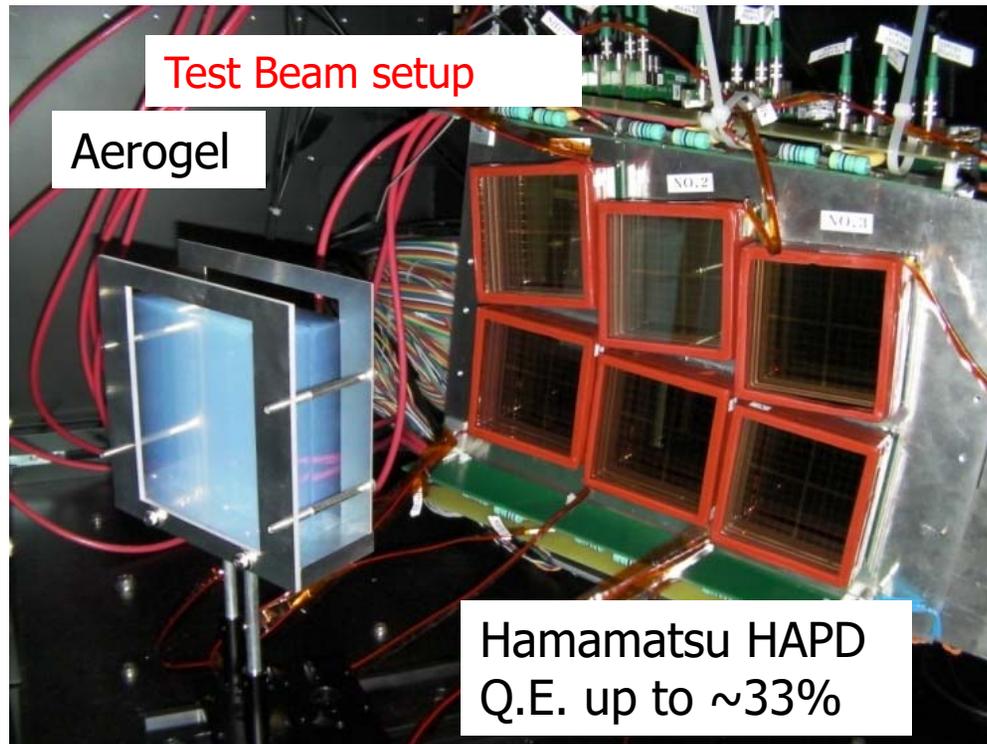


- Very fast and compact readout electronics with waveform sampling for a precise time measurement →
- Production of large quartz pieces, construction of modules, mechanics and installation methods
- Analytic expressions for the very complex 2D likelihood functions.

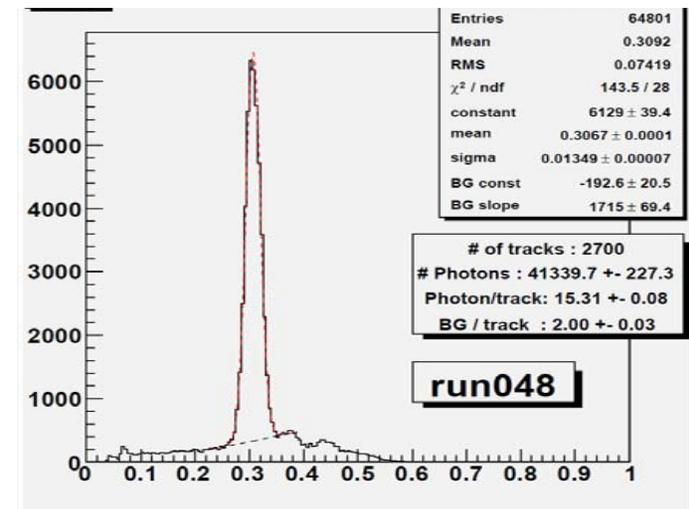
# HAPD as the Aerogel RICH photon detector



Clear Cherenkov image observed



Cherenkov angle distribution

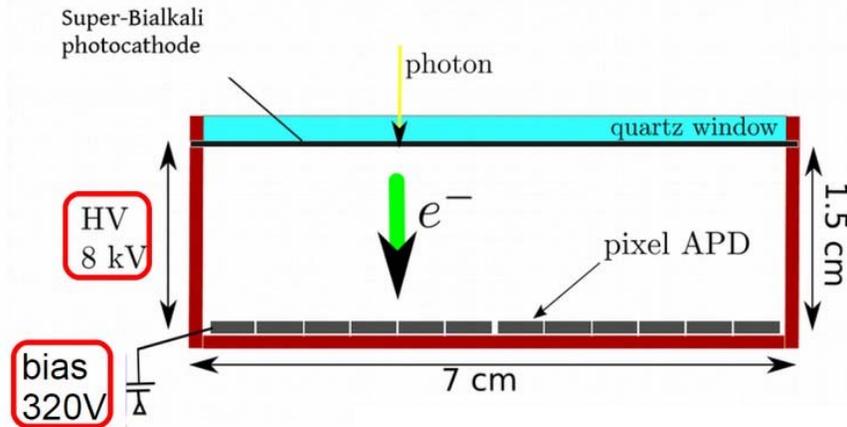


**6.6  $\sigma$  p/K at 4GeV/c !**

→ NIM A595 (2008) 180

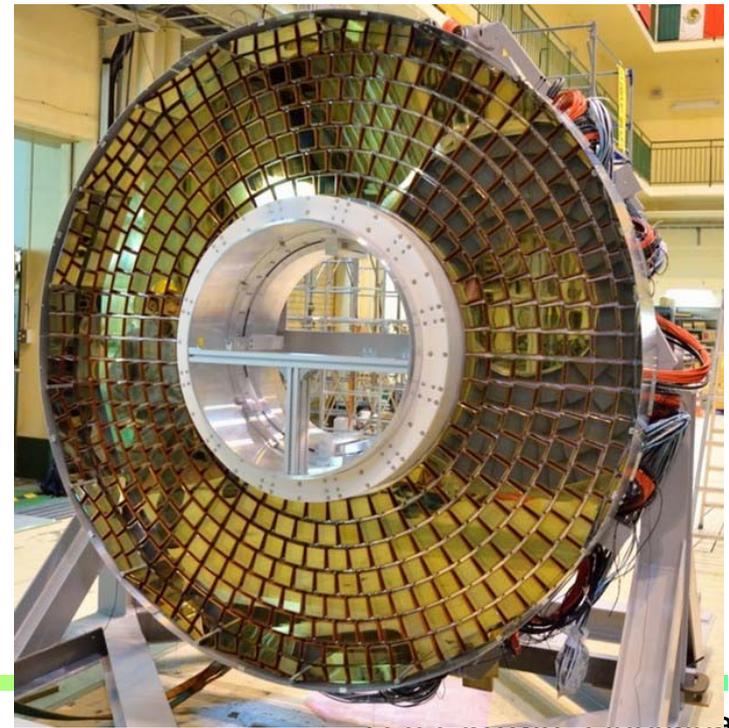
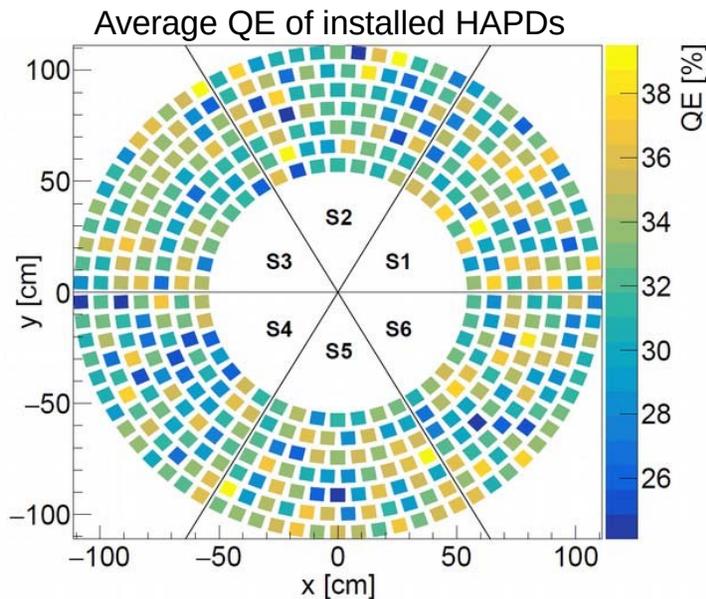
# ARICH photo sensor: HAPD

- **HAPD – Hybrid Avalanche Photo-Detector**



Size	73x73 mm
# of channels	144 (36-ch APDx4)
Total gain	>60000 (1500 x 40)
Peak QE	~30%
Active area	64%
Weight	220g

- 420 modules to cover the detector plane



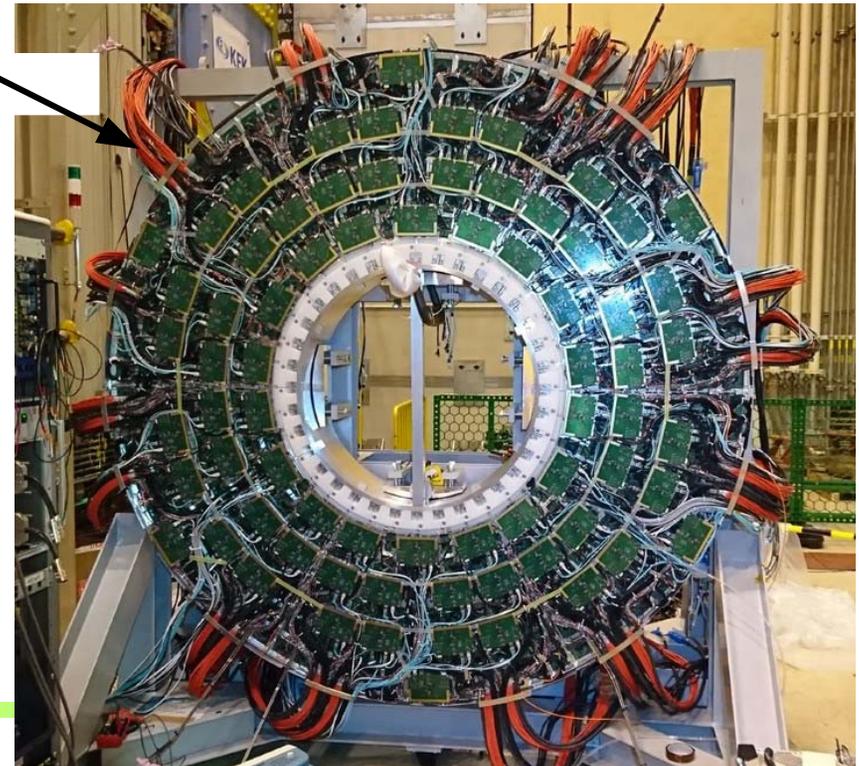
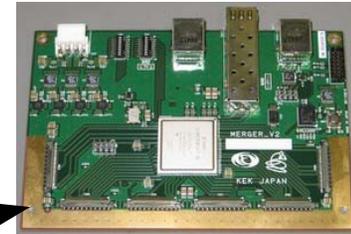
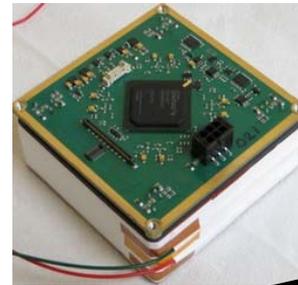
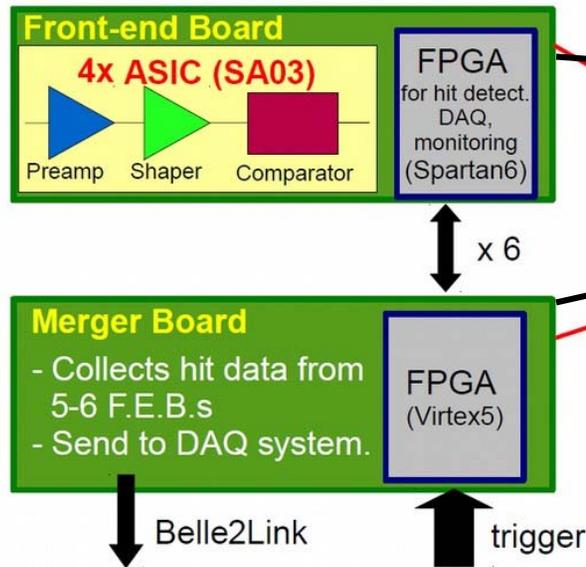
Feb. 20

SY

TECHNICAL, Ljubljana

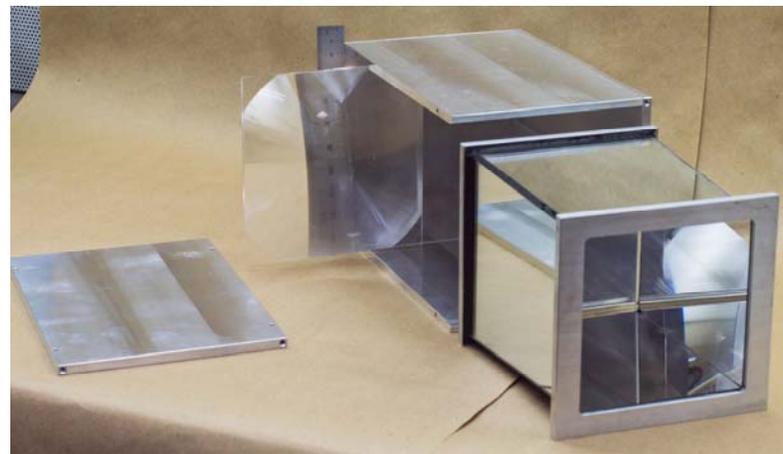
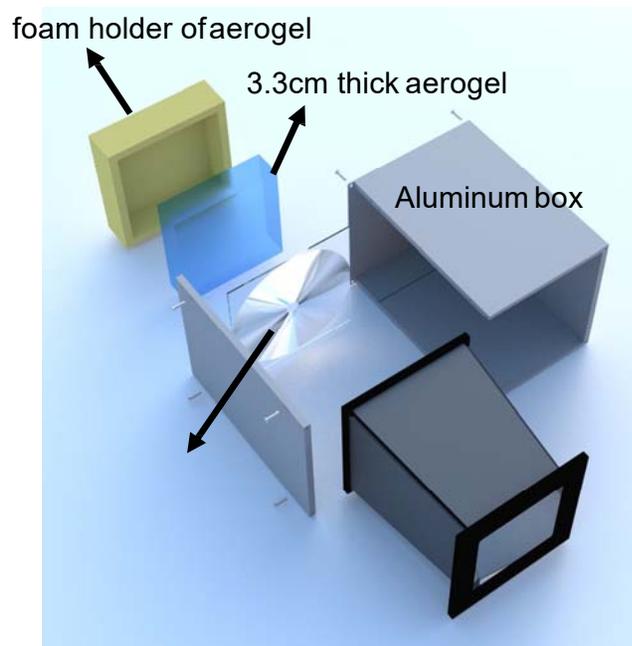
# ARICH read-out electronics

- In total ~60k channels
- Limited space of 5cm behind HAPDs



- Variable gain (3.1-12.5 V/pC),  
shaping time (100-200 ns)  
→ optimization for  
increased noise levels

## 2nd mRICH Prototype - Verify the PID Capability



6" focal length Fresnel lens  
mirror set



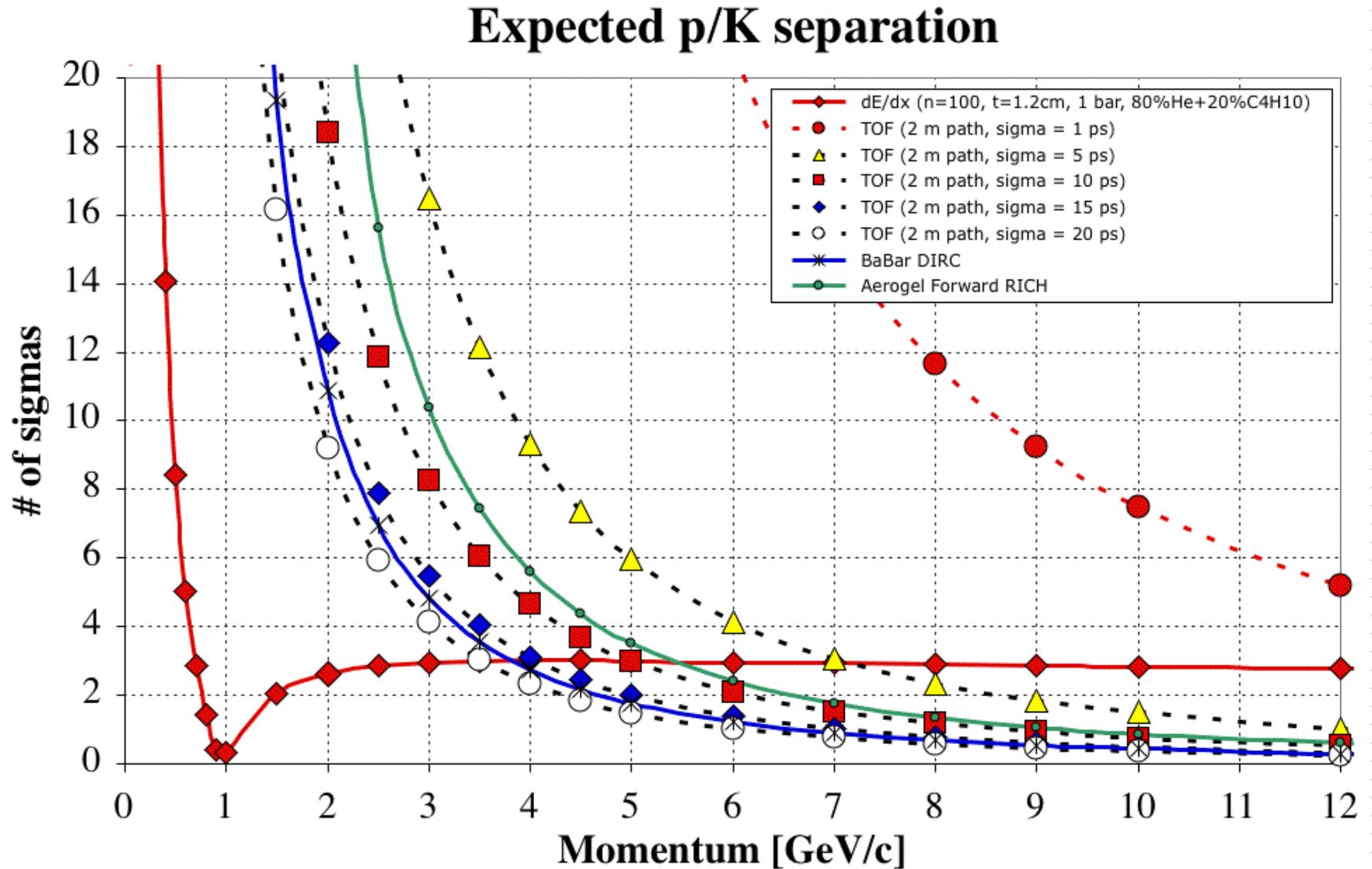
### FEATURES

- High quantum efficiency: 33 % typ.
- High collection efficiency: 80 % typ.
- Single photon peaks detectable at every anode (pixel)
- Wide effective area: 48.5 mm × 48.5 mm
- 16 × 16 multianode, pixel size: 3 mm × 3 mm / anode



1. Longer focal length (Fresnel lens)
2. Smaller pixel size sensors

# Time-of-flight with fast photon detectors

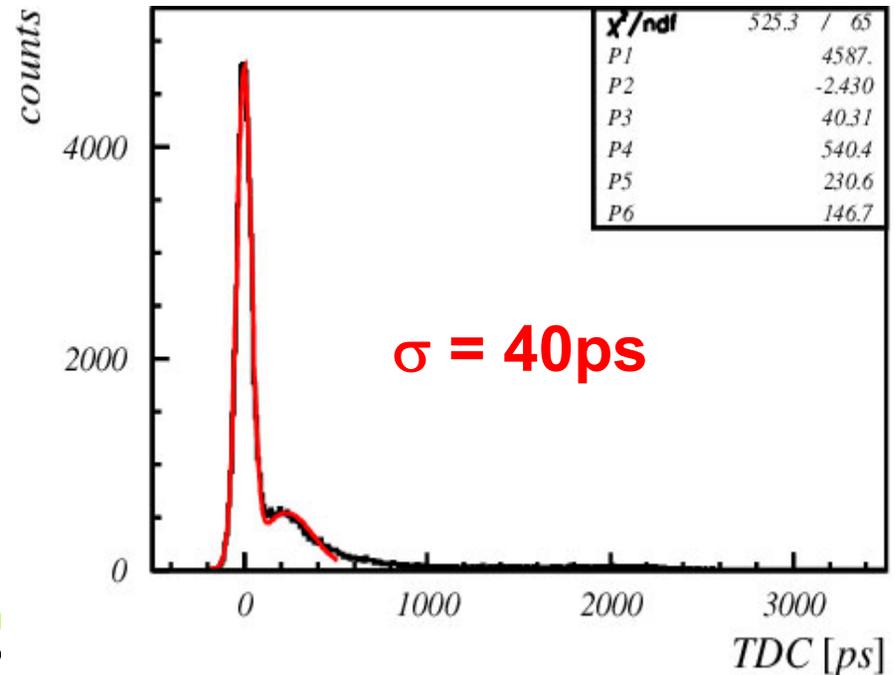
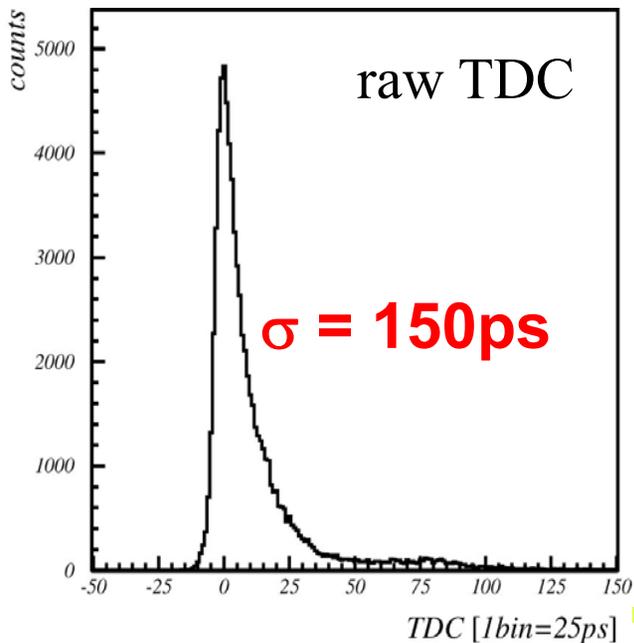
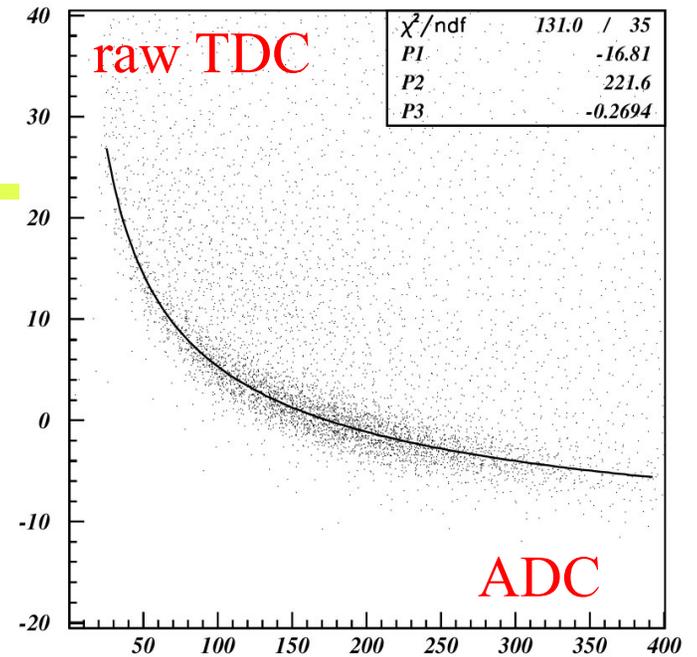


# Time walk correction

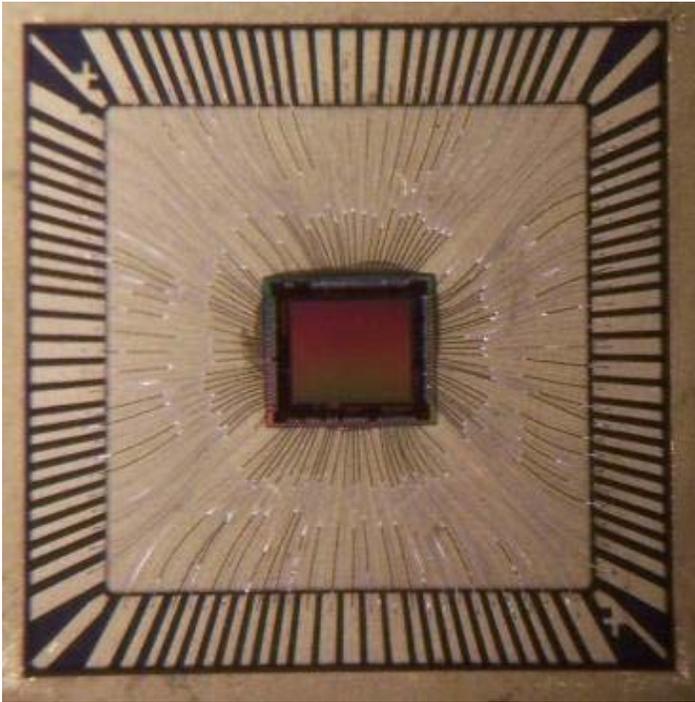
TDC vs. ADC correlation is fitted with

$$TDC = P1 + \sqrt{\frac{P2}{ADC - P3}}$$

and used for TDC correction



# Time walk correction 3: waveform sampling



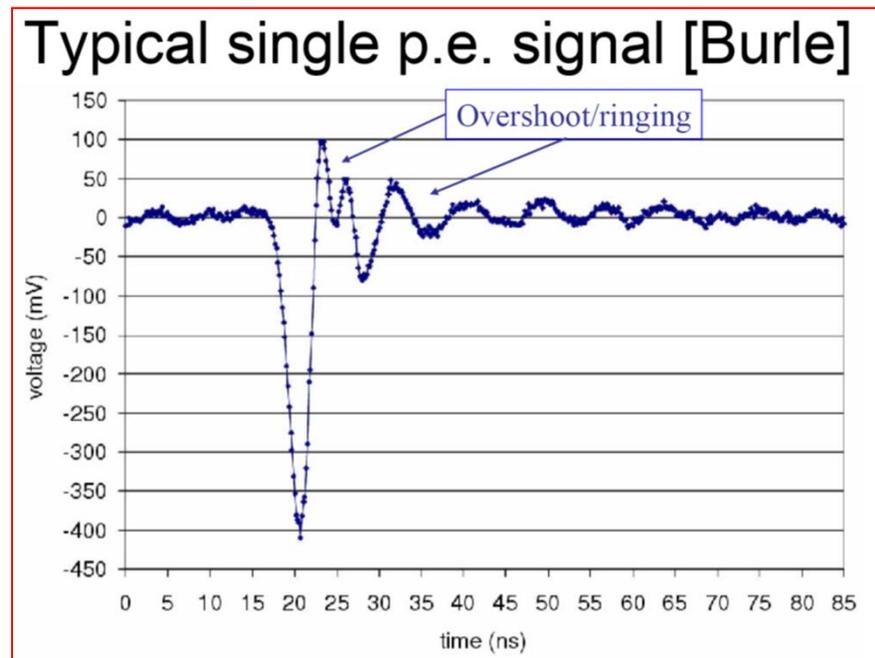
3mm x 2.8mm, TSMC 0.25um

- 64k samples deep
- Multi-MSa/s to Multi-GSa/s

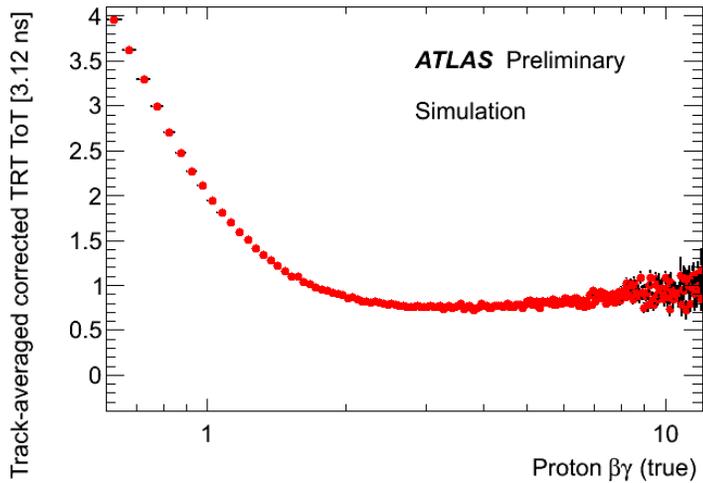
Gary Varner (Hawaii)

Variant of the LABRADOR 3

Successfully flew on ANITA in Dec 06/Jan 07 ( $\leq 50$ ps timing)

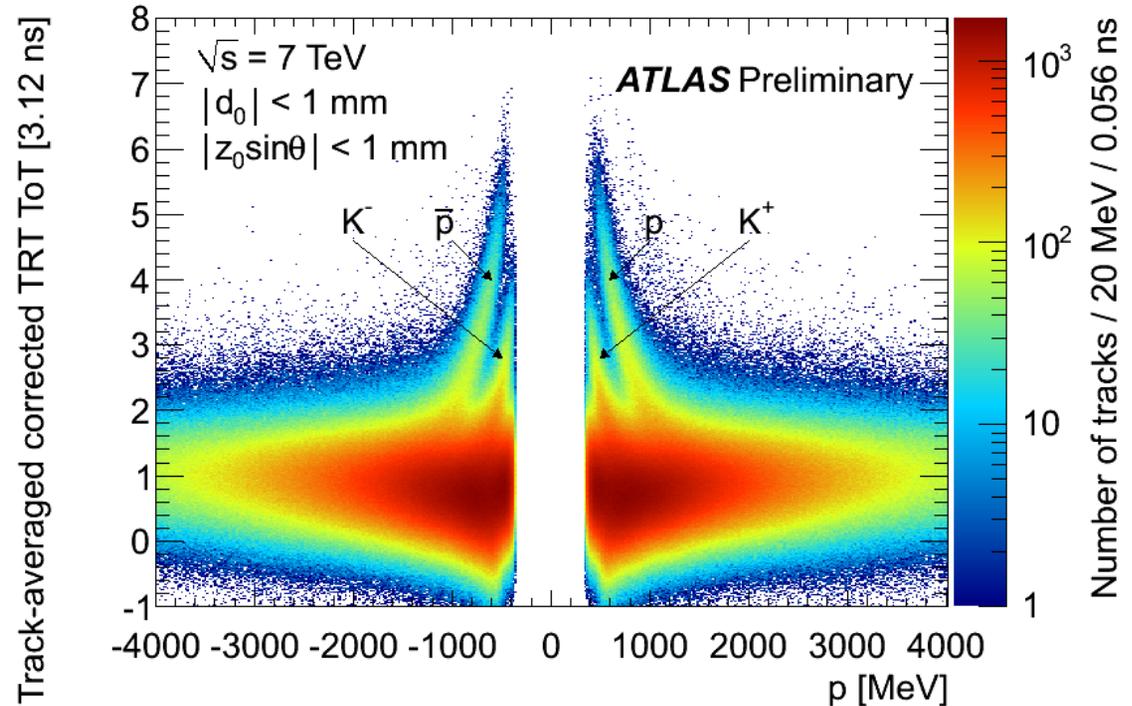


# Time-over-Threshold (ToT): dE/dx in ATLAS TRT



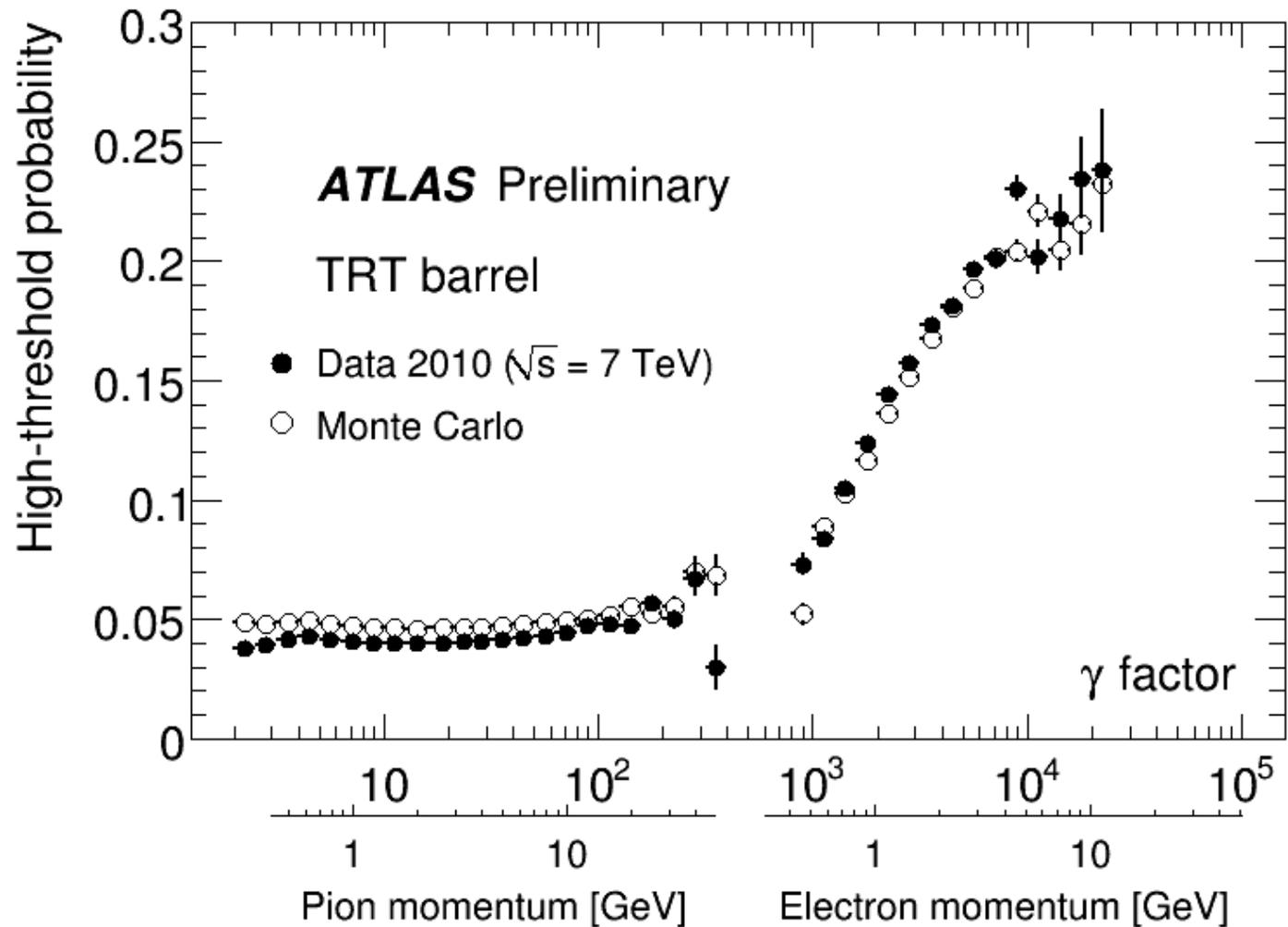
Track-averaged ToT distribution as a function of the track momentum.

The relation between the track ToT measurement and the track  $\beta\gamma$ , obtained from MC studies.

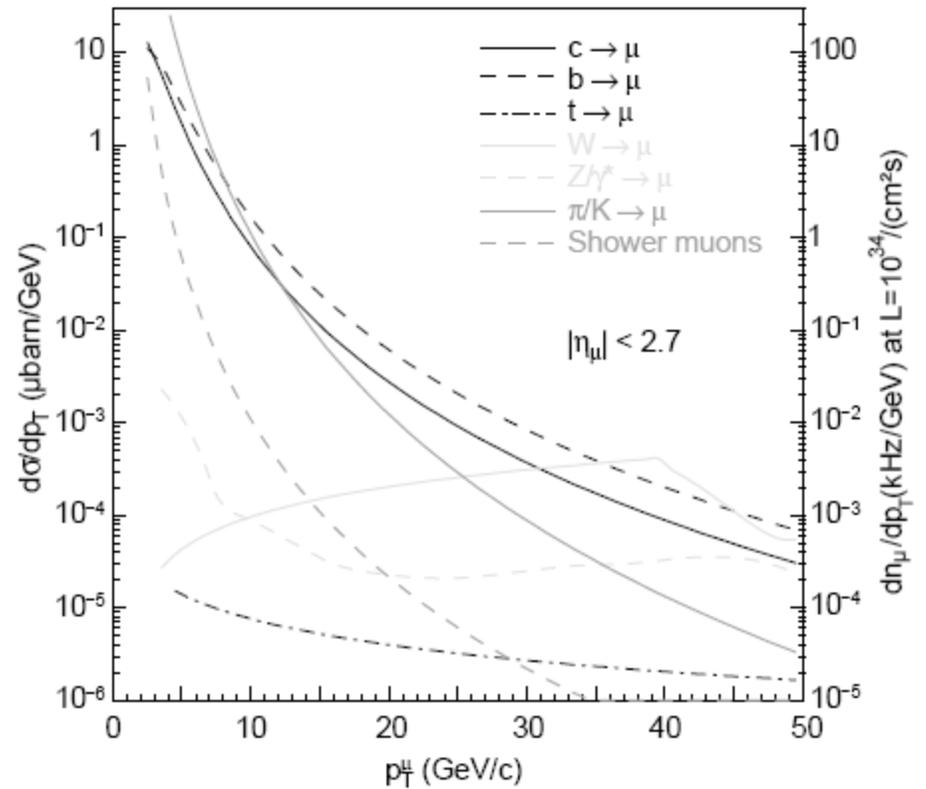
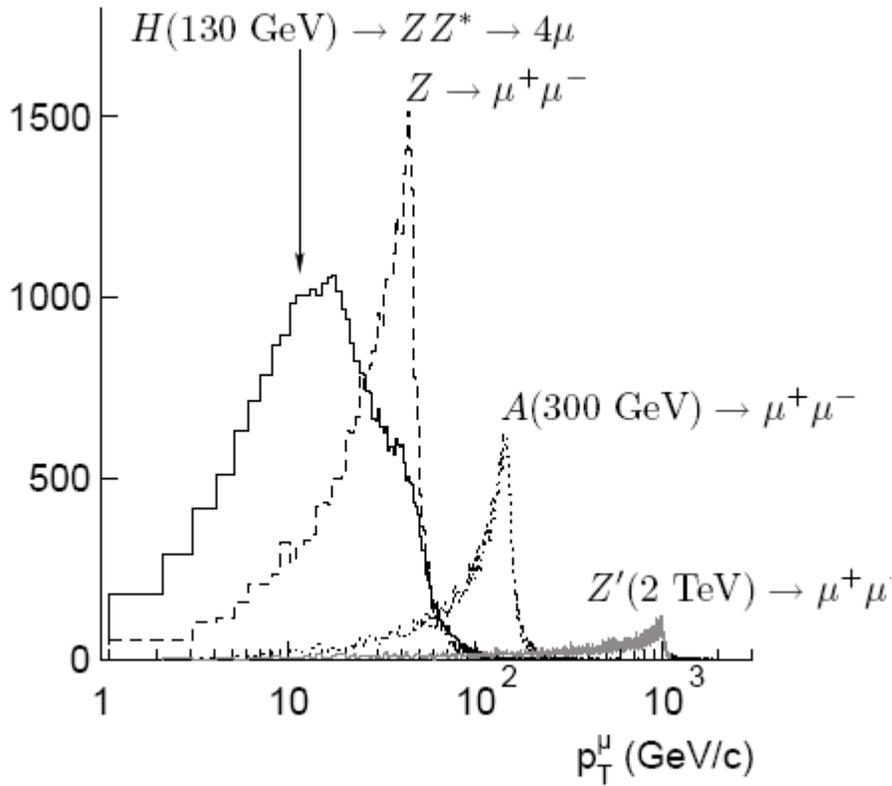


# TRT performance in 2010 data

e/pion separation: high threshold hit probability per straw



# Muon spectrum



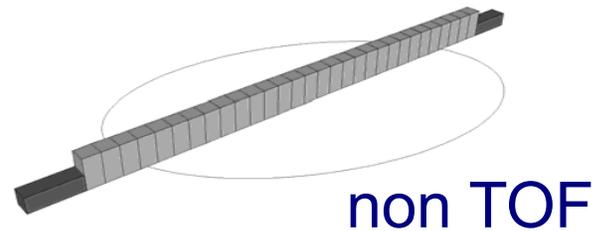
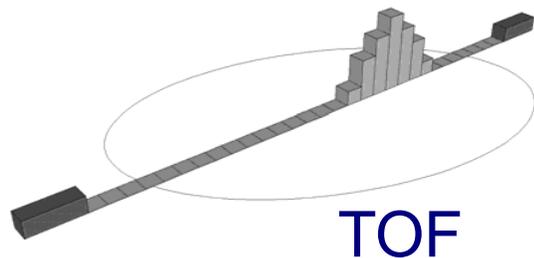
# Applications in medical imaging: advances in TOF-PET

Time-of-Flight difference of annihilation gammas is used to improve the contrast of images obtained with PET

Localization of source position along the line of response:

$$\Delta t \sim 66\text{ps} \rightarrow \Delta x = c_0 \Delta t / 2 \sim 1\text{cm}$$

$\Delta t$  = coincidence resolving time, CRT



However, PET systems based on SiPM readout are reaching CRT of  $\sim 300$  ps, and only with small crystals  $\sim 3 \times 3 \times 3$  mm<sup>3</sup> CRT < 100 ps

Novel photon detectors – MCP-PMTs and SiPMs – have excellent timing resolution  
→ TOF resolution limited by the spread in photon emission and arrival time

Faster annihilation gamma detection method → a faster light emission mechanism

# Annihilation gamma detection with Cherenkov light

Cherenkov light is promptly produced by a charge particle traveling through the medium with velocity higher than the speed of light  $c_0/n$ . Photoelectron emits Cherenkov light in  $\sim 1\text{ps}$ .

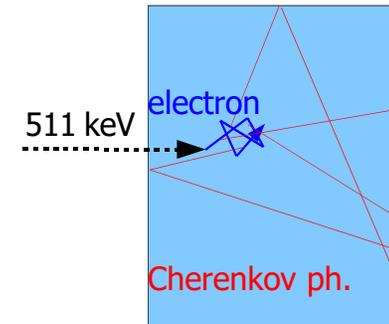
Disadvantage of Cherenkov light is the small number of Cherenkov photons produced per interaction

$$N \approx \frac{370}{eV\text{ cm}} l \Delta E \sin^2 \vartheta \approx 370 \times 0.01 \times 2 \times 0.75 \approx 8$$

→ detection at a single photon level!

Cherenkov radiator:  $\text{PbF}_2$  an excellent candidate

- high gamma stopping power
- high fraction of gamma interactions via photoeffect → electrons with maximal kinetic energy → more Cherenkov photons



	$\rho$ (g/cm <sup>3</sup> )	$n$	e <sup>-</sup> Cherenkov threshold (keV)	Cutoff wavelength (nm)	Attenuation length (cm)	Photofraction
<b>PbF<sub>2</sub></b>	<b>7.77</b>	<b>1.82</b>	<b>101</b>	<b>250</b>	<b>0.91</b>	<b>46%</b>
LYSO	7.4				1.14	32%
LaBr <sub>3</sub>	5.1				2.23	15%

+ high transmission in visible and near UV

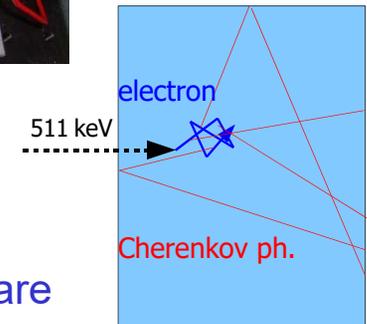
# Excellent TOF PET timing with MCP PMTs

Pioneering experiment, two detectors in a back-to-back configuration:  
PbF<sub>2</sub> 25x25x15 mm<sup>3</sup> with  
MCP-PMT as photodetectors

- single photon timing ~ 50 ps FWHM
- active surface 22.5x22.5 mm<sup>2</sup>



black painted, Teflon wrapped, bare



Timing resolution (black painted):

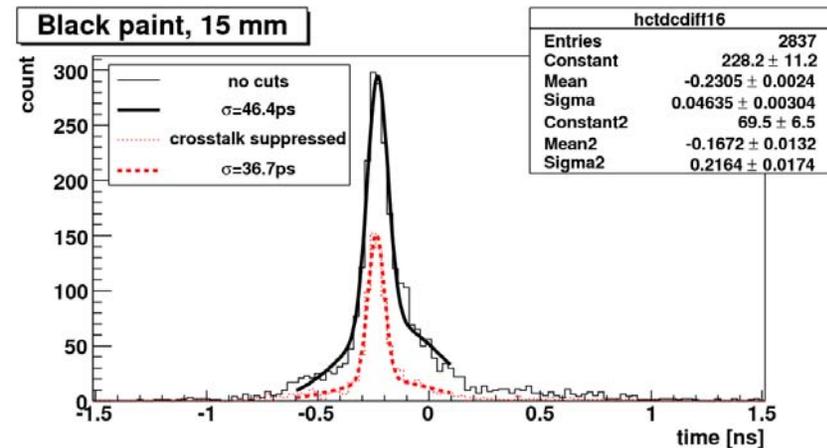
~ 70 ps FWHM, 5mm crystal

~100 ps FWHM 15mm crystal

Efficiency (Teflon wrapped):

~ 6%, single side

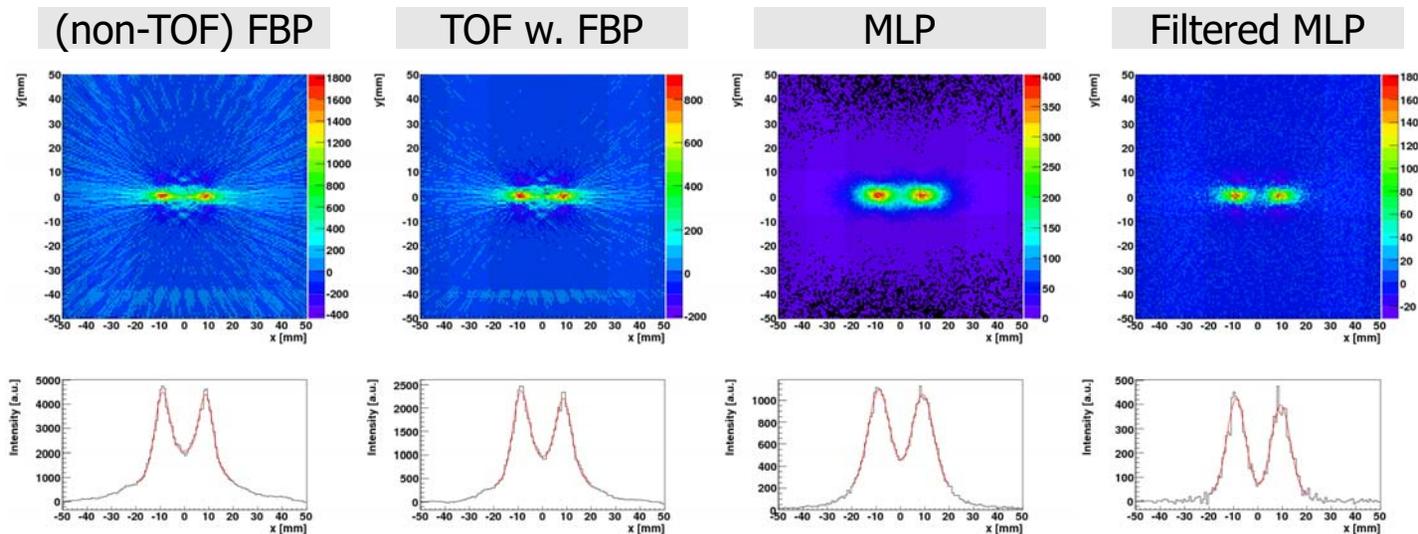
(typically ~ 30% for LSO)



**NIM A654(2011)532**

# Reconstruction - experiment

Two  $^{22}\text{Na}$  point sources at +10 mm and -10 mm 4x4 segmented, black painted  $\text{PbF}_2$  radiators



→ A simple, very fast Most-likely-point (MLP) method ( $\sim$ histogramming of points) already gives a reasonable image

→ NIM A732 (2013) 595

# Cherenkov based PET scanner?

PbF<sub>2</sub> not a scintillator → considerably **cheaper!**

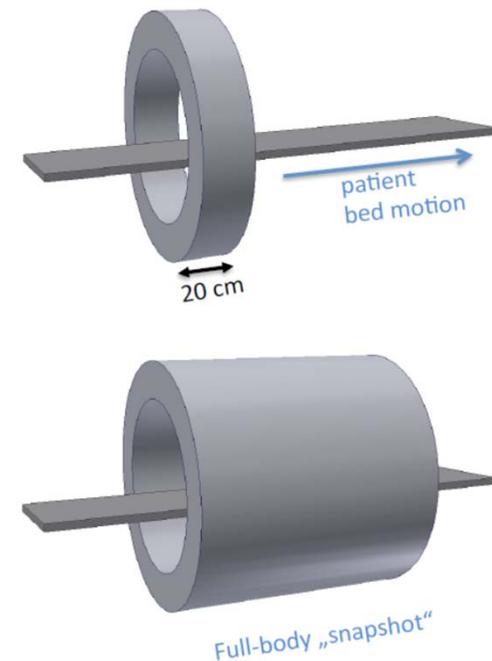
Smaller attenuation length than LYSO – **small parallax error**

→ **Cheaper normal** scanner or

→ **Total/half body** device

Extending axial FOV 20 cm → 200 cm:  
estimated 6-fold increase in SNR →

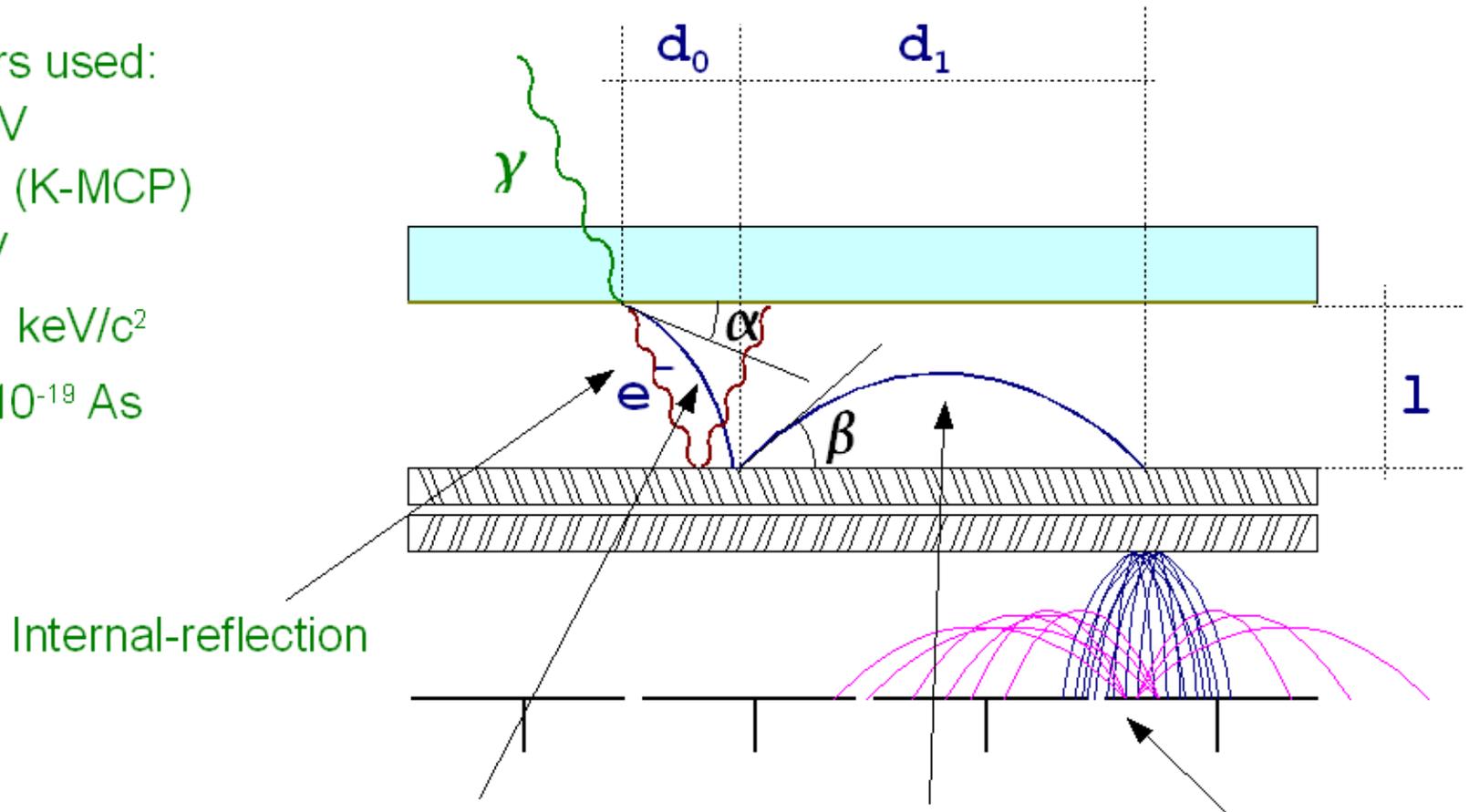
- Better image quality
- OR Shorter scanning time
- OR Less injected activity: 8 mSv → 0.2 mSv



# MCP PMT: processes involved in photon detection

Parameters used:

- $U = 200 \text{ V}$
- $l = 6 \text{ mm (K-MCP)}$
- $E_0 = 1 \text{ eV}$
- $m_e = 511 \text{ keV}/c^2$
- $e_0 = 1.6 \cdot 10^{-19} \text{ As}$



Internal-reflection

Photo-electron:

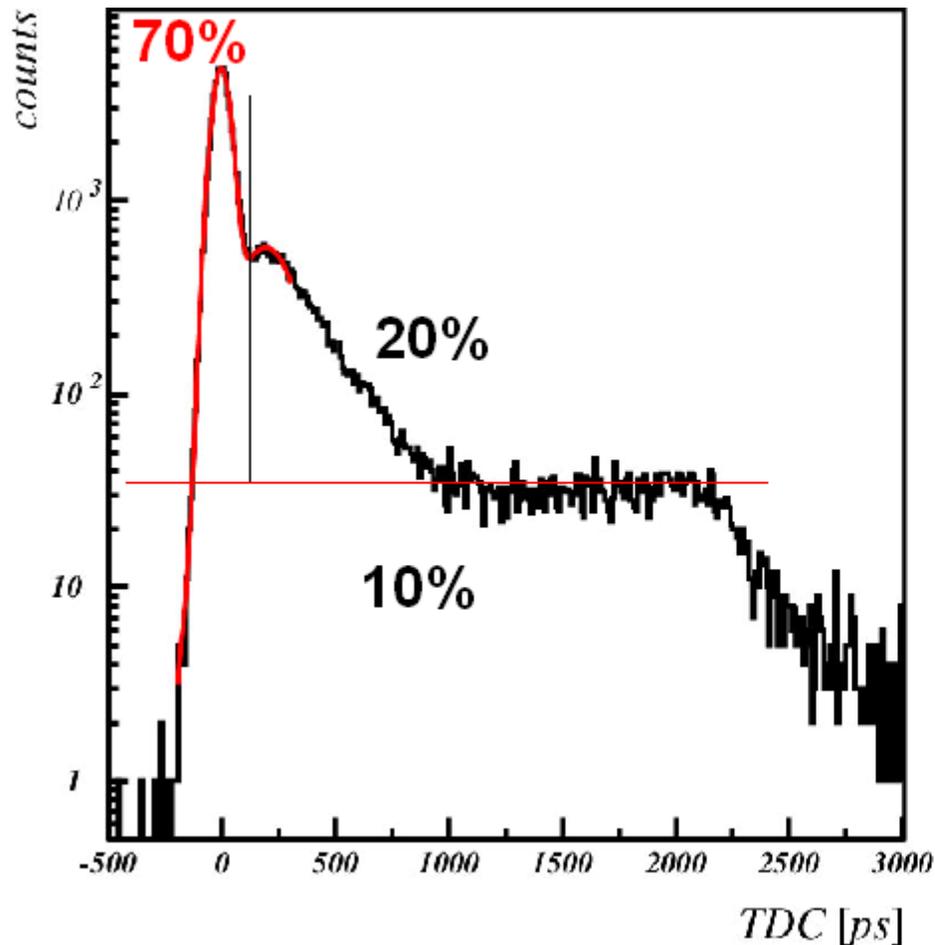
- $d_{0,\max} \sim 0.8 \text{ mm}$
- $t_0 \sim 1.4 \text{ ns}$
- $\Delta t_0 \sim 100 \text{ ps}$

Backscattering:

- $d_{1,\max} \sim 12 \text{ mm}$
- $t_{1,\max} \sim 2.8 \text{ ns}$

Charge sharing

# MCP PMT timing

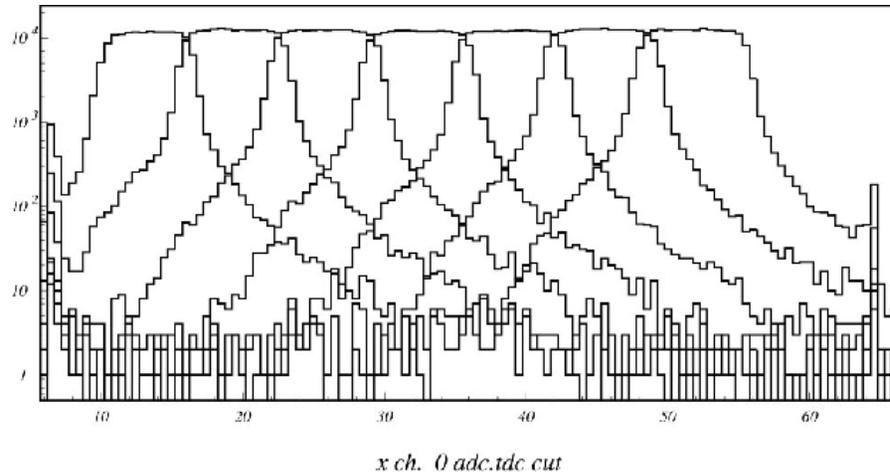
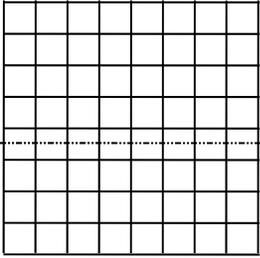


Tails understood (scattering of photoelectrons off the MCP), can be significantly reduced by:

- decreased photocathode-MCP distance and
- increased voltage difference

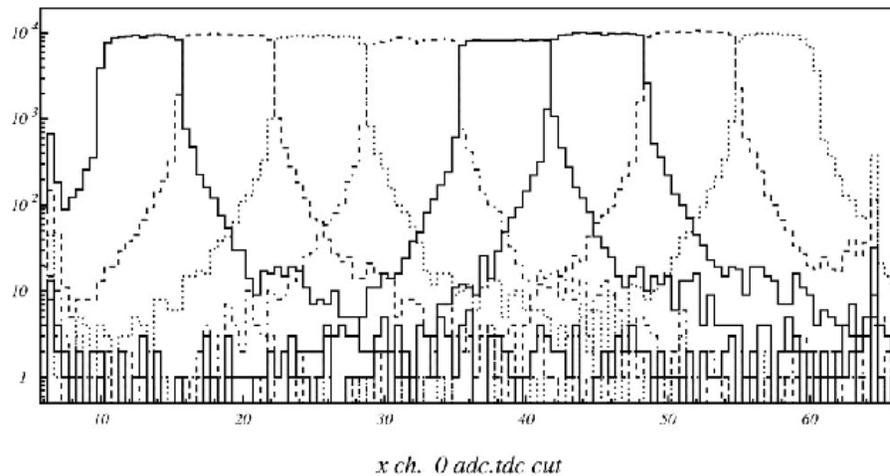
- **prompt signal ~ 70%**
- **short delay ~ 20%**
- **~ 10% uniform distribution**

# MCP PMT: sensitivity



Number of detected hits on individual channels as a function of light spot position.

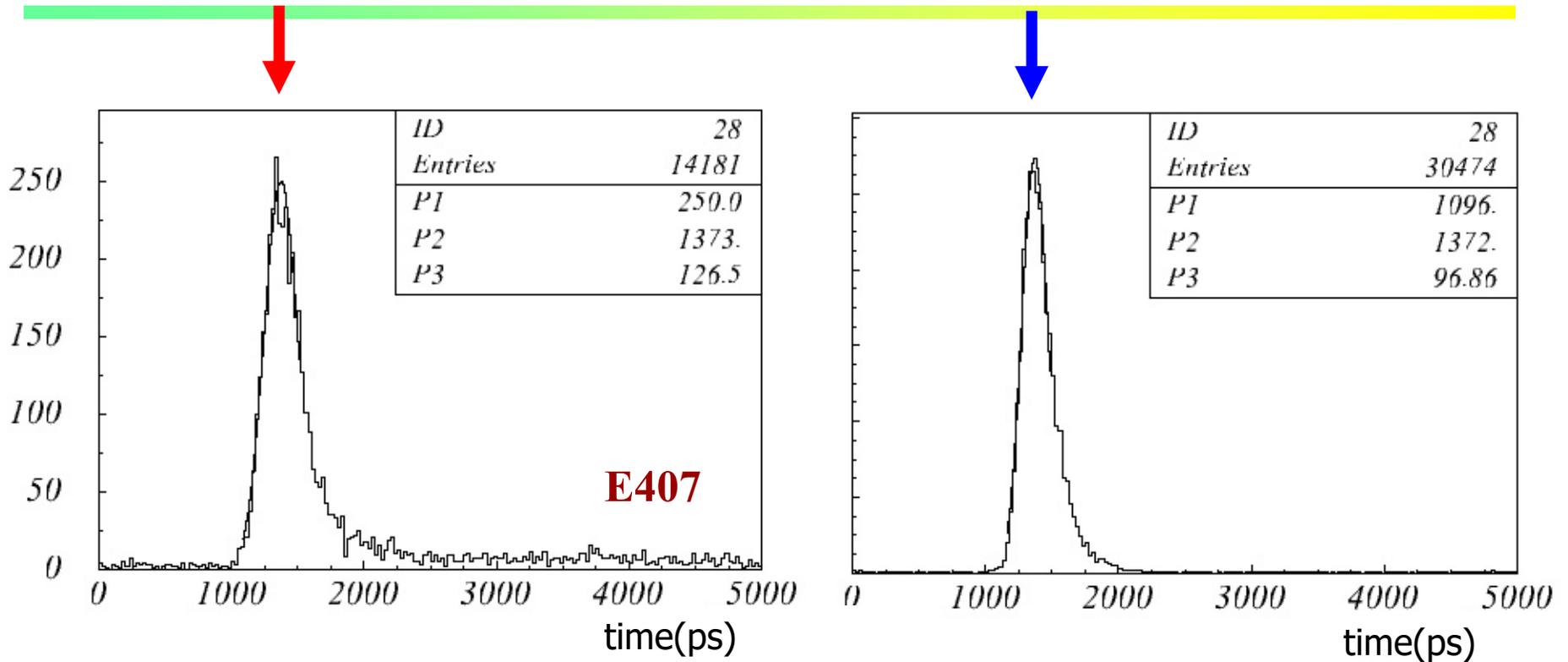
$B = 0 \text{ T}$ ,  
 $HV = 2400 \text{ V}$



$B = 1.5 \text{ T}$ ,  
 $HV = 2500 \text{ V}$

In the presence of magnetic field, charge sharing and cross talk due to long range photoelectron back-scattering are considerably reduced.

# Time resolution: blue vs red

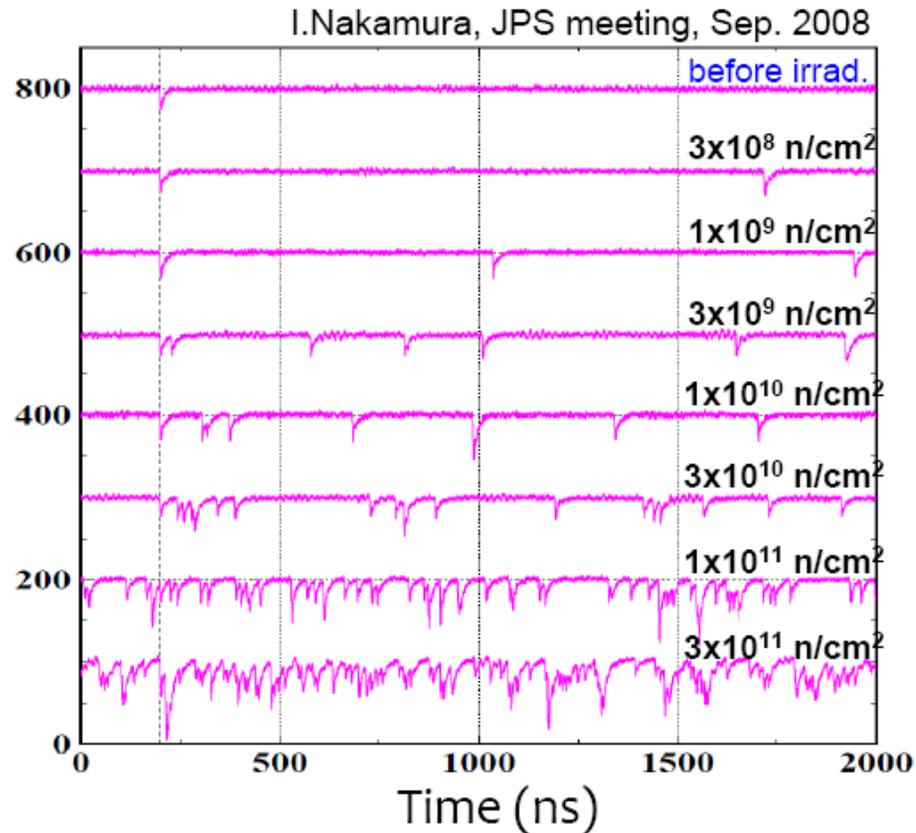


	E407	S137	H100C	H050C	H025C
$\sigma_{\text{red}}$ (ps)	<b>127</b>	<b>182</b>	<b>145</b>	<b>212</b>	<b>154</b>
$\sigma_{\text{blue}}$ (ps)	<b>97</b>	<b>151</b>	<b>136</b>	<b>358</b>	<b>135</b>

•  $\sigma \approx 100$  ps

•  $\sigma_{\text{red}} > \sigma_{\text{blue}}$

# Radiation damage



Expected fluence at 50/ab at  
Belle II:  $2-20 \cdot 10^{11} \text{ n cm}^{-2}$   
→ Worst than the lowest line

→ Very hard to use present SiPMs as single photon detectors in Belle II because of radiation damage by neutrons

→ Also: could only be used with a sophisticated electronics – wave-form sampling

# COMPASS RICH-1 upgrade

## Performance:

$\sim 60$  detected photons per ring at saturation ( $\beta = 1$ )  $\rightarrow N_0 \sim 66 \text{ cm}^{-1}$

$\sigma_\theta \sim 0.3 \text{ mrad} \rightarrow 2 \sigma \pi\text{-K}$  separation at  $\sim 60 \text{ GeV/c}$

K-ID efficiency ( $K^\pm$  from  $\Phi$  decay)  $> 90\%$

$\pi \rightarrow K$  misidentification ( $\pi^\pm$  from  $K_S$  decay)  $\sim 1\%$

